

Ab Initio Nuclear Structure Theory with Chiral NN+3N Interactions

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From QCD to Nuclear Structure

Nuclear Structure

Low-Energy QCD

From QCD to Nuclear Structure

Nuclear Structure

**NN+3N Interaction
from Chiral EFT**

Low-Energy QCD

- chiral EFT based on the relevant degrees of freedom & symmetries of QCD
- provides consistent NN, 3N,... interaction plus currents

From QCD to Nuclear Structure

Nuclear Structure

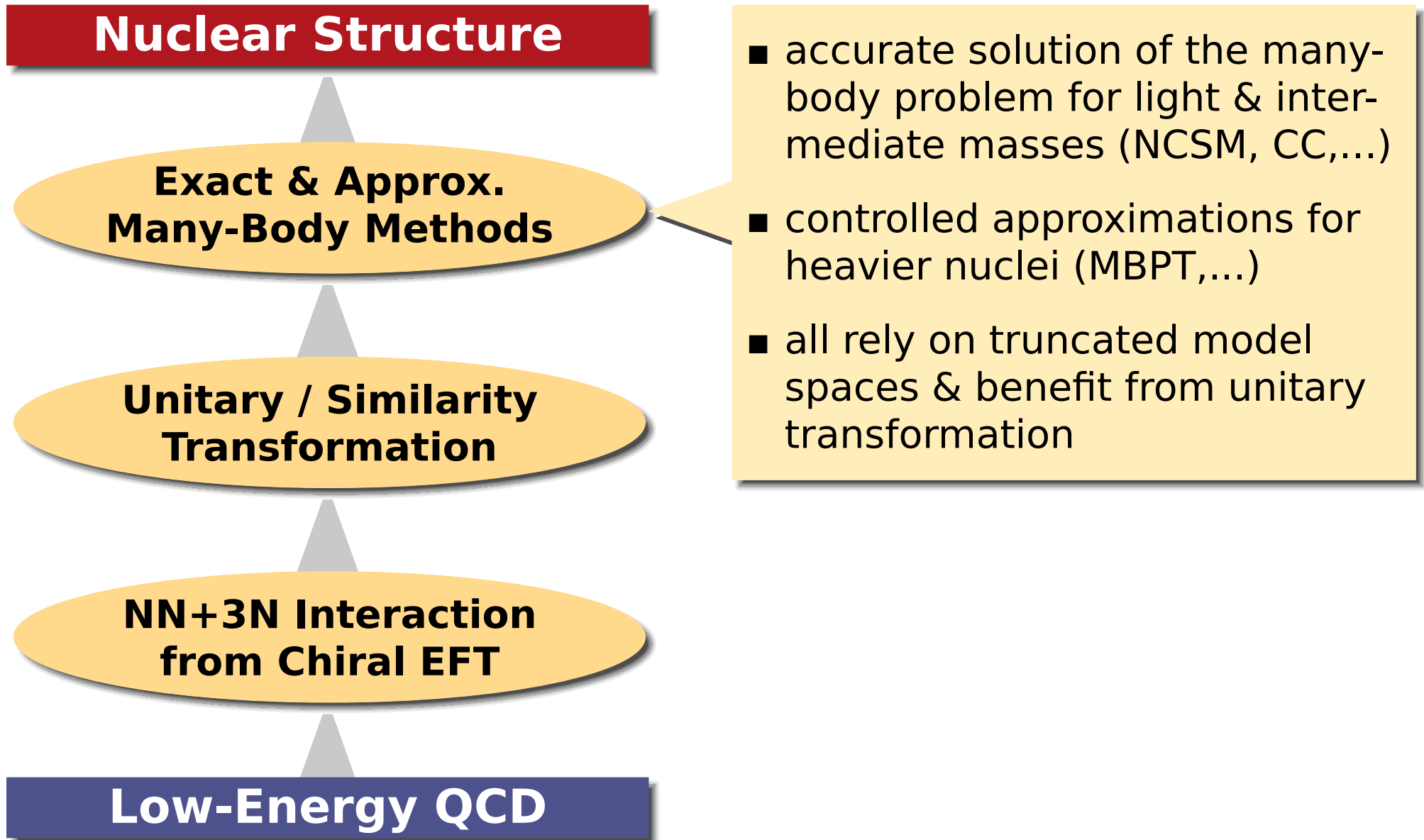
Unitary / Similarity Transformation

**NN+3N Interaction
from Chiral EFT**

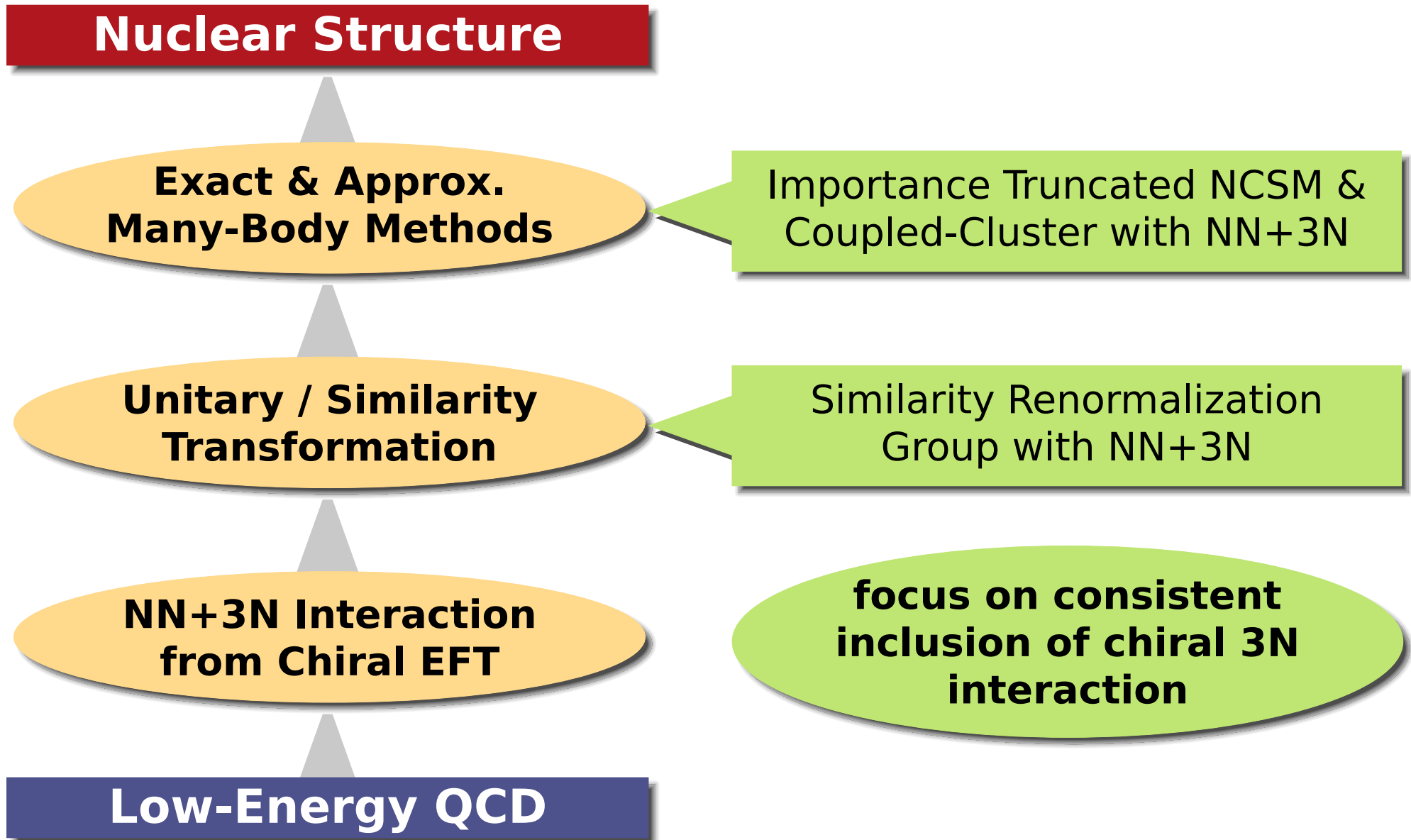
Low-Energy QCD

- adapt Hamiltonian to truncated low-energy model space
 - tame short-range correlations
 - improve convergence behavior
- transform Hamiltonian & observables consistently

From QCD to Nuclear Structure



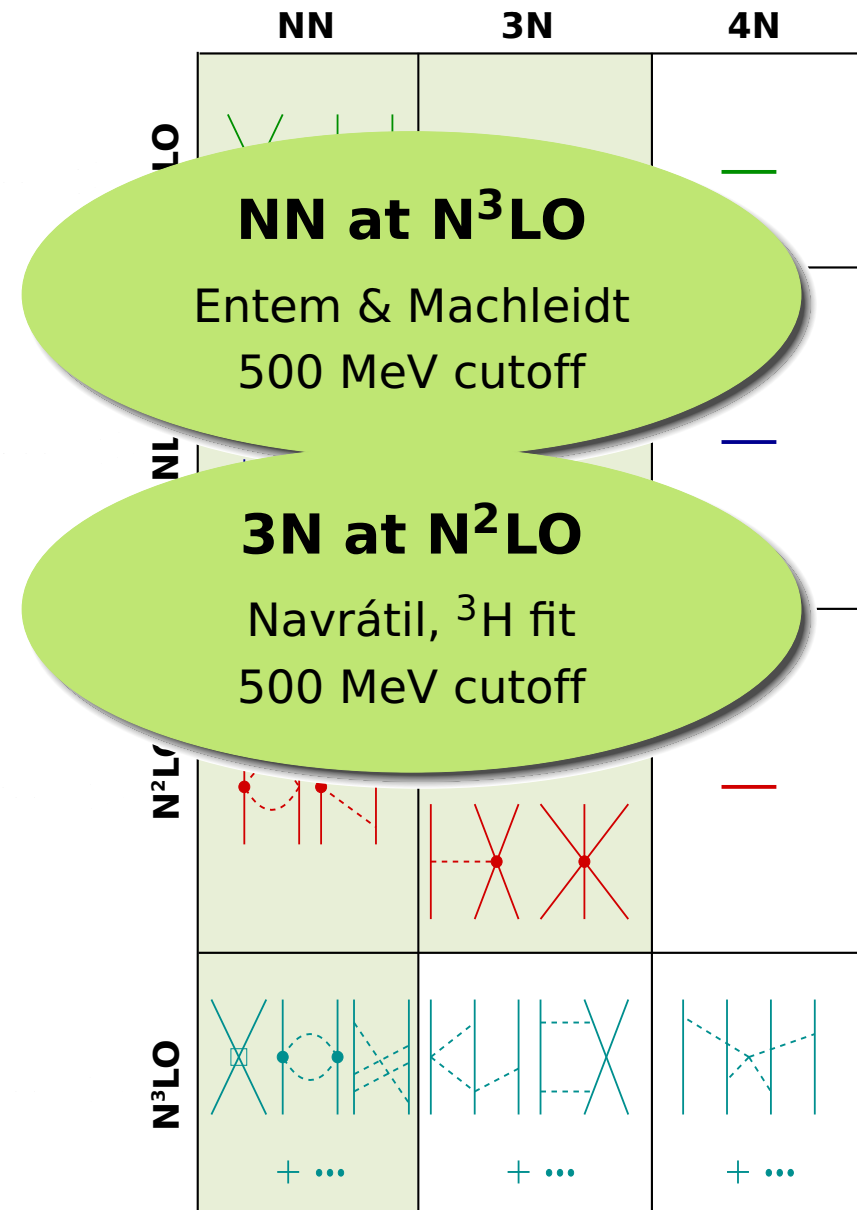
From QCD to Nuclear Structure



Nuclear Interactions from Chiral EFT

Nuclear Interactions from Chiral EFT

- low-energy **effective field theory** for relevant degrees of freedom (π, N) based on symmetries of QCD
- long-range **pion dynamics** explicitly
- short-range physics absorbed in **contact terms**, low-energy constants fitted to experiment ($NN, \pi N, \dots$)
- hierarchy of **consistent NN, 3N, ... interactions** (plus currents)
- many **ongoing developments**
 - 3N interaction at N^3LO
 - explicit inclusion of Δ -resonance
 - formal issues: power counting, renormalization, cutoff choice, ...



