

Frontiers in Nuclear Structure Theory from a FAIR Perspective

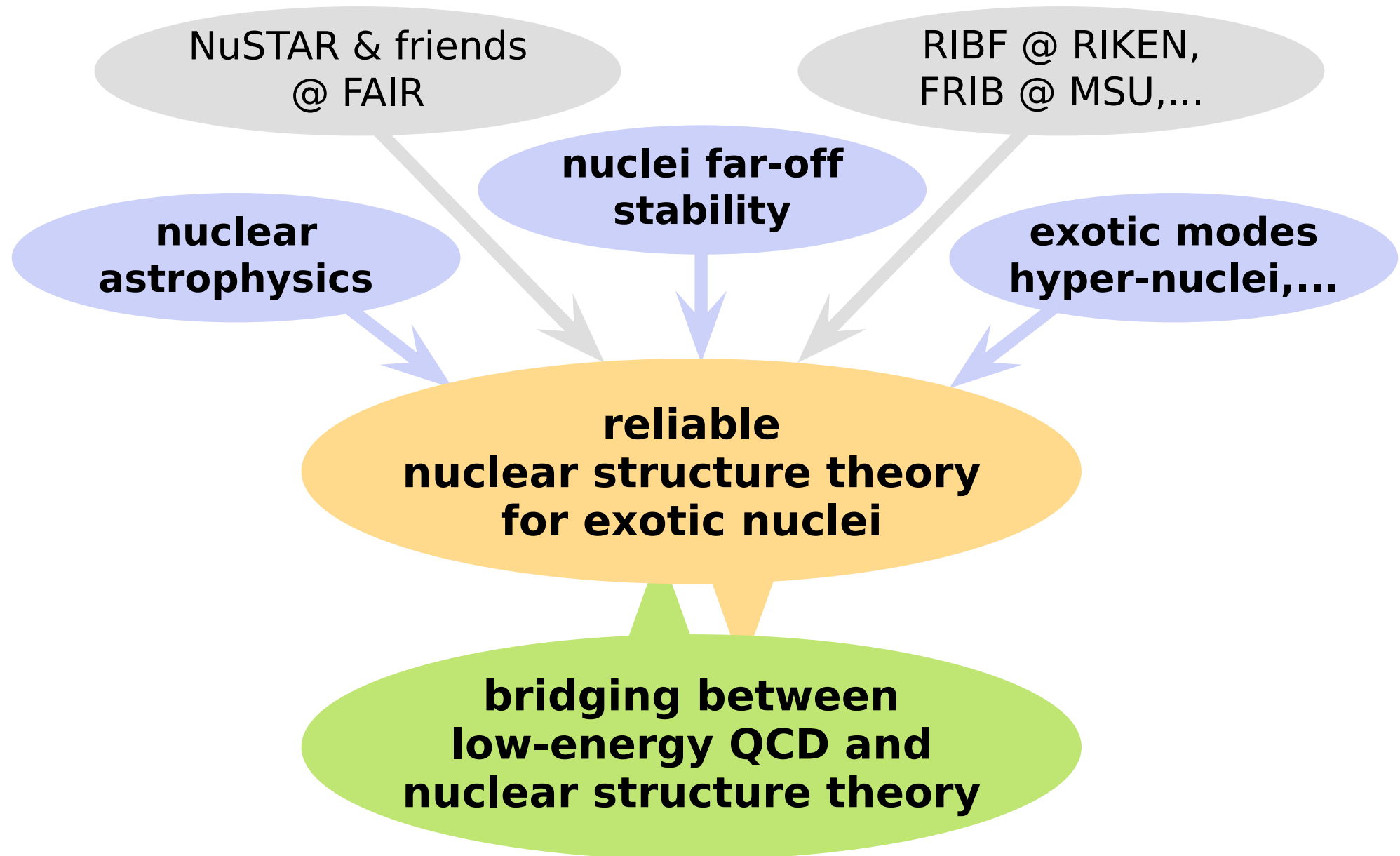
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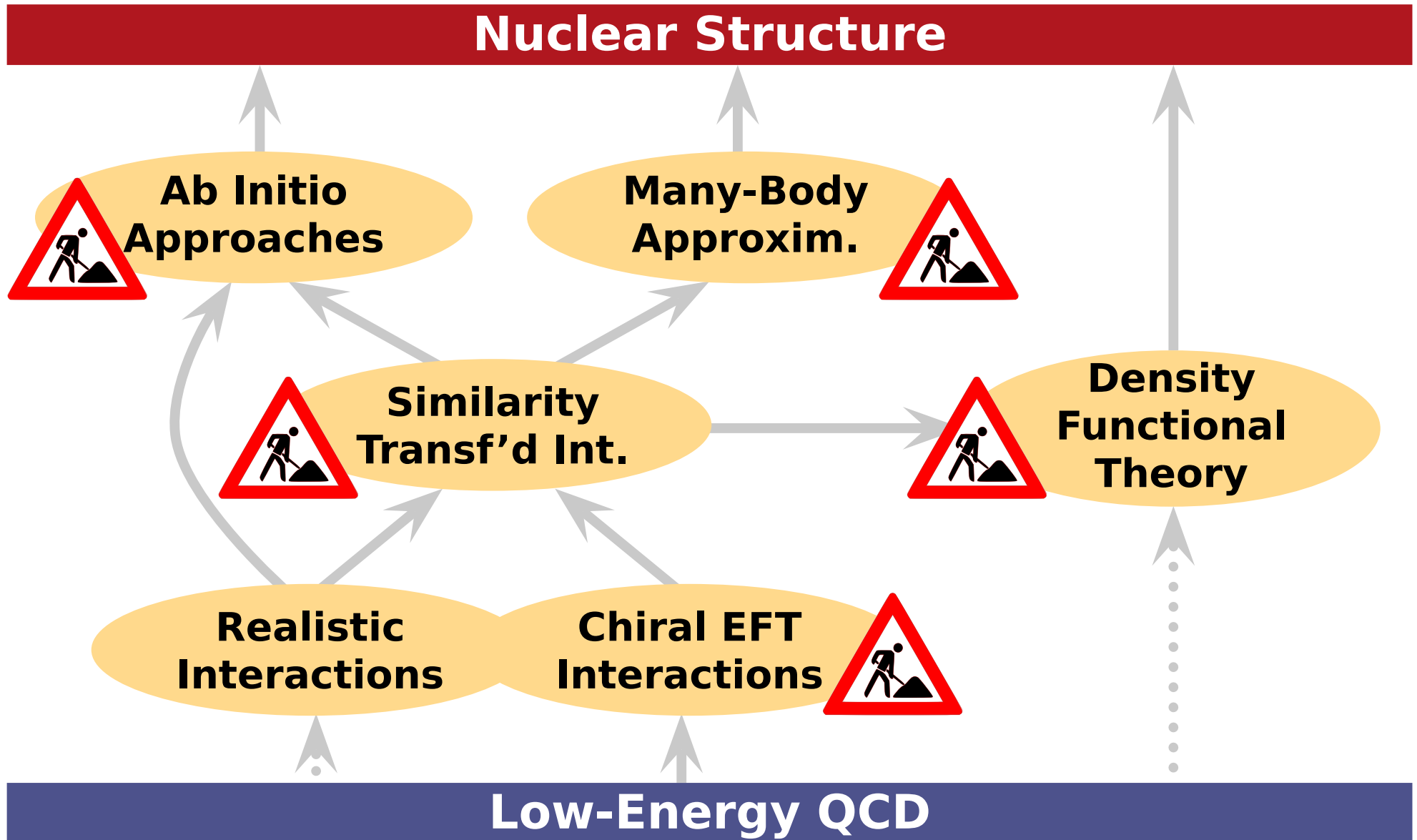


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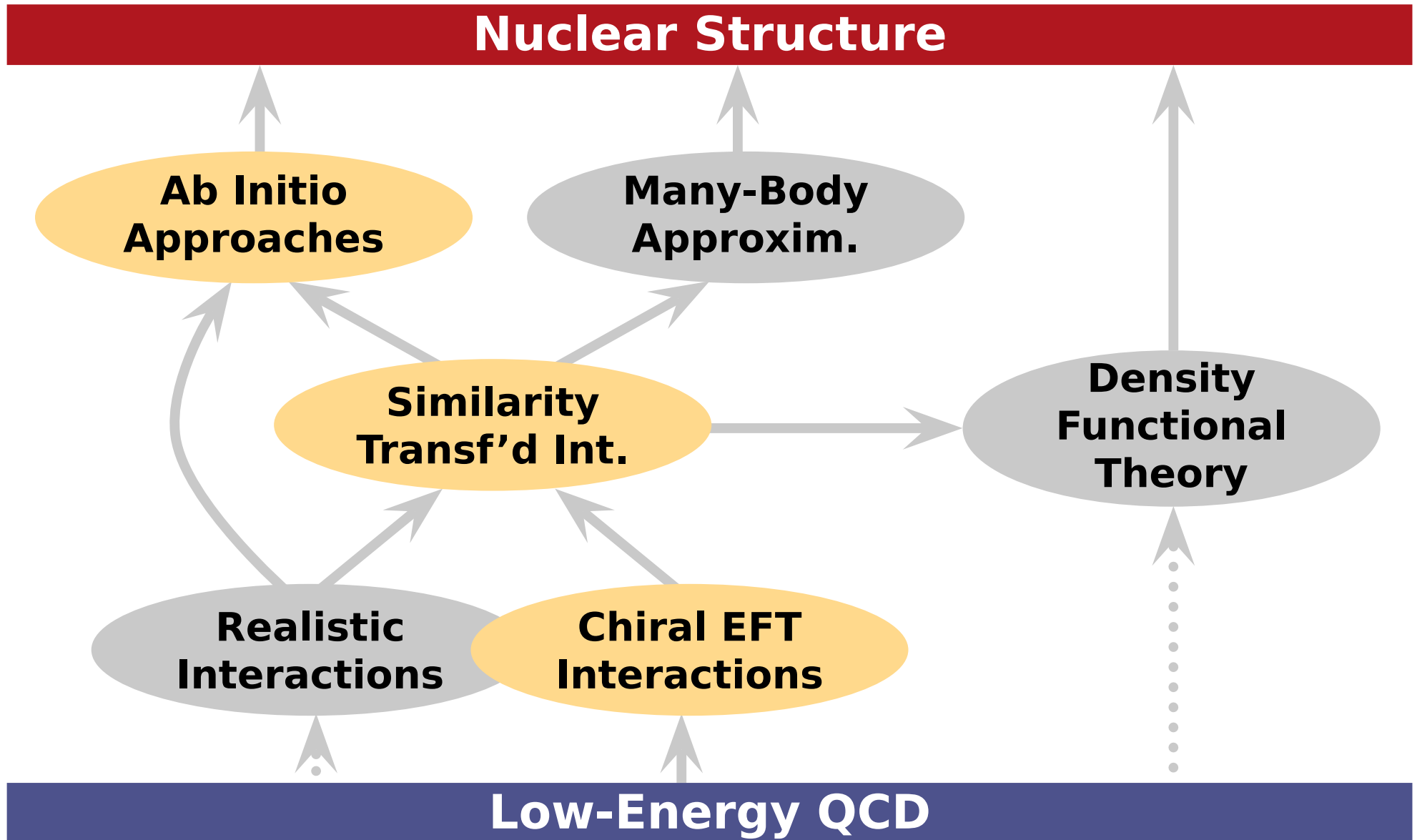
Nuclear Structure from a FAIR Perspective



