

Ab Initio Nuclear Structure Theory

with

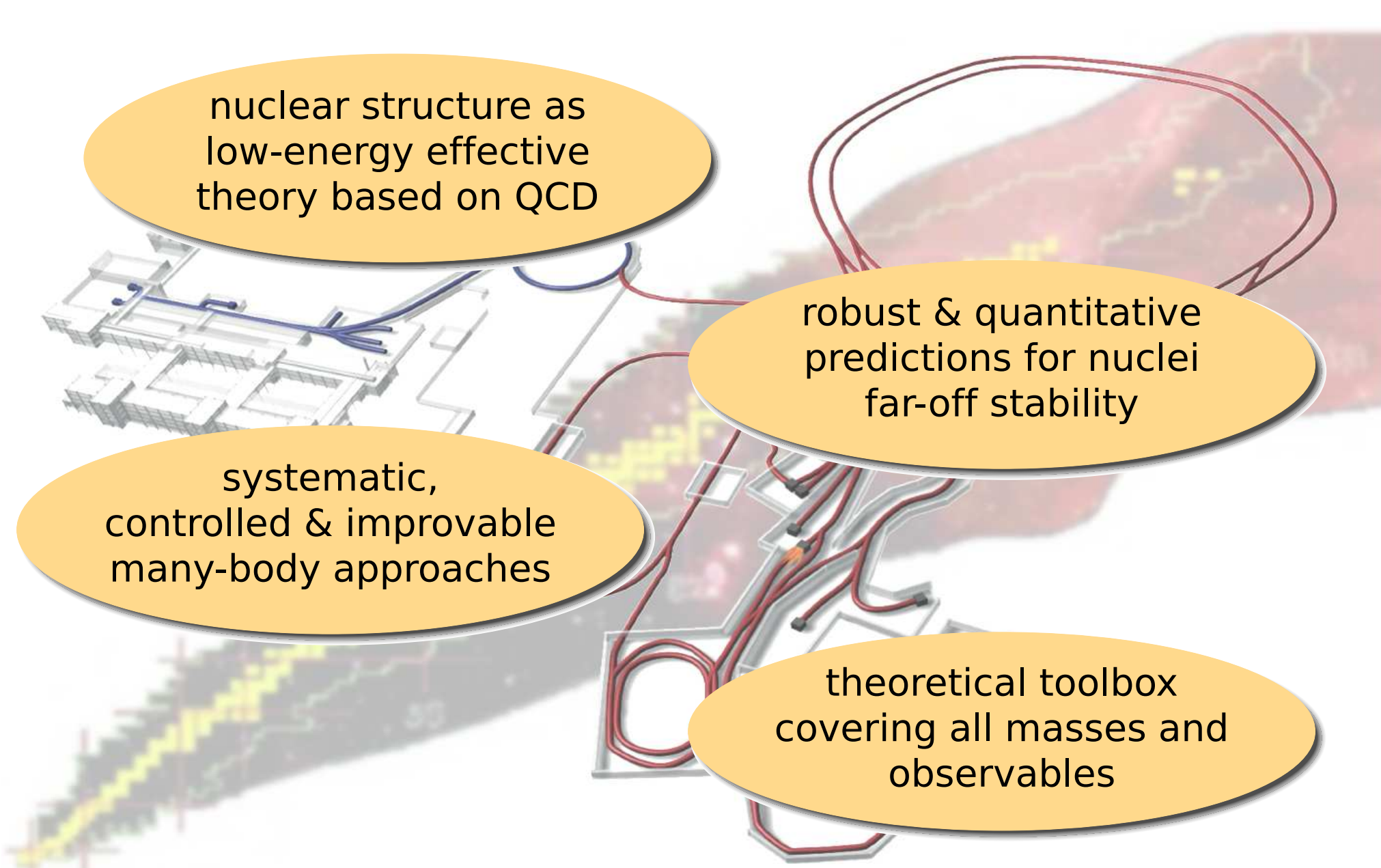
Chiral NN plus 3N Interactions

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Institut für Kernphysik



TECHNISCHE
UNIVERSITÄT
DARMSTADT

Nuclear Theory — Wish List



nuclear structure as
low-energy effective
theory based on QCD

robust & quantitative
predictions for nuclei
far-off stability

systematic,
controlled & improvable
many-body approaches

theoretical toolbox
covering all masses and
observables

Ab Initio Building Blocks

Nuclear Structure Observables

NEW
Nuclear Lattice Simulations
chiral EFT on lattice

IMPROVED
Exact Ab-Initio Solutions
few-body et al.

IMPROVED
Direct Ab-Initio Solutions
few-body, no-core shell model, etc.

IMPROVED
Approx. Many-Body Methods
controlled & improvable schemes

NEW
Similarity Transformations
physics-conserving transform. of observables

NEW
Chiral Interactions
consistent & improvable NN, 3N,... interactions

NEW
Chiral Effective Field Theory
systematic low-energy effective theory of QCD

IMPROVED
Energy-Density Functional Theor
guided by chiral EFT

Quantum Chromodynamics

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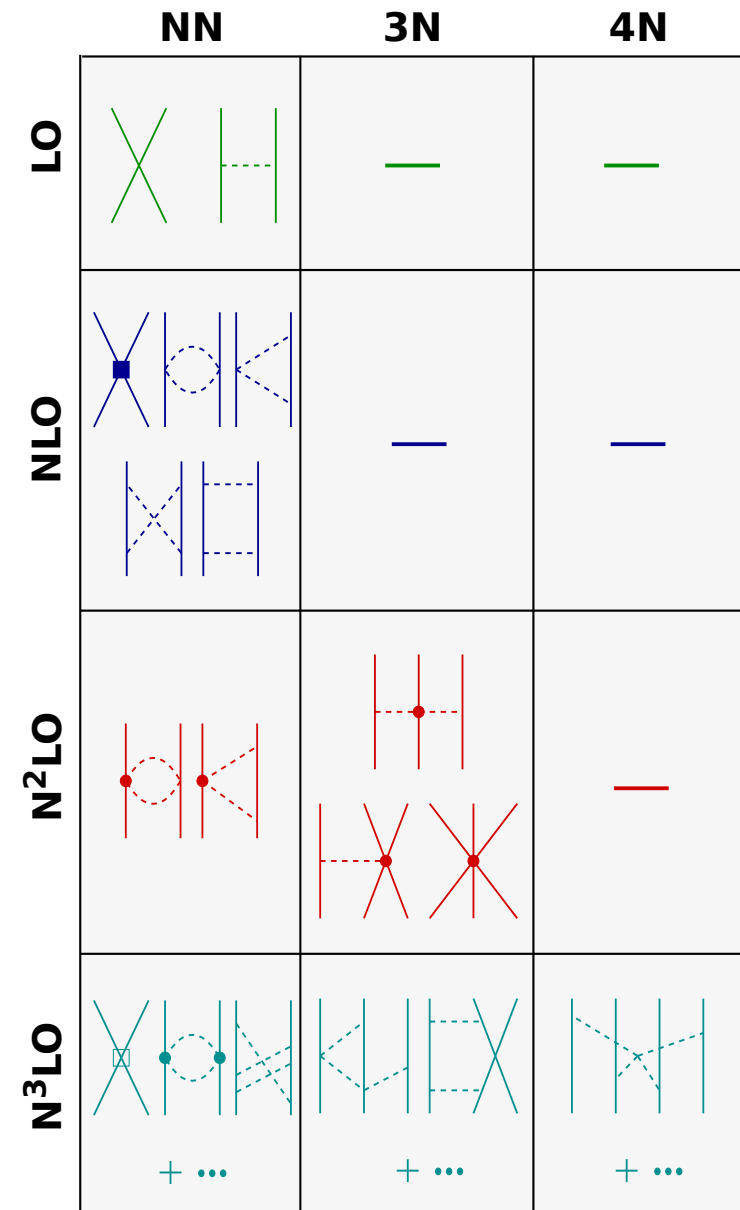
Chiral Effective Field Theory

systematic low-energy effective theory of QCD

Quantum Chromodynamics

Chiral EFT for Nuclear Interactions

- low-energy **effective field theory** for relevant degrees of freedom (π, N) based on symmetries of QCD
- long-range **pion dynamics** explicitly
- short-range physics absorbed in **contact terms**, low-energy constants fitted to experiment ($NN, \pi N, \dots$)
- hierarchy of **consistent NN, 3N, ... interactions** (plus currents)
 - 3N interaction at $N^3\text{LO}$
 - explicit inclusion of Δ -resonance
 - formal issues: power counting, renormalization, cutoff choice, ...



adapted from Meißner (2005)

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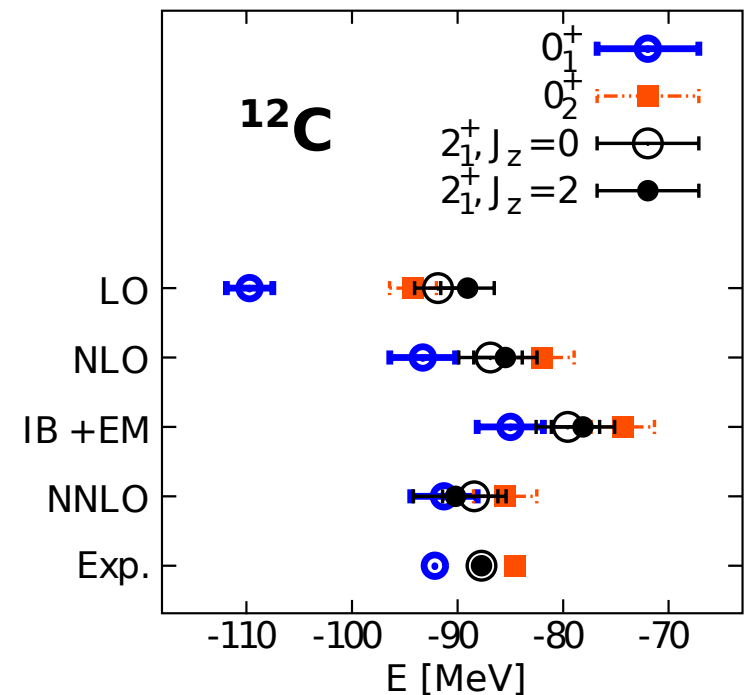
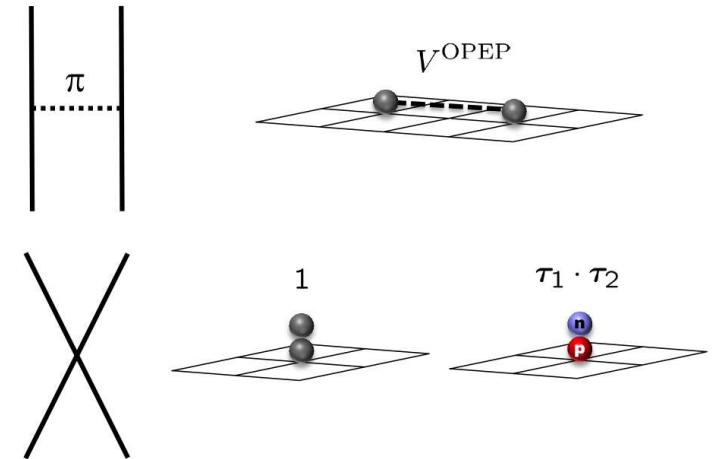
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Nuclear Lattice Simulations

- put **chiral EFT on a space-time lattice** and use Lattice-QCD technology
- lattice defines IR and UV cutoffs
- fit LECs to scattering and ground-state observables on the lattice
- **Euclidean time projection** to extract ground and excited states
- ground states of ${}^4\text{He}$, ${}^8\text{Be}$, ${}^{12}\text{C}$,...
- highlight: **Hoyle state in ${}^{12}\text{C}$**
- **statistical and extrapolation errors** still too large for spectroscopy



Epelbaum, Krebs, Lee, Meißner (2011)