Similarity Renormalization Group

Roth, Langhammer, Calci et al. — Phys. Rev. Lett. 107, 072501 (2011)

Roth, Neff, Feldmeier — Prog. Part. Nucl. Phys. 65, 50 (2010)

Roth, Reinhardt, Hergert — Phys. Rev. C 77, 064033 (2008)

Hergert, Roth — Phys. Rev. C 75, 051001(R) (2007)

Similarity Renormalization Group

continuous transformation driving
Hamiltonian to band-diagonal form
with respect to a chosen basis

- **unitary transformation** of Hamiltonian

$$\tilde{H}_\alpha = U_\alpha^\dagger H U_\alpha$$

simplicity and flexibility
are great advantages of
the SRG approach

- **evolution equations** for \tilde{H}_α and U_α

$$\frac{d}{d\alpha} \tilde{H}_\alpha = [\eta_\alpha, \tilde{H}_\alpha]$$

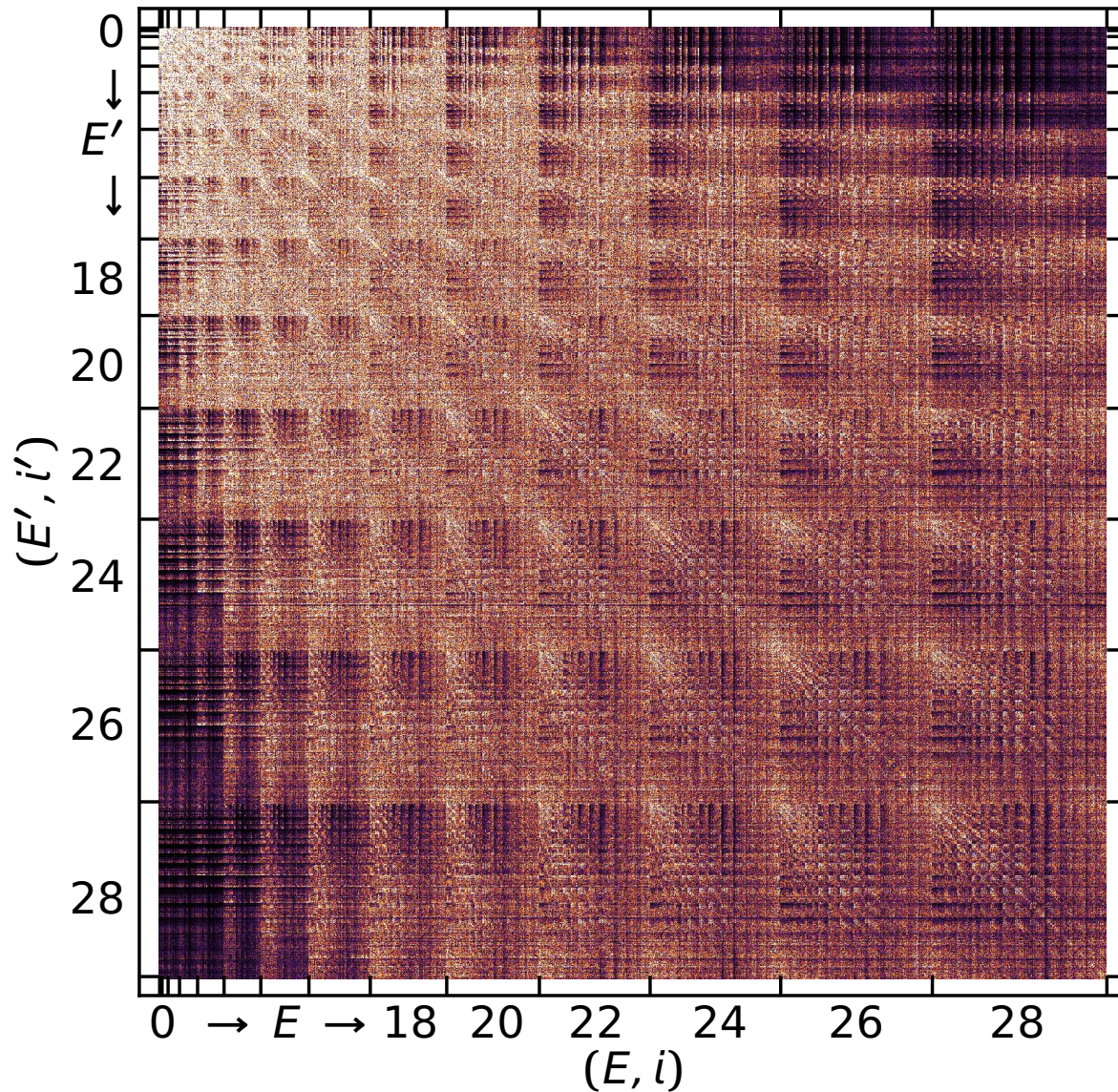
solve SRG evolution
equations using two- &
three-body Jacobi HO
representation

- **dynamic generator**: commutator with the operator in whose
eigenbasis H shall be diagonalized

$$\eta_\alpha = (2\mu)^2 [T_{\text{int}}, \tilde{H}_\alpha]$$

SRG Evolution in Three-Body Space

3B-Jacobi HO matrix elements

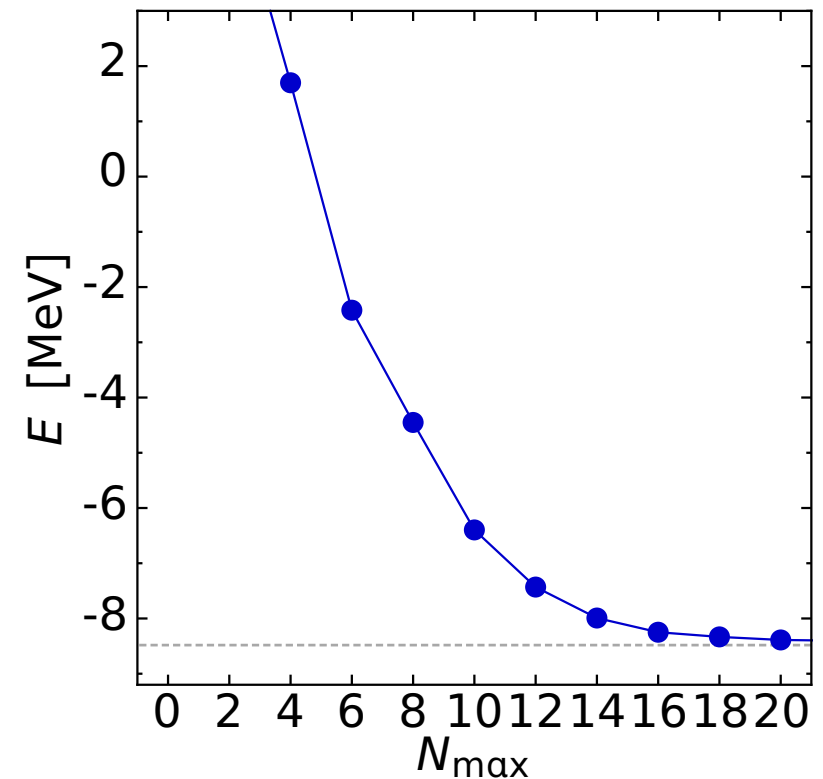


$$\alpha = 0.000 \text{ fm}^4$$

$$\Lambda = \infty \text{ fm}^{-1}$$

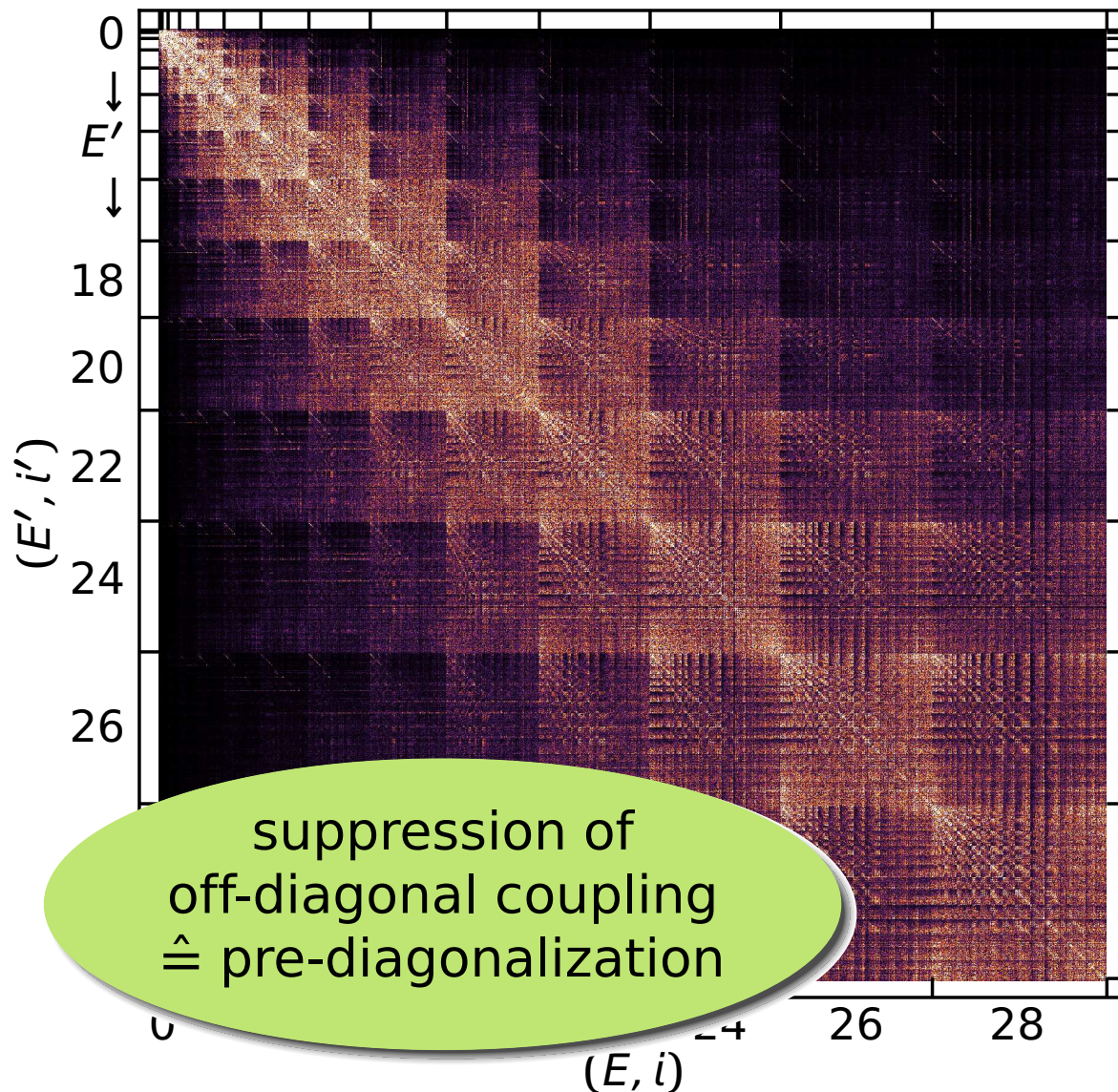
$$J^\pi = \frac{1}{2}^+, T = \frac{1}{2}, \hbar\Omega = 28 \text{ MeV}$$