Modern Nuclear Structure Theory



Modern Nuclear Structure Theory



Nuclear Interactions from QCD

Realistic Nuclear Interactions

■ QCD ingredients

- chiral effective field theory
- meson-exchange theory

■ short-range phenomenology

- contact terms or parameterization of short-range potential

■ experimental two-body data

- scattering phase-shifts & deuteron properties reproduced with high precision

■ supplementary 3N interaction

- adjusted to spectra of light nuclei

Argonne
V18

CD Bonn

Nijmegen
I/II

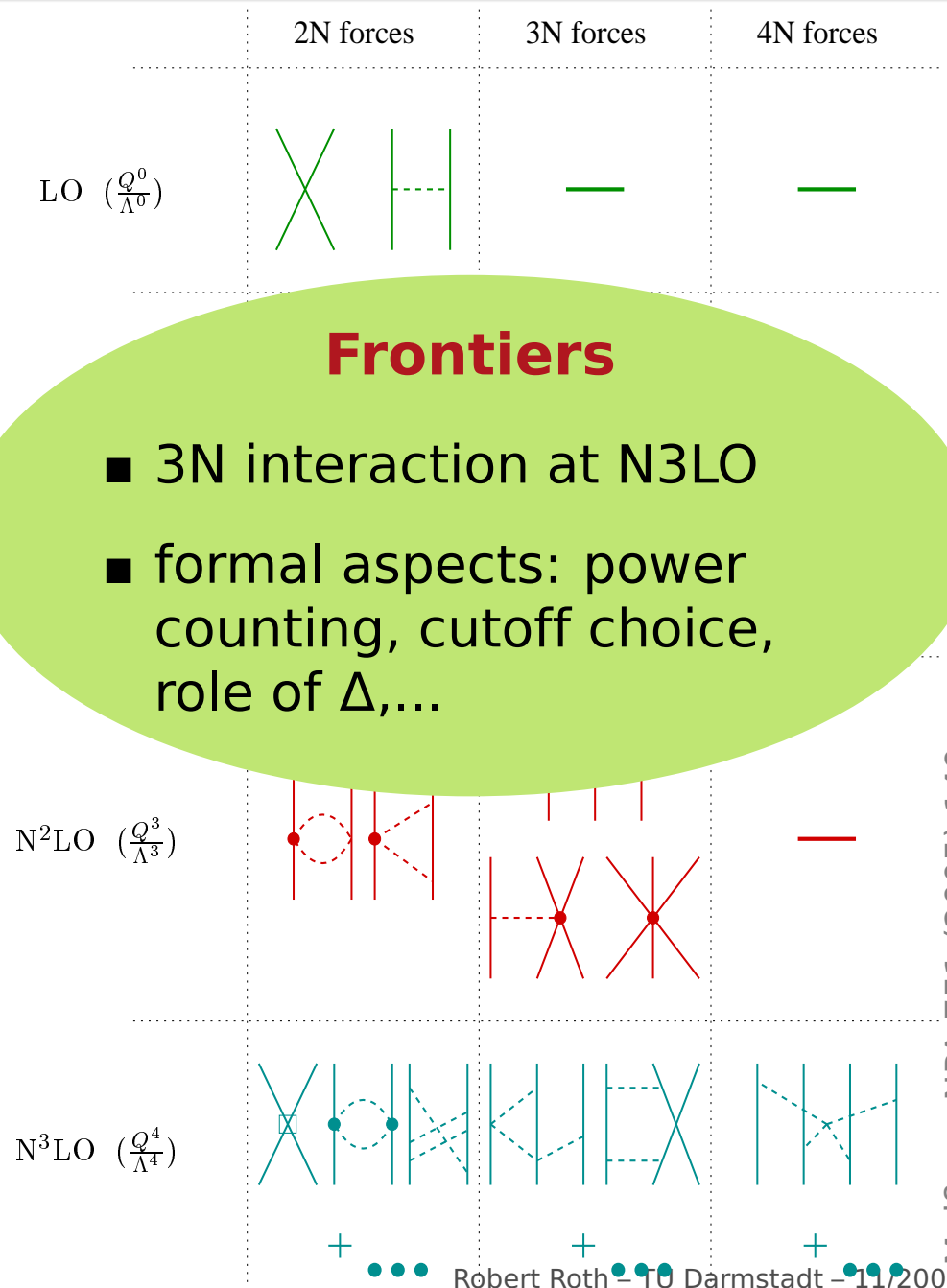
Chiral
N3LO

Argonne V18
+ Illinois 2

Chiral N3LO
+ N2LO

Nuclear Interaction from Chiral EFT

- **EFT for relevant degrees of freedom (π, N)** based on symmetries of QCD
- long-range **pion dynamics** treated explicitly
- short-range physics absorbed in **contact terms**
- low-energy constants fitted to experimental data ($NN, \pi N$)
- hierarchy of **consistent NN, 3N, ... interactions** (including current operators)



Similarity Transformed Interactions

Why Transformed Interactions?

Realistic Interactions

- generate strong correlations in many-body states
- short-range central & tensor correlations most important

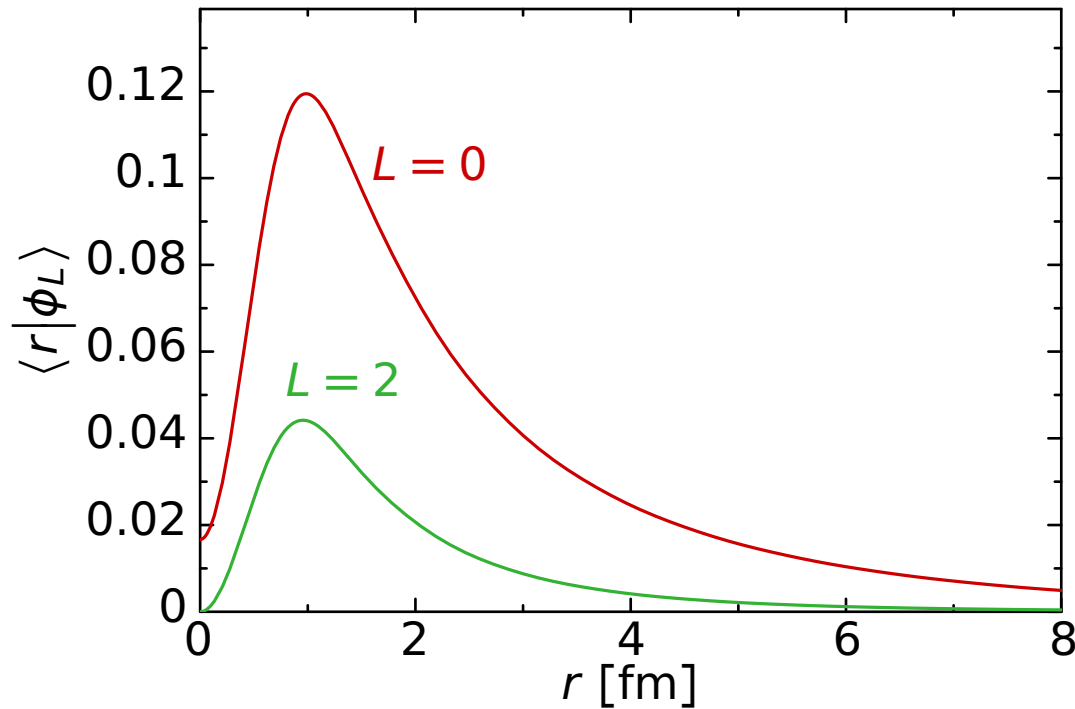
Many-Body Methods

- rely on truncated many-nucleon Hilbert spaces
- not capable of describing short-range correlations
- extreme: Hartree-Fock based on single Slater determinant

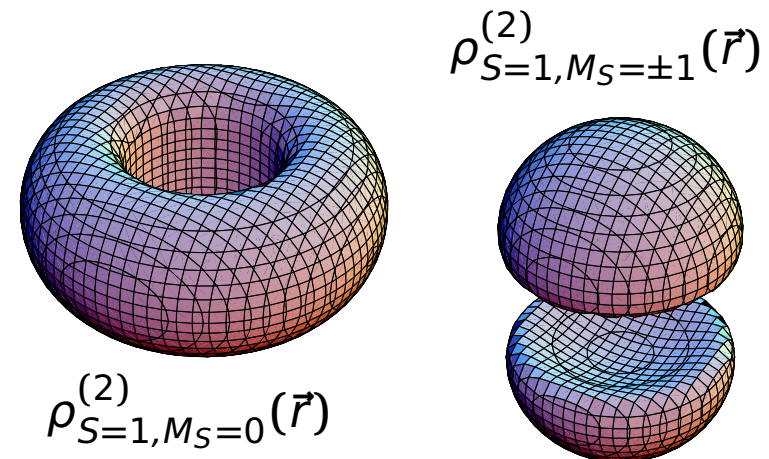
Similarity Transformation

- adapt realistic potential to the available model space
- conserve experimentally constrained properties (phase shifts)

Deuteron: Manifestation of Correlations



- **exact deuteron solution** for Argonne V18 potential



short-range repulsion
suppresses wavefunction
at small distances r

central correlations

tensor interaction
generates $L=2$ admixture
to ground state

tensor correlations

Similarity Transformed Interactions

Unitary Correlation Operator Method (UCOM)

H. Feldmeier et al. — Nucl. Phys. A 632 (1998) 61

T. Neff et al. — Nucl. Phys. A713 (2003) 311

R. Roth et al. — Nucl. Phys. A 745 (2004) 3

R. Roth et al. — Phys. Rev. C 72, 034002 (2005)

Unitary Correlation Operator Method

Correlation Operator

define a unitary operator C to describe the effect of short-range correlations

$$C = \exp[-iG] = \exp\left[-i \sum_{i<j} g_{ij}\right]$$

Correlated States

imprint short-range correlations onto uncorrelated many-body states

$$|\tilde{\psi}\rangle = C |\psi\rangle$$

Correlated Operators

adapt Hamiltonian to uncorrelated states (pre-diagonalization)

$$\tilde{O} = C^\dagger O C$$

$$\langle \tilde{\psi} | O | \tilde{\psi}' \rangle = \langle \psi | C^\dagger O C | \psi' \rangle = \langle \psi | \tilde{O} | \psi' \rangle$$

Unitary Correlation Operator Method

explicit ansatz for unitary transformation operator **motivated by the physics of short-range correlations**

Central Correlator C_r

- radial distance-dependent shift in the relative coordinate of a nucleon pair

$$g_r = \frac{1}{2} [s(r) q_r + q_r s(r)]$$

$$q_r = \frac{1}{2} \left[\frac{\vec{r}}{r} \cdot \vec{q} + \vec{q} \cdot \frac{\vec{r}}{r} \right]$$

Tensor Correlator C_Ω

- angular shift depending on the orientation of spin and relative coordinate of a nucleon pair

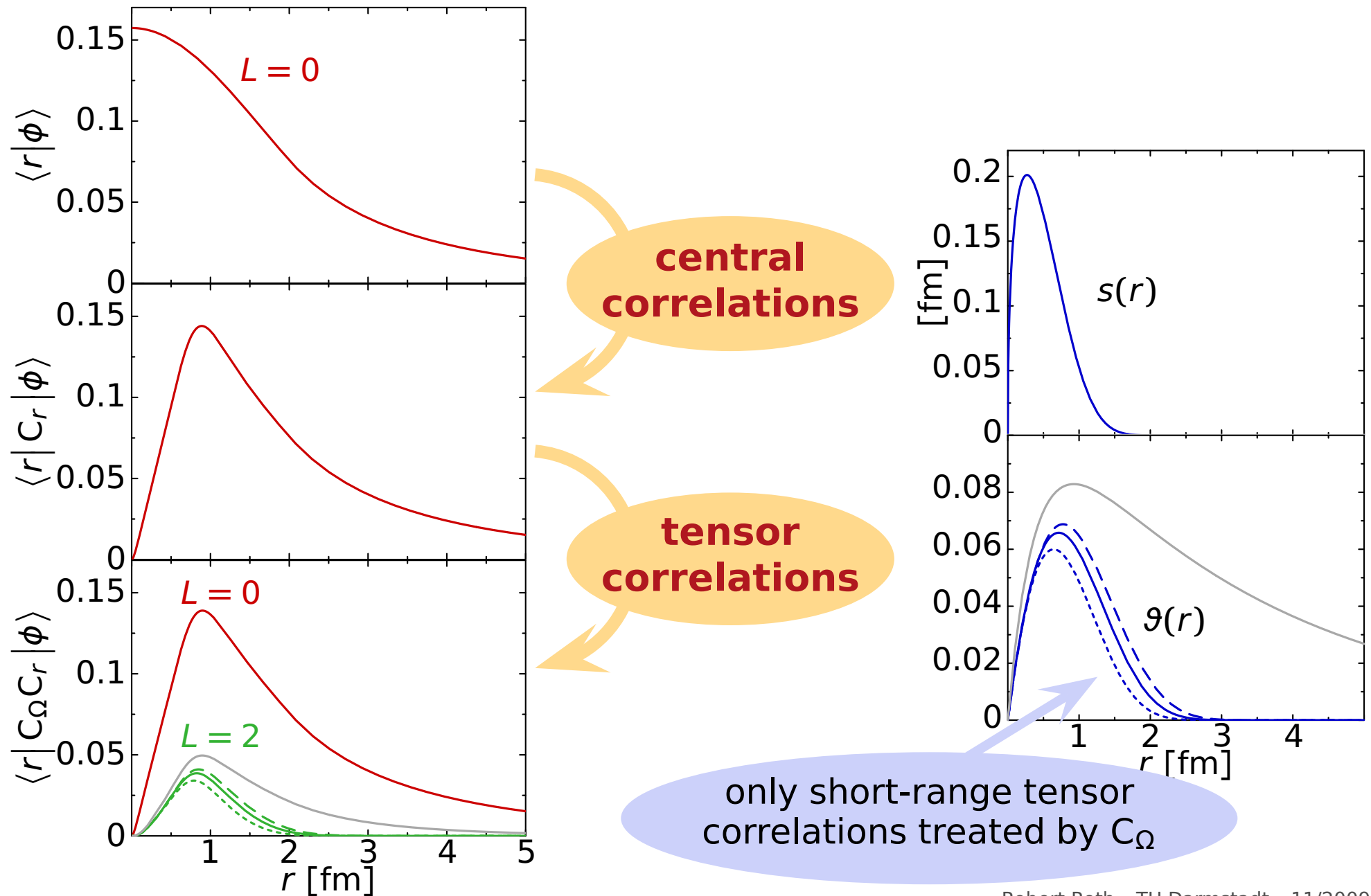
$$g_\Omega = \frac{3}{2} \vartheta(r) [(\vec{\sigma}_1 \cdot \vec{q}_\Omega)(\vec{\sigma}_2 \cdot \vec{r}) + (\vec{r} \leftrightarrow \vec{q}_\Omega)]$$

