Hartree-Fock and RPA for Axially Deformed Nuclei with Realistic Interactions

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Outline

Introduction

Basics RPA Formalism for Deformed Nuclei

Results ²⁰Ne Angular Momentum Projection Three-Body Forces ²⁸Si Deformation ³²S Transition Densities

Basics

Interaction

Unitarily transformed interaction, based on realistic AV18 NN interaction

- suitable for restricted Hilbert spaces
- tamed short-range correlations
- improved convergence behavior
- conservation of phase shifts
- phenomenological 3-body contact term one parameter

Many body method

Hartree-Fock and Random Phase Approximation (RPA) with axial deformation

- Spherical Harmonic Oscillator basis, truncated at e = 2 n + l ≤ 14, l ≤ 10
- ► HF variation over n,l,j,m_j m_j and parity of l are good quantum numbers (m_j is usually reffered to as k)
- exact angular momentum projection
- approximative variation after projection with β constraint

RPA Formalism for Axially Deformed Nuclei

Vibration creation operator

ph selection rules

$$\begin{split} \hat{Q}_{\omega}^{\dagger} &= \sum_{ph} X_{ph}^{\omega} \, \hat{a}_{p}^{\dagger} \hat{a}_{h} - \sum_{ph} Y_{ph}^{\omega} \, \hat{a}_{h}^{\dagger} \hat{a}_{p} & \pi_{p} - \pi_{h} = \pi_{\omega} \\ & |m_{p} - m_{h}| = |m_{\omega}| \\ \hat{Q}_{\omega} |\mathsf{RPA}\rangle &= 0 & , \qquad \hat{Q}_{\omega}^{\dagger} |\mathsf{RPA}\rangle = |\omega\rangle \end{split}$$

Quasi-Boson-Approximation

$$\langle \mathsf{RPA} | \hat{H} \, \hat{Q}_\omega^\dagger | \mathsf{RPA} \rangle \rightarrow \langle \mathsf{HF} | [\hat{H}, \hat{Q}_\omega^\dagger] | \mathsf{HF} \rangle$$

RPA equations in *ph*-space

$$\begin{pmatrix} A & B \\ -B^* & -A^* \end{pmatrix} \begin{pmatrix} X^{\nu} \\ Y^{\nu} \end{pmatrix} = \hbar \omega_{\nu} \begin{pmatrix} X^{\nu} \\ Y^{\nu} \end{pmatrix}$$
$$A_{ph,p'h'} = \delta_{pp'} \delta_{hh'} (e_p - e_h) + H_{hp',ph'} \quad , \quad B_{ph,p'h'} = H_{hh',pp'}$$

RPA Transition Amplitudes

Unprojected transition amplitudes

$$(\mathsf{RPA} \| \hat{T}^{\lambda} \| \omega) = \sum_{\mu, ph} \left(X^{\omega}_{ph} + (-1)^{K_{ph}} Y^{\omega}_{ph}
ight) \langle h | \hat{T}^{\lambda}_{\mu} | p
angle$$

Projected transition amplitudes

$$(\operatorname{RPA} \| \hat{T}^{\lambda} \| \omega) = (2 J_0 + 1) N_0 N_\omega (-1)^{J_0 - K_0} \sum \begin{pmatrix} J_0 & \lambda & J_\omega \\ -K_0 & \mu & K_0 - \mu \end{pmatrix}$$

no additional approximations for
angular momentum projection
$$\hat{P}^J_{MK} = \frac{2 J + 1}{2} \int_{-1}^{1} d^J_{MK}(\beta) e^{i\beta J_y} d\cos(\beta)$$

HF Ground State - ²⁰Ne



Effect of HF angular momentum projection - ²⁰Ne



Effect of RPA angular momentum projection - ²⁰Ne



Effect of 3-body forces - ²⁰Ne



HF Ground State - ²⁸Si



Effect of prolate and oblate deformation - ²⁸Si



HF Ground State - ³²S



K contributions - ³²S

0.6

 $R(E1IV)~(e^{2}fm^{2}/MeV)$



► K defines the shape of the oscillation

Isovector Dipole





Transition densities - ${}^{32}S$ J = 0 K = 0 E = 8.2MeV



Transition densities - ${}^{32}S$ J = 0 K = 0 E = 11.5MeV



E (MeV)

Transition densities - ${}^{32}S$ J = 0 K = 0 E = 11.6MeV



Transition densities - ${}^{32}S$ J = 0 K = 0 E = 13.9MeV



Transition densities - ${}^{32}S$ J = 2 K = 0 E = 13.9MeV



Transition densities - ${}^{32}S$ J = 2 K = 0 E = 16.0 MeV



Transition densities - ${}^{32}S$ J = 2 K = 1 E = 2.3MeV



Transition densities - ${}^{32}S$ J = 2 K = 1 E = 12.4 MeV



Transition densities - ${}^{32}S$ J = 2 K = 1 E = 18.0MeV



Transition densities - ${}^{32}S$ J = 2 K = 1 E = 18.1MeV



Transition densities - ${}^{32}S$ J = 2 K = 2 E = 20.2MeV



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