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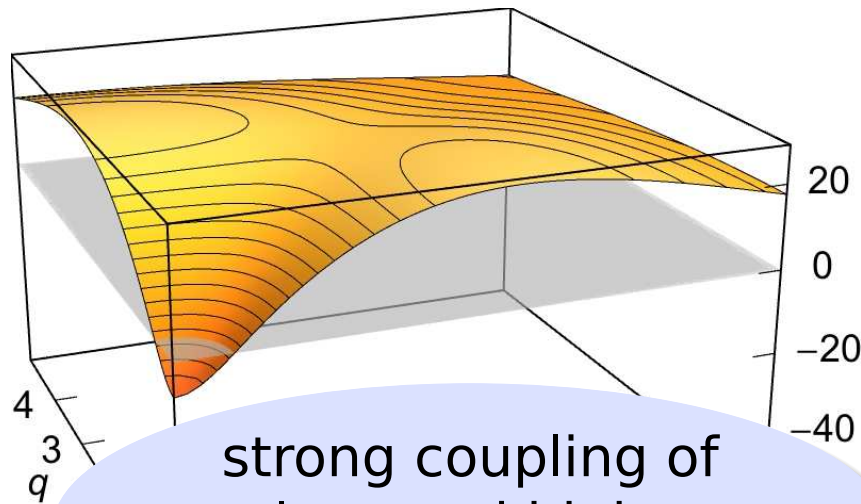
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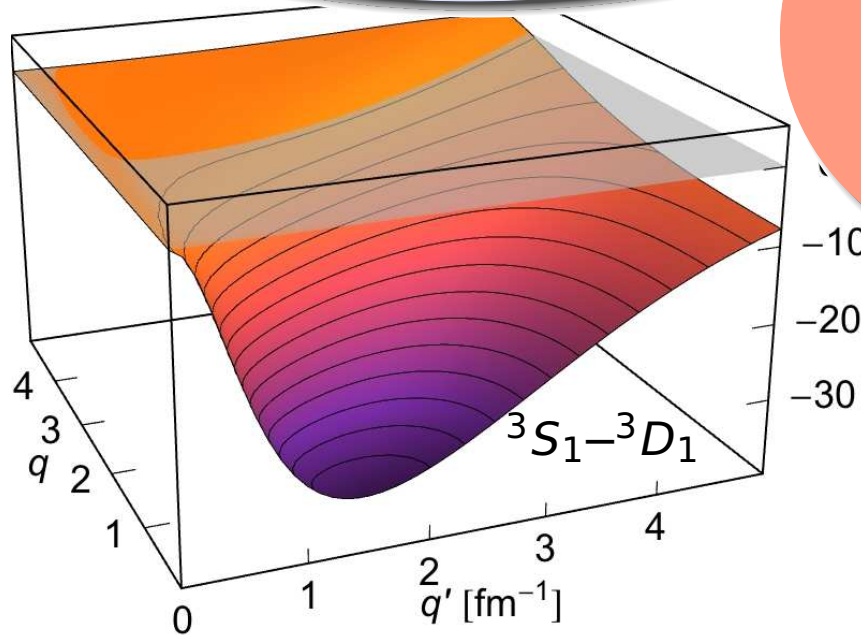
Quantum Chromodynamics

Why Similarity Transformations?

momentum-space matrix elements



strong coupling of low- and high-momentum modes



$^3S_1-^3D_1$

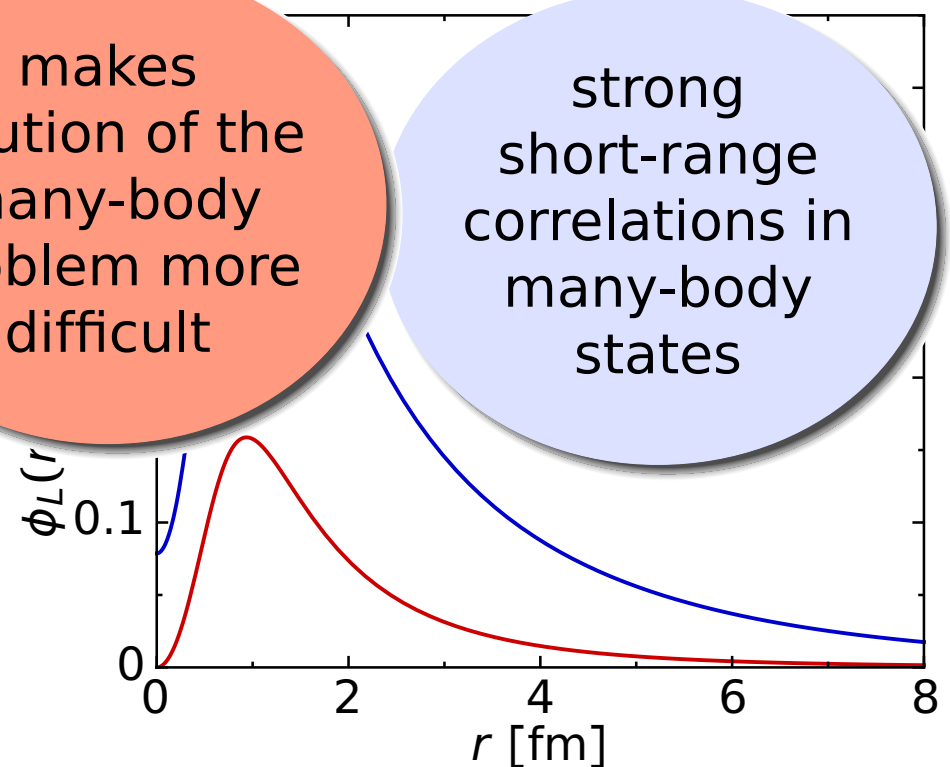
Argonne V18

$$J^\pi = 1^+, T = 0$$

deuteron wave-function

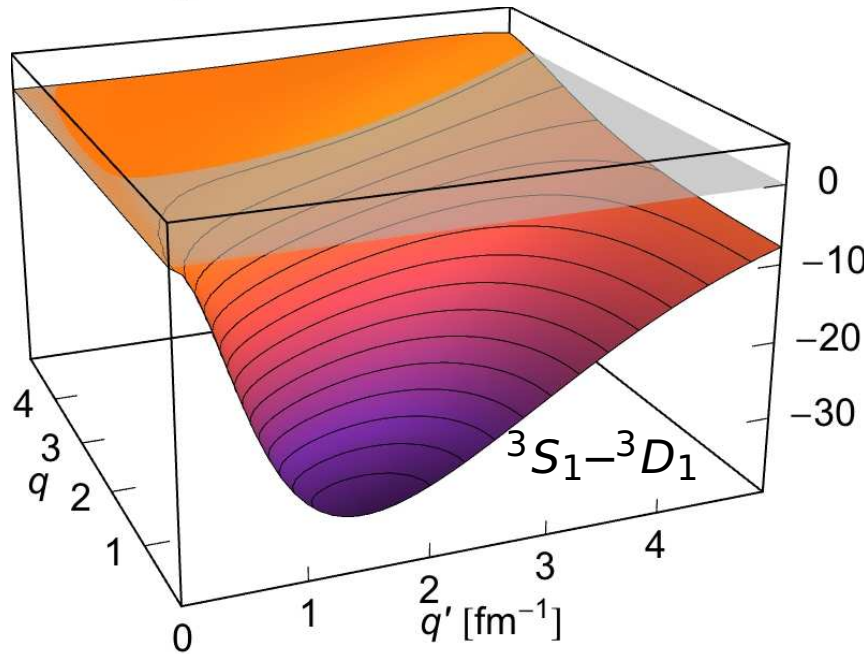
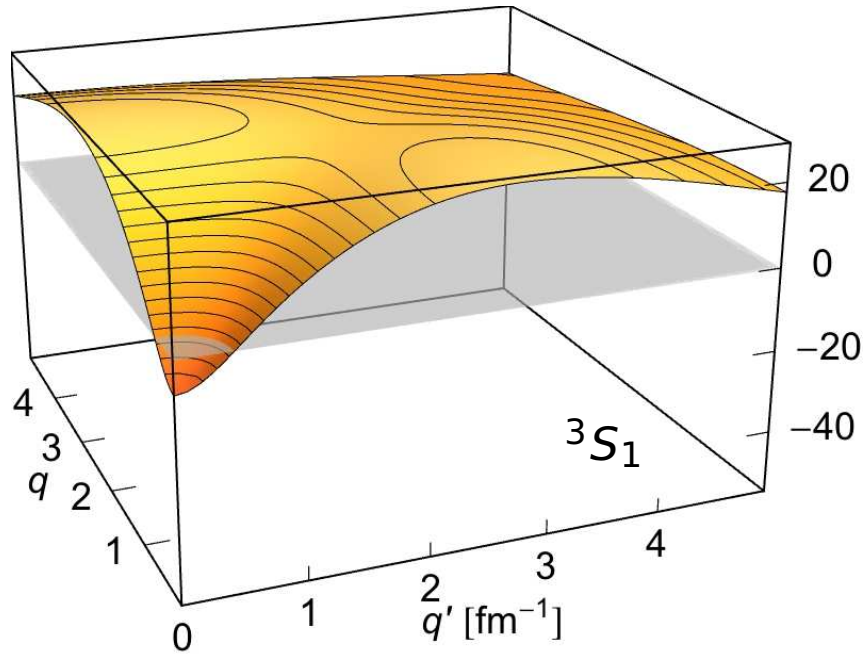
makes solution of the many-body problem more difficult

strong short-range correlations in many-body states



Why Similarity Transformations?