$$\vec{q}_\Omega = \vec{q} - \frac{\vec{r}}{r} q_r$$

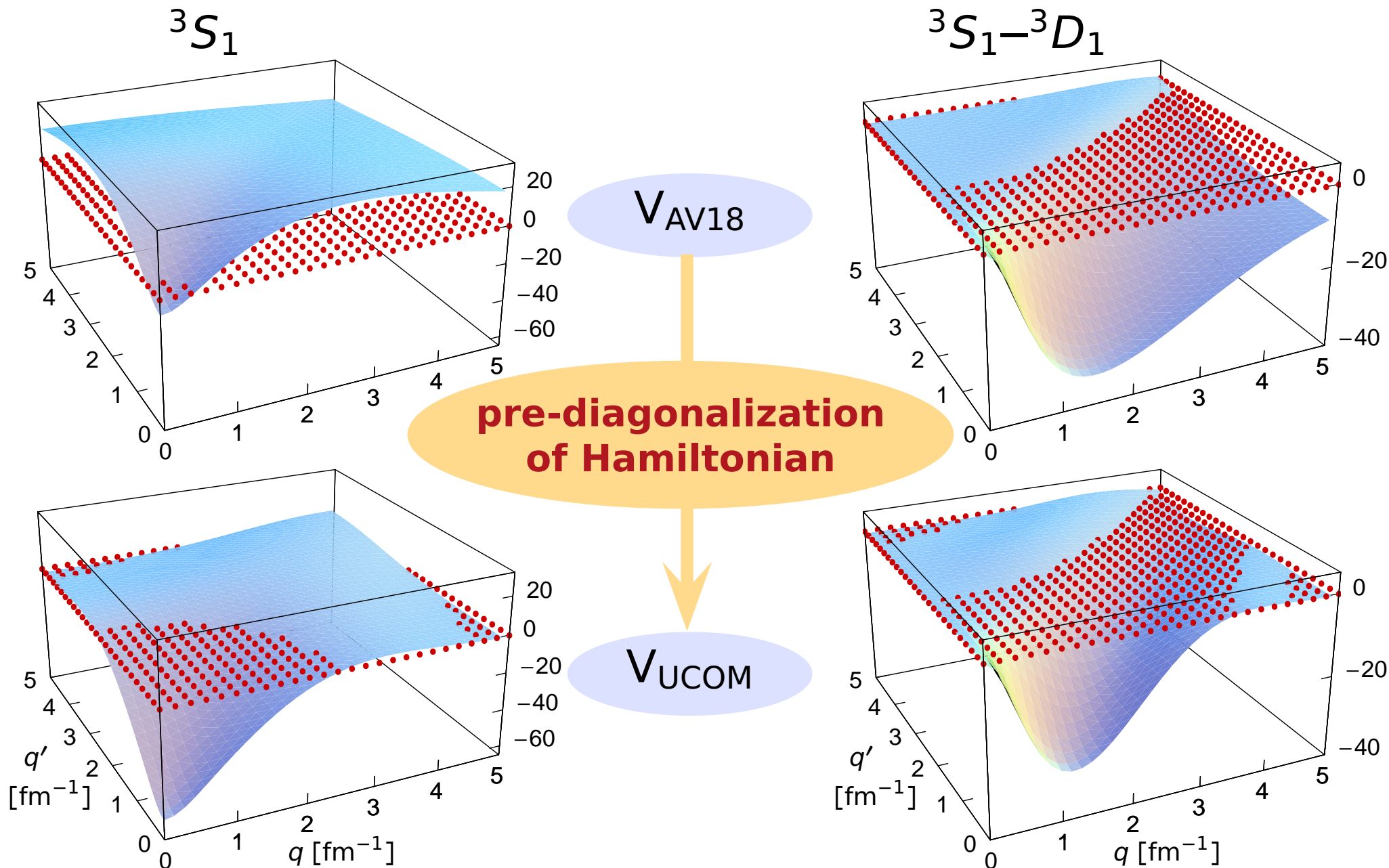
$$C = C_\Omega C_r = \exp\left(-i \sum_{i < j} g_{\Omega,ij}\right) \exp\left(-i \sum_{i < j} g_{r,ij}\right)$$

- $s(r)$ and $\vartheta(r)$ are optimized for the initial potential

Correlated States: The Deuteron



Correlated Interaction: V_{UCOM}



Similarity Transformed Interactions

Similarity Renormalization Group (SRG)

Hergert & Roth — Phys. Rev. C 75, 051001(R) (2007)

Bogner et al. — Phys. Rev. C 75, 061001(R) (2007)

Roth, Reinhardt, Hergert — Phys. Rev. C 77, 064033 (2008)

Similarity Renormalization Group

flow evolution of the **Hamiltonian to band-diagonal form** with respect to uncorrelated many-body basis

Flow Equation for Hamiltonian

- evolution equation for Hamiltonian

$$\tilde{H}(\alpha) = C^\dagger(\alpha) H C(\alpha) \quad \rightarrow \quad \frac{d}{d\alpha} \tilde{H}(\alpha) = [\eta(\alpha), \tilde{H}(\alpha)]$$

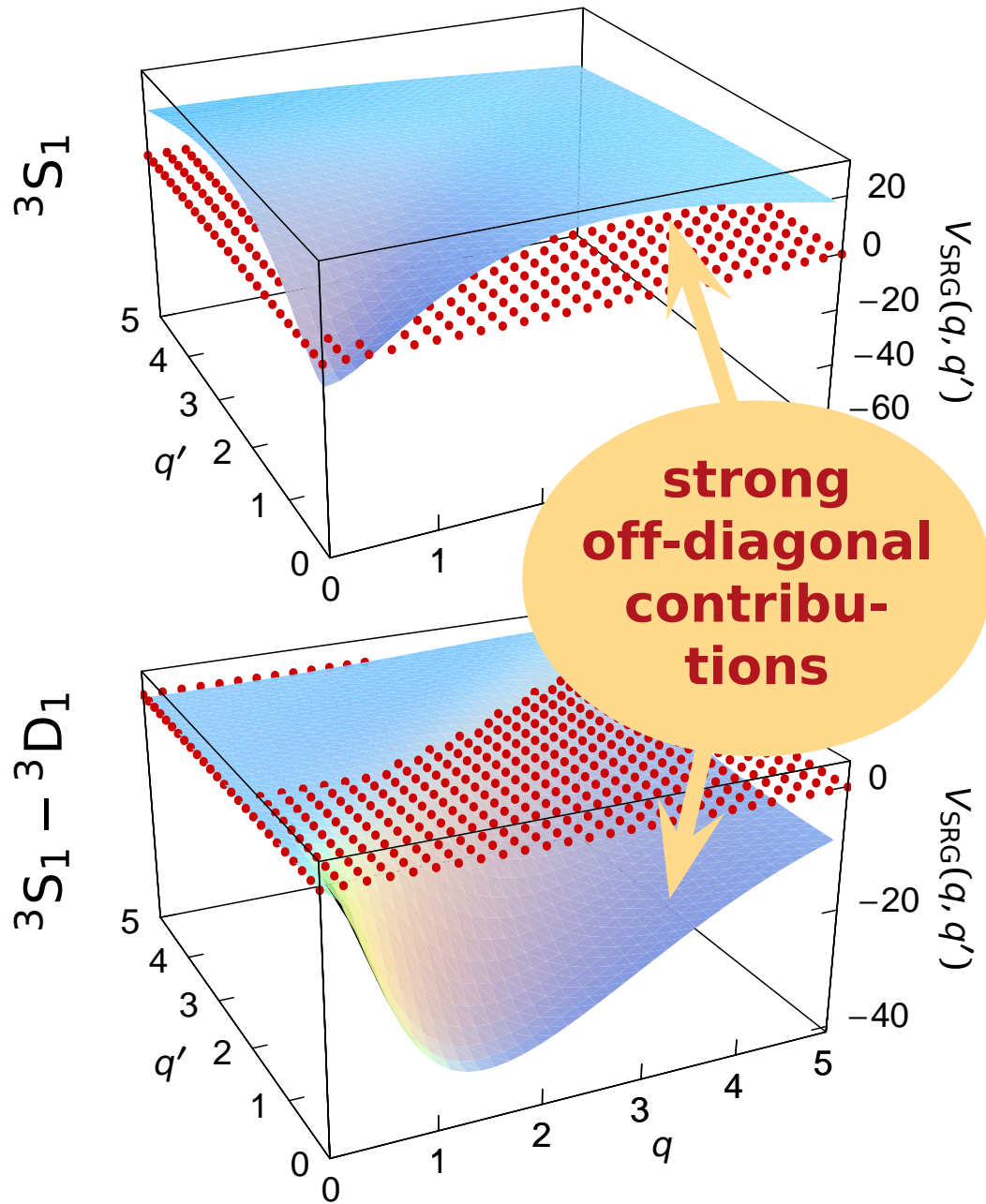
- dynamical generator defined as commutator with the operator in whose eigenbasis H shall be diagonalized

$$\eta(\alpha) \stackrel{2B}{=} \frac{1}{2\mu} [\vec{q}^2, \tilde{H}(\alpha)]$$

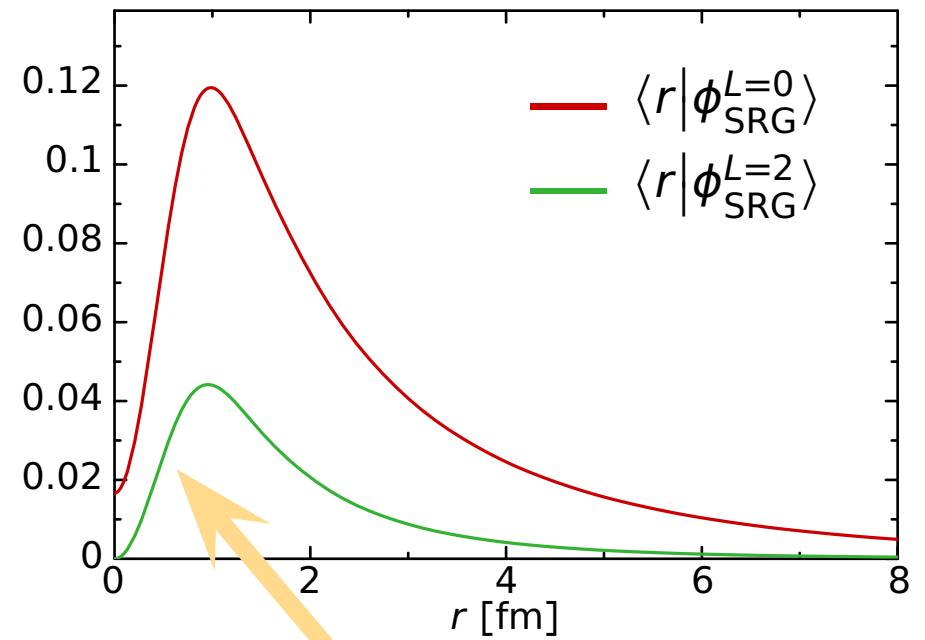
UCOM vs. SRG

$\eta(0)$ has the same structure as UCOM generators g_r & g_Ω

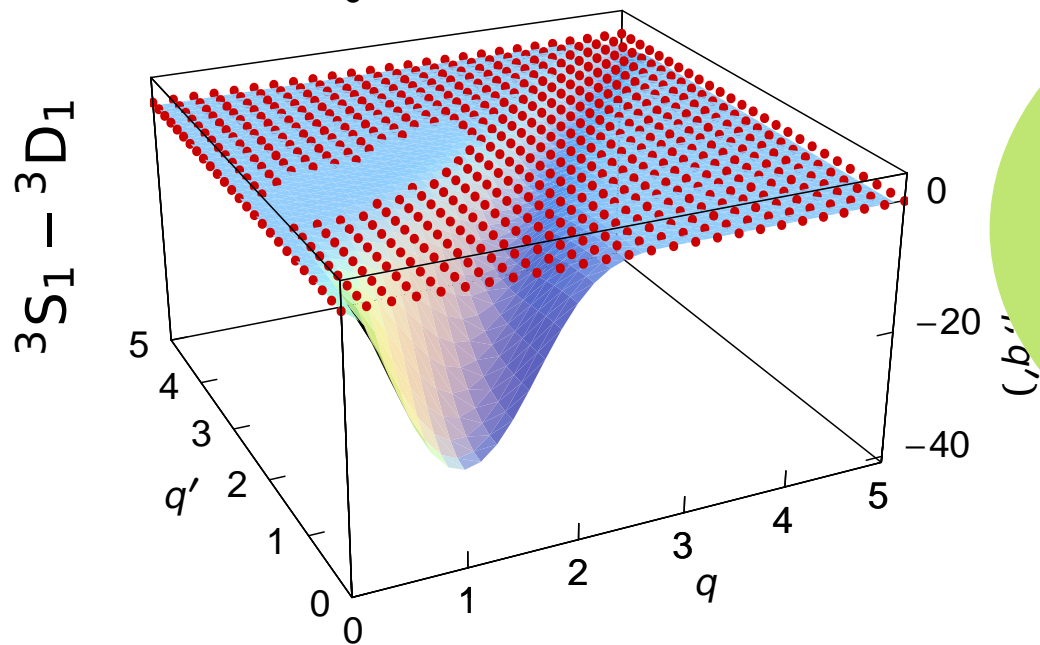
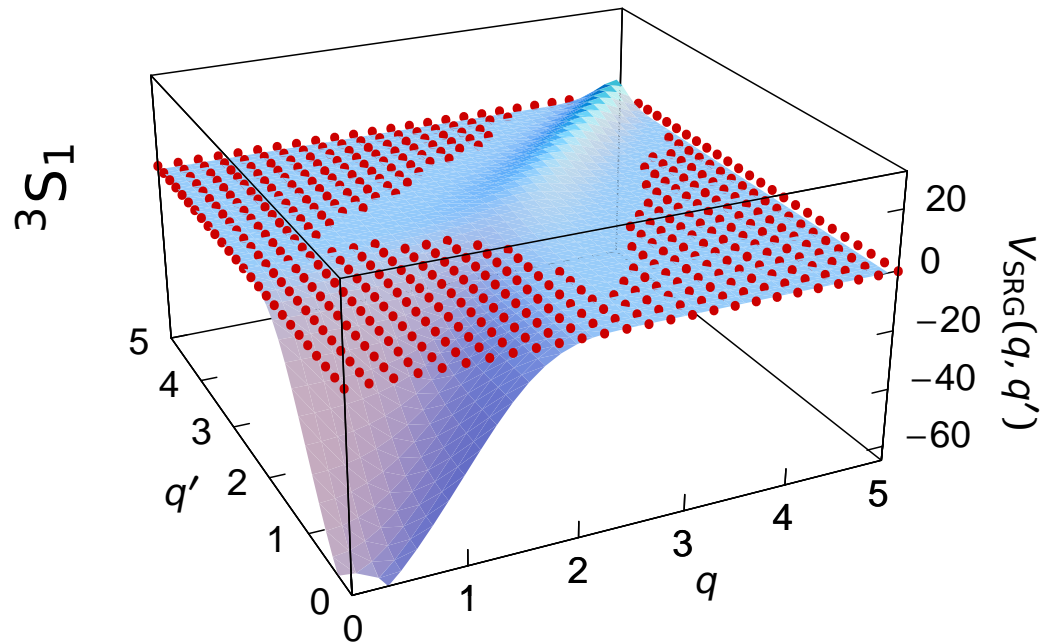
SRG Evolution: The Deuteron



