

Ab Initio Nuclear Structure beyond the p-Shell

Importance Truncated No-Core Shell Model



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- Motivation
- Modern Effective Interactions
 - Unitary Correlation Operator Method
 - Similarity Renormalization Group
- Innovative Many-Body Methods
 - No-Core Shell Model
 - Importance Truncated NCSM
- Perspectives

From QCD to Nuclear Structure

Nuclear Structure

Low-Energy QCD

From QCD to Nuclear Structure

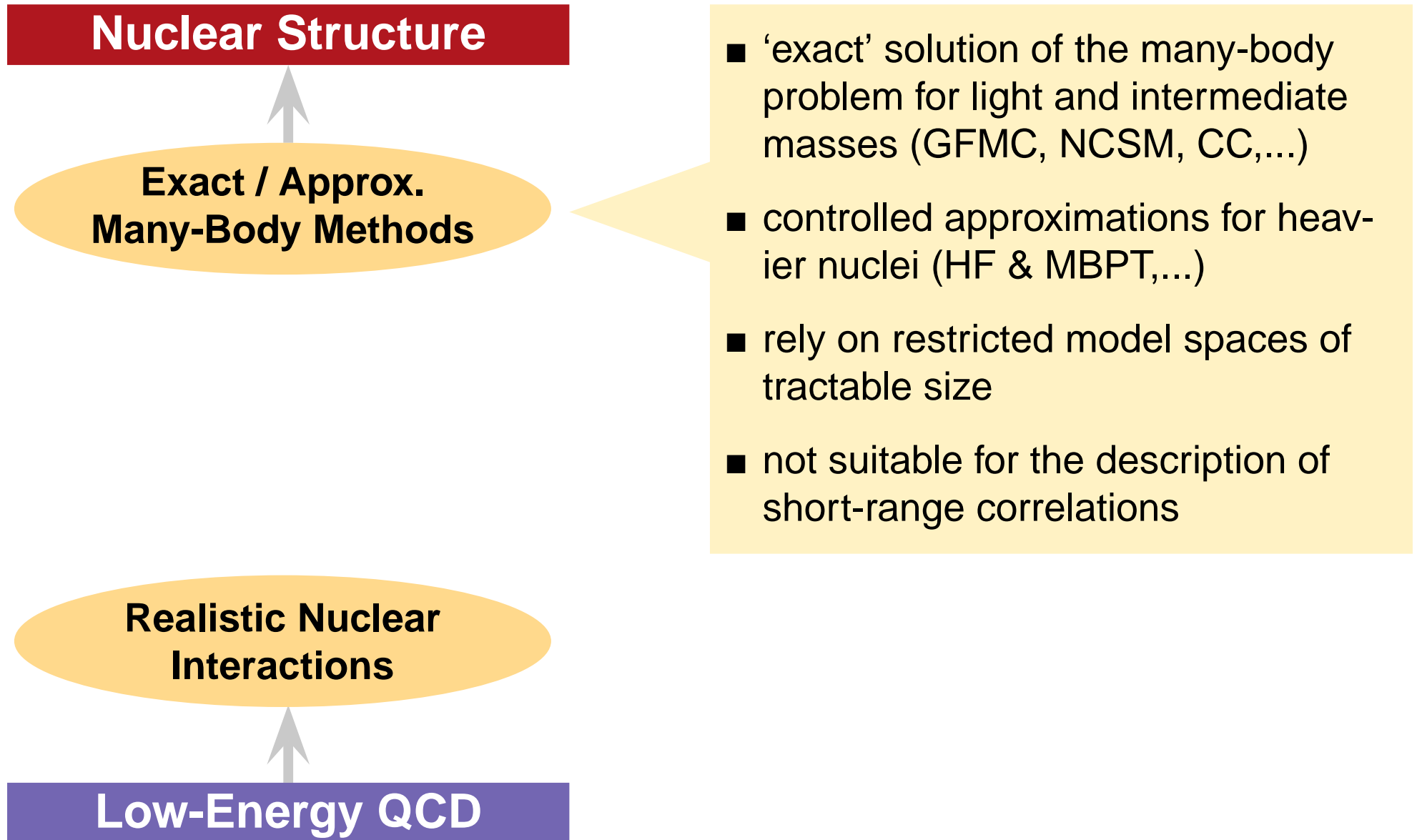
Nuclear Structure

Realistic Nuclear Interactions

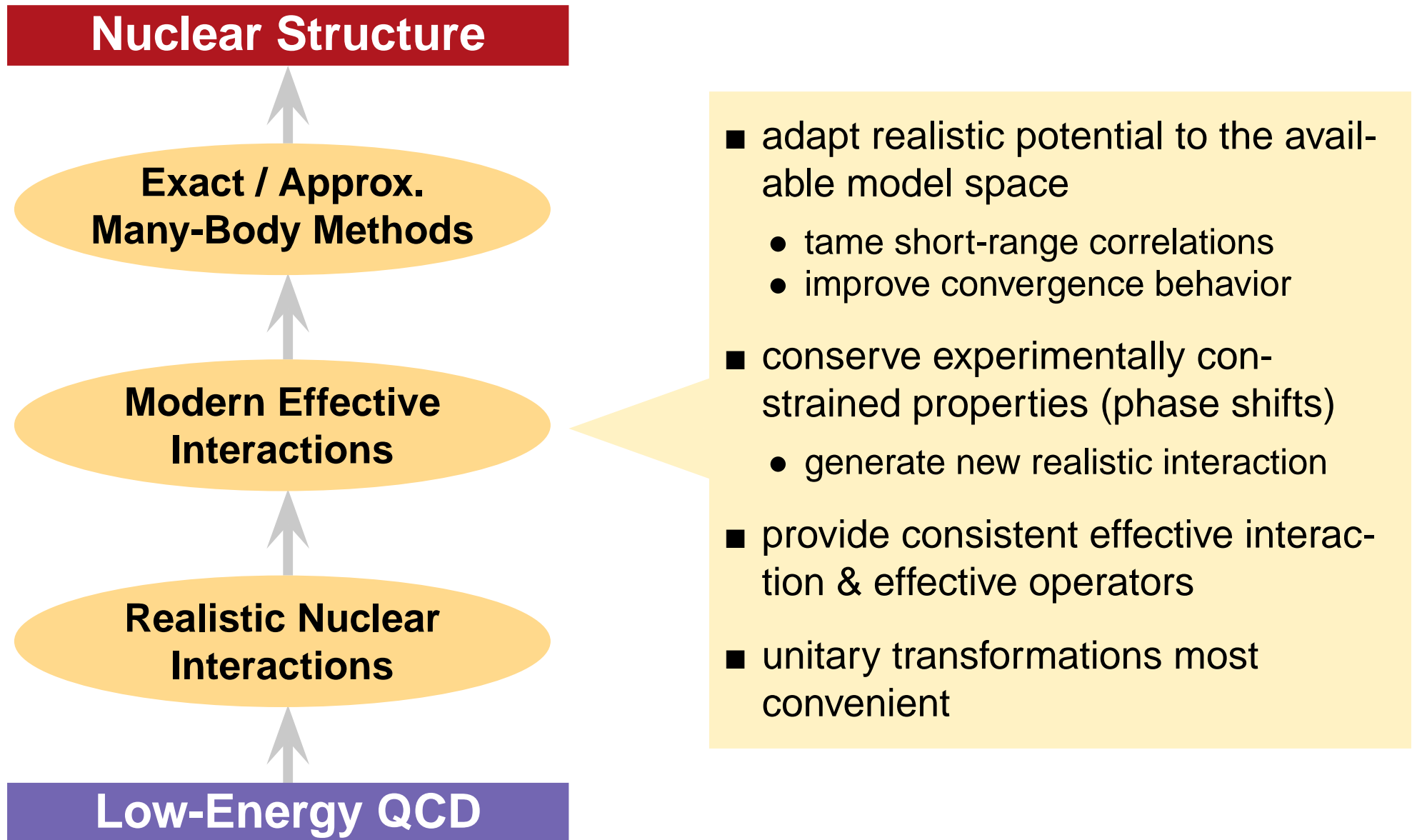
Low-Energy QCD

- chiral interactions: consistent NN & 3N interaction derived within χ EFT
- traditional NN-interactions: Argonne V18, CD Bonn,...
- reproduce experimental NN phase-shifts with high precision
- induce strong short-range central & tensor correlations

From QCD to Nuclear Structure



From QCD to Nuclear Structure



Modern Effective Interactions

Unitary Correlation Operator Method (UCOM)