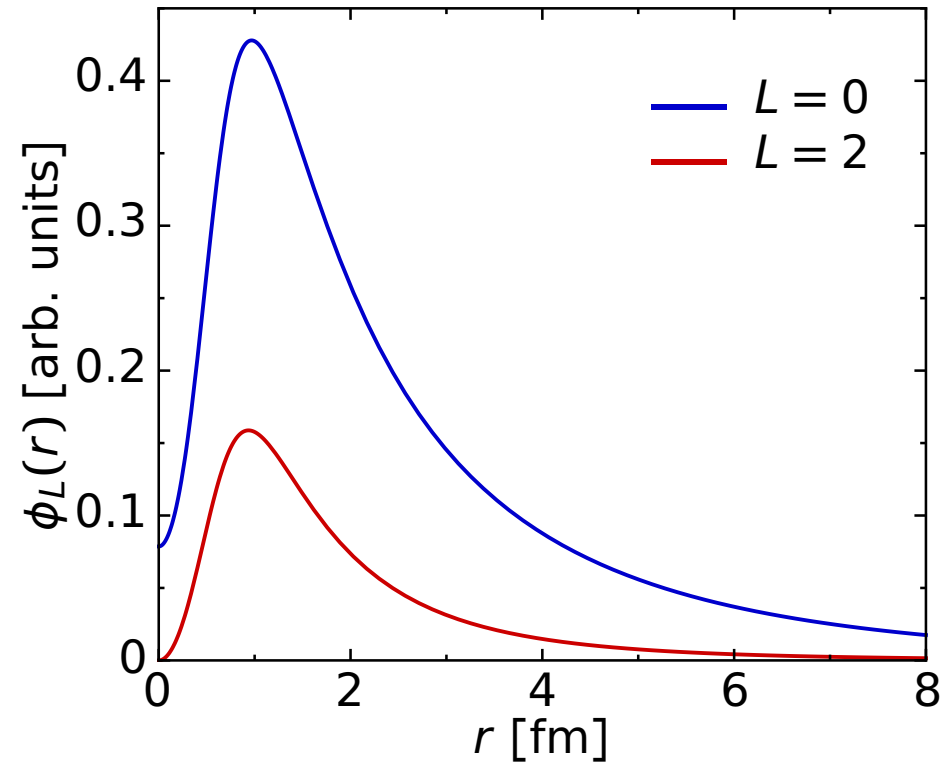
momentum-space matrix elements



Argonne V18

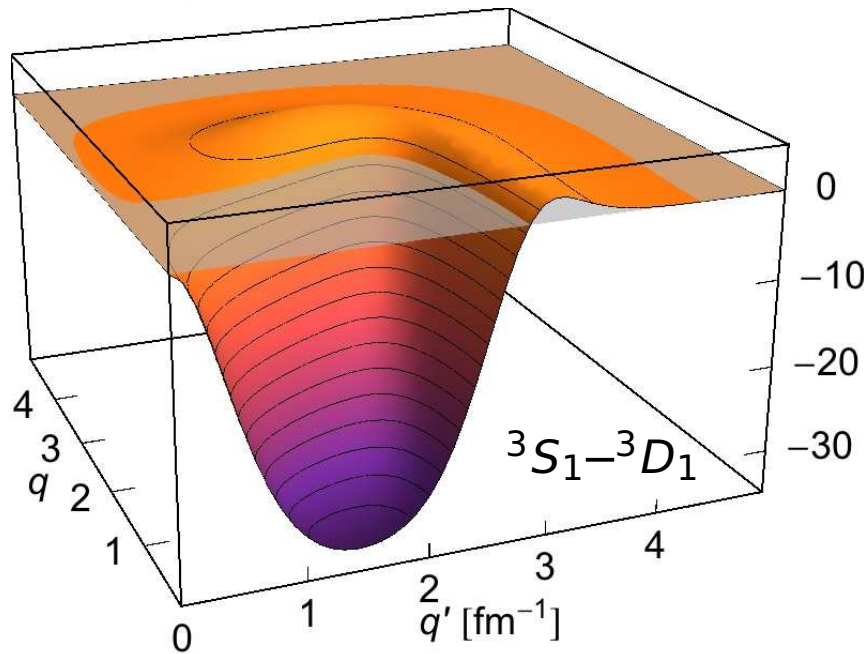
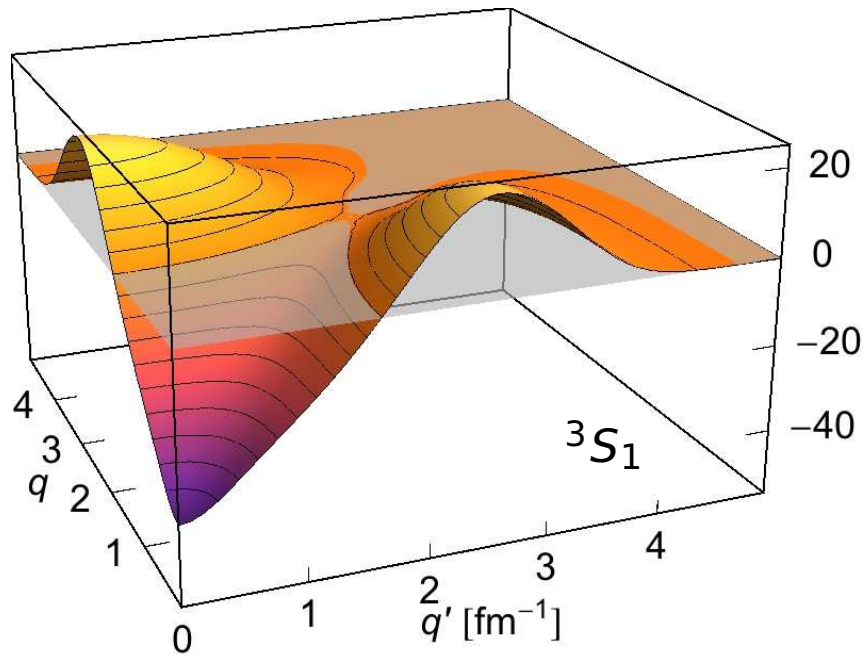
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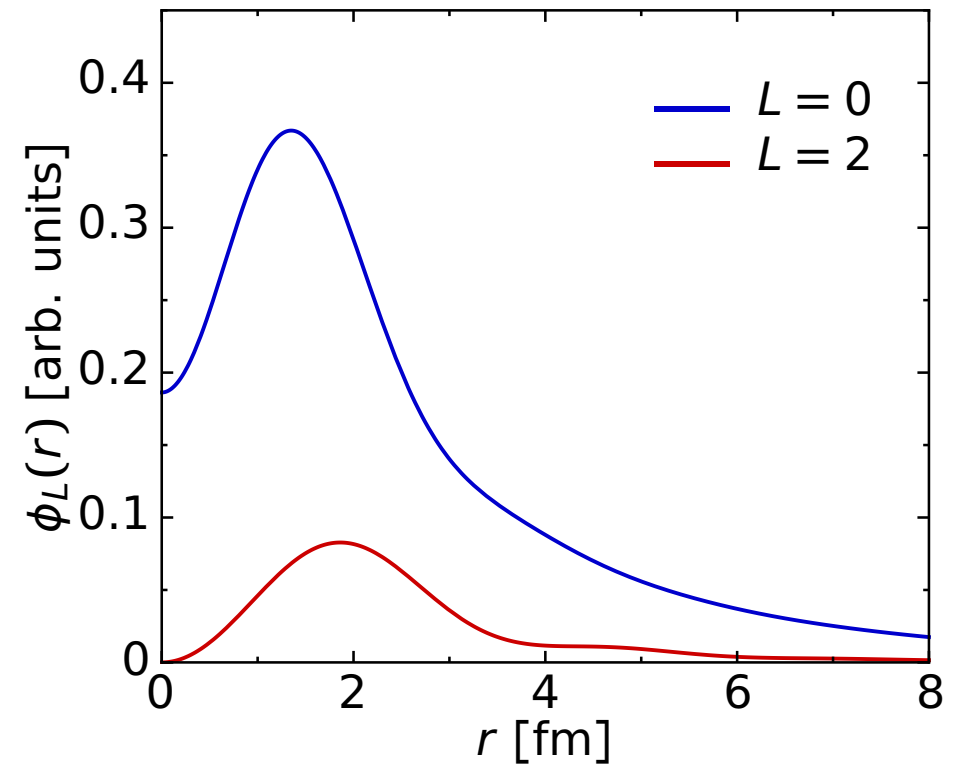


chiral $N^3\text{LO}$

Entem & Machleidt, 500 MeV

$$J^\pi = 1^+, T = 0$$

deuteron wave-function



Similarity Renormalization Group

continuous transformation driving
Hamiltonian to band-diagonal form
with respect to a chosen basis

- **unitary transformation** of Hamiltonian

$$\tilde{H}_\alpha = U_\alpha^\dagger H U_\alpha$$

simplicity and flexibility
are great advantages of
the SRG approach

- **evolution equations** for \tilde{H}_α and U_α

$$\frac{d}{d\alpha} \tilde{H}_\alpha = [\eta_\alpha, \tilde{H}_\alpha]$$

$$\frac{d}{d\alpha} U_\alpha = \dots$$

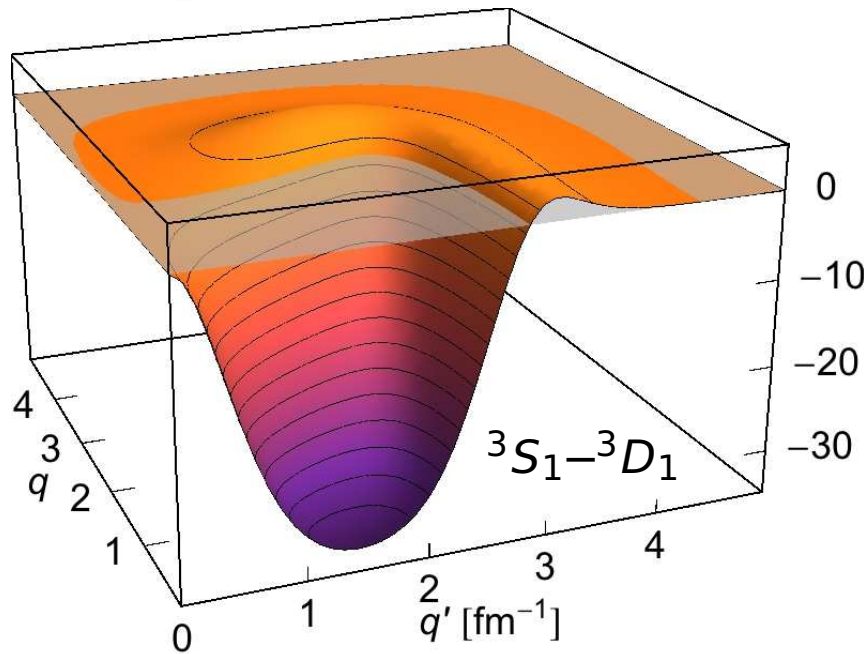
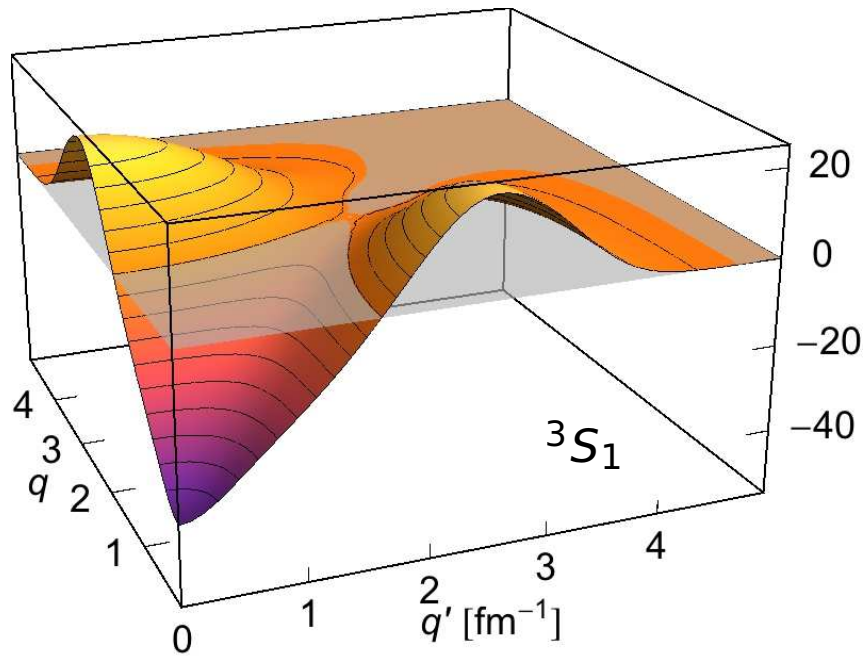
other transformation
approaches (UCOM, V_{lowk})
follow as special cases

- **dynamic generator**: commutator with the operator in whose eigenbasis H shall be diagonalized