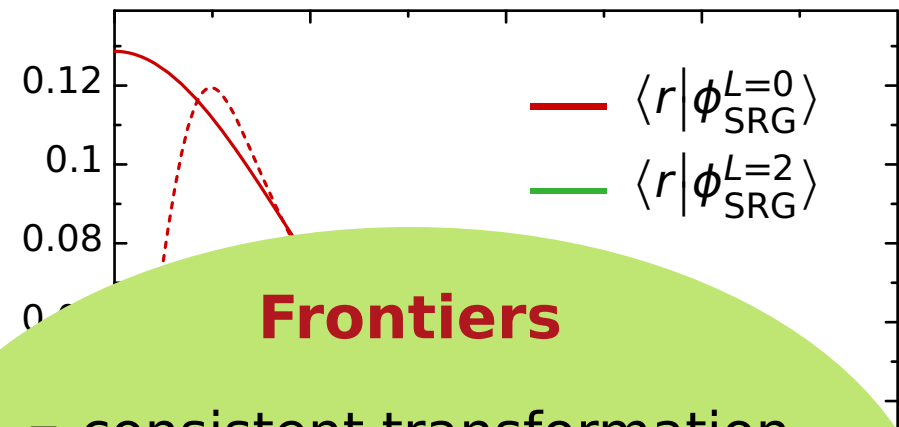
Argonne V18



SRG Evolution: The Deuteron



$$\alpha = 0.1000 \text{ fm}^4$$



Frontiers

- consistent transformation of NN+3N Hamiltonian
- transformed observables
- alternative SRG & UCOM generators

Ab Initio Approaches

No-Core Shell Model

Roth et al. — Phys. Rev. C 72, 034002 (2005)

Roth & Navrátil — in preparation

No-Core Shell Model: Basics

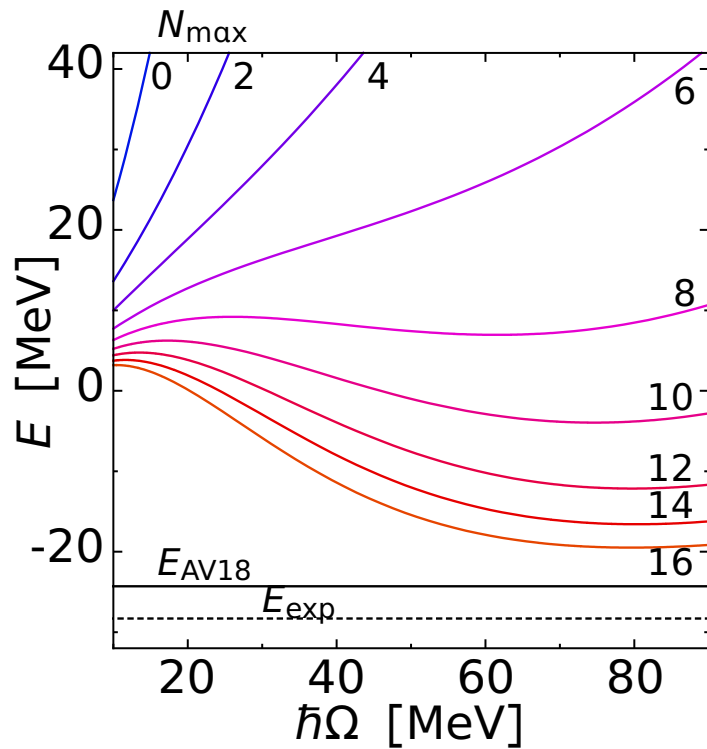
- special case of a **full configuration interaction (CI)** scheme
- **many-body basis**: Slater determinants $|\Phi_\nu\rangle$ composed of harmonic oscillator single-particle states

$$|\Psi\rangle = \sum_\nu C_\nu |\Phi_\nu\rangle$$

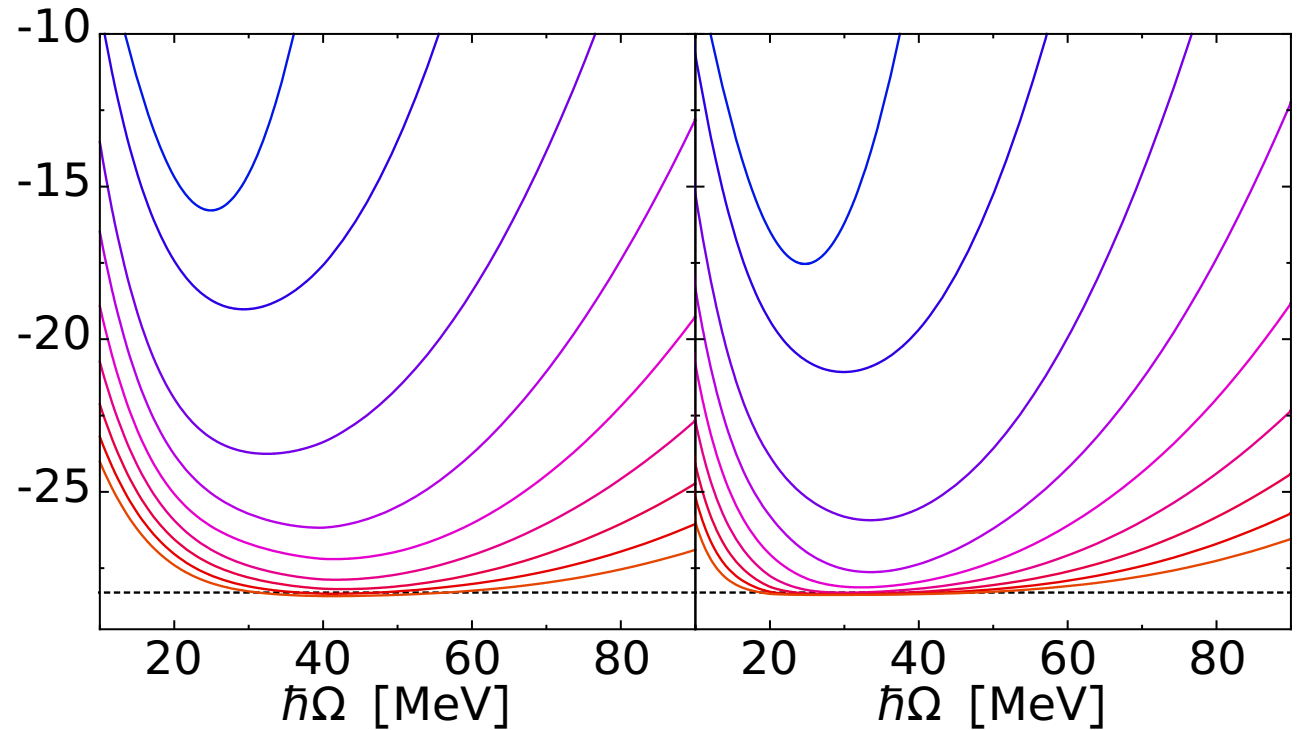
- **model space**: spanned by basis states $|\Phi_\nu\rangle$ with unperturbed excitation energies of up to $N_{\max}\hbar\Omega$
 - ▶ **exact factorization** of intrinsic and CM component is possible
- numerical solution of **eigenvalue problem** for H_{int} within $N_{\max}\hbar\Omega$ model space via Lanczos methods
 - ▶ model spaces of **up to 10^9 basis states** are used routinely
- increase N_{\max} until **convergence** is observed