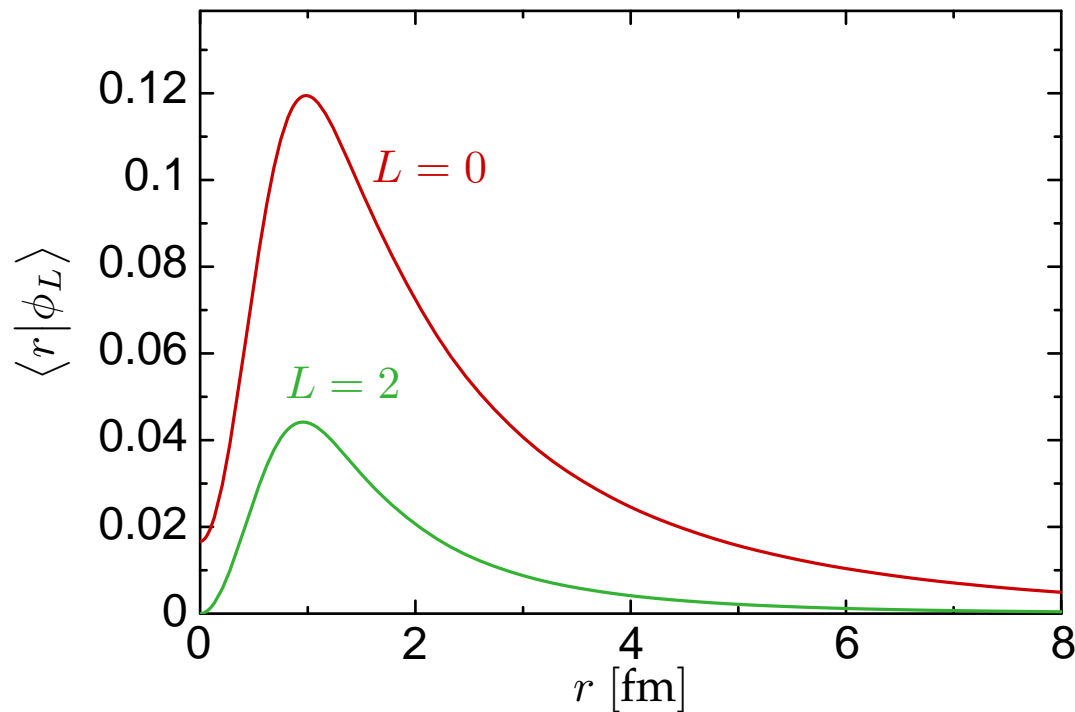
H. Feldmeier et al. — Nucl. Phys. A 632 (1998) 61

T. Neff et al. — Nucl. Phys. A713 (2003) 311

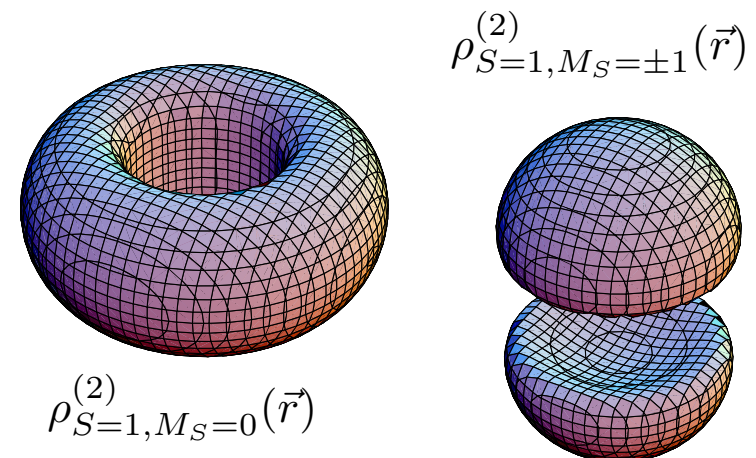
R. Roth et al. — Nucl. Phys. A 745 (2004) 3

R. Roth et al. — Phys. Rev. C 72, 034002 (2005)

Deuteron: Manifestation of Correlations



■ **exact deuteron solution**
for Argonne V18 potential



short-range repulsion
suppresses wavefunction at
small distances r

central correlations

tensor interaction
generates D-wave admixture
in the ground state

tensor correlations

Unitary Correlation Operator Method

explicit ansatz for unitary transformation operator **motivated by the physics of short-range central & tensor correlations**

Central Correlator C_r

- radial distance-dependent shift in the relative coordinate of a nucleon pair

$$\mathbf{g}_r = \frac{1}{2} [s(\mathbf{r}) \mathbf{q}_r + \mathbf{q}_r s(\mathbf{r})]$$

$$\mathbf{q}_r = \frac{1}{2} \left[\frac{\vec{r}}{r} \cdot \vec{q} + \vec{q} \cdot \frac{\vec{r}}{r} \right]$$

Tensor Correlator C_Ω

- angular shift depending on the orientation of spin and relative coordinate of a nucleon pair

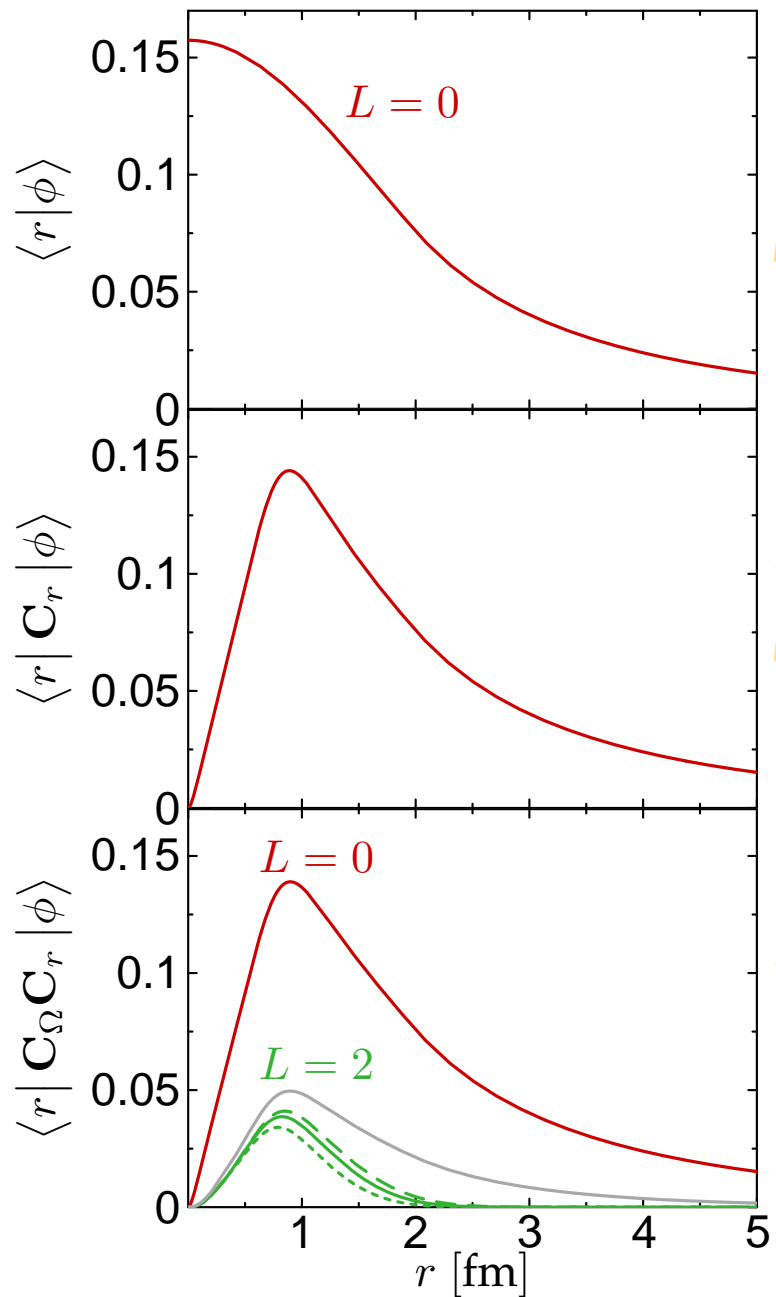
$$\mathbf{g}_\Omega = \frac{3}{2} \vartheta(\mathbf{r}) [(\vec{\sigma}_1 \cdot \vec{q}_\Omega)(\vec{\sigma}_2 \cdot \vec{r}) + (\vec{r} \leftrightarrow \vec{q}_\Omega)]$$

$$\vec{q}_\Omega = \vec{q} - \frac{\vec{r}}{r} \mathbf{q}_r$$

$$\mathbf{C} = \mathbf{C}_\Omega \mathbf{C}_r = \exp\left(-i \sum_{i < j} \mathbf{g}_{\Omega,ij}\right) \exp\left(-i \sum_{i < j} \mathbf{g}_{r,ij}\right)$$

- $s(r)$ and $\vartheta(r)$ for given potential determined by constrained energy minimization in the two-body system (for each S, T)