NCSM ground state ${}^3\text{H}$



SRG Evolution in Three-Body Space

3B-Jacobi HO matrix elements

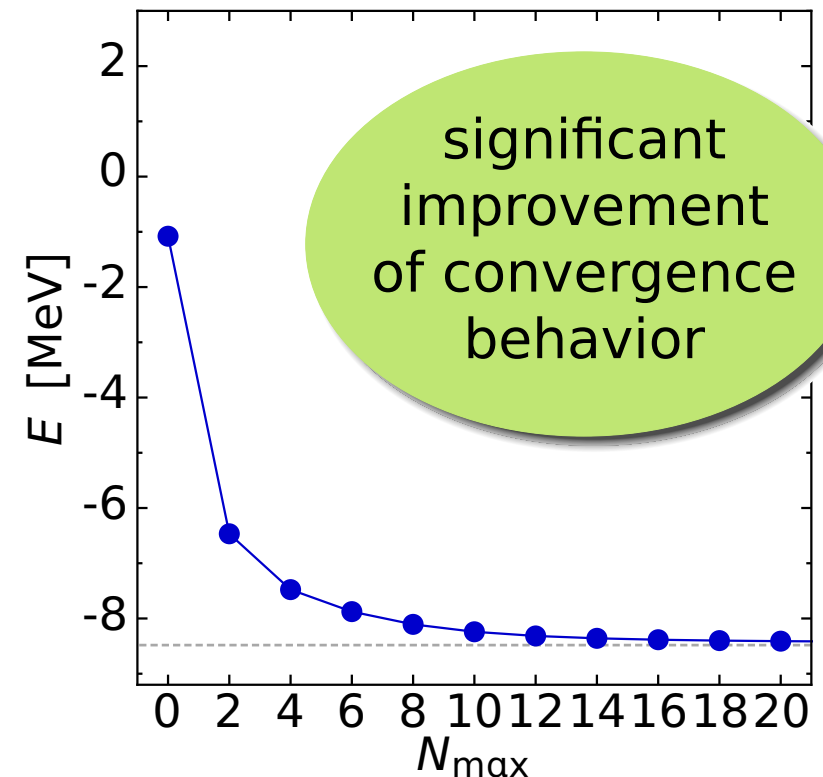


$$\alpha = 0.320 \text{ fm}^4$$

$$\Lambda = 1.33 \text{ fm}^{-1}$$

$$J^\pi = \frac{1}{2}^+, T = \frac{1}{2}, \hbar\Omega = 28 \text{ MeV}$$

NCSM ground state ${}^3\text{H}$



Calculations in A-Body Space

- evolution **induces n -body contributions** $\tilde{H}_\alpha^{[n]}$ to Hamiltonian

$$\tilde{H}_\alpha = \tilde{H}_\alpha^{[1]} + \tilde{H}_\alpha^{[2]} + \tilde{H}_\alpha^{[3]} + \tilde{H}_\alpha^{[4]} + \dots$$

- truncation of cluster series inevitable — formally destroys unitarity and invariance of energy eigenvalues (independence of α)

Three SRG-Evolved Hamiltonians

- **NN only**: start with NN initial Hamiltonian and keep two-body terms only
- **NN+3N-induced**: start with NN initial Hamiltonian and keep two- and three-body terms
- **NN+3N-full**: start with NN+3N initial Hamiltonian and keep two- and three-body terms

α -variation provides a **diagnostic tool** to assess the contributions of omitted many-body interactions

Sounds easy, but...

- ❶ computation of initial 2B/3B-Jacobi HO matrix elements of chiral NN+3N interactions
 - we use Petr Navratil's ManyEff code for computing 3B-Jacobi matrix elements and corresponding CFPs
- ❷ SRG evolution in 2B/3B space and cluster decomposition
 - efficient implementation using adaptive ODE solver & BLAS; largest block takes a few hours on single node
- ❸ transformation of 2B/3B Jacobi HO matrix elements into JT-coupled representation
 - formulated transformation directly into JT-coupled scheme; highly efficient implementation; can handle $E_{3\max} = 16$ in JT-coupled scheme
- ❹ data management and on-the-fly decoupling in many-body codes
 - invented optimized storage scheme for fast on-the-fly decoupling; can keep all matrix elements up to $E_{3\max} = 16$ in memory

Importance Truncated NCSM

Roth, Langhammer, Calci et al. — Phys. Rev. Lett. 107, 072501 (2011)

Navrátil, Roth, Quaglioni — Phys. Rev. C 82, 034609 (2010)

Roth — Phys. Rev. C 79, 064324 (2009)

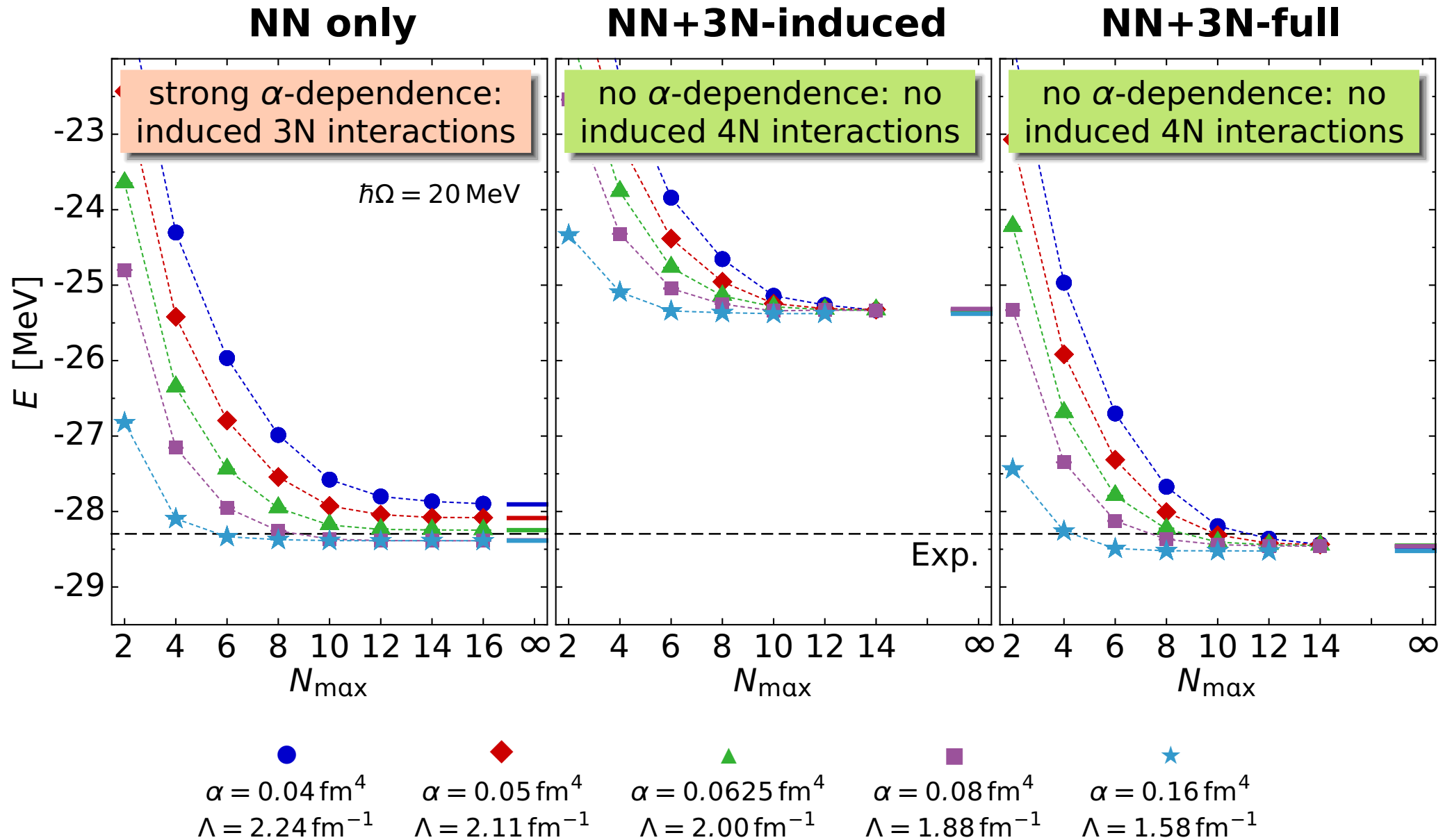
Roth, Navrátil — Phys. Rev. Lett. 99, 092501 (2007)