^4He : NCSM Convergence

V_{AV18}



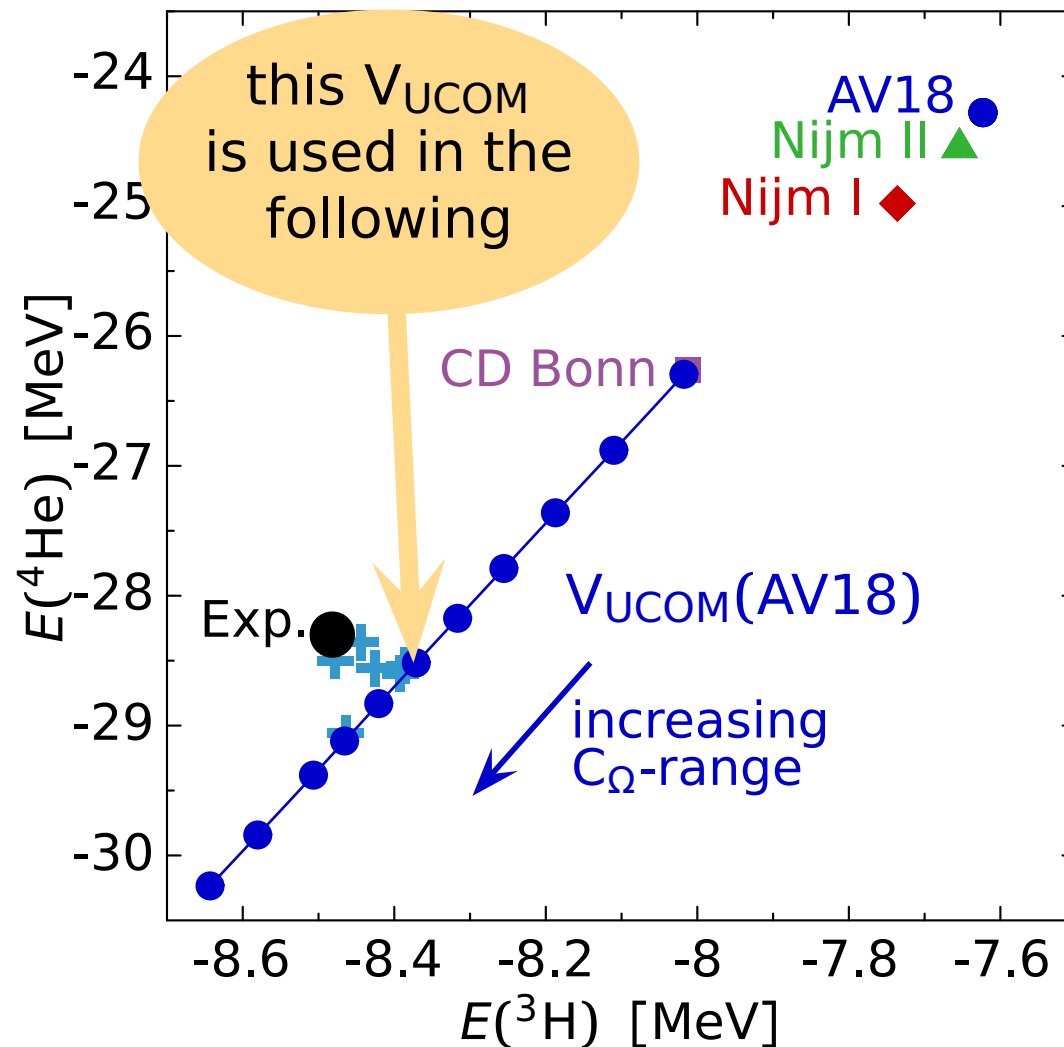
V_{UCOM}
MIN, $I_9 = 0.09 \text{ fm}^3$



V_{SRG}
 $\bar{\alpha} = 0.03 \text{ fm}^4$

- I_9 or $\bar{\alpha}$ adjusted such that ^4He binding energy is reproduced

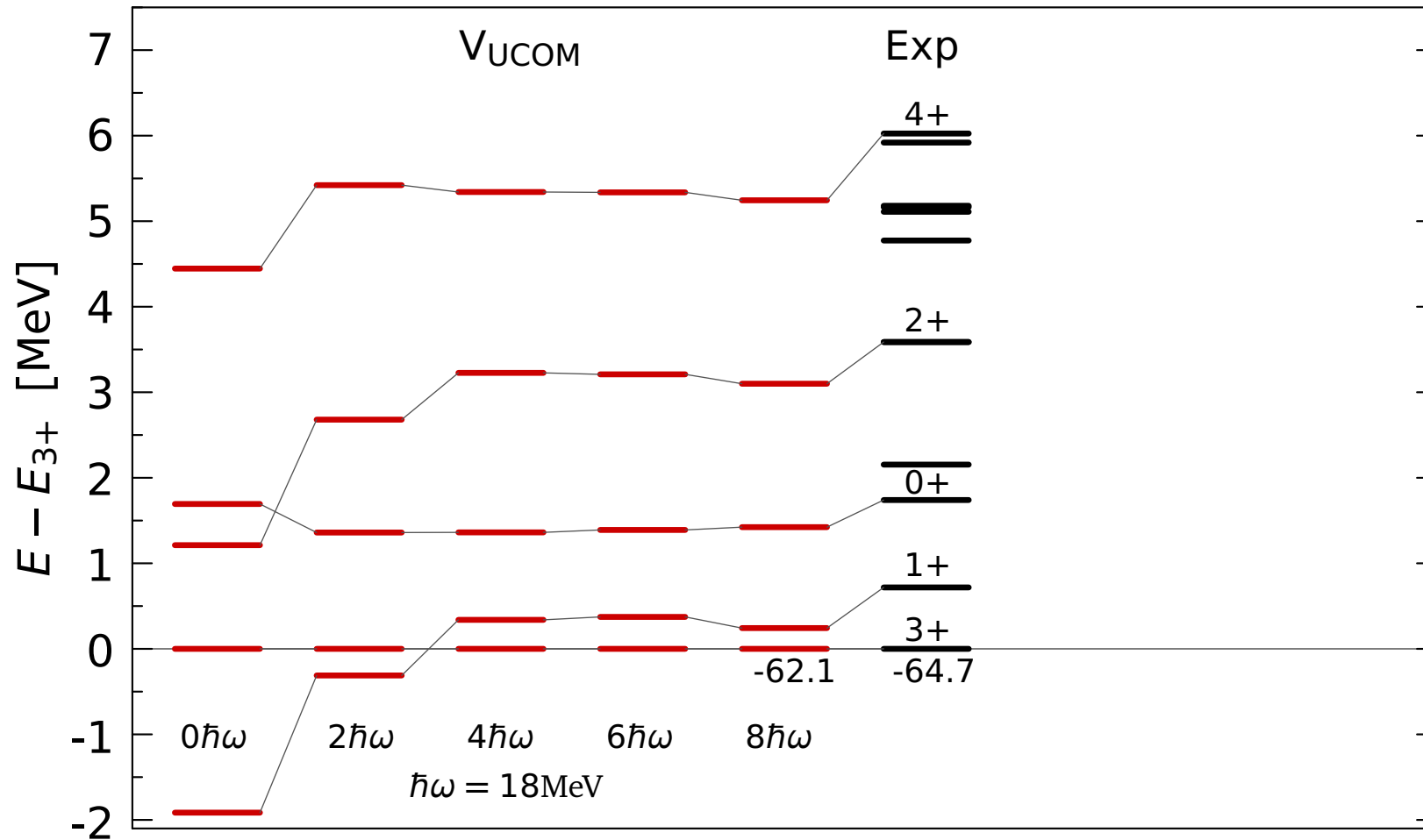
Tjon-Line and 3N Interactions



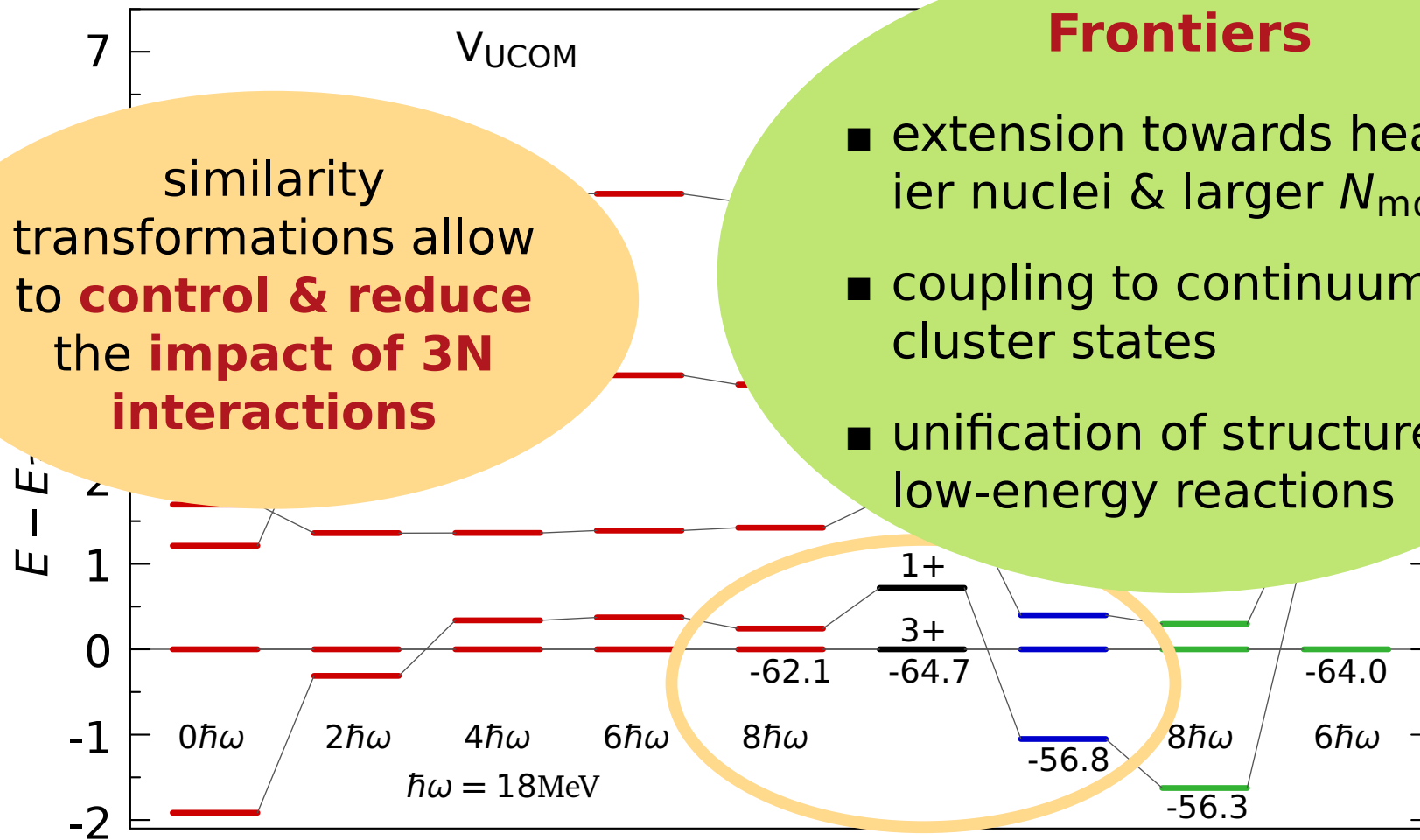
- **Tjon-line:** $E(^4\text{He})$ vs. $E(^3\text{H})$ for phase-shift equivalent NN-interactions
- change of C_{Ω} -correlator range results in shift along Tjon-line

**minimize net
3N interaction**
by choosing
correlator close to
experimental point

^{10}B : Hallmark of a 3N Interaction?



^{10}B : Hallmark of a 3N Interaction?



Ab Initio Approaches

Importance Truncated No-Core Shell Model

Roth — Phys. Rev. C 79, 064324 (2009)

Roth, Gour & Piecuch — Phys. Lett. B 679, 334 (2009)

Roth, Gour & Piecuch — Phys. Rev. C 79, 054325 (2009)

Roth & Navrátil — Phys. Rev. Lett. 99, 092501 (2007)

Importance Truncated NCSM

- converged NCSM calculations are essentially restricted to p-shell

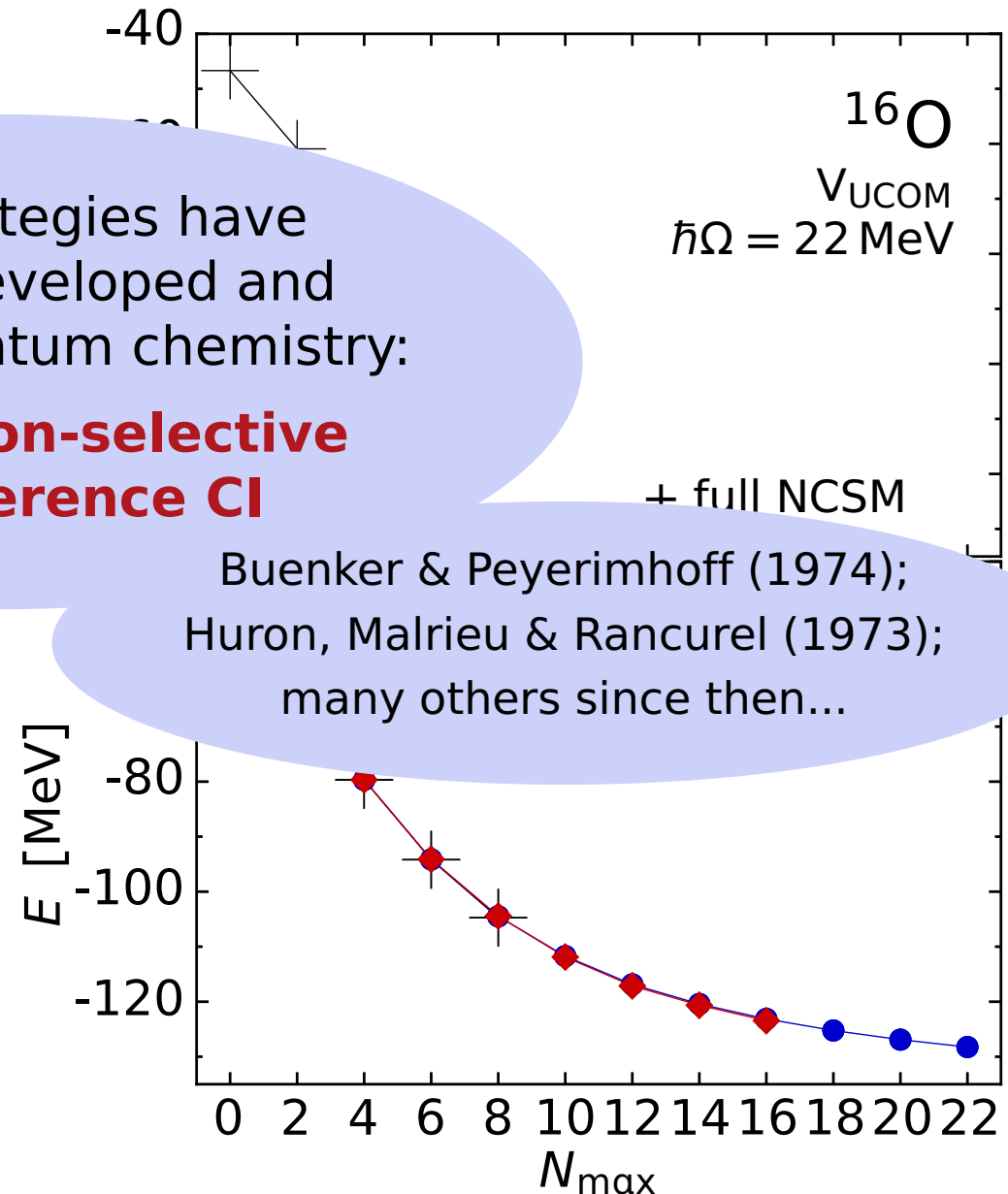
- full 10 oscillator shells for ^{16}O (basis dimension $\sim 10^{10}$)

similar strategies have first been developed and applied in quantum chemistry:

configuration-selective multireference CI

Importance Truncation

reduce NCSM space to the relevant basis states using an **a priori importance measure** derived from MBPT



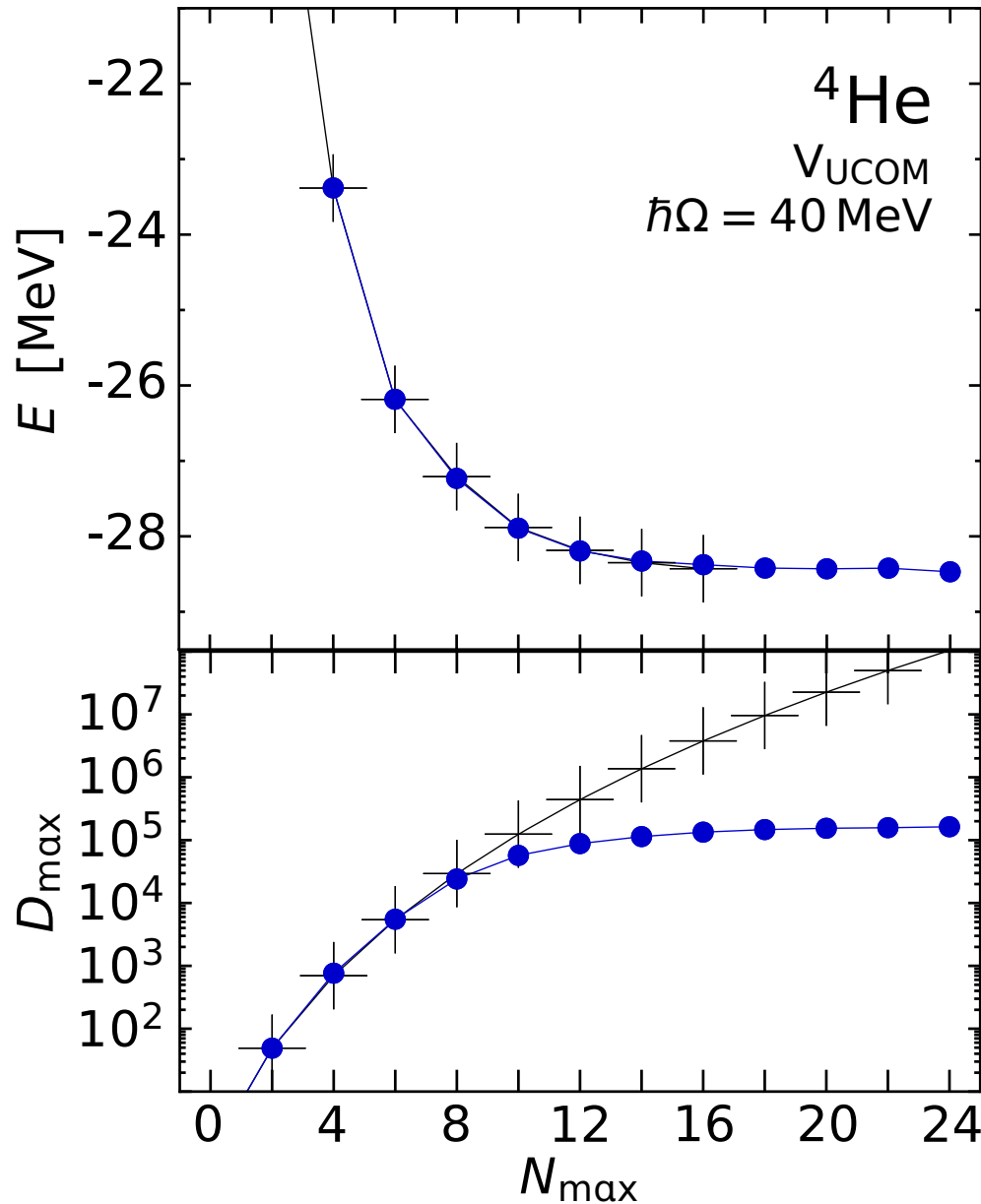
Importance Truncation: General Idea

- given an initial approximation $|\Psi_{\text{ref}}\rangle$ for the **target state**
- **measure the importance** of individual basis state $|\Phi_\nu\rangle$ via first-order multiconfigurational perturbation theory

$$K_\nu = -\frac{\langle \Phi_\nu | H | \Psi_{\text{ref}} \rangle}{\epsilon_\nu - \epsilon_{\text{ref}}}$$