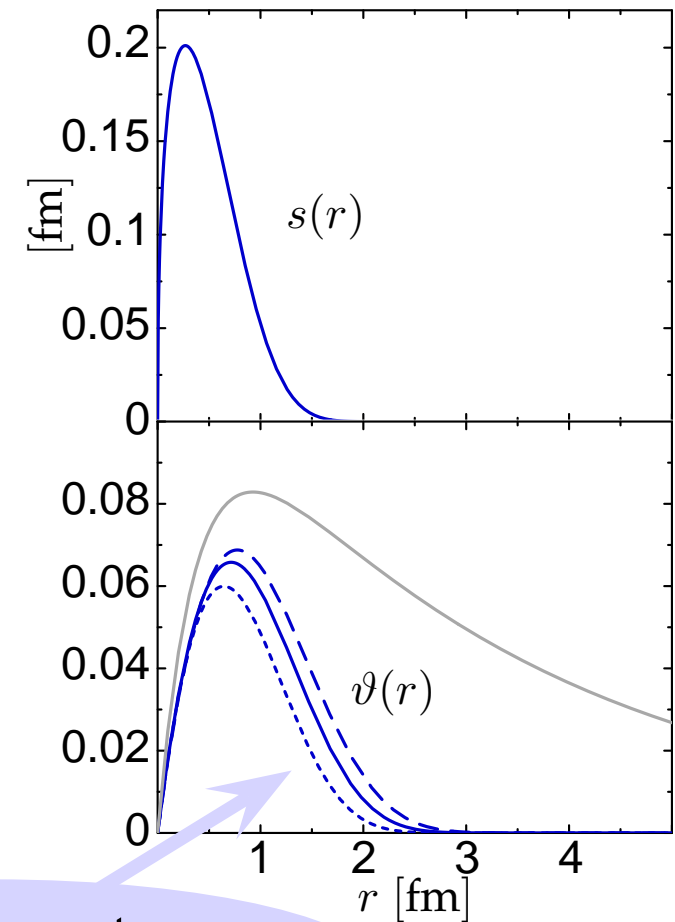
Correlated States: The Deuteron



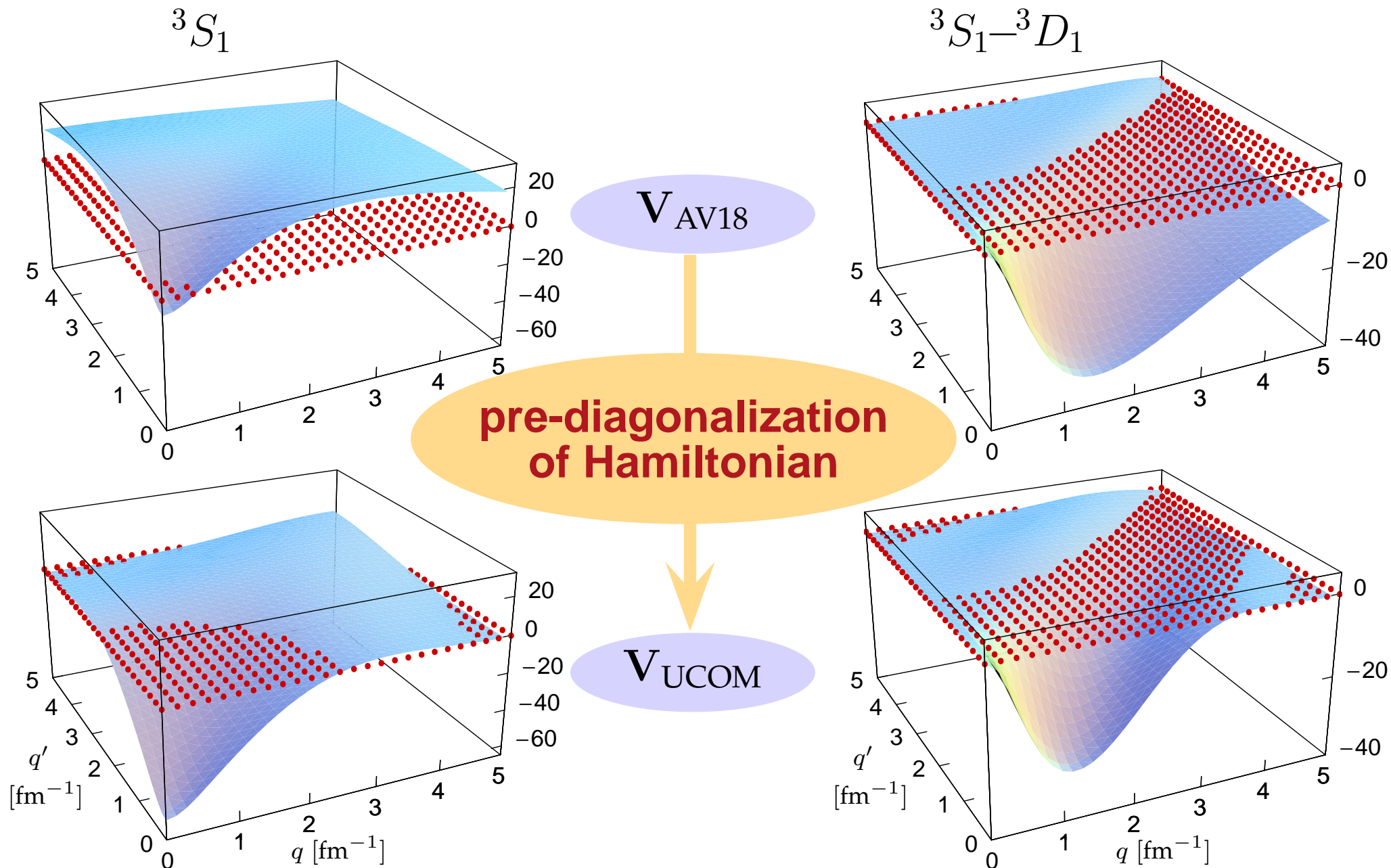
central correlations

tensor correlations

only short-range tensor correlations treated by C_{Ω}



Correlated Interaction: V_{UCOM}



Modern Effective Interactions

Similarity Renormalization Group (SRG)

Hergert & Roth — Phys. Rev. C 75, 051001(R) (2007)

Bogner et al. — Phys. Rev. C 75, 061001(R) (2007)

Roth, Reinhardt, Hergert — arXiv:0802.4239

Similarity Renormalization Group

unitary transformation of the **Hamiltonian to a band-diagonal form** with respect to a given uncorrelated many-body basis

Flow Equation for Hamiltonian

- evolution equation for Hamiltonian

$$\tilde{\mathbf{H}}(\alpha) = \mathbf{C}^\dagger(\alpha) \mathbf{H} \mathbf{C}(\alpha) \quad \rightarrow \quad \frac{d}{d\alpha} \tilde{\mathbf{H}}(\alpha) = [\boldsymbol{\eta}(\alpha), \tilde{\mathbf{H}}(\alpha)]$$

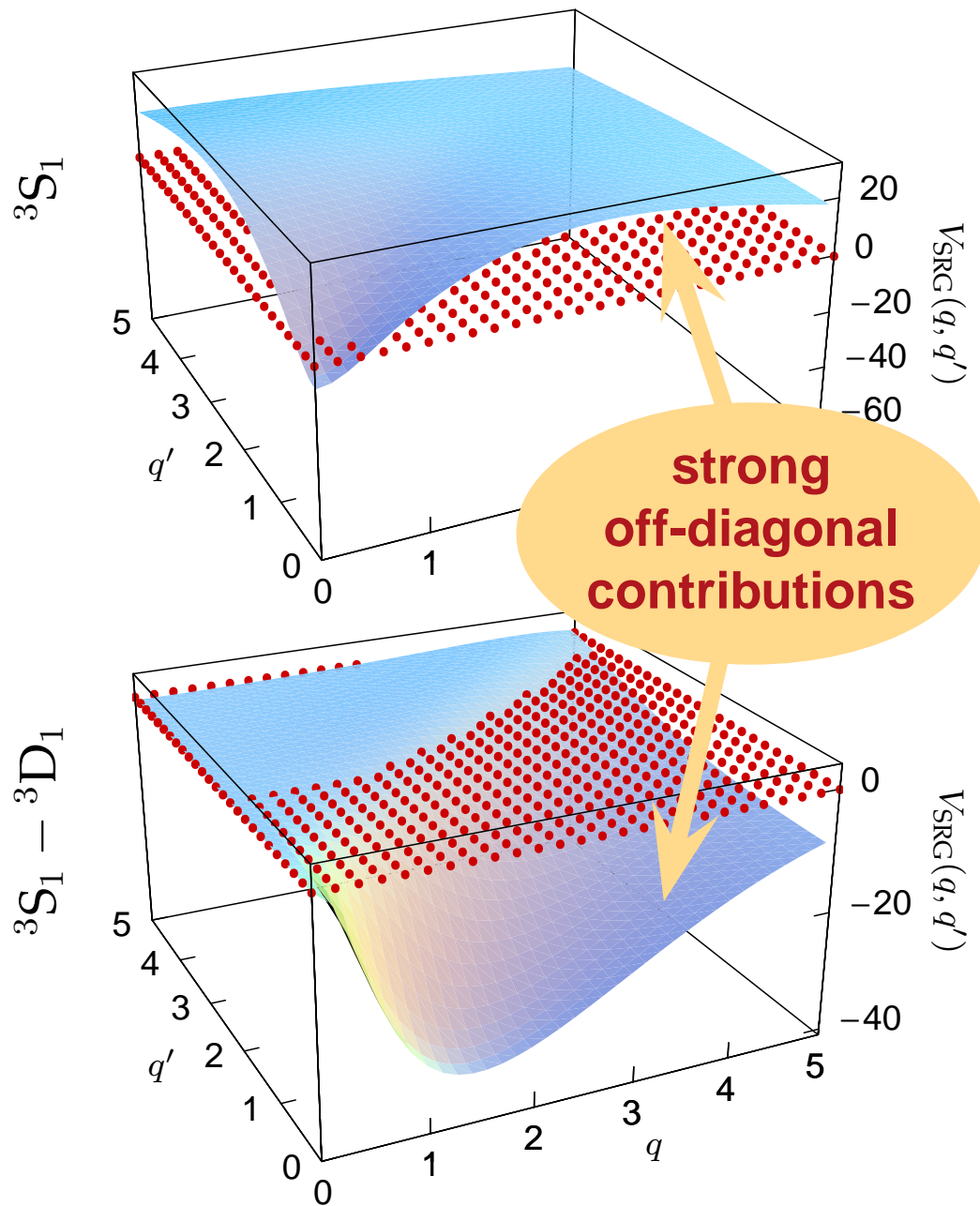
- dynamical generator defined as commutator with the operator in whose eigenbasis \mathbf{H} shall be diagonalized