Importance Truncated NCSM

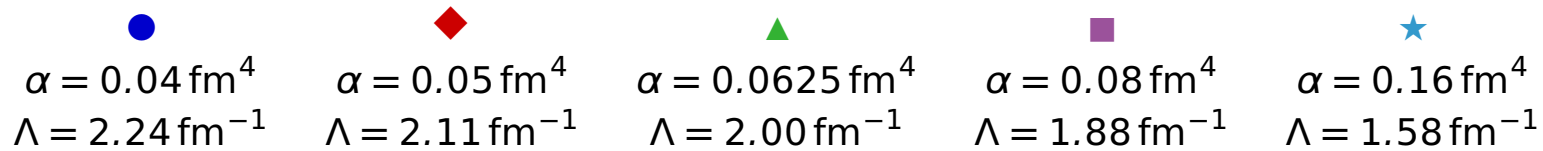
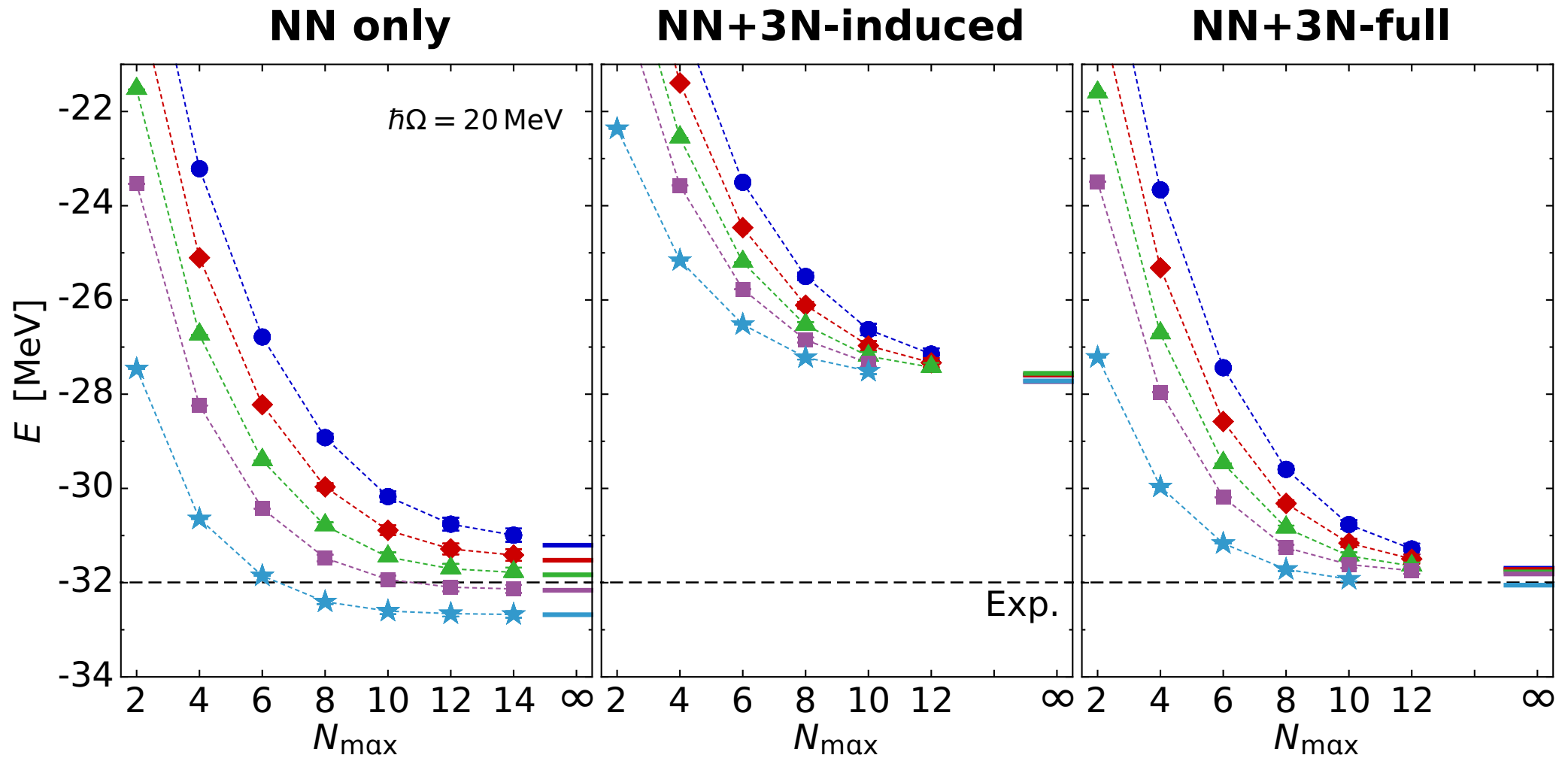
NCSM is one of the most powerful and universal exact ab-initio methods

- construct matrix representation of Hamiltonian using a **basis of HO Slater determinants** truncated w.r.t. HO excitation energy $N_{\max}\hbar\Omega$
- solve **large-scale eigenvalue problem** for a few extremal eigenvalues
- **all relevant observables** can be computed from the eigenstates
- range of applicability limited by **factorial growth** of basis with N_{\max} & A
- adaptive **importance truncation** extends the range of NCSM by reducing the model space to physically relevant states
- we have developed a **parallelized IT-NCSM/NCSM code** capable of handling $3N$ matrix elements up to $E_{3\max} = 16$