- construct **importance truncated space** spanned by basis states with $|K_\nu| \geq K_{\text{min}}$ and solve eigenvalue problem
- **iterative scheme**: repeat construction of importance truncated model space using eigenstate as improved reference $|\Psi_{\text{ref}}\rangle$
- **threshold extrapolations** and **perturbative corrections** can be used to account for discarded basis states

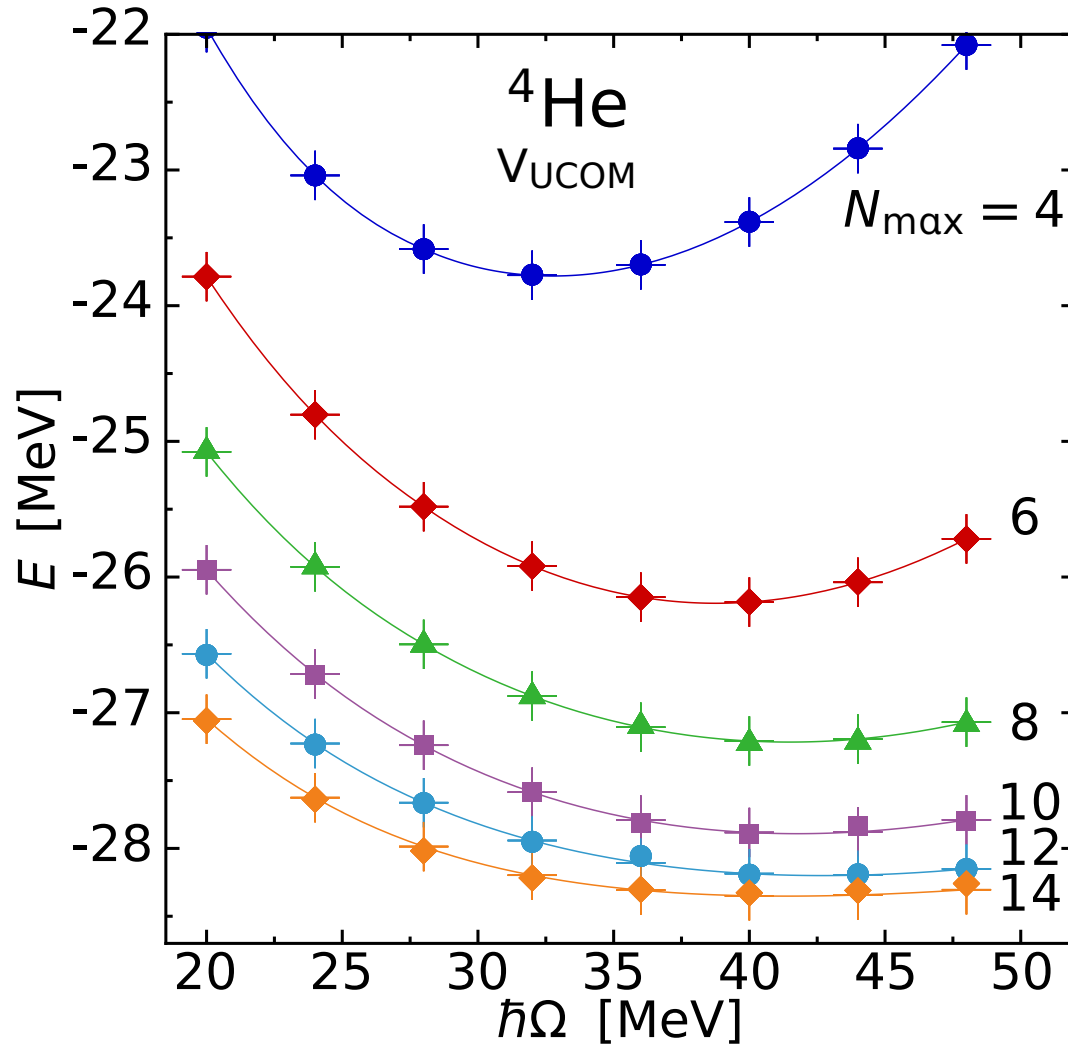
^4He : Importance-Truncated NCSM



- **sequential IT-NCSM(seq)**: single importance update using $(N_{\text{max}} - 2)\hbar\Omega$ eigenstate as reference
- **reproduces exact NCSM result** for all N_{max}
- reduction of basis by more than two orders of magnitude w/o loss of precision

+ full NCSM
● IT-NCSM(seq)

^4He : Importance-Truncated NCSM



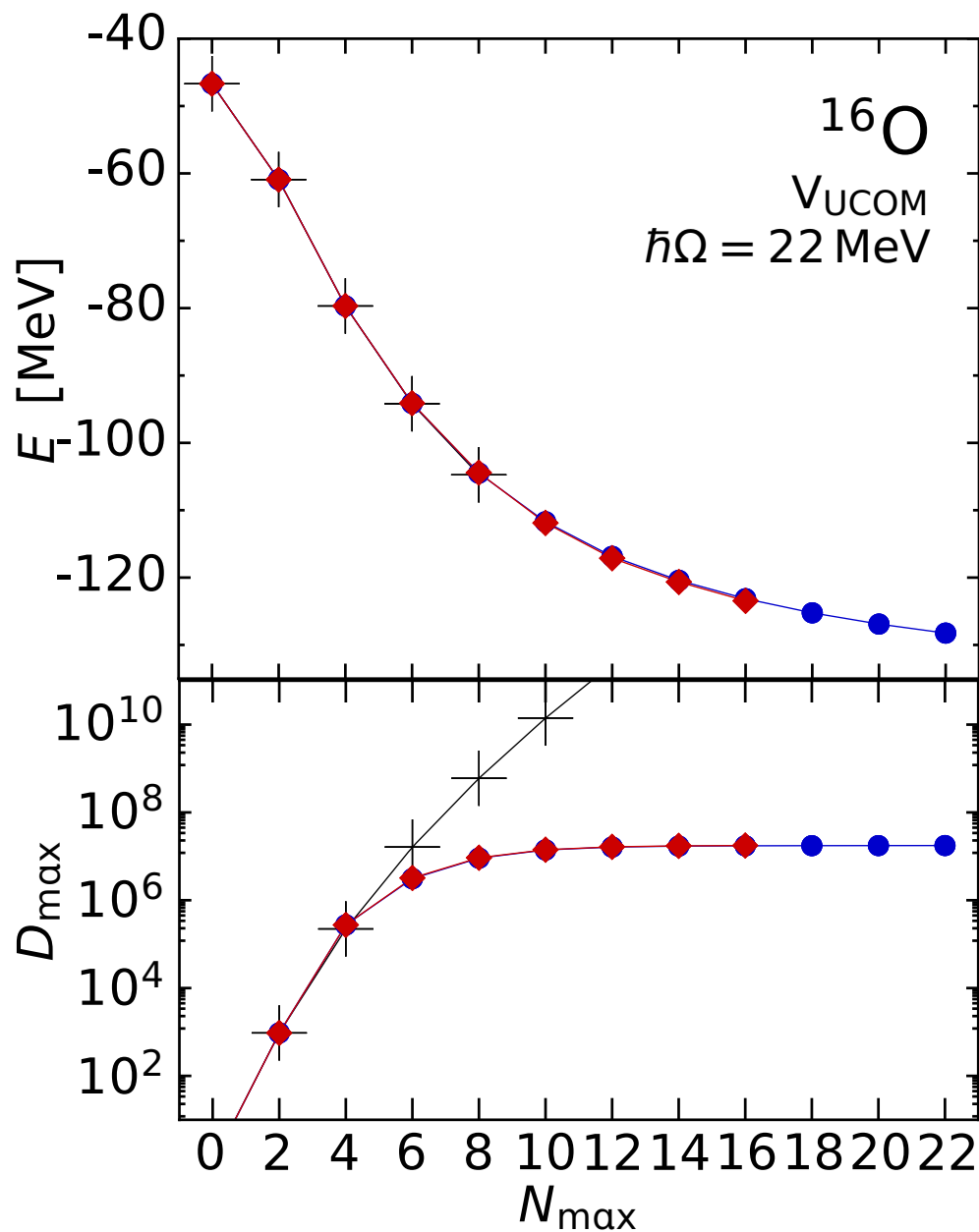
- **reproduces exact NCSM result** for all $\hbar\Omega$ and N_{max}

- importance truncation & threshold extrapolation is robust

- **no center-of-mass contamination** for any N_{max} and $\hbar\Omega$

- + full NCSM
- IT-NCSM(seq)

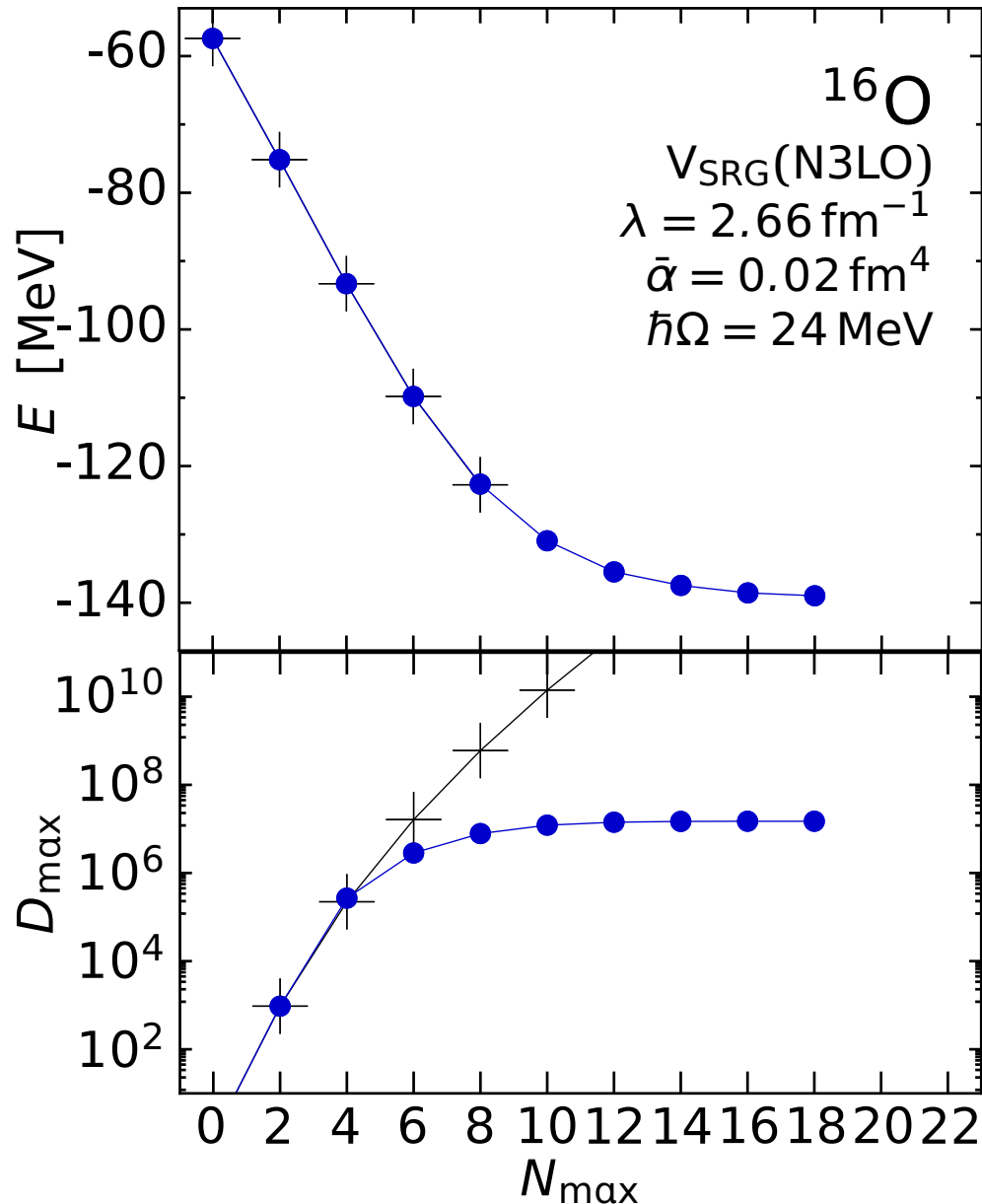
^{16}O : Importance-Truncated NCSM



- IT-NCSM(seq) provides **excellent agreement with full NCSM** calculation
- dimension reduced by **several orders of magnitude**
- possibility to go **way beyond** the domain of the full NCSM