$$\boldsymbol{\eta}(\alpha) \stackrel{2B}{=} \frac{1}{2\mu} [\vec{\mathbf{q}}^2, \tilde{\mathbf{H}}(\alpha)]$$

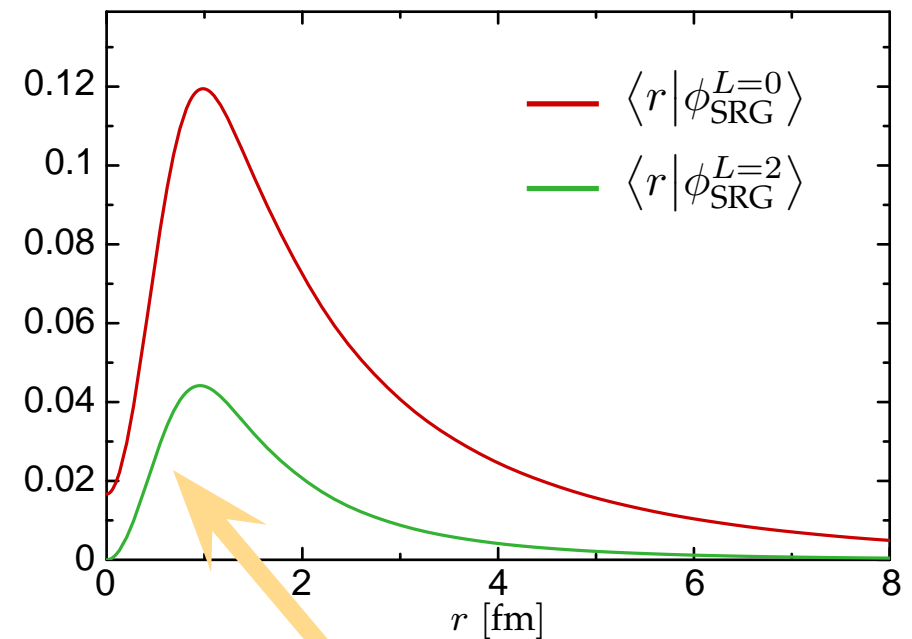
UCOM vs. SRG

$\boldsymbol{\eta}(0)$ has the same structure as the UCOM generators \mathbf{g}_r and \mathbf{g}_Ω

SRG Evolution: The Deuteron

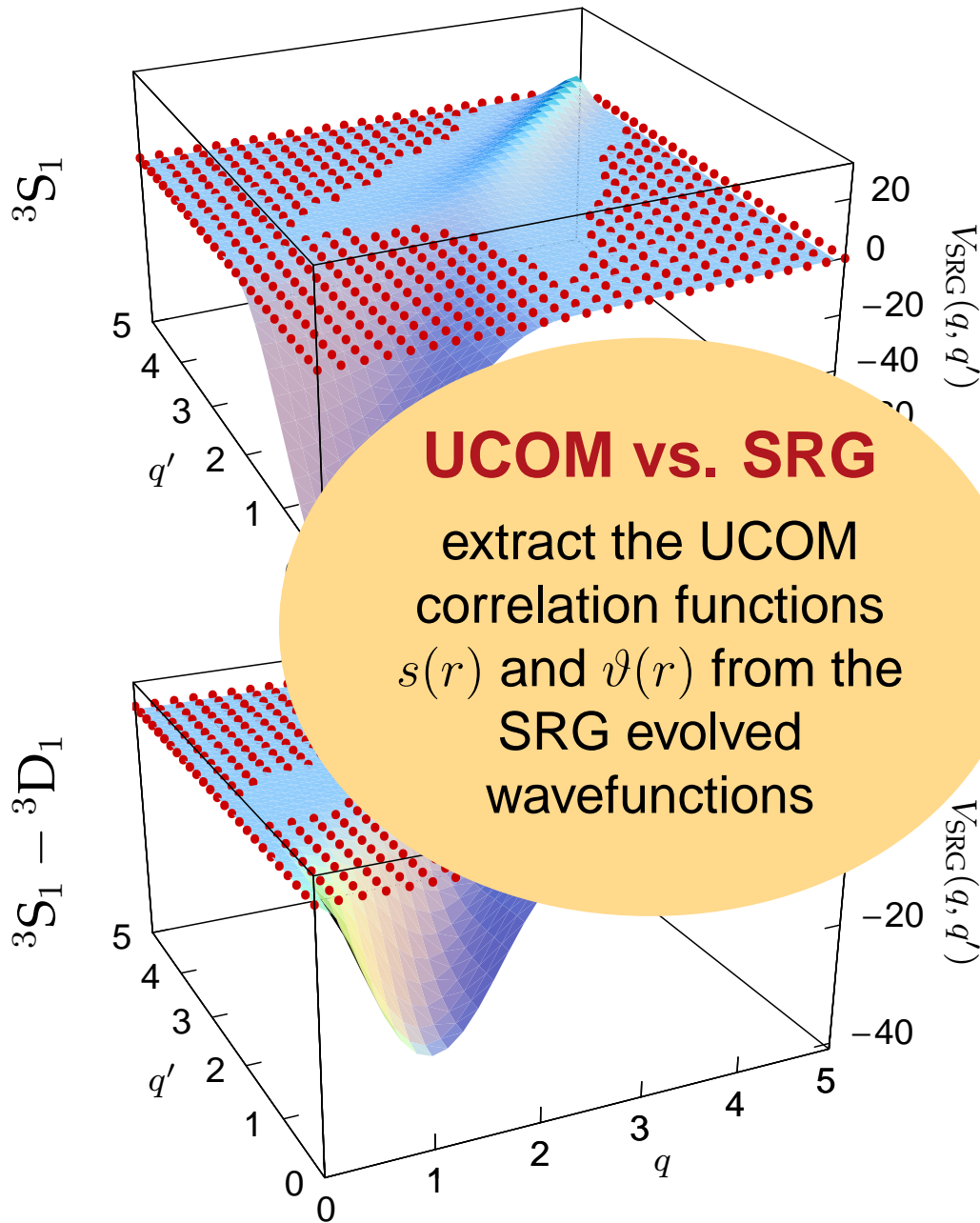


Argonne V18

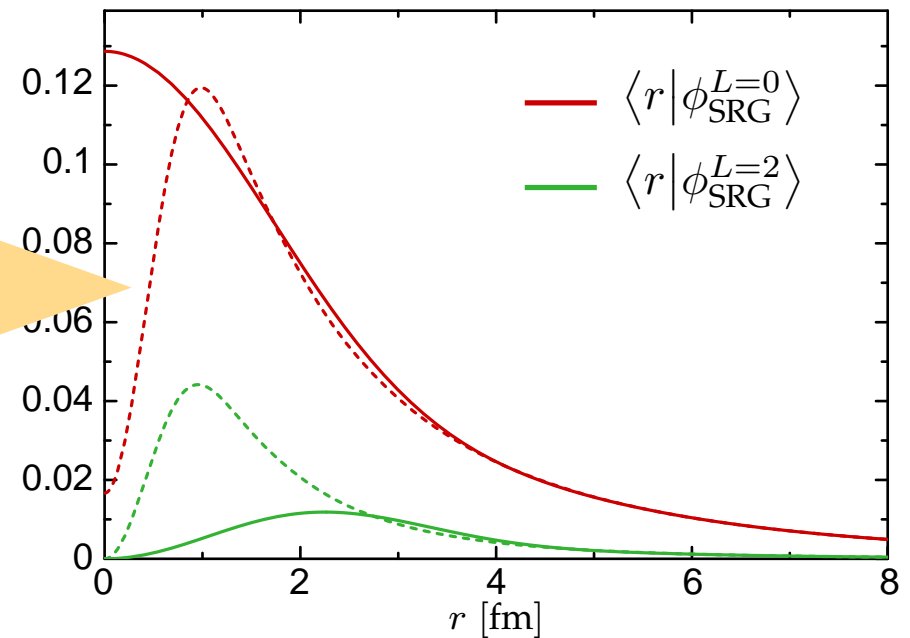


short-range central & tensor correlations

SRG Evolution: The Deuteron



$$\alpha = 0.1000 \text{ fm}^4$$



Exact Many-Body Methods

No-Core Shell Model

Roth et al. — Phys. Rev. C 72, 034002 (2005)