$$\eta_\alpha = (2\mu)^2 [T_{int}, \tilde{H}_\alpha]$$

SRG Evolution in Two-Body Space

momentum-space matrix elements

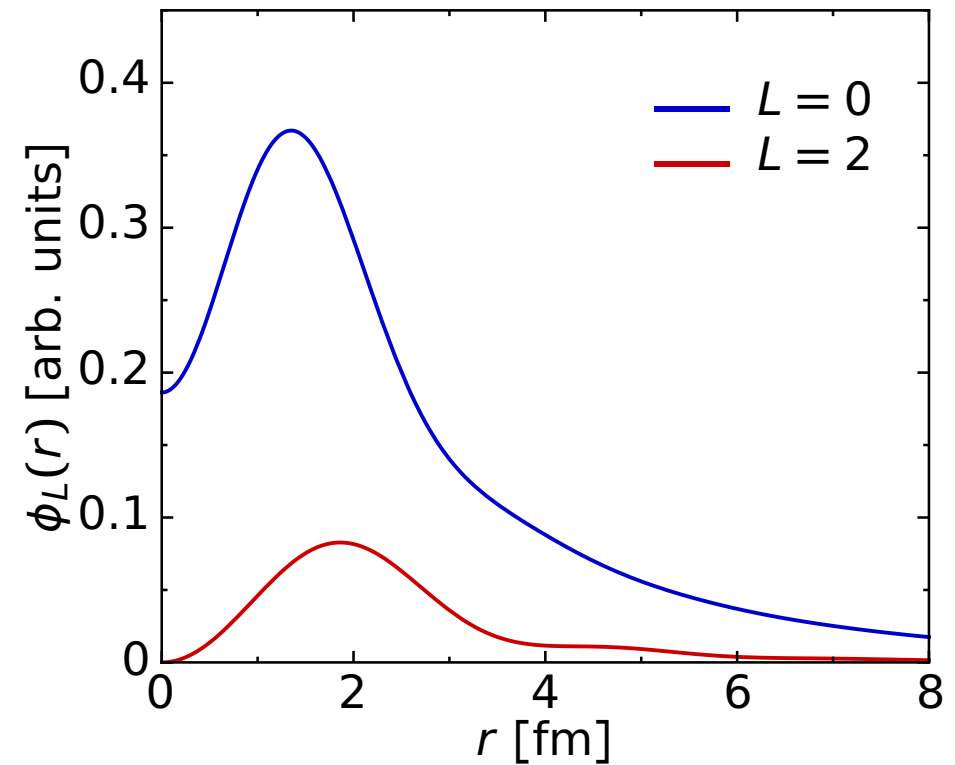


$$\alpha = 0.000 \text{ fm}^4$$

$$\Lambda = \infty \text{ fm}^{-1}$$

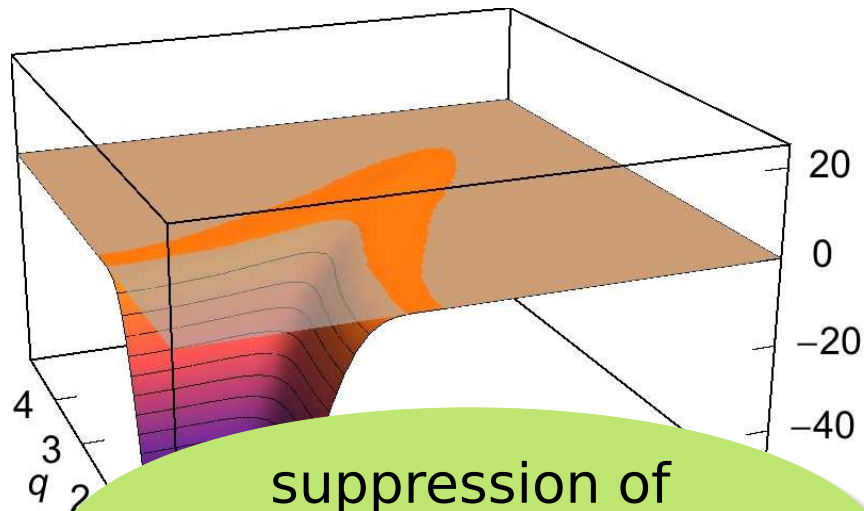
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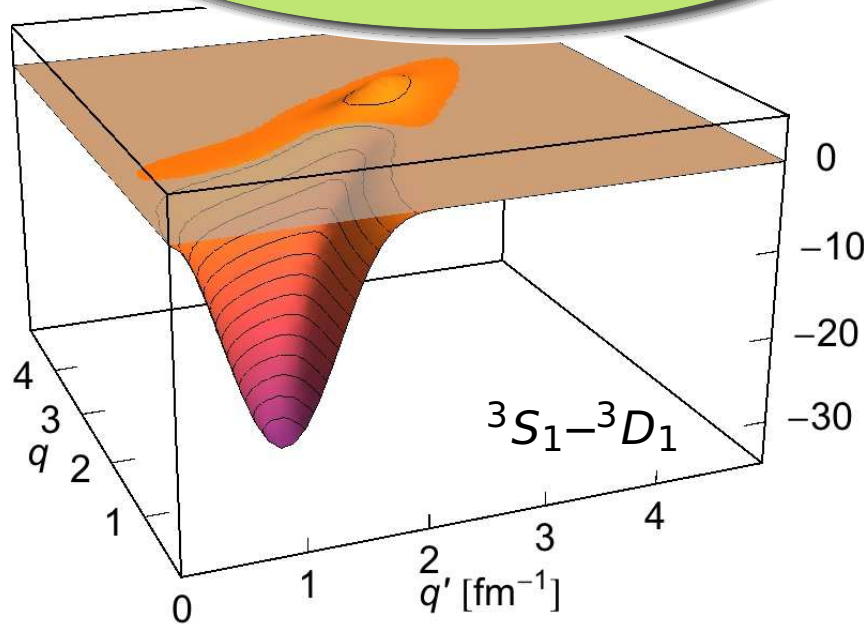