- + full NCSM
- IT-NCSM(seq), $C_{\text{min}} = 0.0005$
- ◆ IT-NCSM(seq), $C_{\text{min}} = 0.0003$

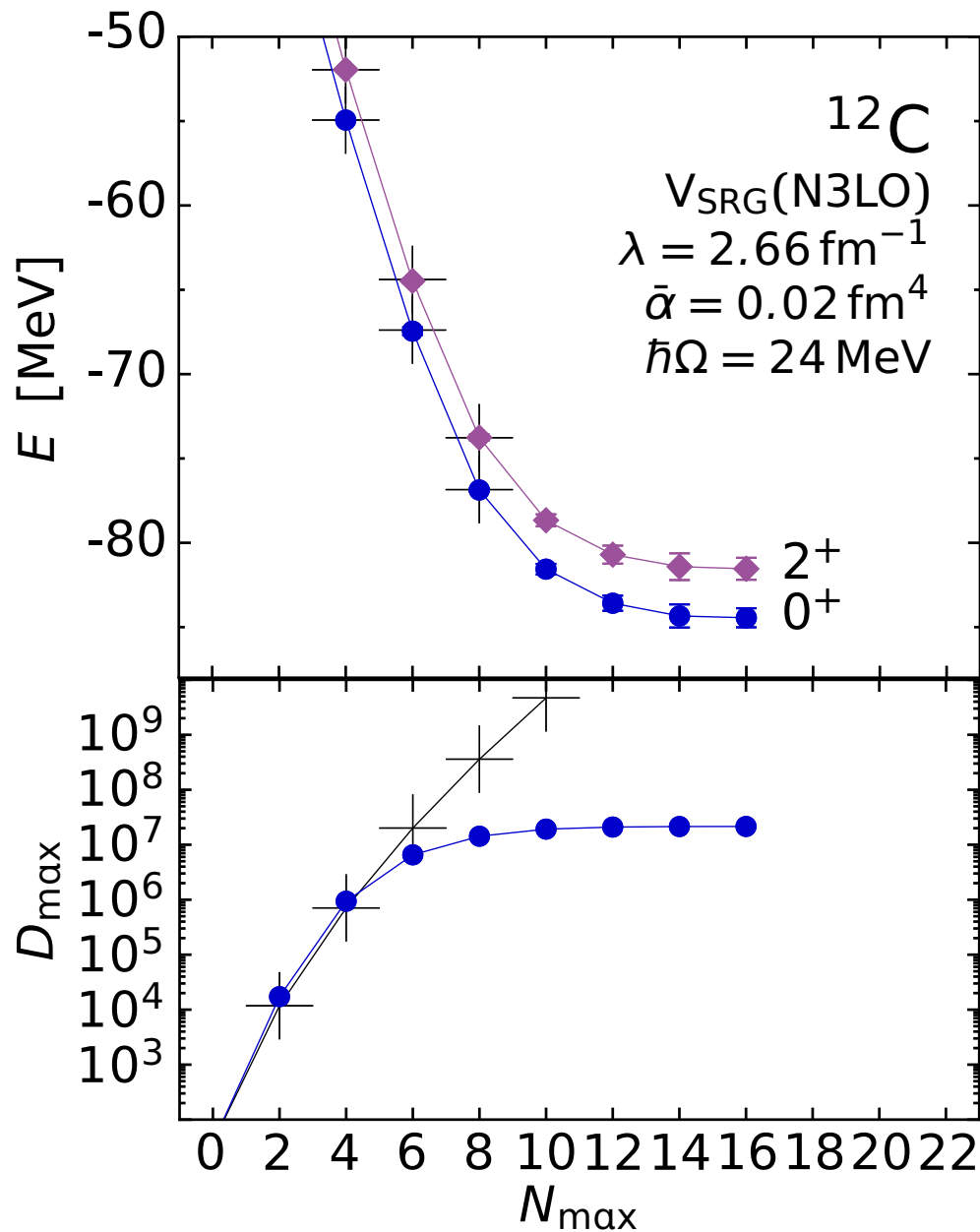
^{16}O : Importance-Truncated NCSM



- **SRG-evolved N3LO potential** provides a much better convergence behavior
- nevertheless, $N_{\text{max}} \leq 8$ calculations are not sufficient
- non-exponential behavior observed with V_{UCOM} is really due to interaction

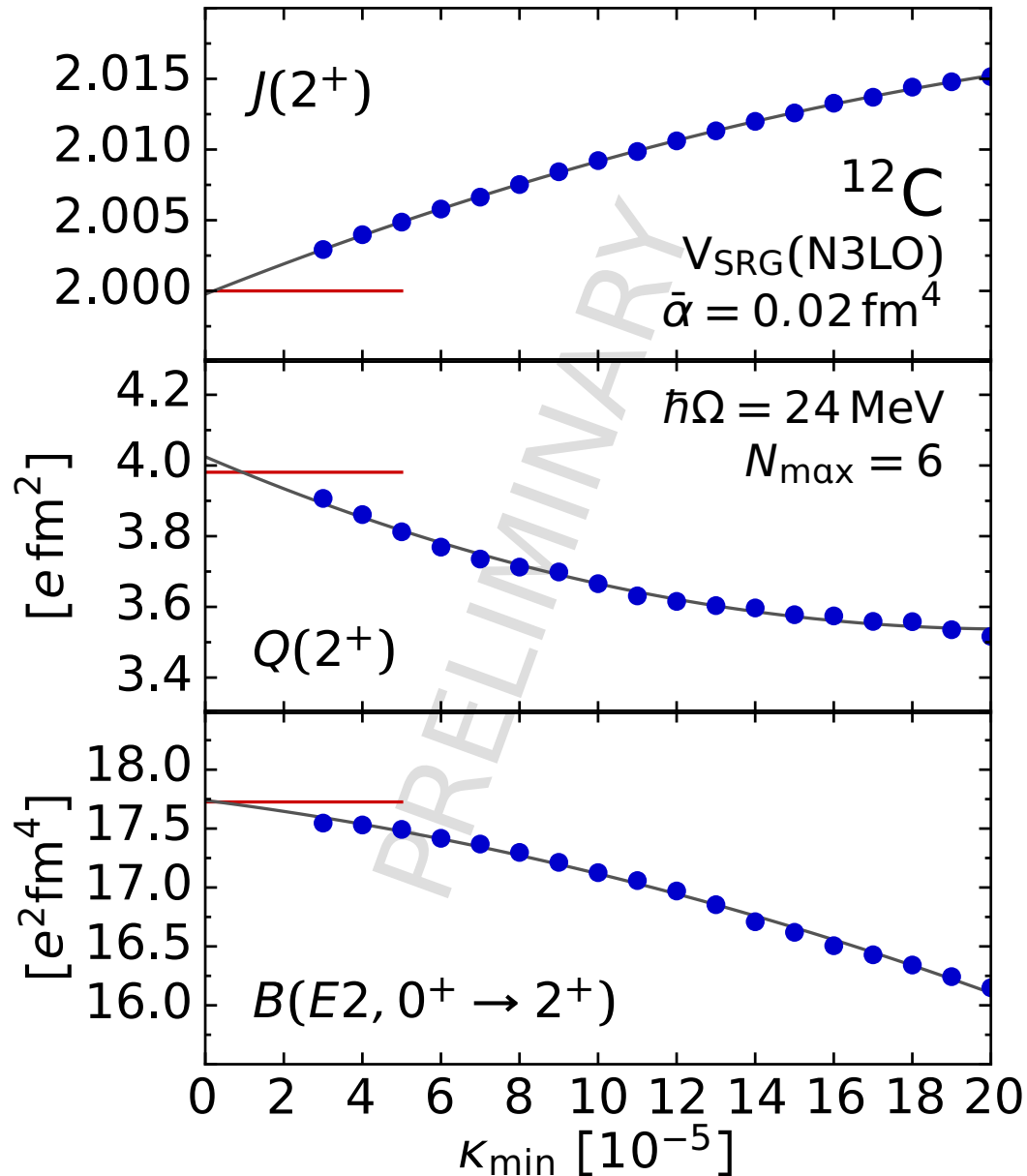
- + full NCSM
- IT-NCSM(seq), $C_{\text{min}} = 0.0005$

^{12}C : IT-NCSM for Excited States



- **target ground & excited states** simultaneously
 - ▶ separate importance measure $\kappa_{\nu}^{(n)}$ for each target state
 - ▶ basis state is included if $|\kappa_{\nu}^{(n)}| \geq \kappa_{\text{min}}$ for any n
 - dimension of importance truncated space **grows linearly** with # of target states
- + full NCSM
 ●◆ IT-NCSM(seq), $C_{\text{min}} = 0.0005$

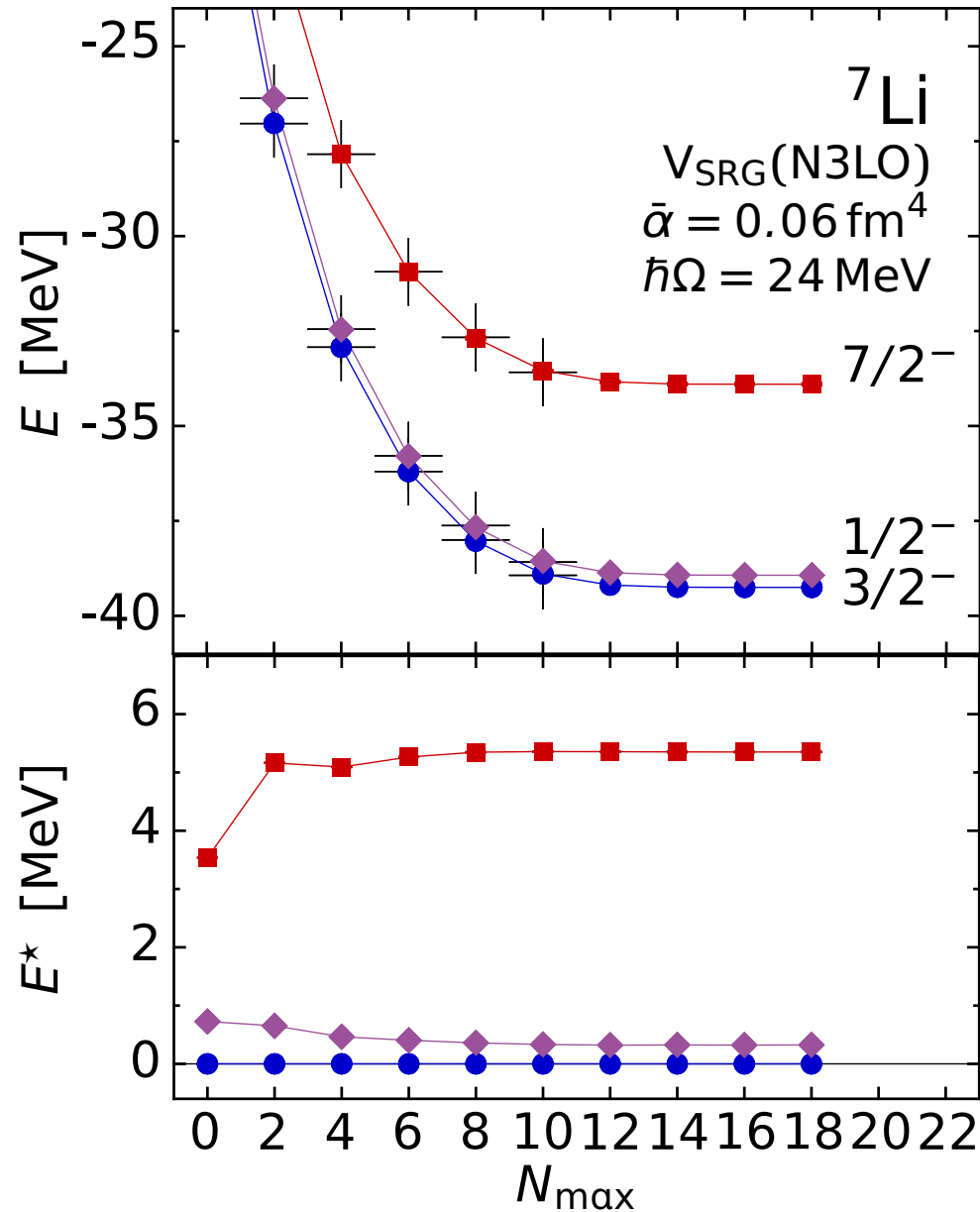
^{12}C : IT-NCSM for Spectroscopy



- access to **spectroscopic observables** via eigenstates
- multipole moments, transition strengths, transition form-factors, densities,...
- simple threshold extrapolation essentially **reproduces full NCSM results**

systematic spectroscopy in p- and sd-shell with large $N_{\max}\hbar\Omega$ spaces