Roth & Navrátil — in preparation

Reminder: No-Core Shell Model

- many-body state is **expanded in Slater determinants** $|\Phi_\nu\rangle$ composed of harmonic oscillator single-particle states

$$|\Psi\rangle = \sum_{\nu} C_{\nu} |\Phi_{\nu}\rangle$$

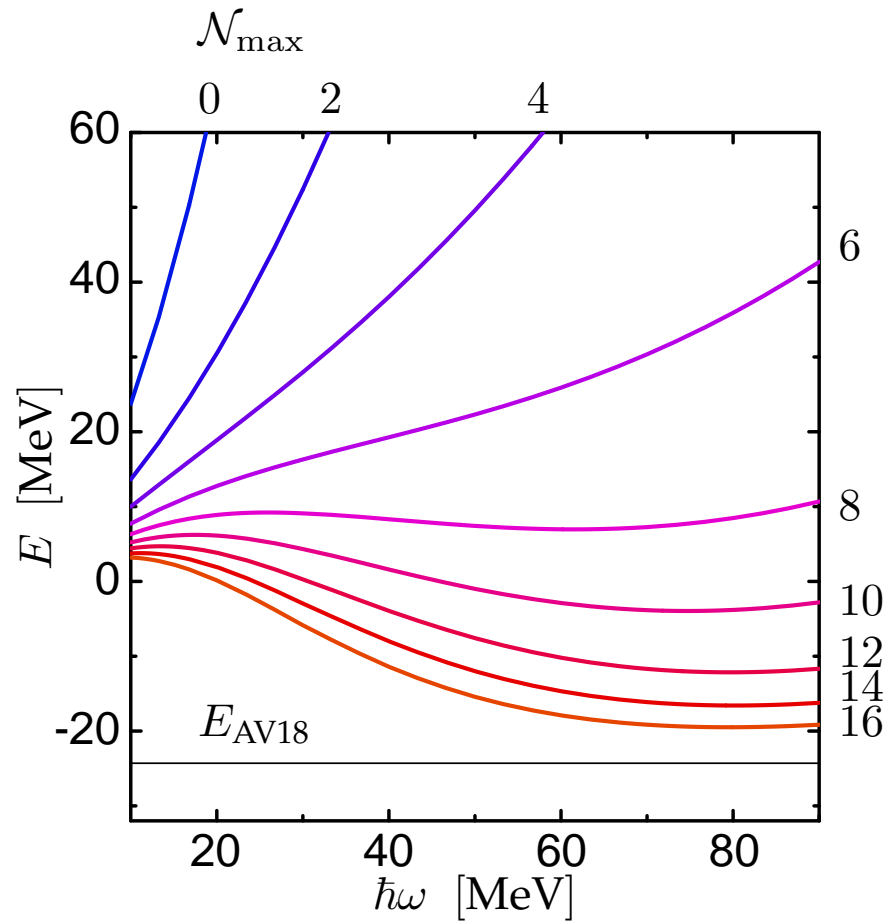
- $\mathcal{N}_{\max} \hbar\omega$ **model space**: truncate basis of Slater determinants with respect to number of oscillator quanta (unperturbed excitation energy)

with increasing model space size more and more **correlations can be described** by the shell model states

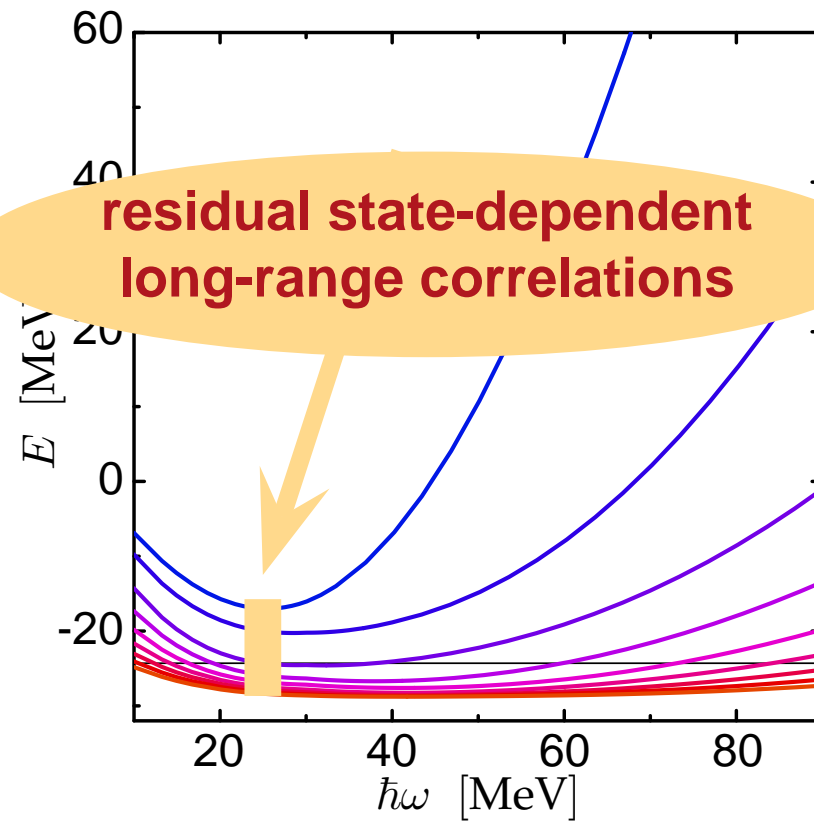
facilitates systematic study of short- and long-range correlations