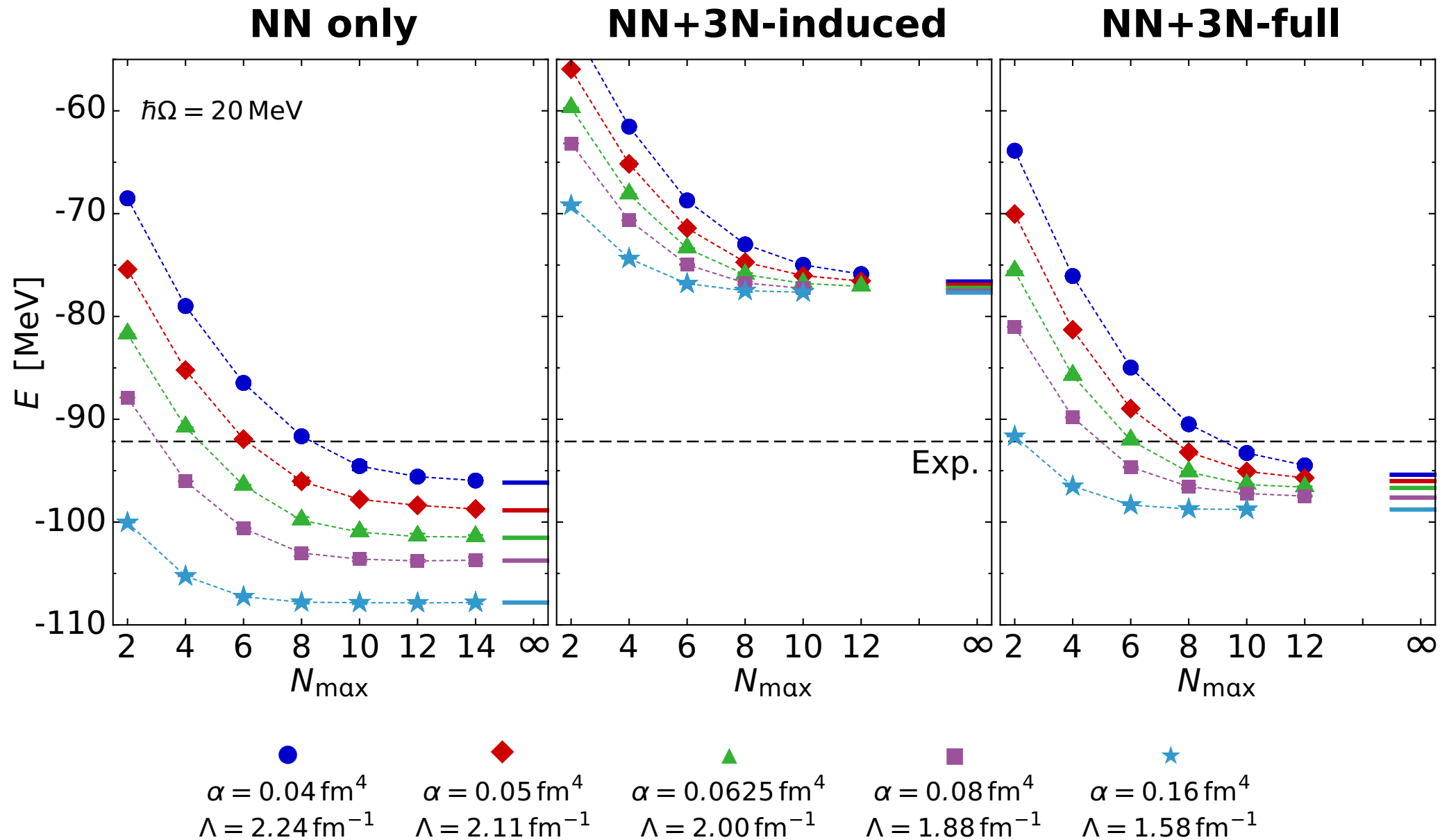
^4He : Ground-State Energies



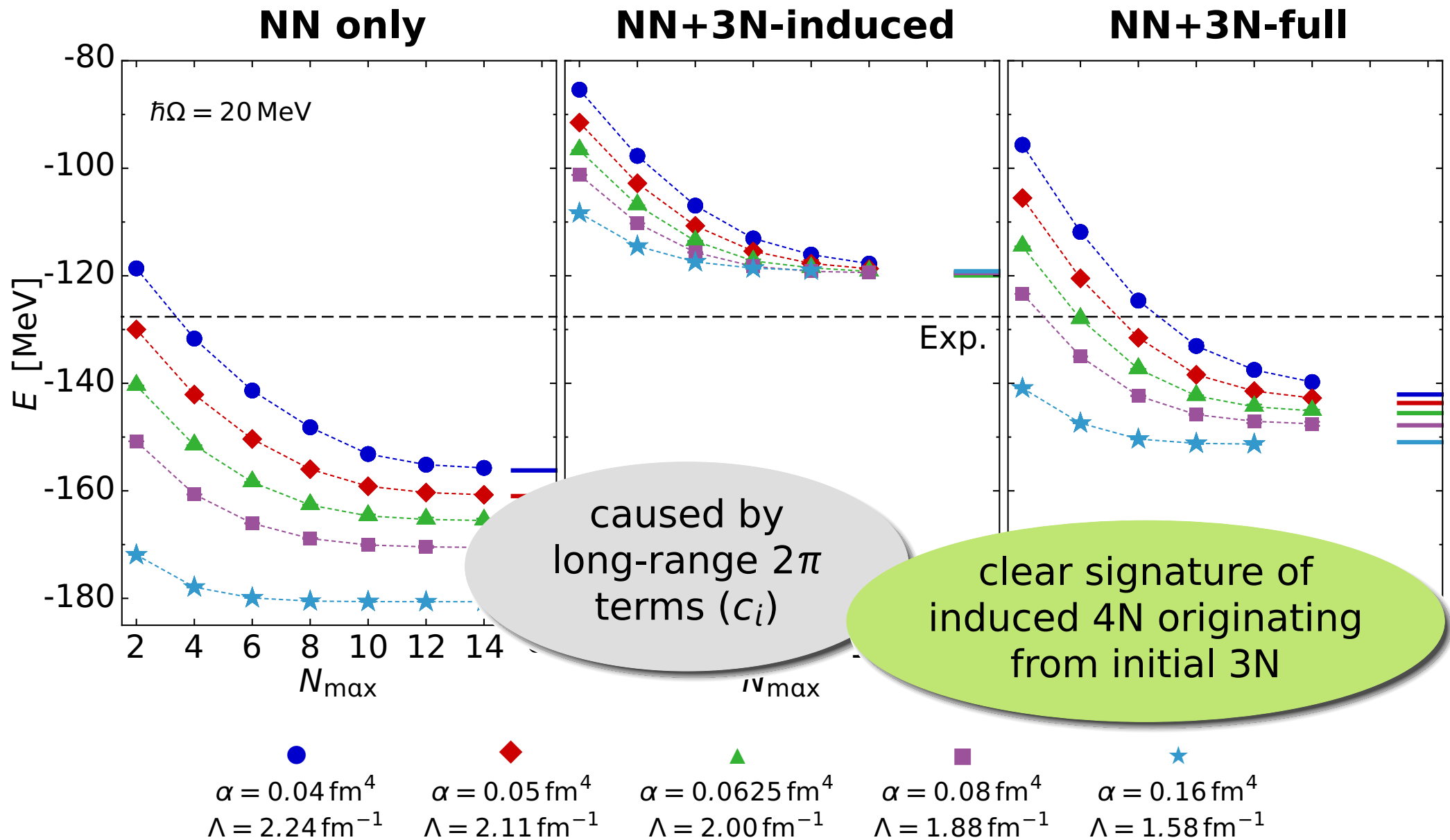
${}^6\text{Li}$: Ground-State Energies



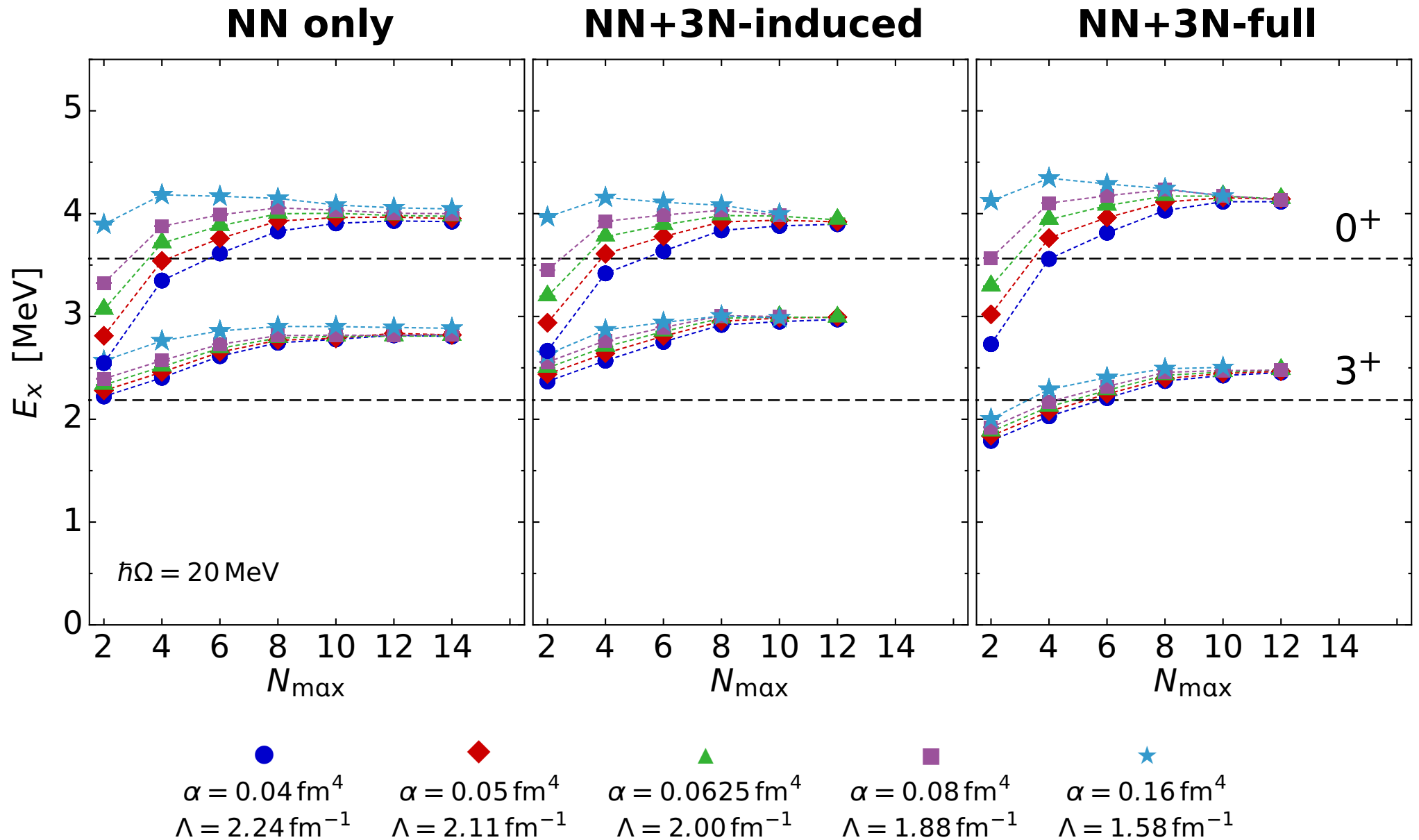
^{12}C : Ground-State Energies



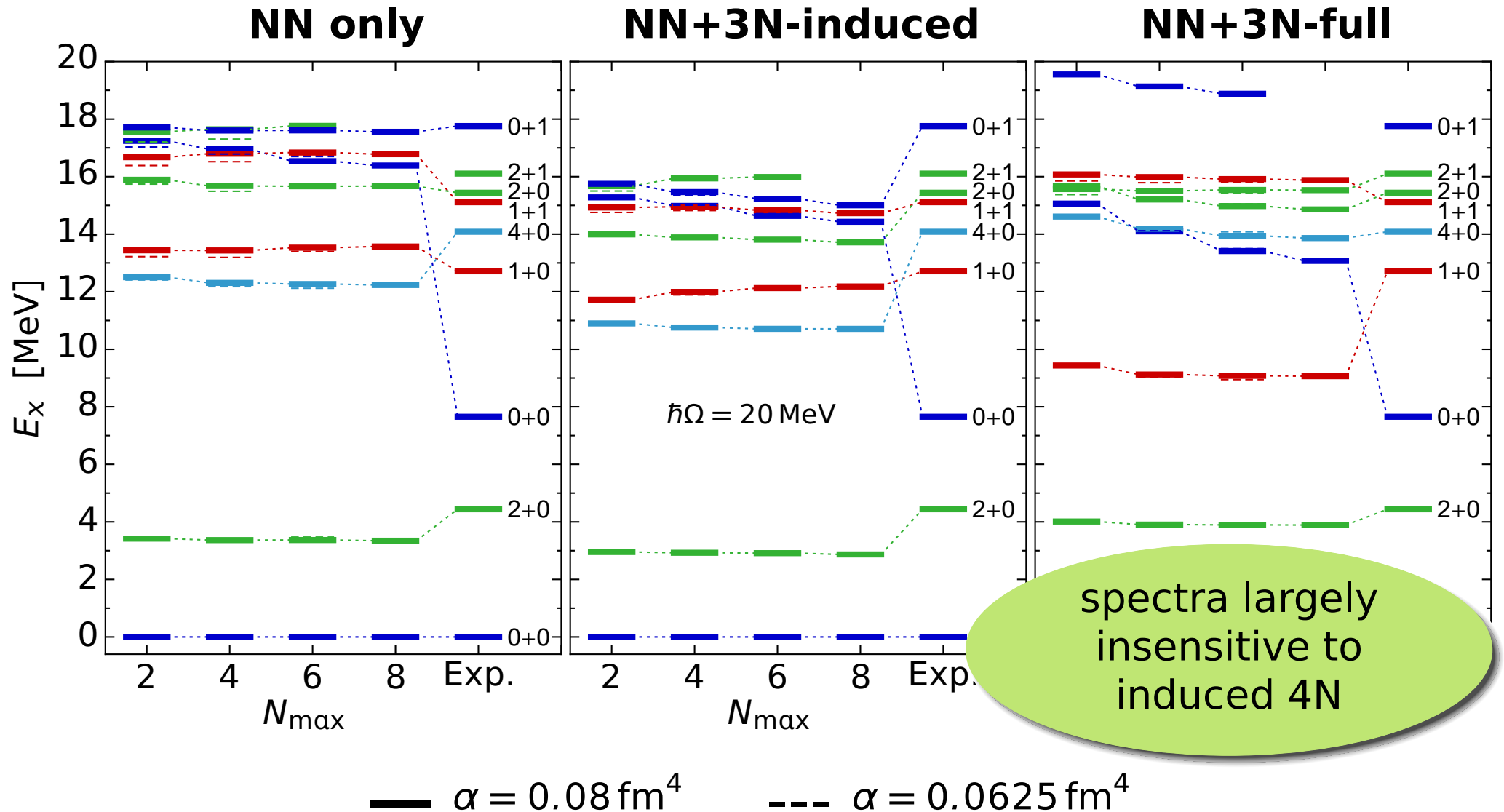
^{16}O : Ground-State Energies



${}^6\text{Li}$: Excitation Energies



Spectroscopy of ^{12}C



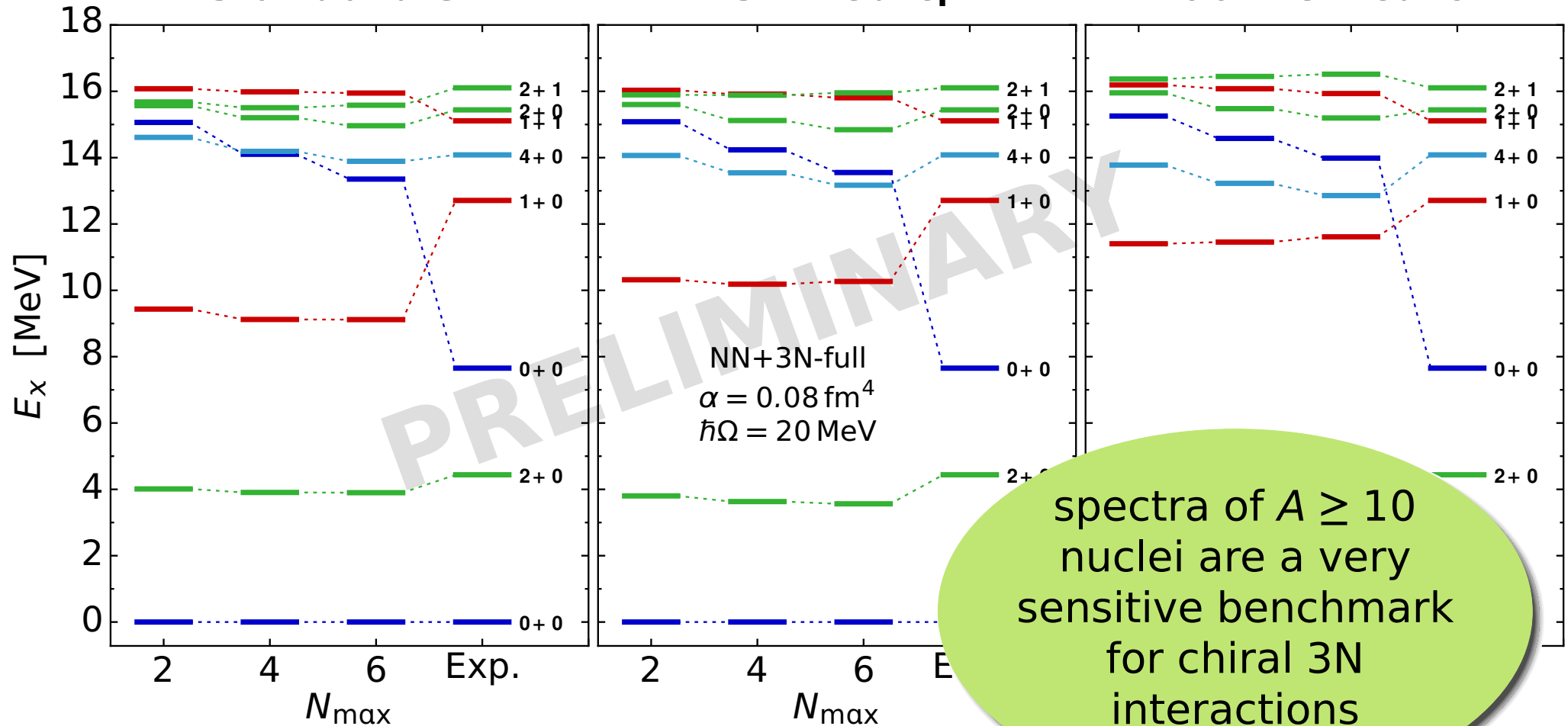
Outlook: Sensitivity on Initial 3N

modified 3N interaction with

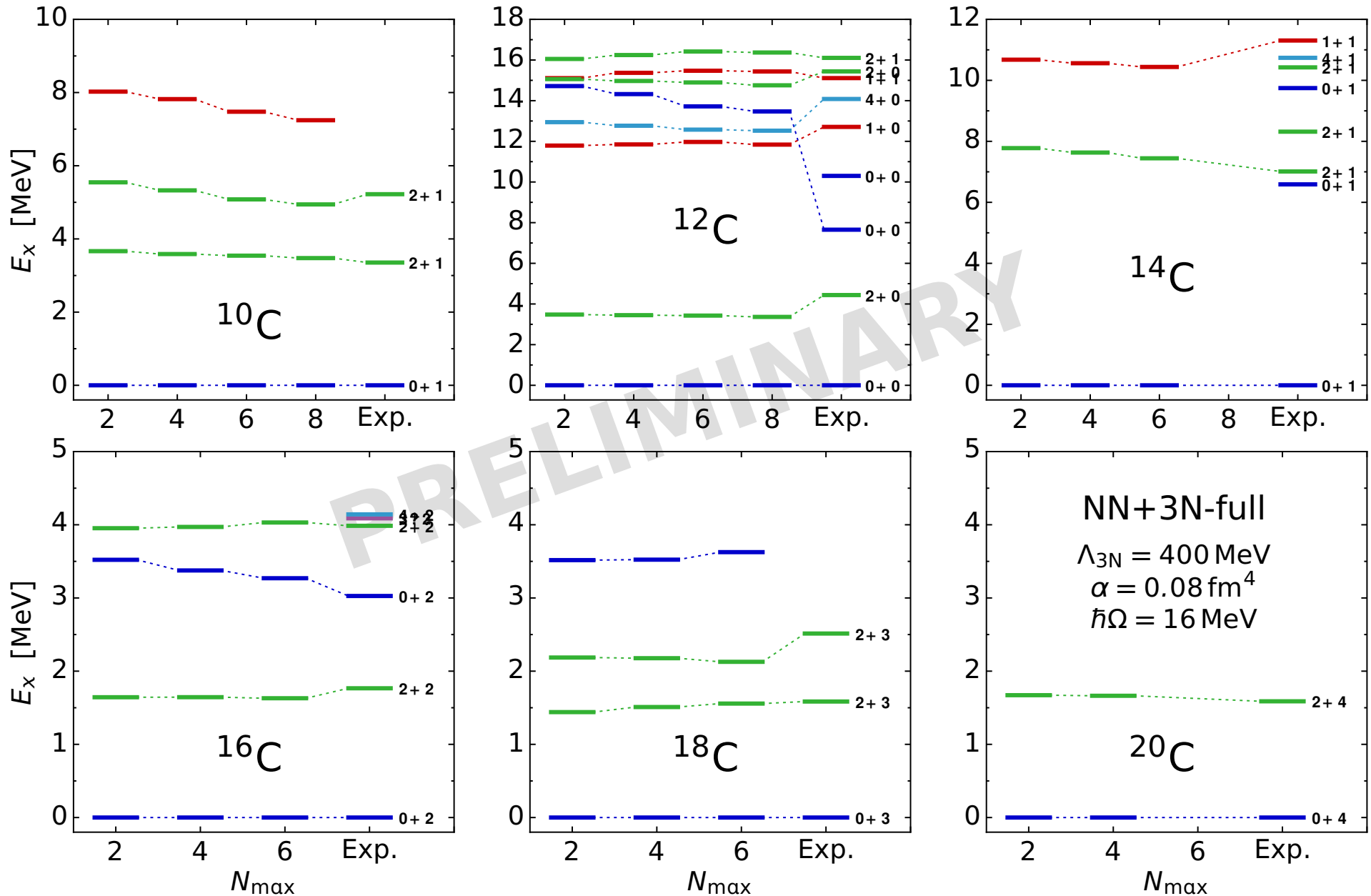
standard 3N

shifted c_i

400 MeV cutoff



Outlook: Carbon Isotopic Chain



Normal-Ordered $3N$ Interaction & Coupled-Cluster Method

Roth, Binder, Vobig et al. — arXiv: 1112.0287 (2011)

Normal-Ordered 3N Interaction

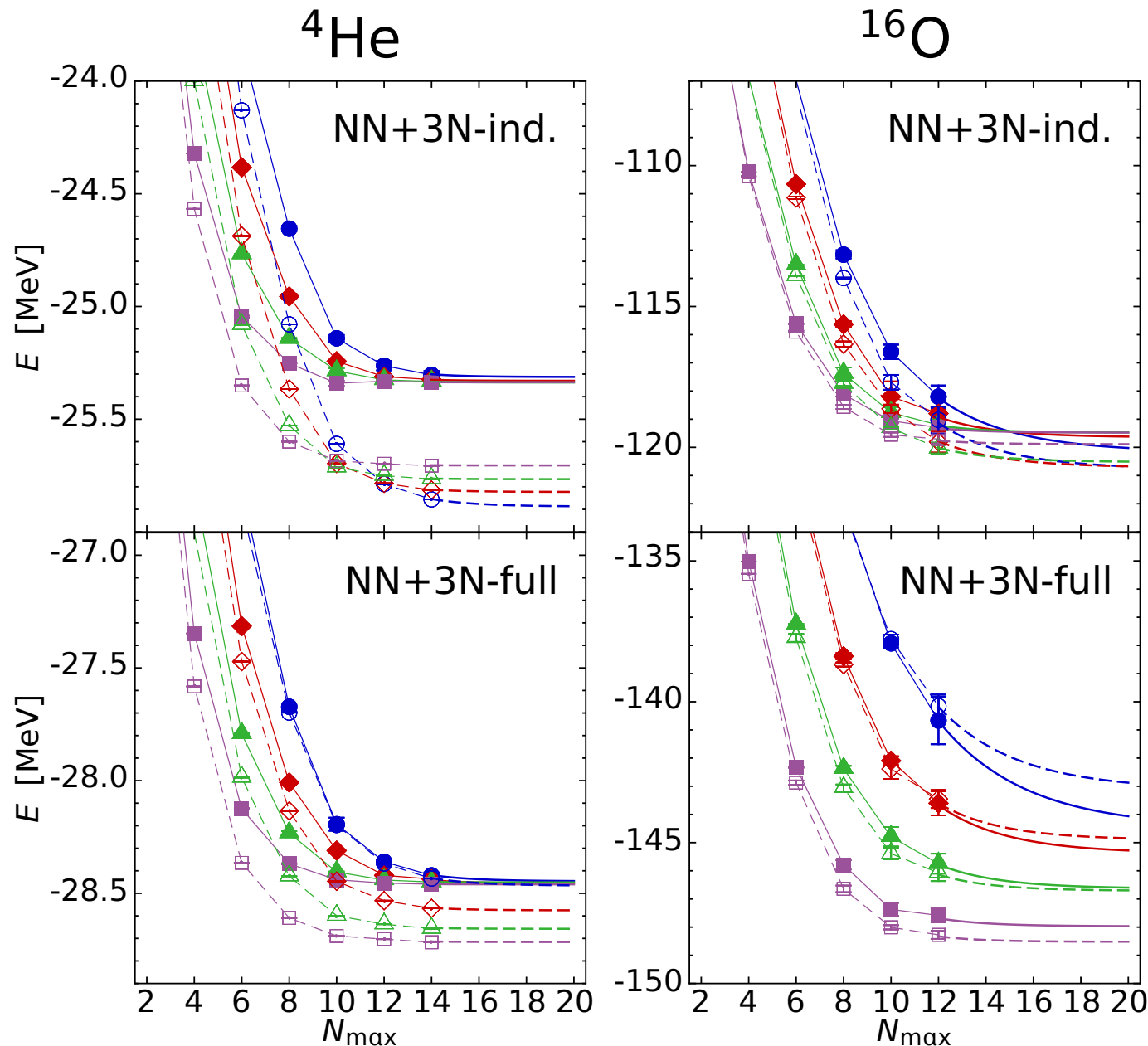
avoid technical challenge of including explicit 3N interactions in many-body calculation

- **idea**: write 3N interaction in normal-ordered form with respect to an A -body reference Slater-determinant ($0\hbar\Omega$ state)