SRG Evolution in Two-Body Space

momentum-space matrix elements



suppression of
off-diagonal coupling
 $\hat{=}$ pre-diagonalization



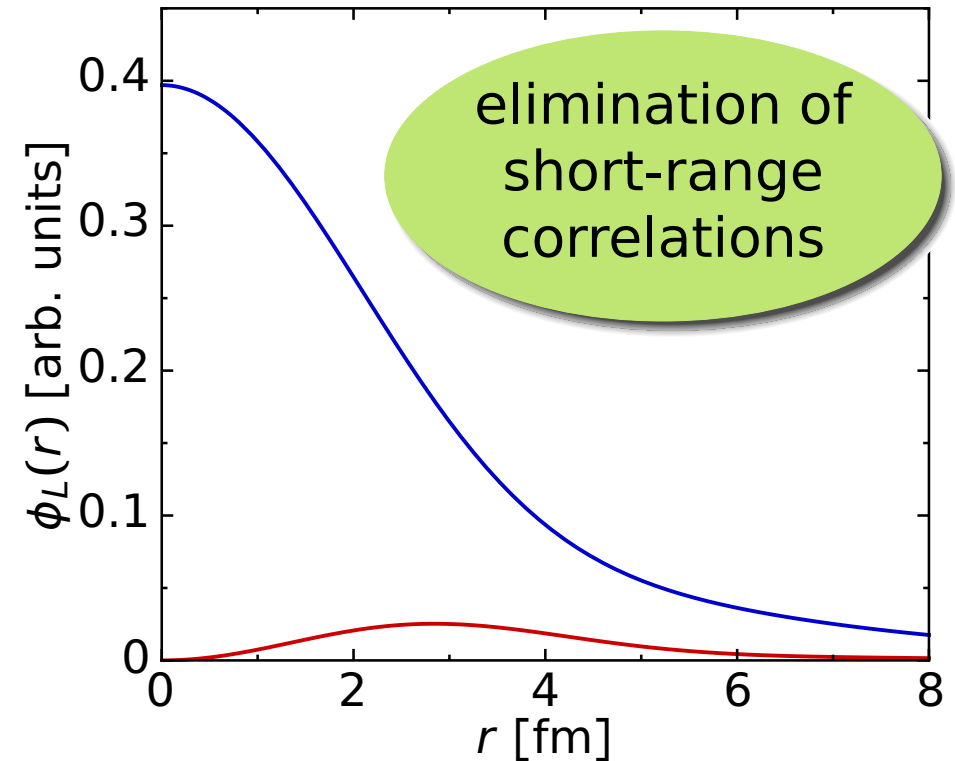
${}^3S_1-{}^3D_1$

$$\alpha = 0.320 \text{ fm}^4$$

$$\Lambda = 1.33 \text{ fm}^{-1}$$

$$J^\pi = 1^+, T = 0$$

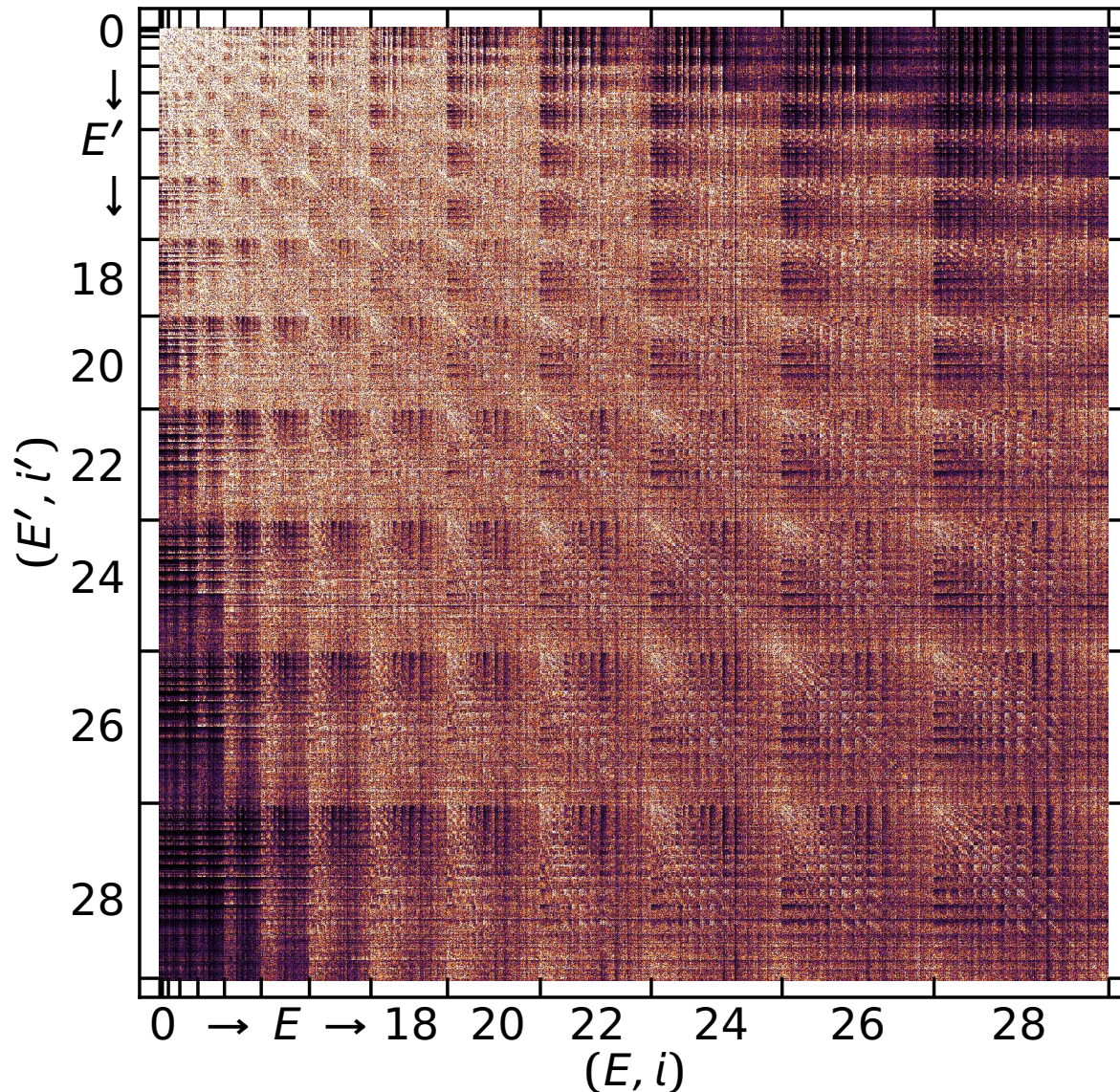
deuteron wave-function



elimination of
short-range
correlations

SRG Evolution in Three-Body Space

3B-Jacobi HO matrix elements

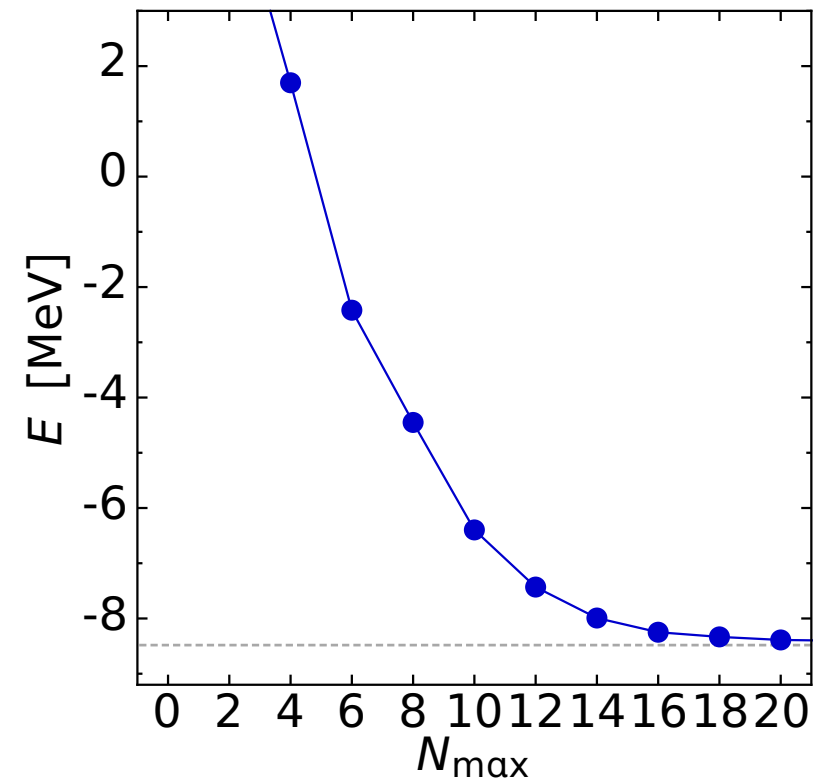


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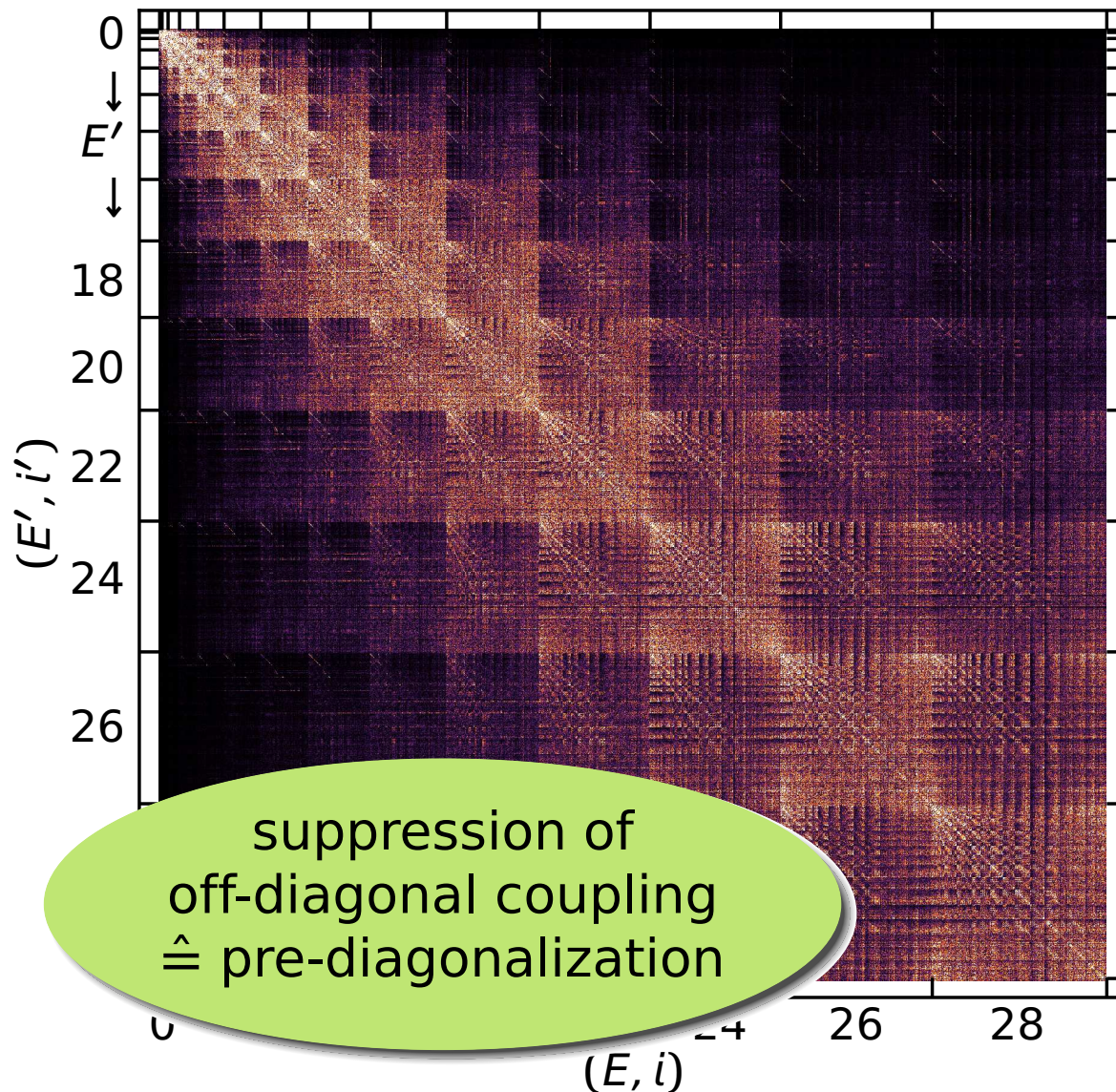
$$J^\pi = \frac{1}{2}^+, T = \frac{1}{2}, \hbar\Omega = 28 \text{ MeV}$$

NCSM ground state ${}^3\text{H}$



SRG Evolution in Three-Body Space

3B-Jacobi HO matrix elements

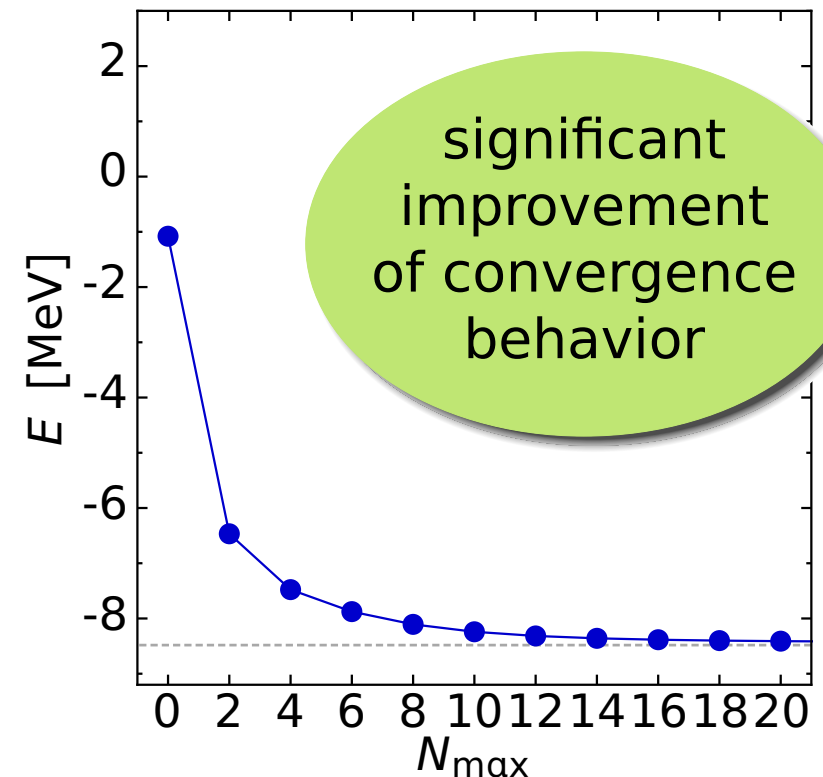


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Hamiltonian in A-Body Space

- unitary transformation **induces n -body contributions** $\tilde{H}_\alpha^{[n]}$ in the Hamiltonian

$$\tilde{H}_\alpha = \tilde{H}_\alpha^{[1]} + \tilde{H}_\alpha^{[2]} + \tilde{H}_\alpha^{[3]} + \dots + \tilde{H}_\alpha^{[A]}$$

- **cluster truncation**: can construct contributions up to $n = 3$ from evolution in 2B and 3B space, but have to discard $n > 3$
 - only the **full evolution in A-body space** conserves A-body energy eigenvalues and, thus, independence of α
 - α -dependence of eigenvalues **Hamiltonian** measures impact of

α -variation provides a **diagnostic tool** to assess the contributions of omitted many-body interactions

Sounds easy, but...

- ❶ computation of initial 2B/3B-Jacobi HO matrix elements of chiral NN+3N interactions
 - we use Petr Navratil's ManyEff code for computing 3B-Jacobi matrix elements and corresponding CFPs
- ❷ SRG evolution in 2B/3B space and cluster decomposition
 - efficient implementation using adaptive ODE solver; largest block takes a few hours on single node
- ❸ transformation of 2B/3B Jacobi HO matrix elements into JT-coupled representation
 - formulated transformation directly into JT-coupled scheme; highly efficient implementation; can handle $E_{3\max} = 16$ in JT-coupled scheme
- ❹ data management and on-the-fly decoupling in many-body codes
 - invented optimized storage scheme for fast on-the-fly decoupling; can keep all matrix elements up to $E_{3\max} = 16$ in memory

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No-Core Shell Model (NCSM)

NCSM is one of the most powerful and universal exact ab-initio methods

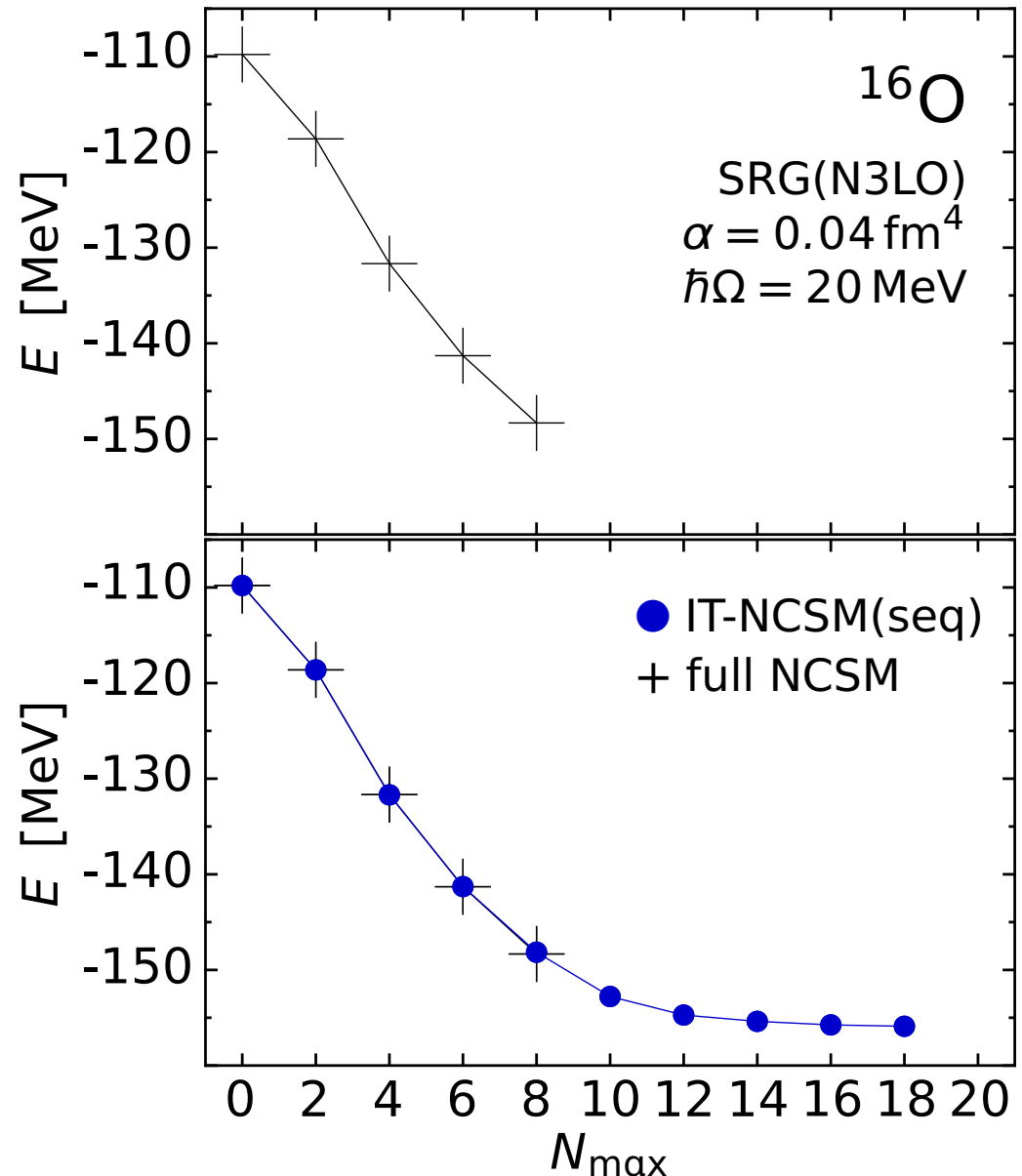
- construct matrix representation of Hamiltonian using a **basis of HO Slater determinants** truncated w.r.t. HO excitation energy $N_{\max}\hbar\Omega$
- solve **large-scale eigenvalue problem** for a few extremal eigenvalues
- **all relevant observables** can be computed from the eigenstates
- range of applicability limited by **factorial growth** of basis with N_{\max} & A
- adaptive **importance truncation** extends the range of NCSM by reducing the model space to physically relevant states
- we have developed a **parallelized IT-NCSM/NCSM code** capable of handling $3N$ matrix elements up to $E_{3\max} = 16$

Importance Truncated NCSM

- converged NCSM calculations essentially restricted to lower/mid p-shell
- full 10 or 12 $\hbar\Omega$ calculation for ^{16}O not really feasible (basis dimension $> 10^{10}$)

Importance Truncation

reduce model space to the relevant basis states using an **a priori importance measure** derived from MBPT



Importance Truncation: General Idea

- given an initial approximation $|\psi_{\text{ref}}^{(m)}\rangle$ for the **target states**
- **measure the importance** of individual basis state $|\Phi_\nu\rangle$ via first-order multiconfigurational perturbation theory

$$K_\nu^{(m)} = -\frac{\langle \Phi_\nu | H | \psi_{\text{ref}}^{(m)} \rangle}{\epsilon_\nu - \epsilon_{\text{ref}}}$$

- construct **importance truncated space** spanned by basis states with $|K_\nu^{(m)}| \geq K_{\text{min}}$ and solve eigenvalue problem

- **sequential scheme**: construct next N_{max} using previous eigenvalues

for $K_{\text{min}} \rightarrow 0$ the full NCSM model space and thus the **exact solution is recovered**

- a posteriori **threshold extrapolation** and **perturbative correction** used to recover contributions from discarded basis states