${}^7\text{Li}$: IT-NCSM for Odd Nuclei



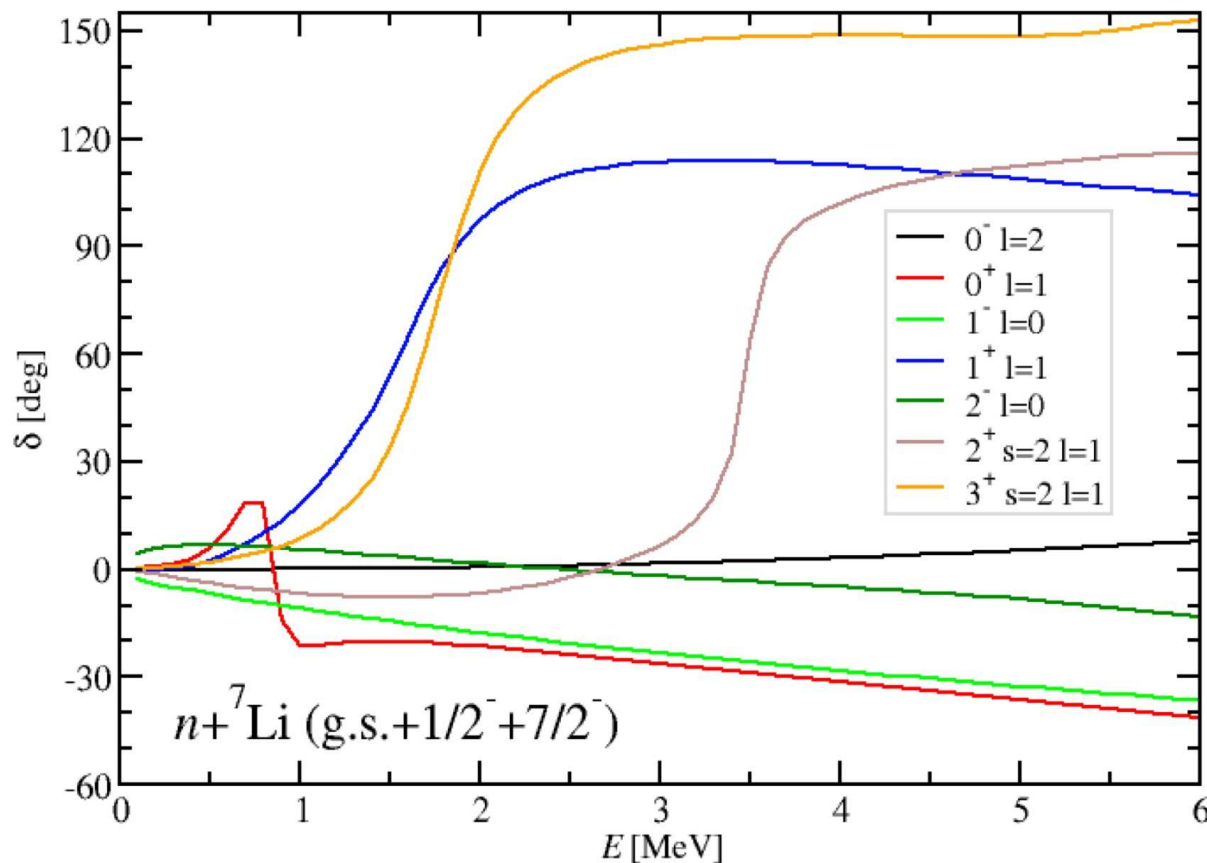
- IT-NCSM(seq) treats a ground state & low-lying excited states for open- and closed-shell nuclei **on the same footing**

- **excellent agreement with full NCSM** calculations in all cases

RGM & IT-NCSM: Ab Initio Reactions

with Petr Navrátil & Sofia Quaglioni (LLNL)

- **IT-NCSM wave function as input for RGM** (Resonating Group Method) calculations of low-energy nucleon-nucleus scattering



- using 3 lowest ${}^7\text{Li}$ states
- so-far up to $N_{\text{max}} = 14$, here $N_{\text{max}} = 8$
- phase-shifts with full NCSM and IT-NCSM input agree
- 2 bound states for ${}^8\text{Li}$
- 4 resonances: 3^+ and 1^+ are known, 0^+ and 2^+ resonances are predictions

Many-Body Approximations

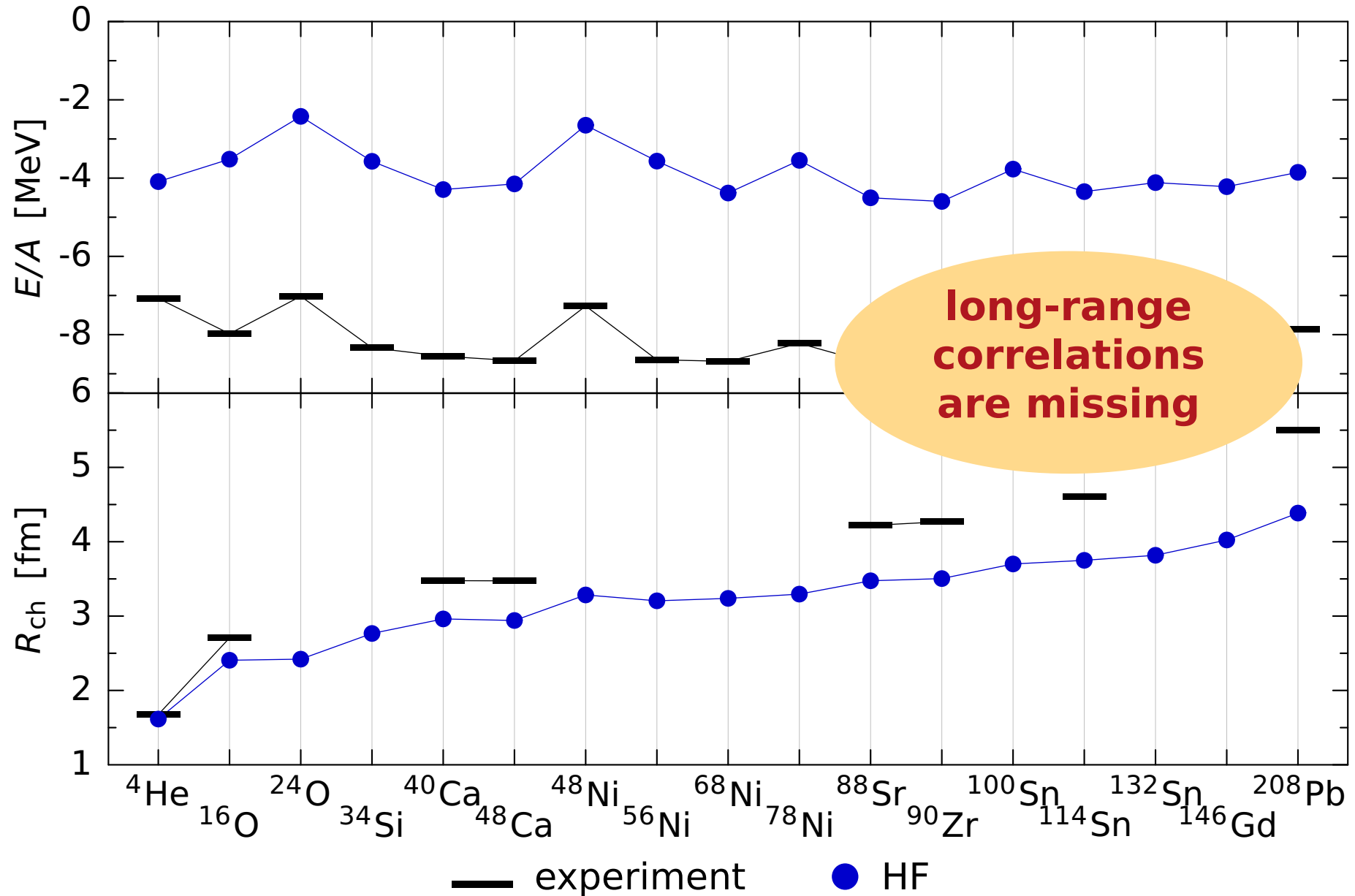
Other Options...

Other Options...

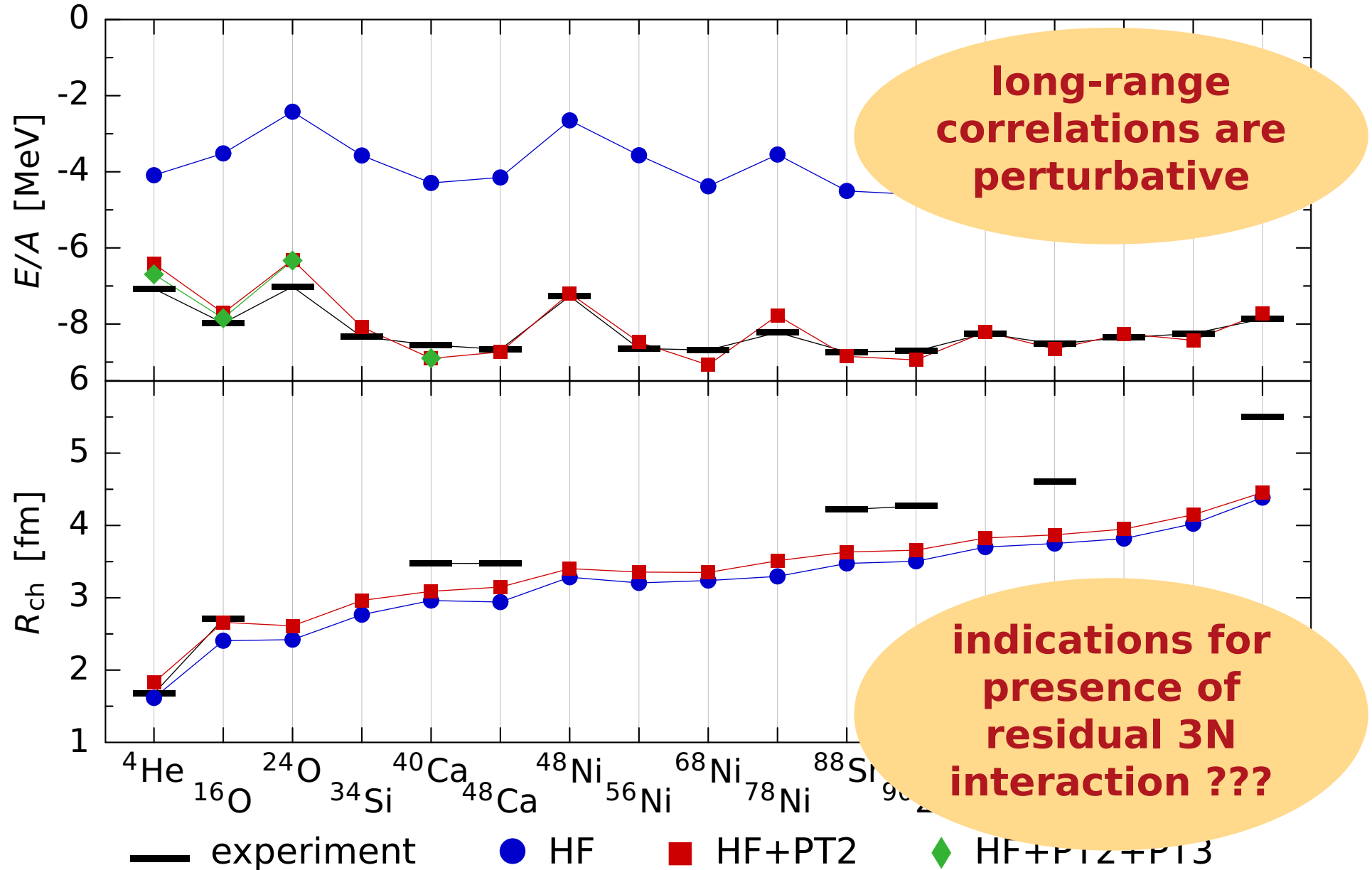
**similarity transformed
interactions (e.g. V_{UCOM} , V_{SRG})
provide universal input for
various many-body methods**

- coupled-cluster method
- Hartree-Fock & many-body perturbation theory
- Hartree-Fock-Bogoliubov & particle-number projection
- RPA & Second-RPA
- FMD with projection & configuration mixing

Hartree-Fock with V_{UCOM}



Perturbation Theory with V_{UCOM}



Conclusions

- three steps from QCD to the nuclear chart
 - QCD-based nuclear interactions
 - similarity transformed interactions (UCOM, SRG,...)
 - exact & approximate many-body methods
- exciting new developments in all three sectors

**QCD-based description of
nuclear structure across
the whole nuclear chart is
within reach**

Epilogue

■ thanks to my group & my collaborators

- S. Binder, A. Calci, B. Erler, A. Günther, M. Hild, H. Krutsch, J. Langhammer, P. Papakonstantinou, S. Reinhardt, F. Schmitt
Institut für Kernphysik, TU Darmstadt
- P. Navrátil, S. Quaglioni
Lawrence Livermore National Laboratory, USA
- H. Hergert, P. Piecuch, J. Gour
Michigan State University, USA
- C. Forssén
Chalmers University of Technology, Sweden
- H. Feldmeier, T. Neff,...
Gesellschaft für Schwerionenforschung (GSI)

Deutsche
Forschungsgemeinschaft
DFG



 **LOEWE** – Landes-Offensive
zur Entwicklung Wissenschaftlich-
ökonomischer Exzellenz