$$\begin{aligned} V_{3N} &= \sum V_{\circ\circ\circ\circ\circ}^{3N} a_{\circ}^{\dagger} a_{\circ}^{\dagger} a_{\circ}^{\dagger} a_{\circ} a_{\circ} a_{\circ} \\ &= W^{0B} + \sum W_{\circ\circ}^{1B} \{a_{\circ}^{\dagger} a_{\circ}\} + \sum W_{\circ\circ\circ}^{2B} \{a_{\circ}^{\dagger} a_{\circ}^{\dagger} a_{\circ} a_{\circ}\} \\ &\quad + \sum W_{\circ\circ\circ\circ\circ}^{3B} \{a_{\circ}^{\dagger} a_{\circ}^{\dagger} a_{\circ}^{\dagger} a_{\circ} a_{\circ} a_{\circ}\} \end{aligned}$$

- **question**: if we neglect the normal-ordered 3B term, how well does this approximation work ?

Benchmark of Normal-Ordered 3N



■ compare IT-NCSM results with full 3N to normal-ord. 3N truncated at the 2B level

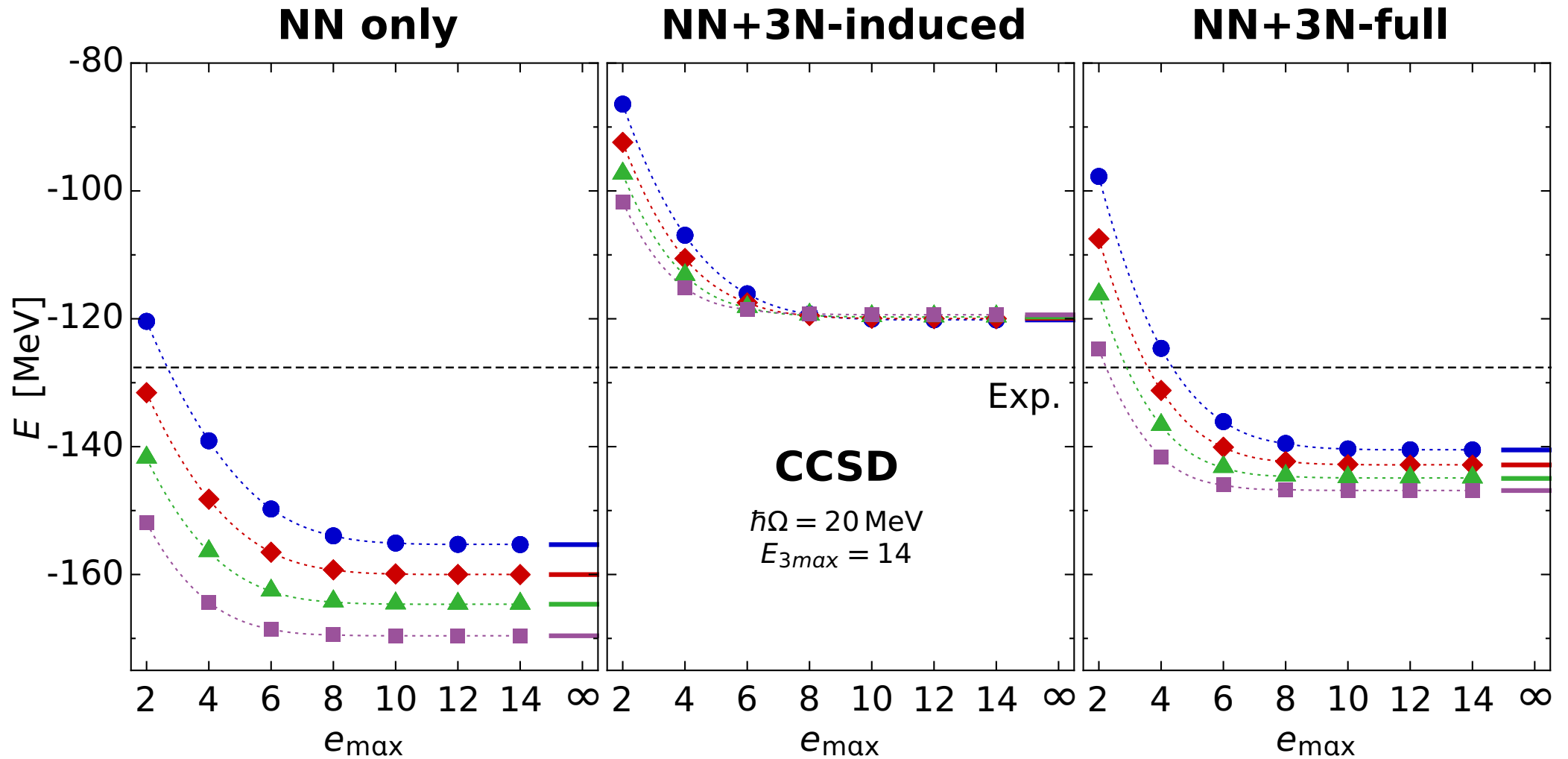
■ typical deviations up to 2% for ${}^4\text{He}$ and 1% for ${}^{16}\text{O}$

full / NO2B

- / ○ $\alpha = 0.04 \text{ fm}^4$
- ◆ / ◇ $\alpha = 0.05 \text{ fm}^4$
- ▲ / △ $\alpha = 0.0625 \text{ fm}^4$
- / □ $\alpha = 0.08 \text{ fm}^4$