A Tale of Three Hamiltonians

Initial Hamiltonian

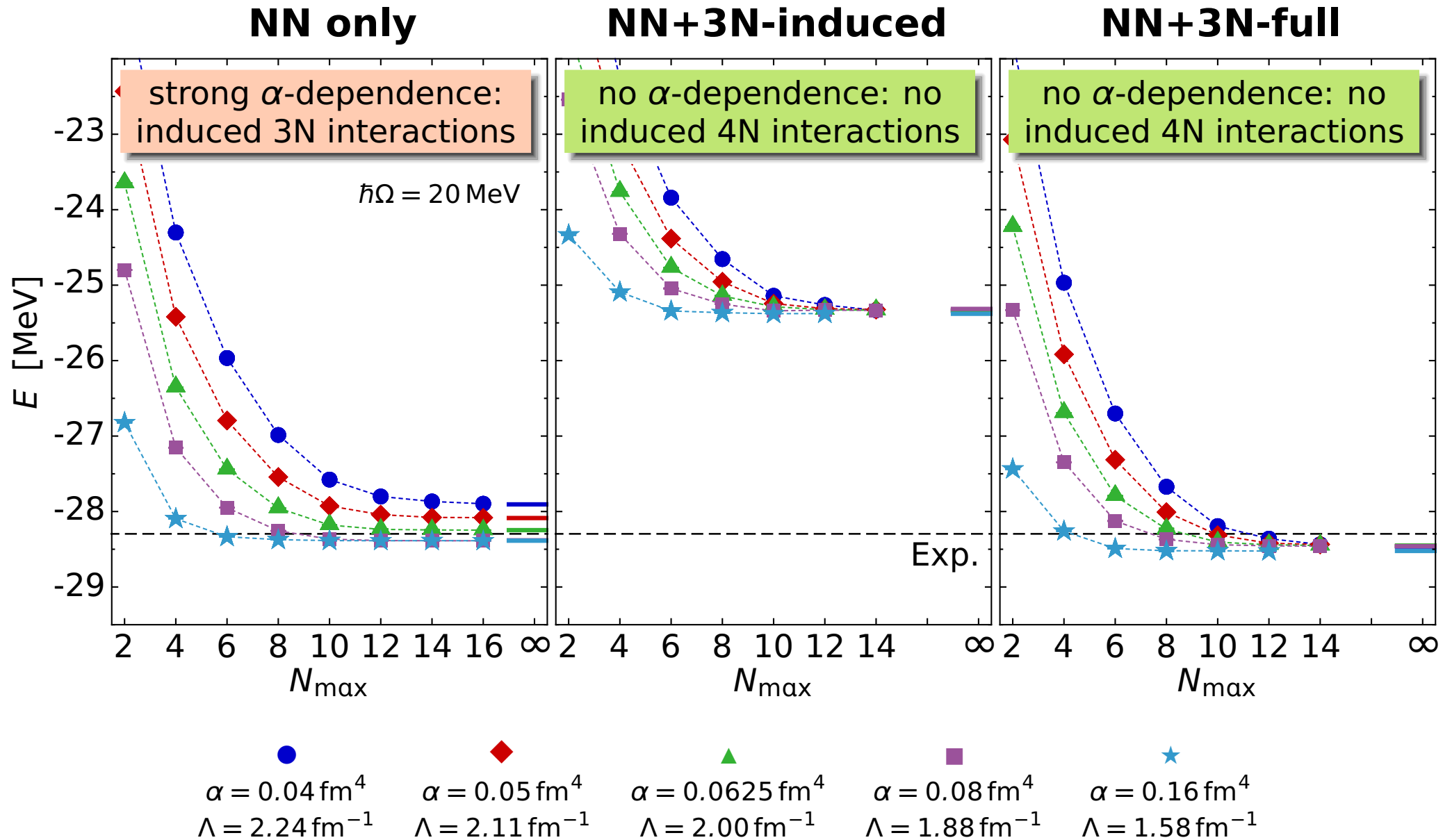
- NN: chiral interaction at $N^3\text{LO}$ (Entem & Machleidt, 500 MeV)
- 3N: chiral interaction at $N^2\text{LO}$ (c_D, c_E from ${}^3\text{H}$ binding & half-life)

SRG-Evolved Hamiltonians

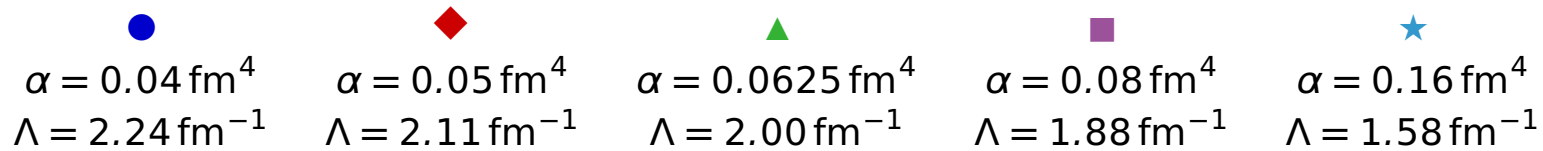
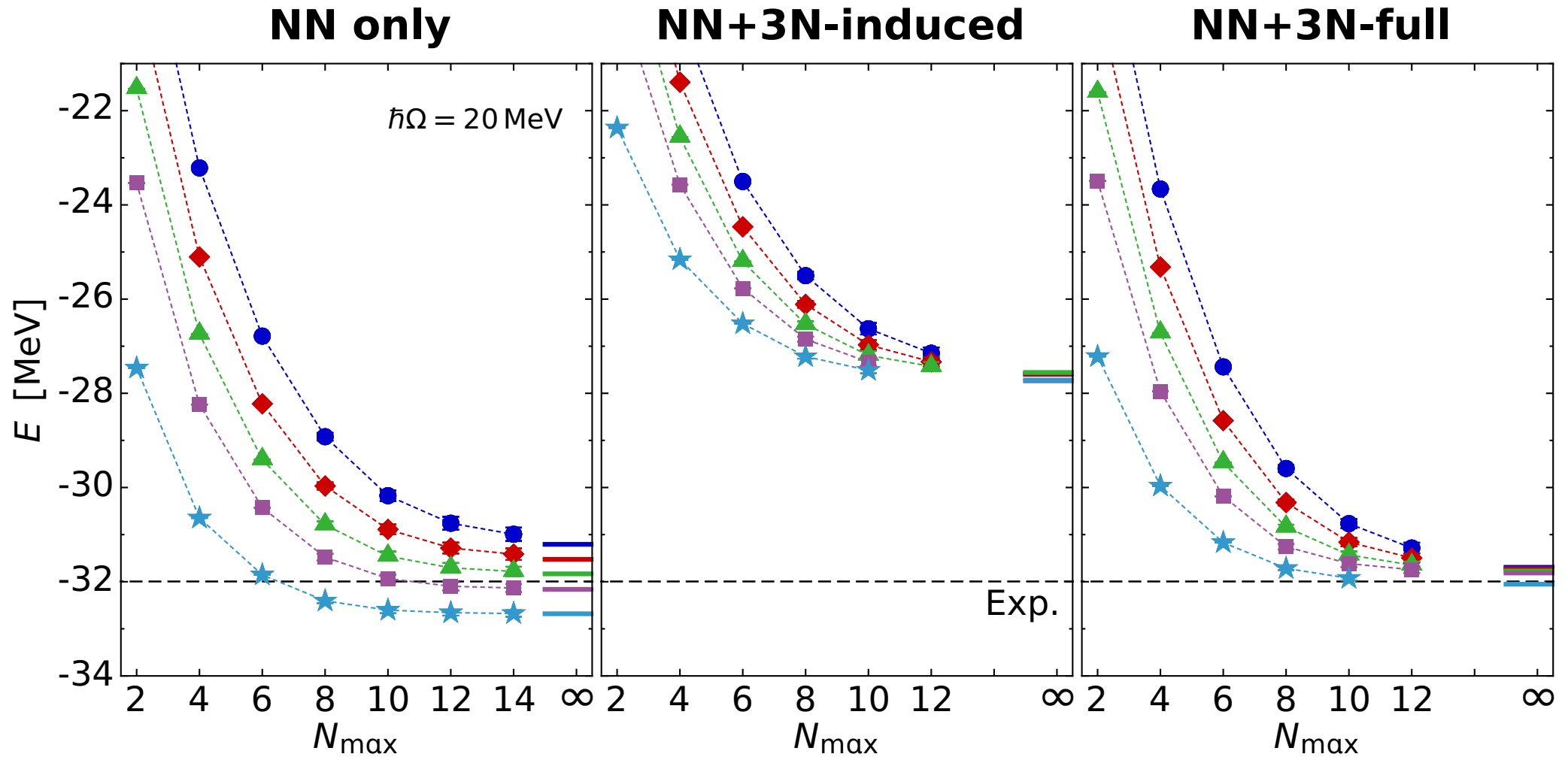
- **NN only**: start with NN initial Hamiltonian and keep two-body terms only
- **NN+3N-induced**: start with NN initial Hamiltonian and keep two- and three-body terms
- **NN+3N-full**: start with NN+3N initial Hamiltonian and keep two- and three-body terms

α -variation provides a **diagnostic tool** to assess the contributions of omitted many-body interactions