^4He : Convergence

V_{AV18}

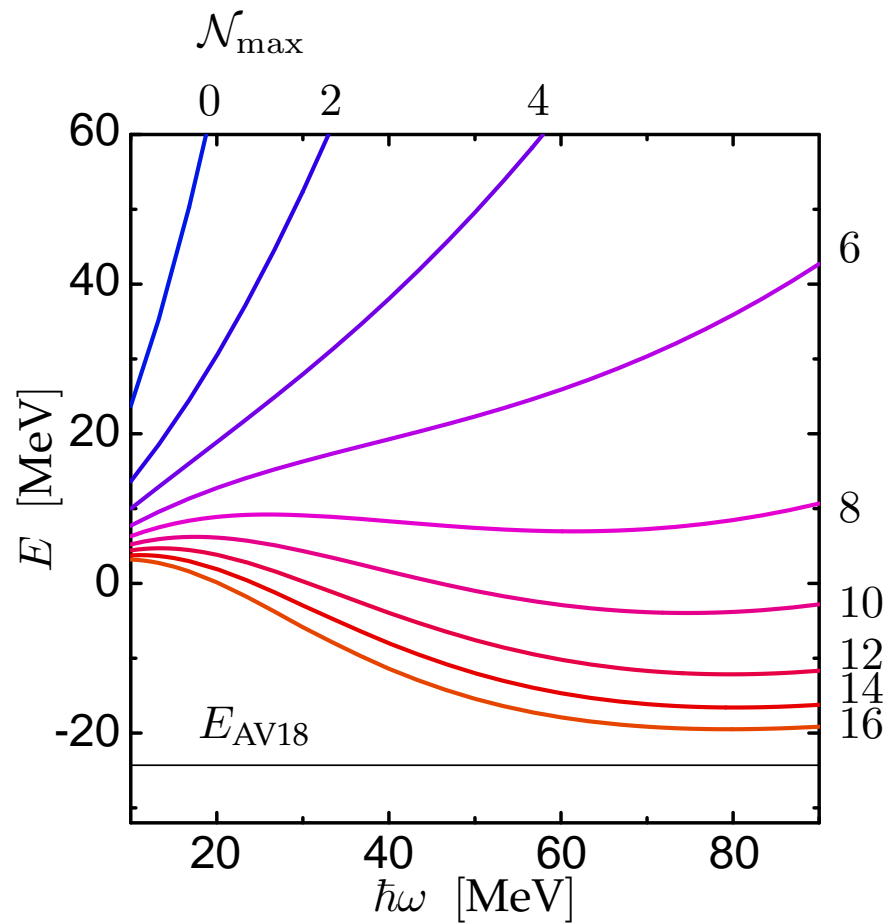


V_{UCOM}

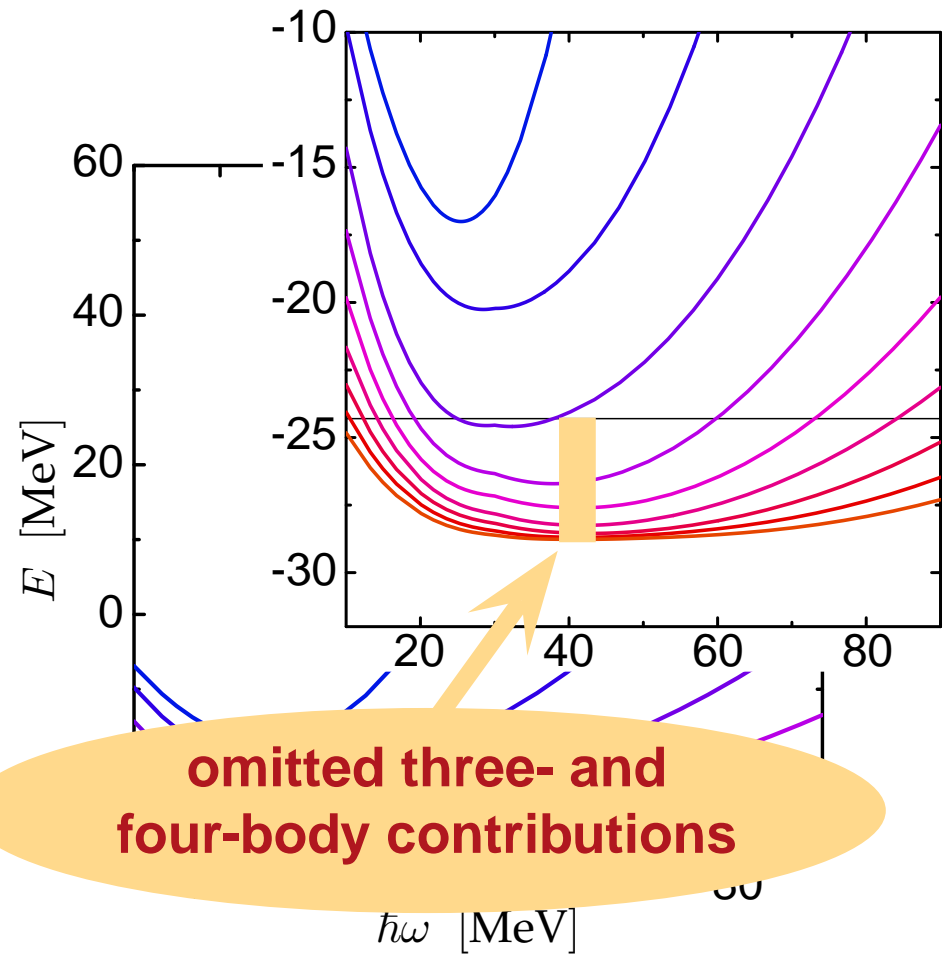


^4He : Convergence

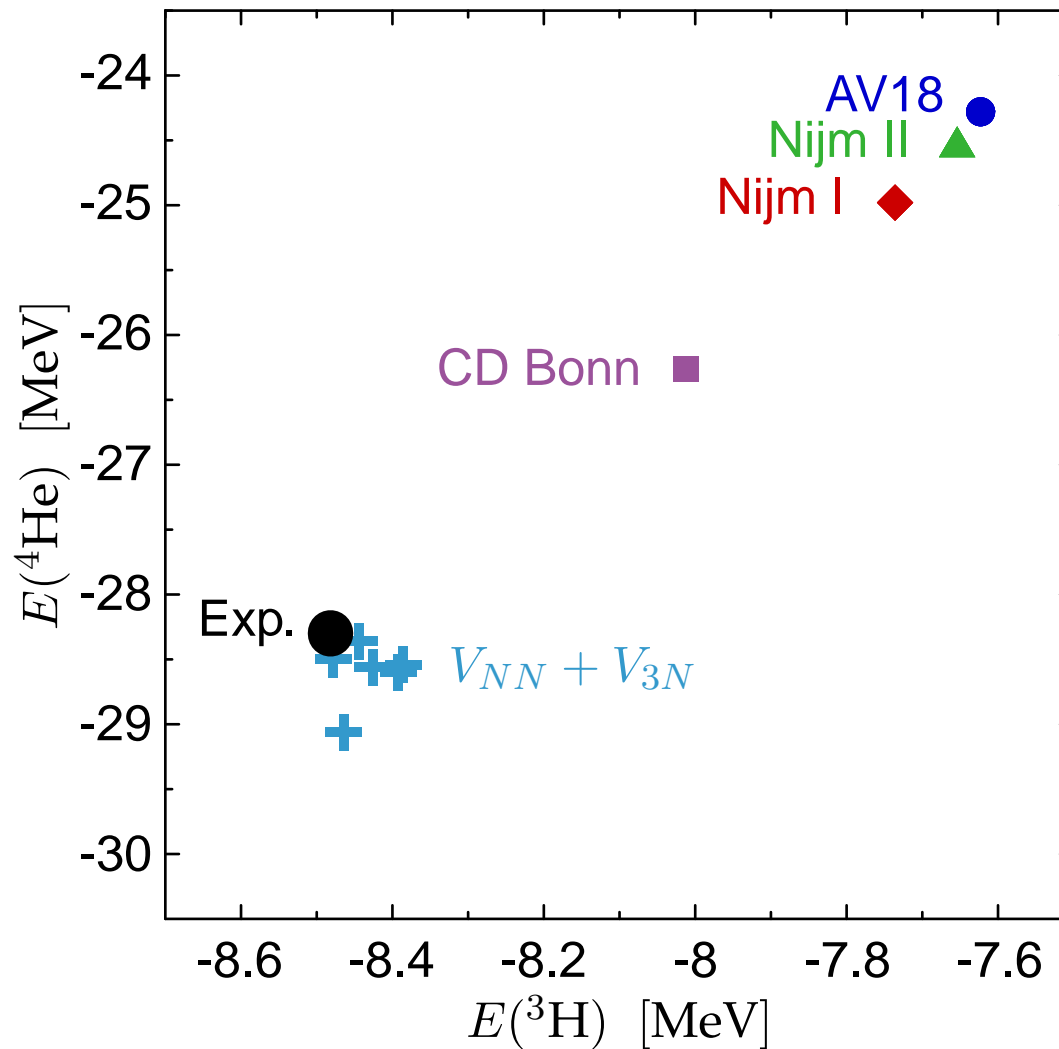
V_{AV18}



V_{UCOM}

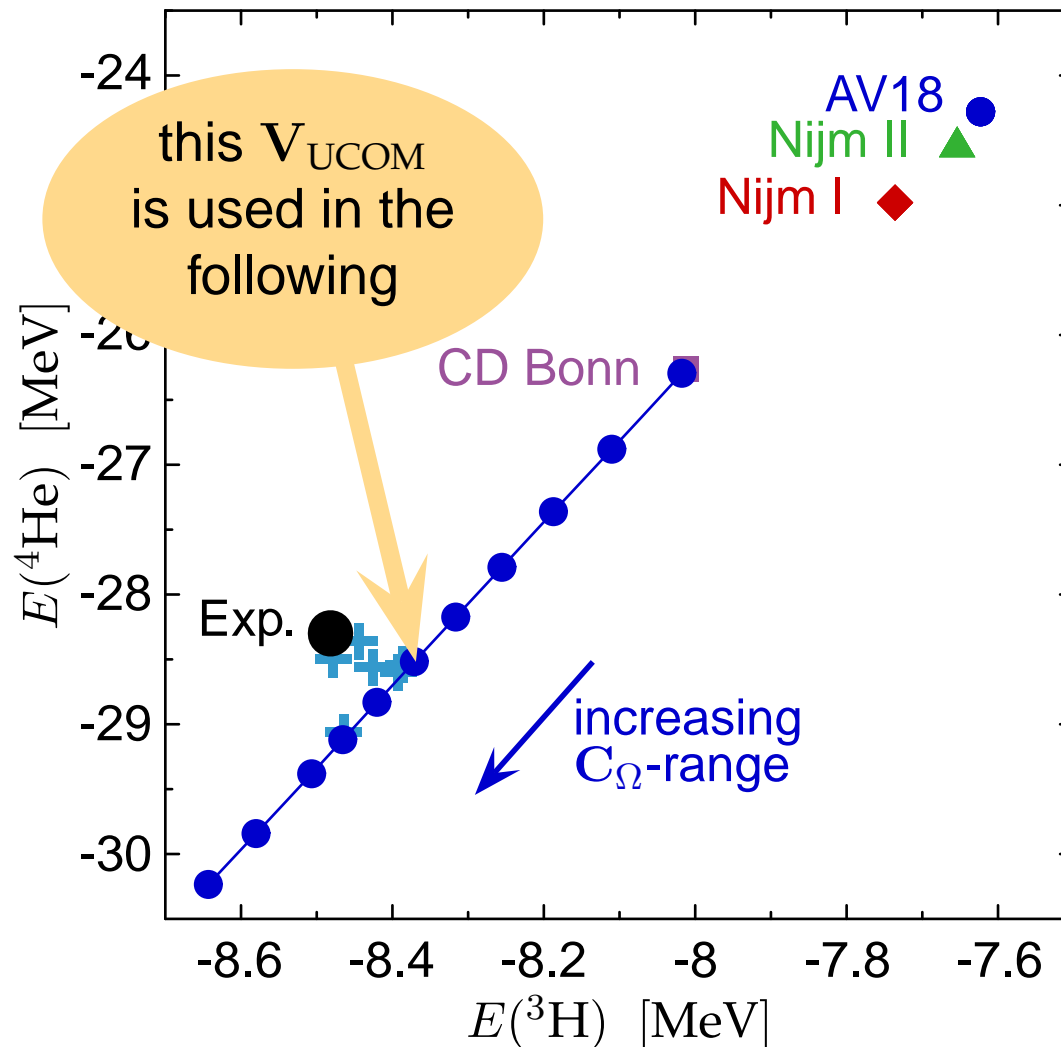


Three-Body Interactions — Tjon Line



- **Tjon-line:** $E({}^4\text{He})$ vs. $E({}^3\text{H})$ for phase-shift equivalent NN-interactions