$\hbar\Omega = 20 \text{ MeV}$

^{16}O : Coupled-Cluster with $3N_{\text{NO2B}}$



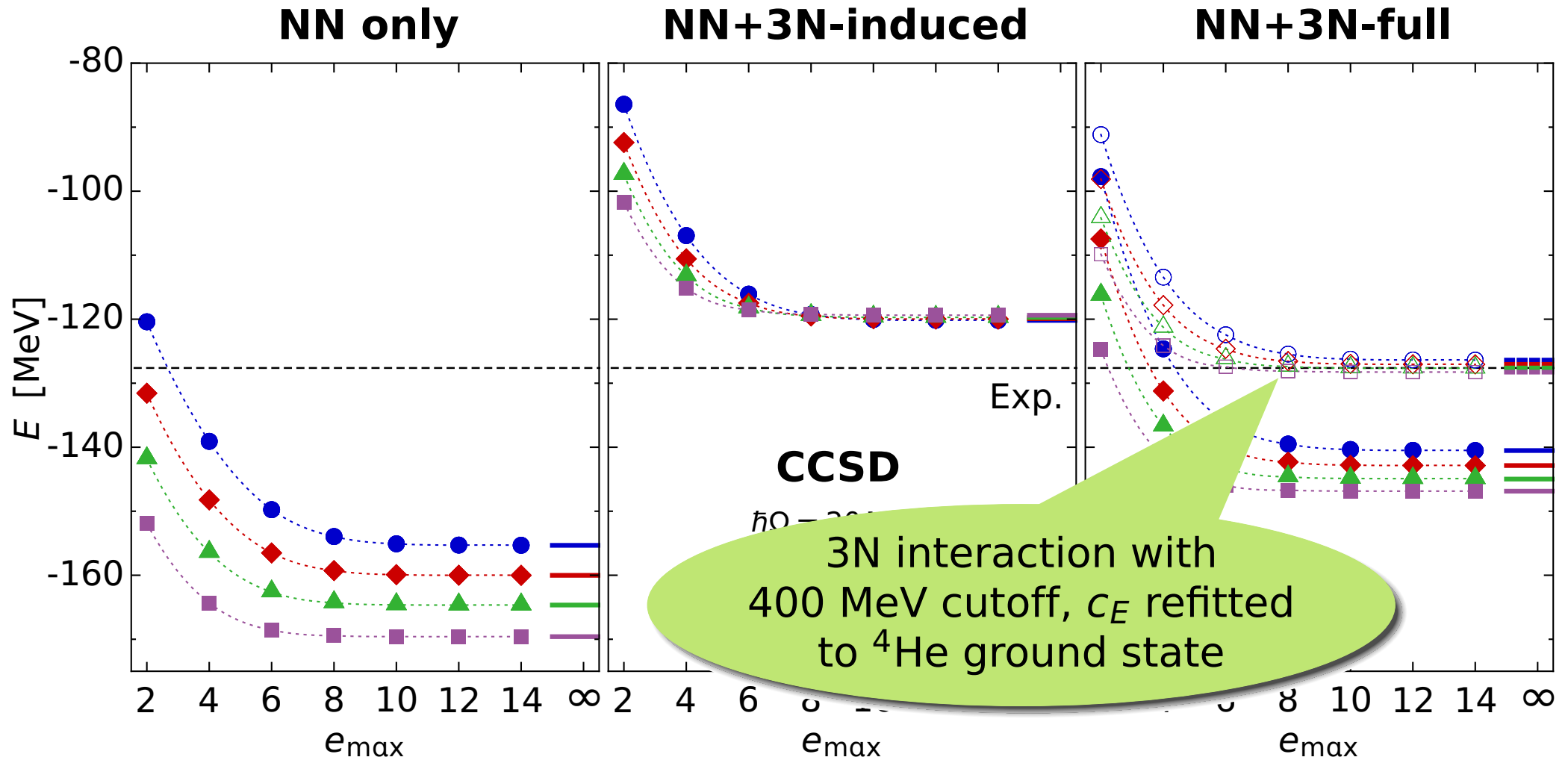
● $\alpha = 0.04 \text{ fm}^4$
 $\Lambda = 2.24 \text{ fm}^{-1}$

◆ $\alpha = 0.05 \text{ fm}^4$
 $\Lambda = 2.11 \text{ fm}^{-1}$

▲ $\alpha = 0.0625 \text{ fm}^4$
 $\Lambda = 2.00 \text{ fm}^{-1}$

■ $\alpha = 0.08 \text{ fm}^4$
 $\Lambda = 1.88 \text{ fm}^{-1}$

^{16}O : Coupled-Cluster with $3N_{\text{NO2B}}$



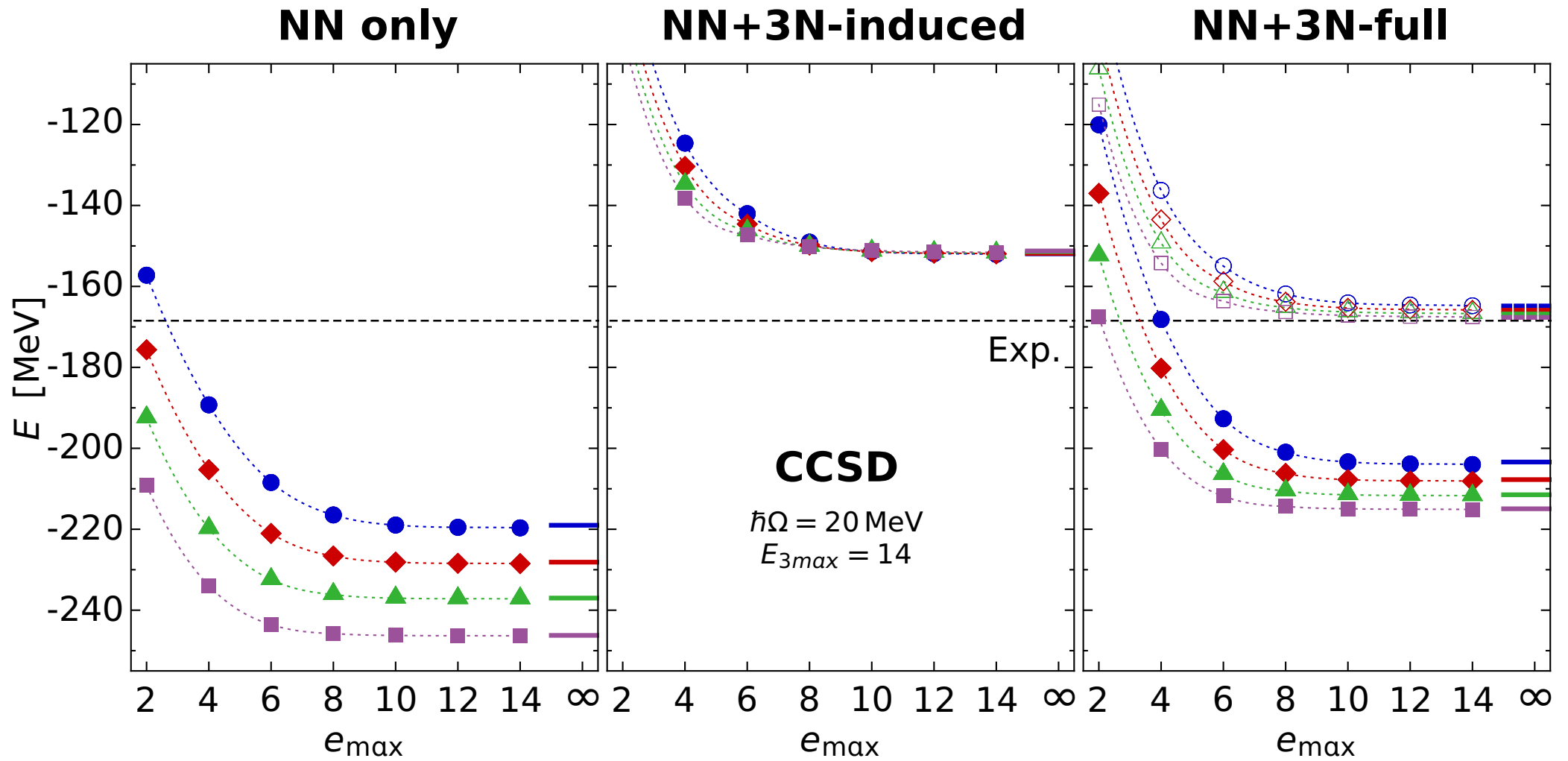
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^{24}O : Coupled-Cluster with $3N_{\text{NO2B}}$