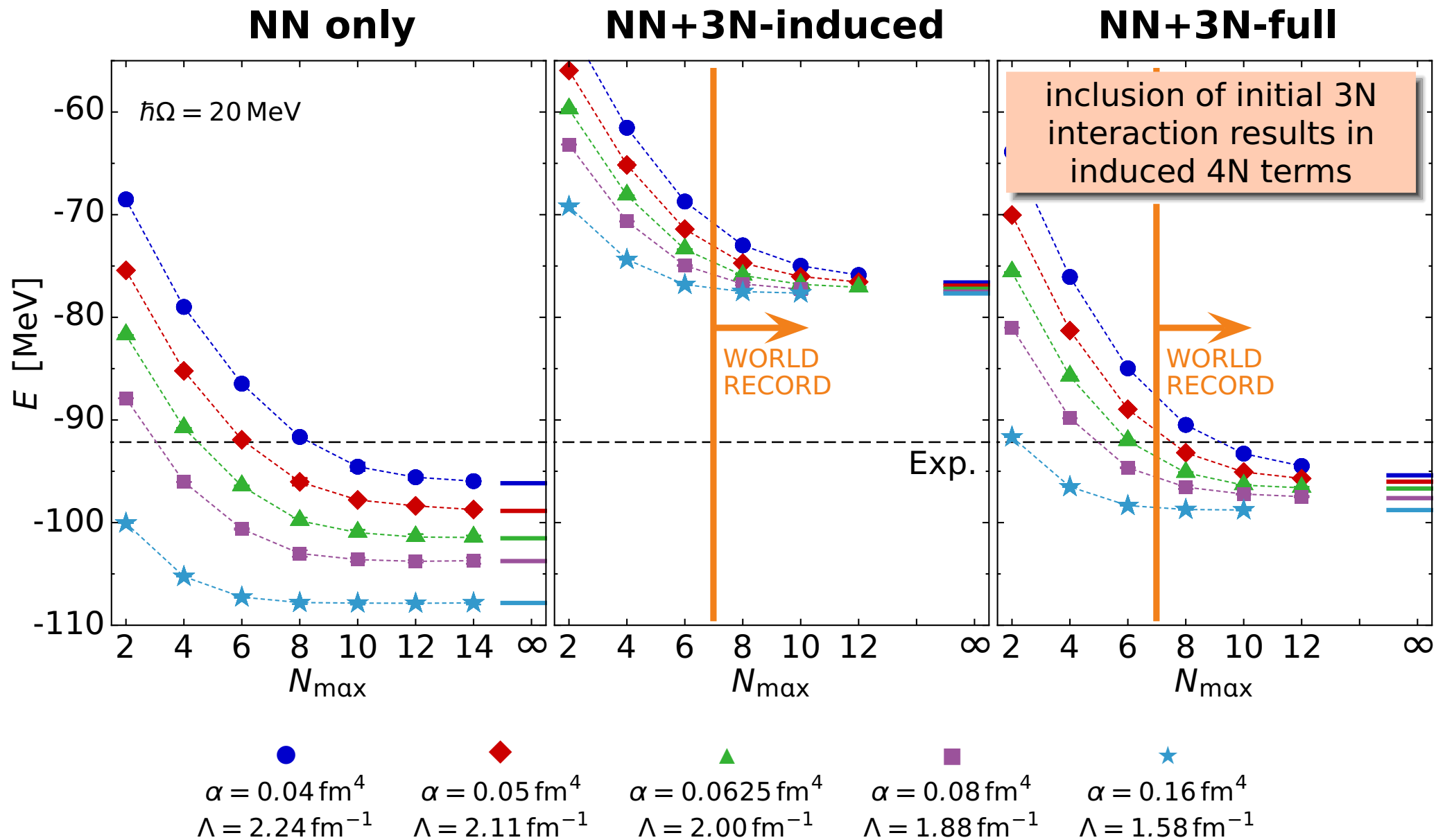
^4He : Ground-State Energies



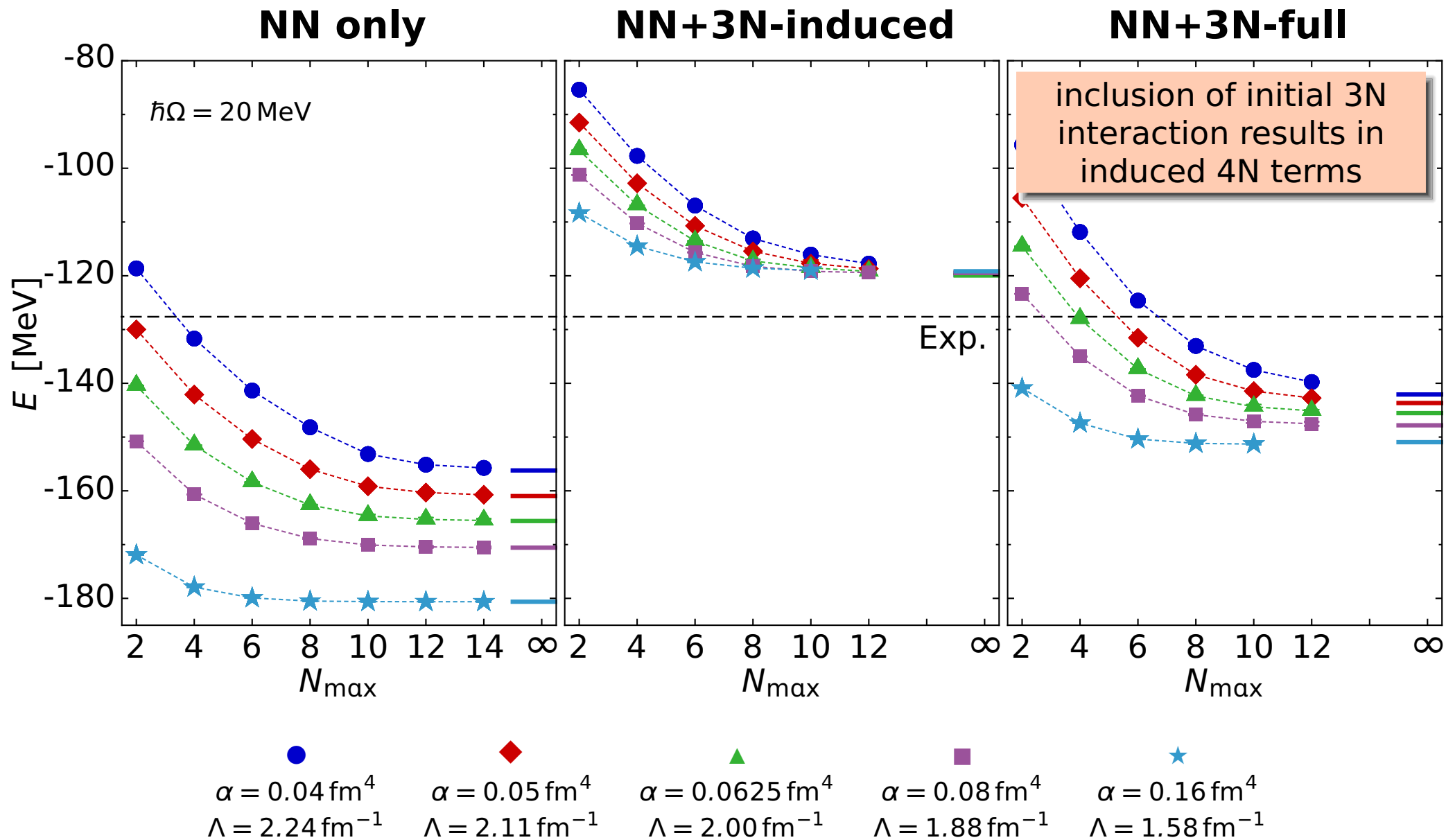
${}^6\text{Li}$: Ground-State Energies



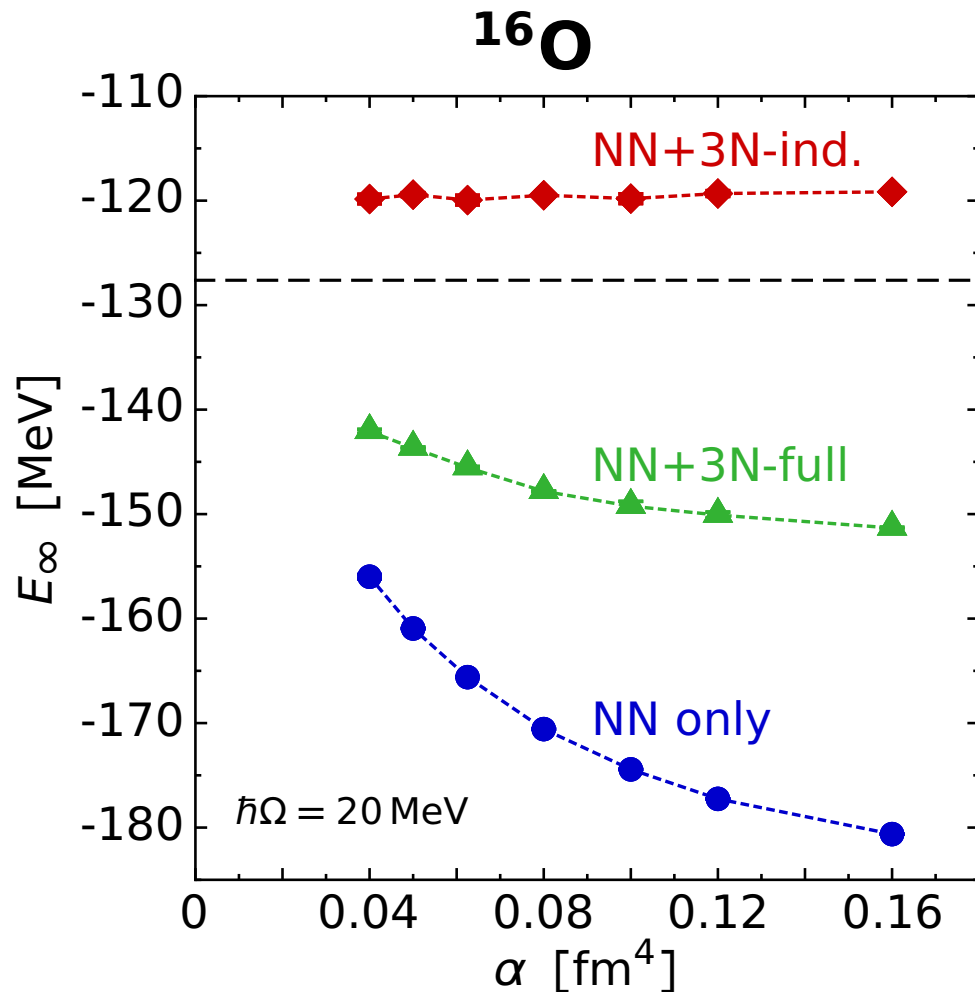
^{12}C : Ground-State Energies



^{16}O : Ground-State Energies



^{16}O : Energy vs. Flow Parameter



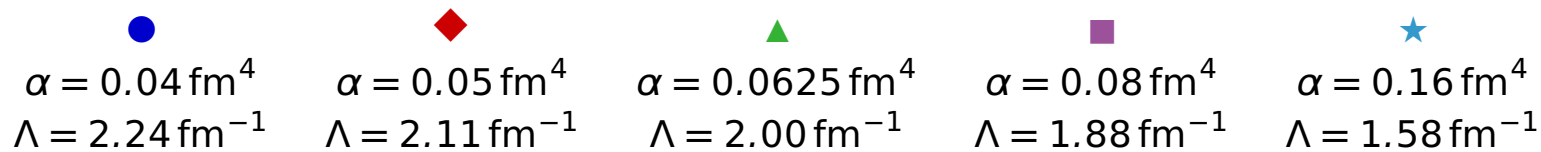
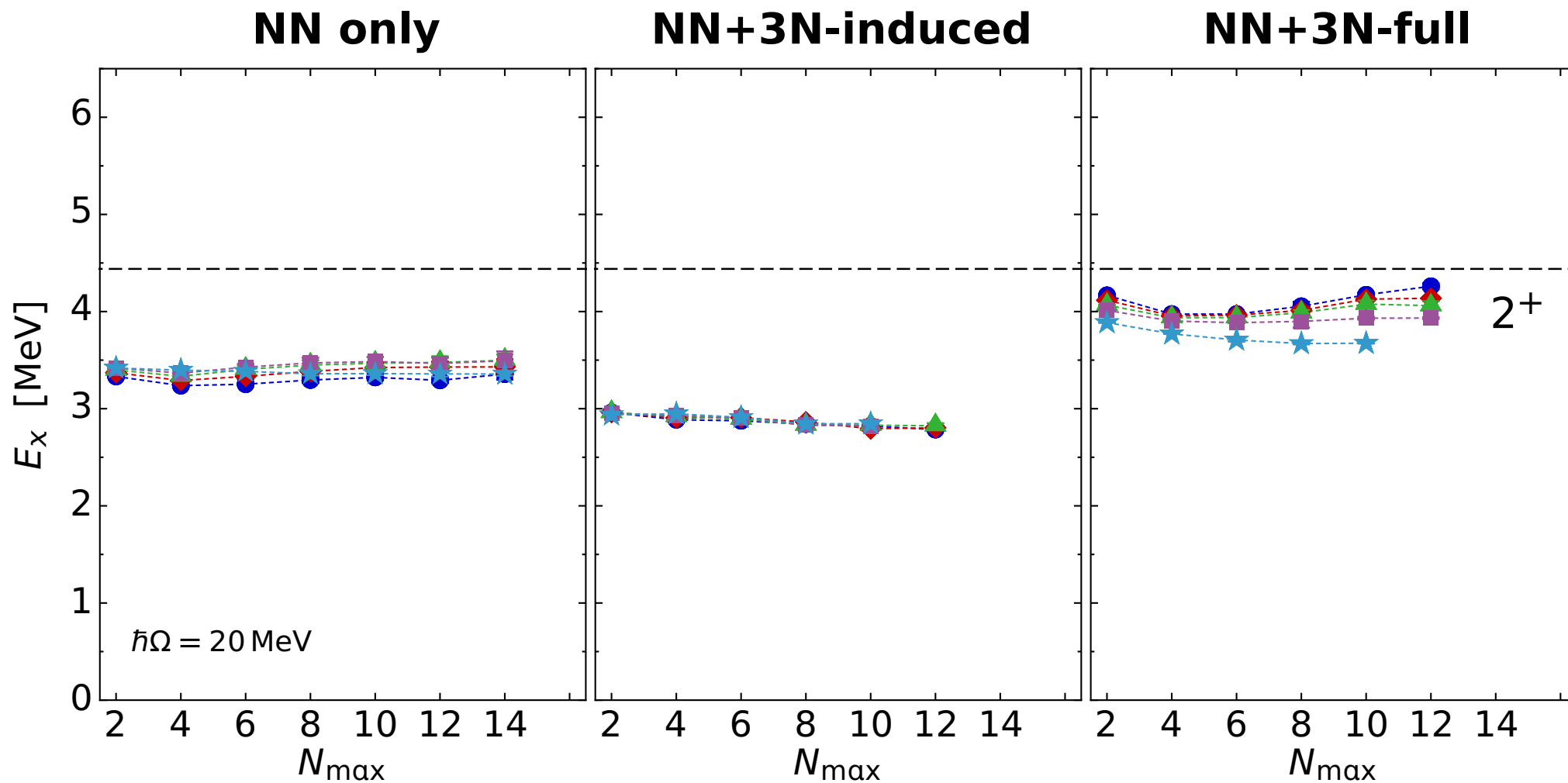
■ initial NN Hamiltonian

- induced 3N interactions are significant
- no indication of induced 4N
- NN+3N-induced unitarily equivalent to initial NN