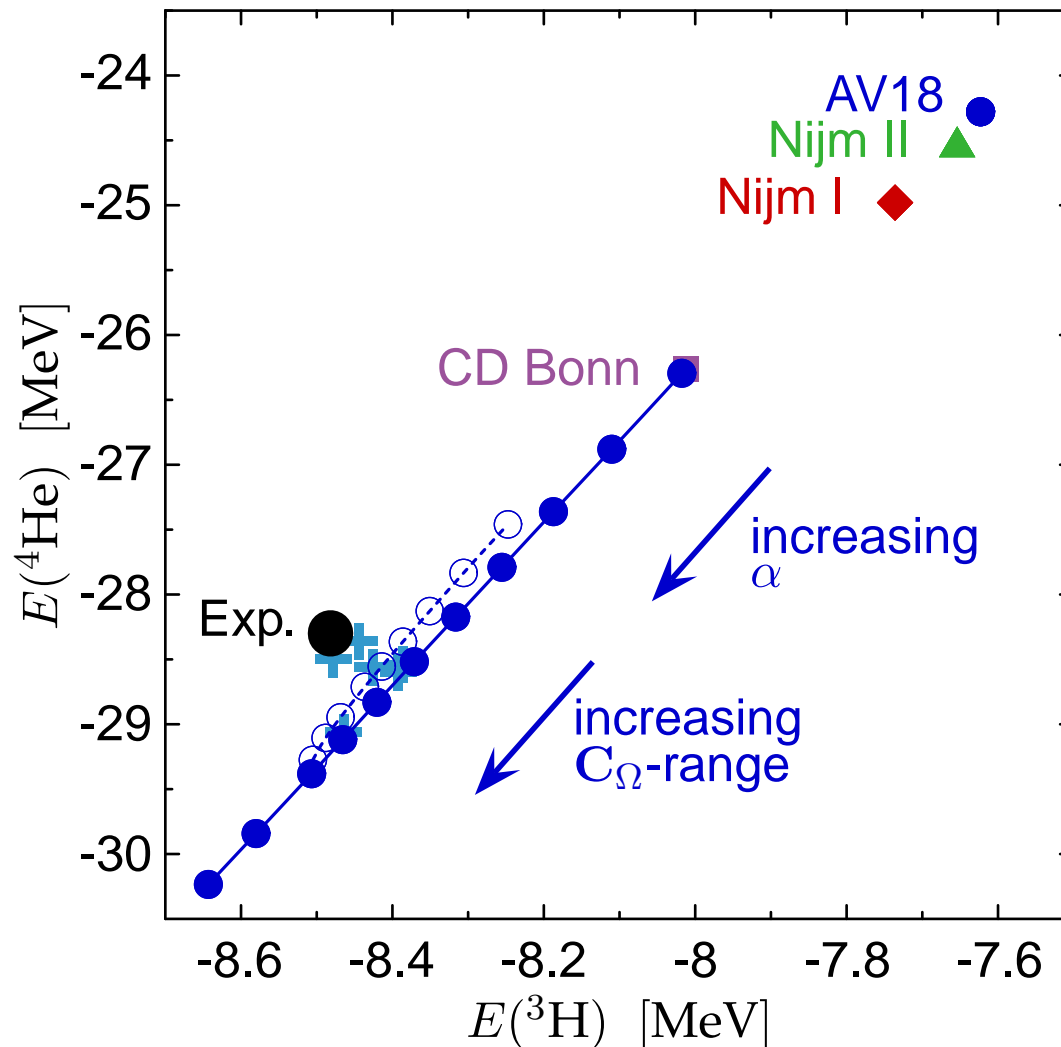
Three-Body Interactions — Tjon Line



- **Tjon-line:** $E(^4\text{He})$ vs. $E(^3\text{H})$ for phase-shift equivalent NN-interactions
- change of C_{Ω} -correlator range results in shift along Tjon-line

minimize net three-body force by choosing correlator with energies close to experimental value