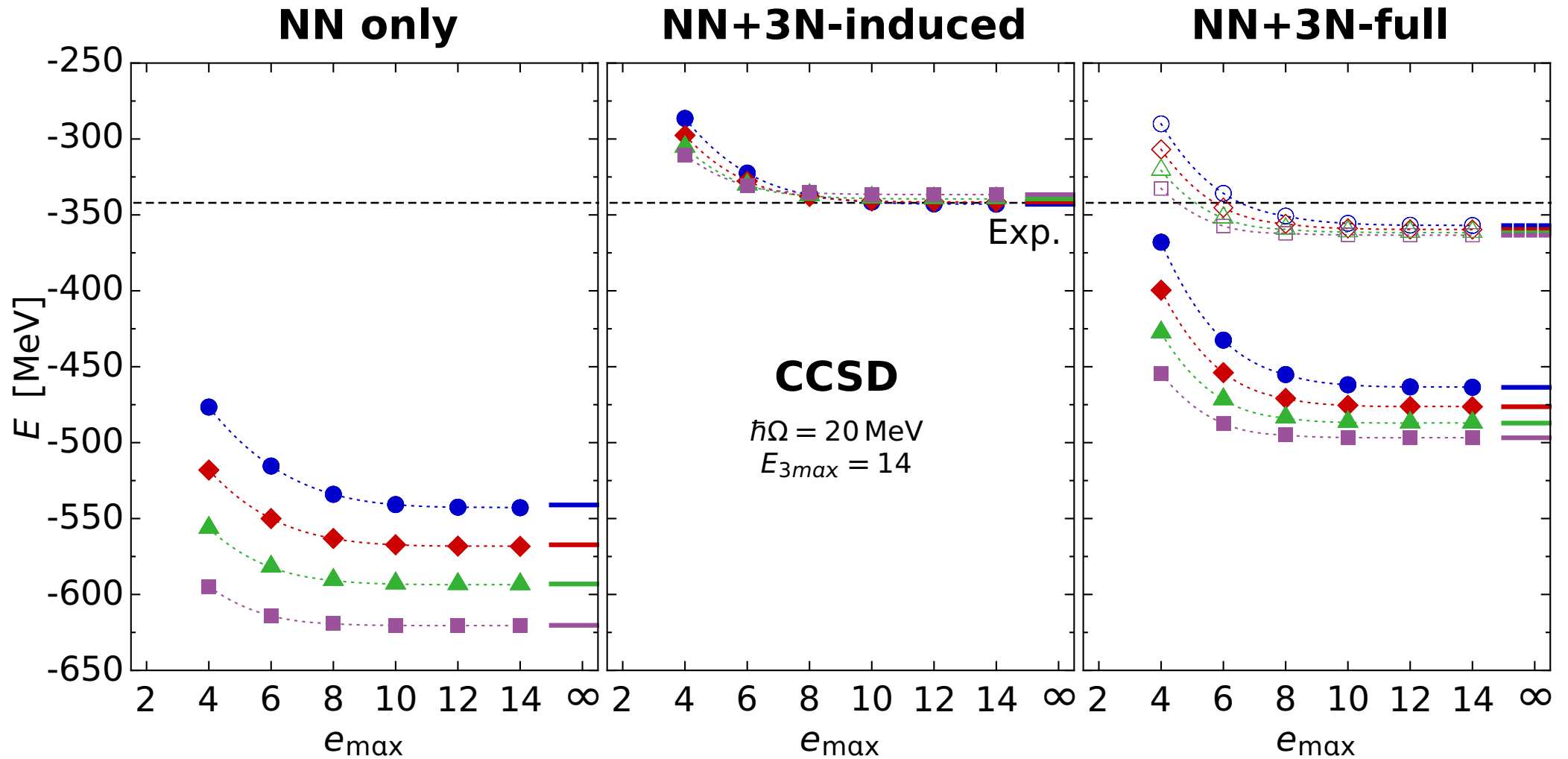
●
 $\alpha = 0.04 \text{ fm}^4$
 $\Lambda = 2.24 \text{ fm}^{-1}$

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 $\alpha = 0.05 \text{ fm}^4$
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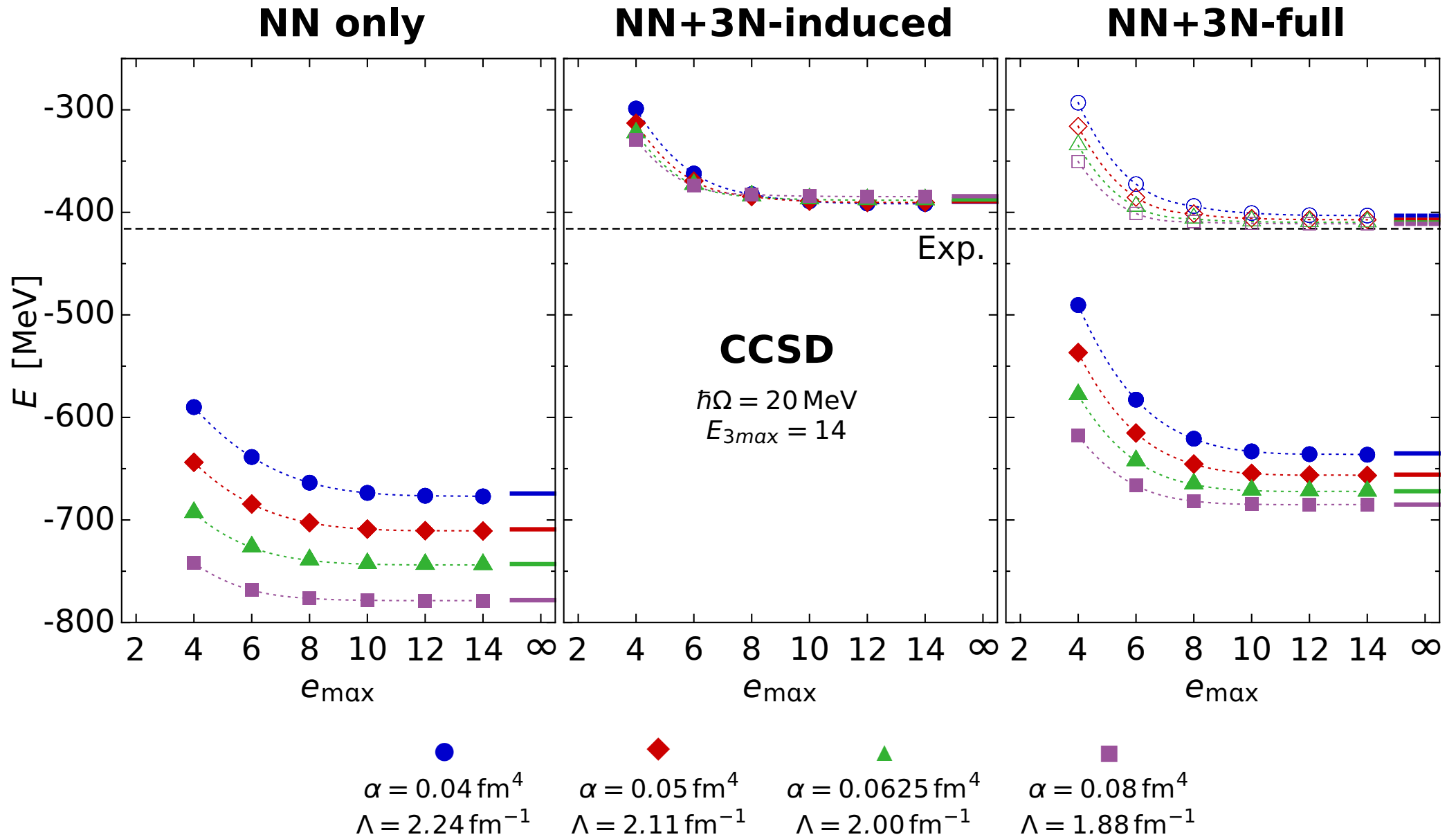
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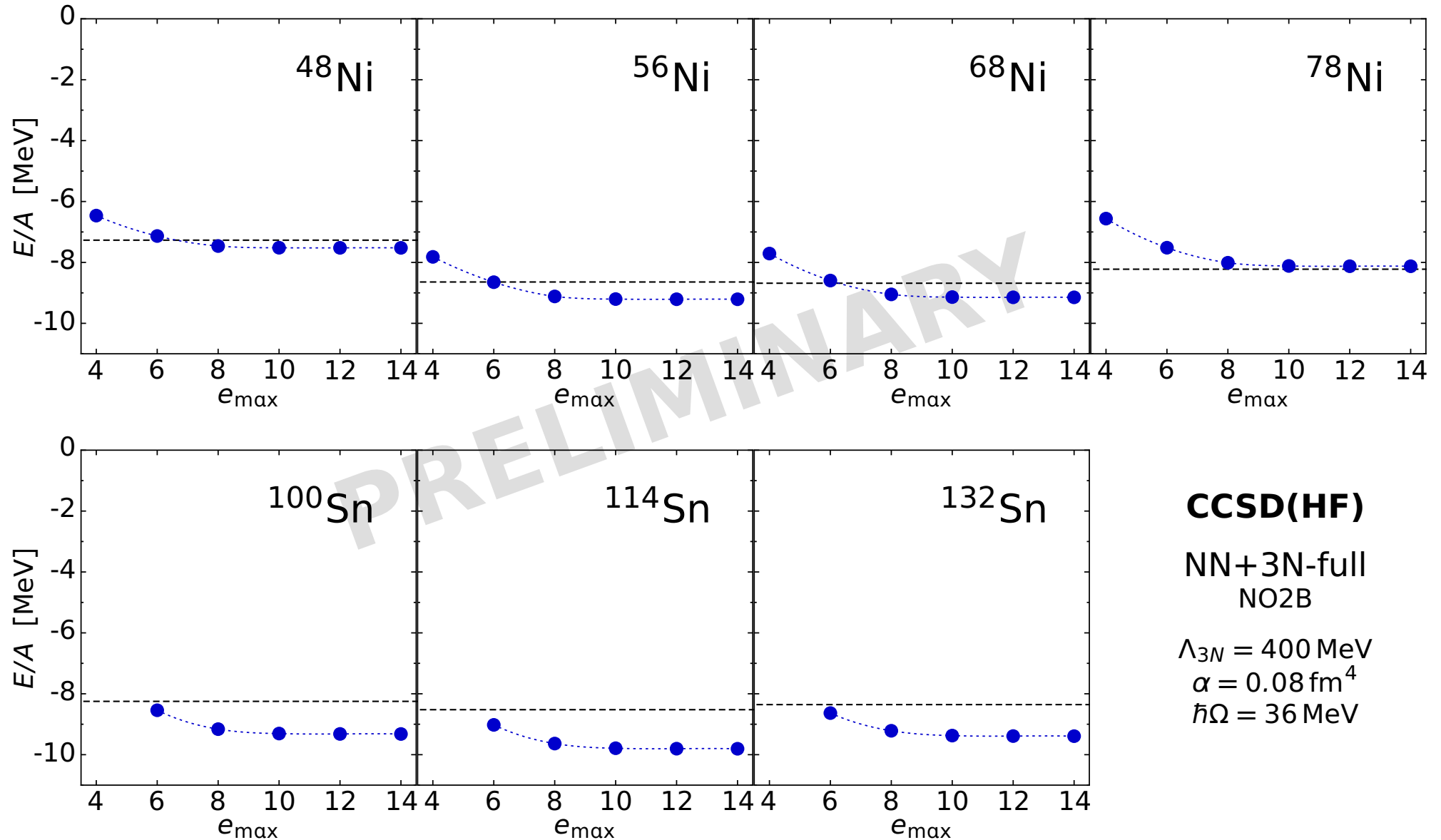
^{40}Ca : Coupled-Cluster with $3N_{\text{NO2B}}$



^{48}Ca : Coupled-Cluster with $3N_{\text{NO2B}}$



Outlook: Chiral 3N for Heavy Nuclei



Conclusions

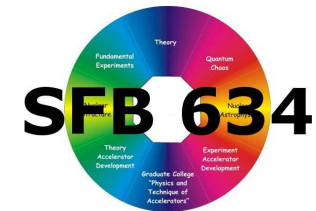
Conclusions

- new era of **ab-initio nuclear structure and reaction theory** connected to QCD via chiral EFT
 - chiral EFT as universal starting point... some issues remain
- consistent **inclusion of 3N interactions** in similarity transformations & many-body calculations
 - breakthrough in computation & handling of 3N matrix elements
- **innovations in many-body theory**: extended reach of exact methods & improved control over approximations
 - versatile toolbox for different observables & mass ranges
- many **exciting applications** ahead...

Epilogue

■ thanks to my group & my collaborators

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