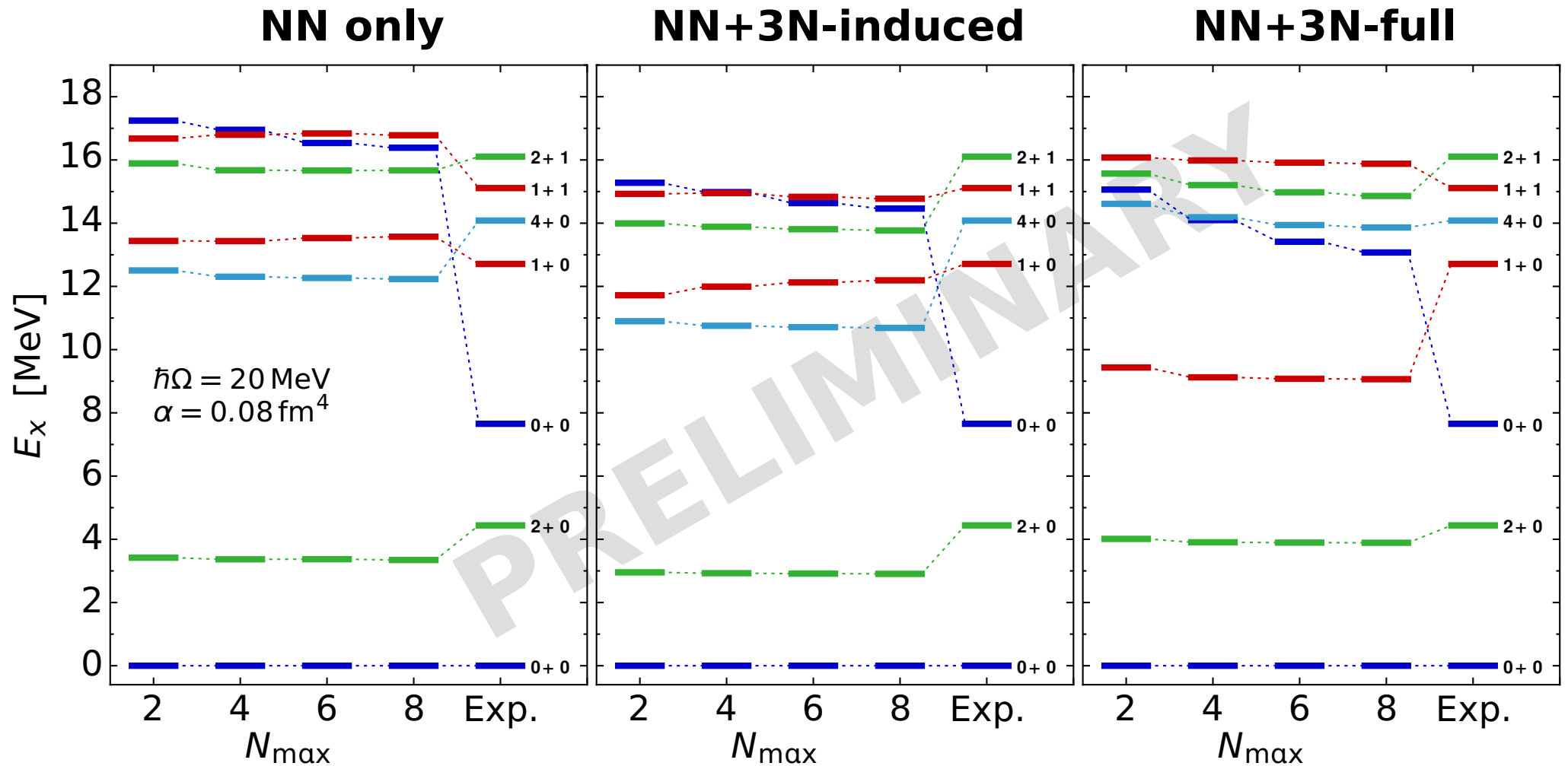
■ initial NN+3N Hamiltonian

- induced 4N interactions are sizable in upper p-shell
- generated by long-range 2π terms of initial 3N interaction
- design modified SRG generator to suppress induced 4N

$^{12}\text{C}: 2^+$ Excitation Energies

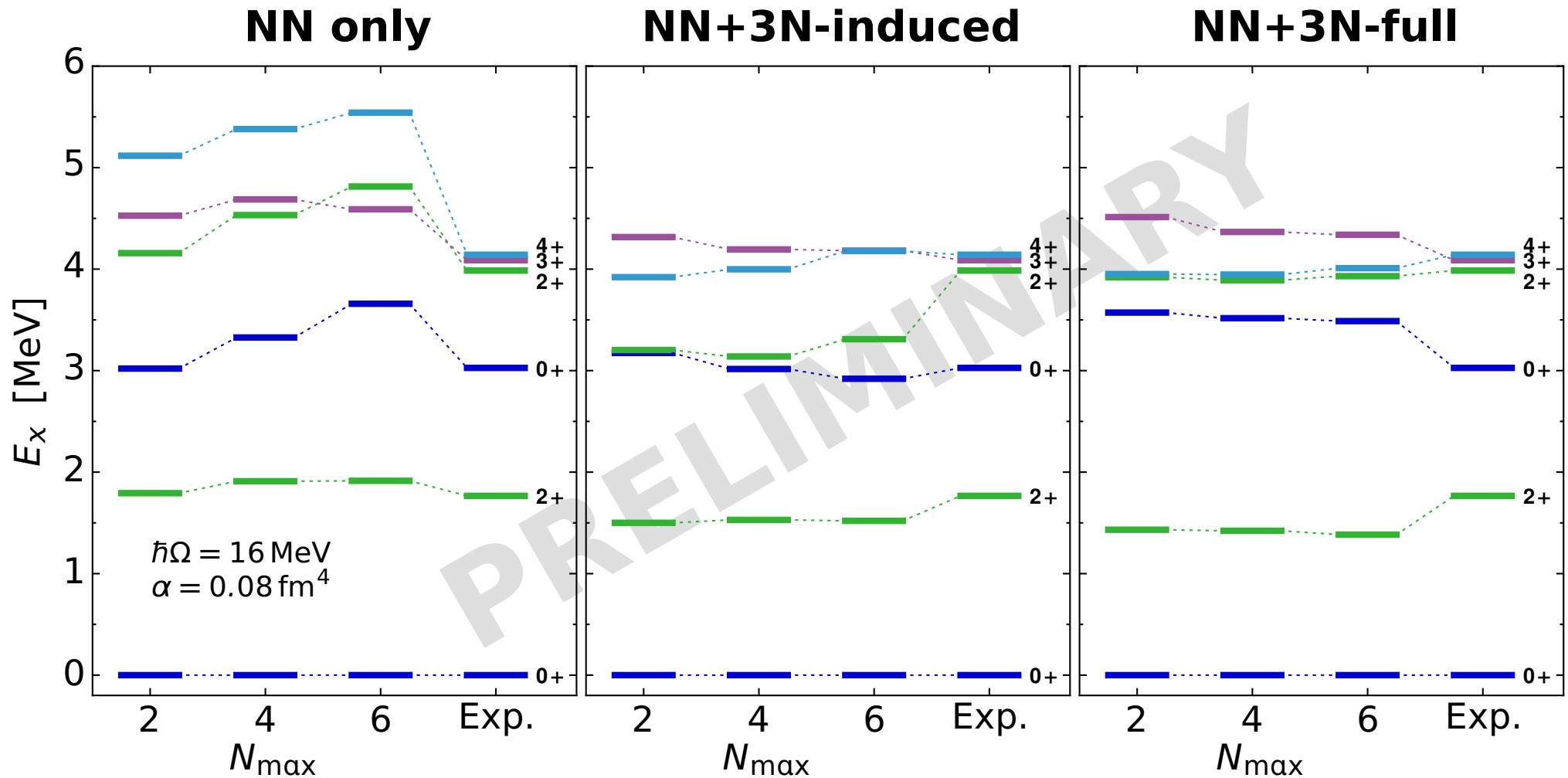


Spectroscopy of ^{12}C



- IT-NCSM gives access to **complete spectroscopy of p- and sd-shell nuclei** starting from chiral NN+3N interactions

Spectroscopy of ^{16}C

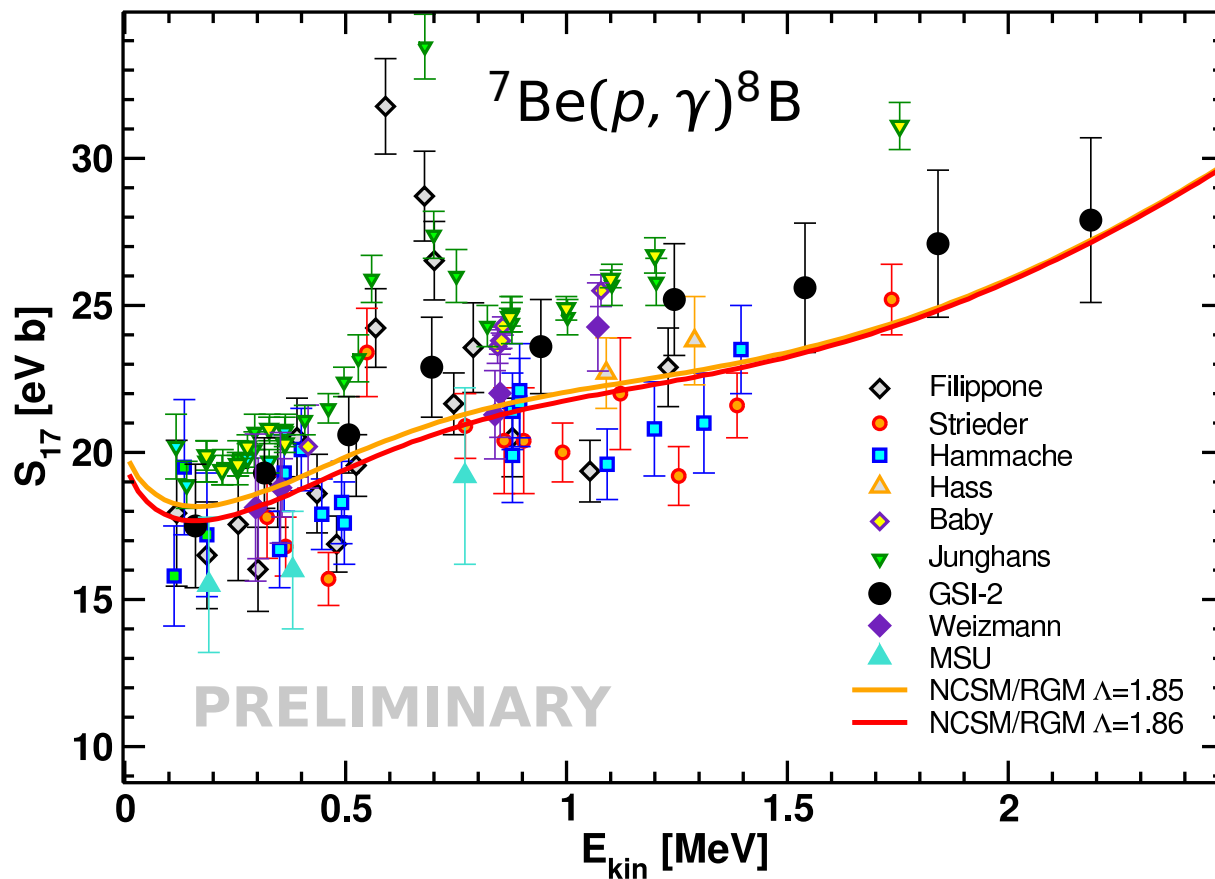


- IT-NCSM gives access to **complete spectroscopy of p- and sd-shell nuclei** starting from chiral NN+3N interactions

Bridge to Ab-Initio Reaction Theory

with P. Navrátil (TRIUMF) & S. Quaglioni (LLNL)

- **NCSM/RGM**: combine Resonating Group Method for description of relative projective-target motion with IT-NCSM for the description of target nucleus



- astrophysical S-factor for proton capture on ${}^7\text{Be}$
- IT-NCSM wave functions for ${}^7\text{Be}$ for up to 8 eigenstates
- solution of the RGM with kernels involving the full many-body information
- SRG-evolved chiral NN interaction with α adjusted to reproduce ${}^8\text{B}$ energy relative to threshold

Conclusions

- new era of **ab-initio nuclear structure and reaction theory** connected to QCD via chiral EFT
 - chiral EFT as universal starting point... some formal issues remain
- consistent **inclusion of 3N interactions** in similarity transformations & many-body calculations
 - breakthrough in computation & handling of 3N matrix elements
- **innovations in many-body theory**: extended reach of exact methods & improved control over approximations
 - versatile toolbox for different observables & mass ranges
- many **exciting applications** ahead...
- see also: HK23.4, HK23.6, HK33.3, HK41.5

Epilogue

■ thanks to my group & my collaborators

- **S. Binder, A. Calci**, B. Erler, A. Günther, M. Hild, H. Krutsch, **J. Langhammer**, P. Papakonstantinou, S. Reinhardt, F. Schmitt, C. Stumpf, K. Vobig, R. Wirth

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- H. Hergert, P. Piecuch

Michigan State University, USA

- C. Forssén

Chalmers University, Sweden

- H. Feldmeier, T. Neff,...

GSI Helmholtzzentrum



Deutsche
Forschungsgemeinschaft

DFG

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for

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LOEWE – Landes-Offensive
zur Entwicklung Wissenschaftlich-
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