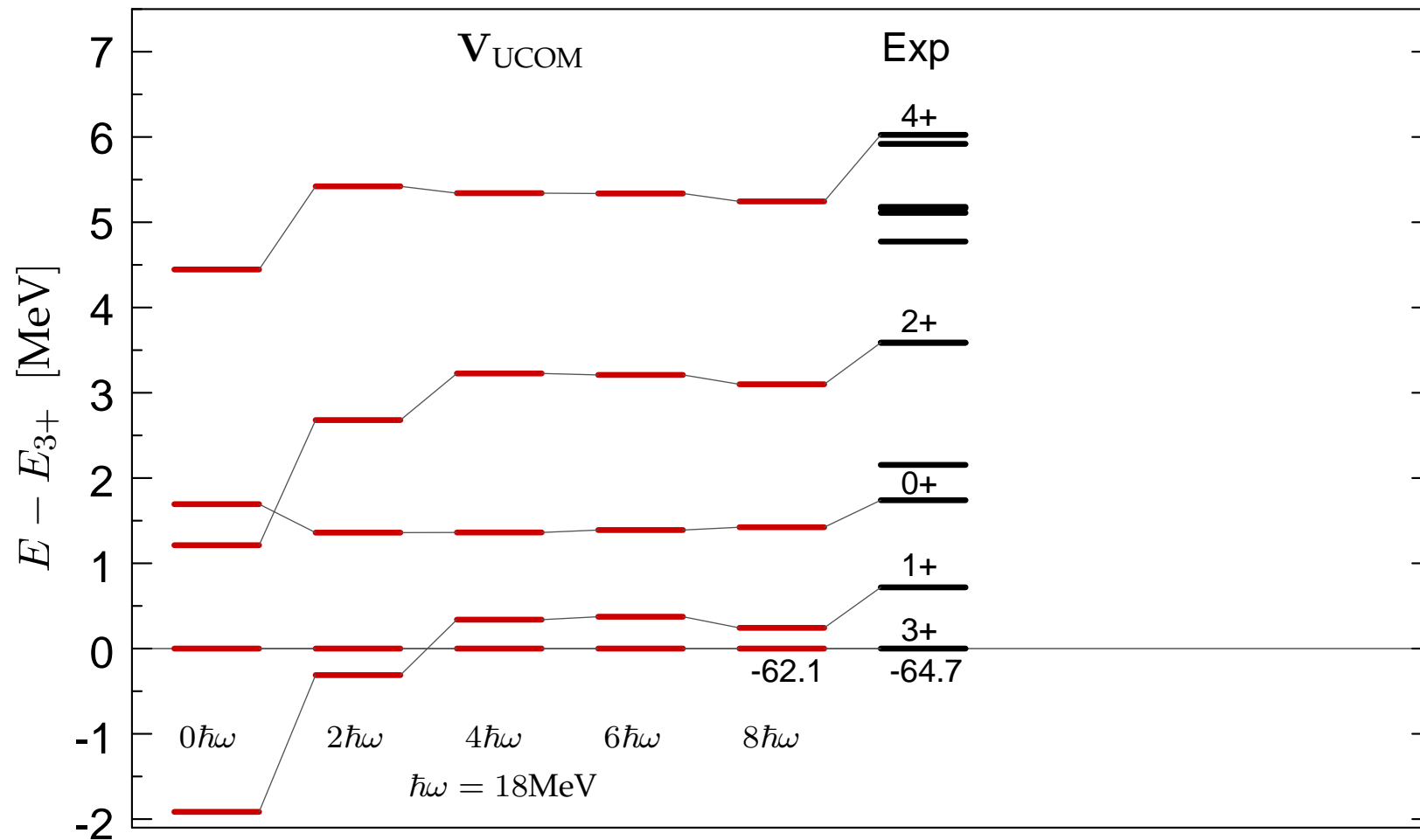
Three-Body Interactions — Tjon Line



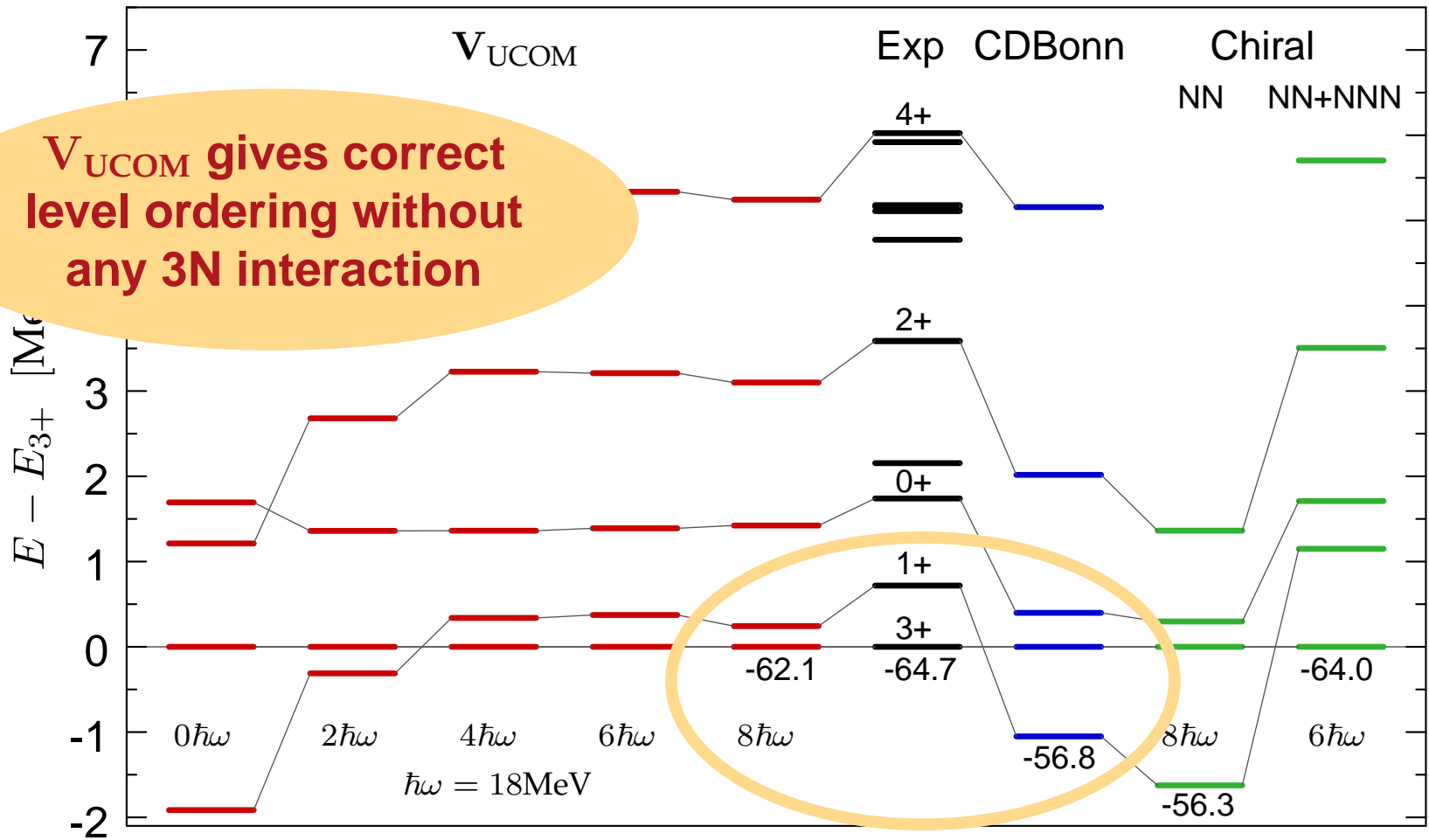
- **Tjon-line:** $E(^4\text{He})$ vs. $E(^3\text{H})$ for phase-shift equivalent NN-interactions
- same behavior for the SRG interaction as function of α

**minimize net
three-body force**
by choosing correlator
with energies close to
experimental value

^{10}B : Hallmark of a 3N Interaction?



^{10}B : Hallmark of a 3N Interaction?



Exact Many-Body Methods

Importance Truncated No-Core Shell Model

Roth & Navrátil — Phys. Rev. Lett. 99, 092501 (2007)

Roth, Piecuch, Gour — in preparation

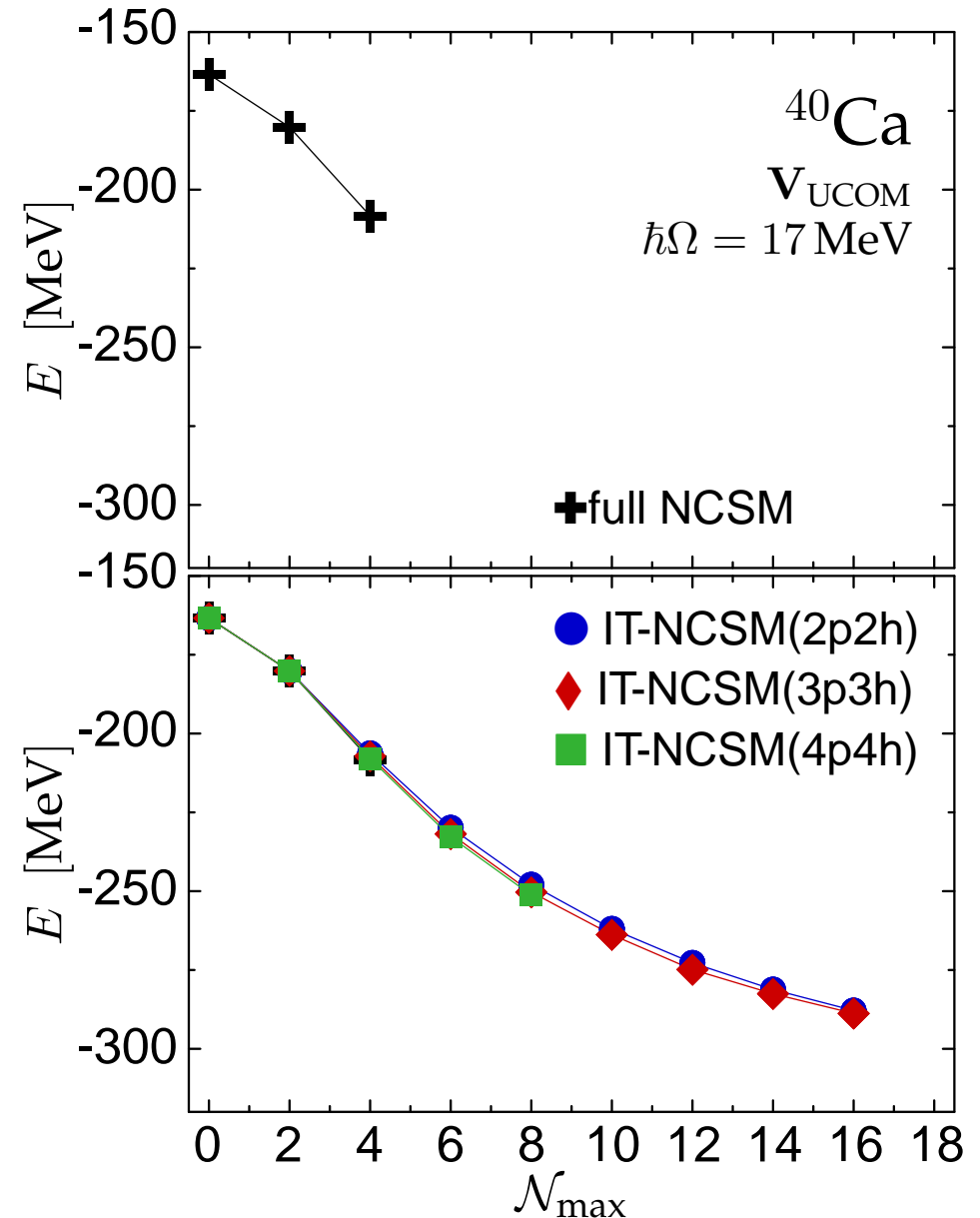
Roth — in preparation

Importance Truncated NCSM

- converged NCSM calculations essentially restricted to p-shell
- full $6\hbar\Omega$ calculation for ^{40}Ca presently not feasible (basis dimension $\sim 10^{10}$)

Importance Truncation

reduce NCSM space to relevant states using an **a priori importance measure** derived from MBPT



Importance Truncation: General Idea

- start with $\mathcal{N}_{\max} \hbar \omega$ **space** of the NCSM

→ separation of intrinsic and center-of-mass component of state

- **importance measure**: identify important basis states $|\Phi_\nu\rangle$ via first-order multiconfigurational perturbation theory

$$\kappa_\nu = - \frac{\langle \Phi_\nu | \mathbf{H} | \Psi_{\text{ref}} \rangle}{\epsilon_\nu - \epsilon_{\text{ref}}}$$

- **importance truncation**: starting from approximation $|\Psi_{\text{ref}}\rangle$ of target state, construct importance truncated space spanned by basis states with $|\kappa_\nu| \geq \kappa_{\text{min}}$

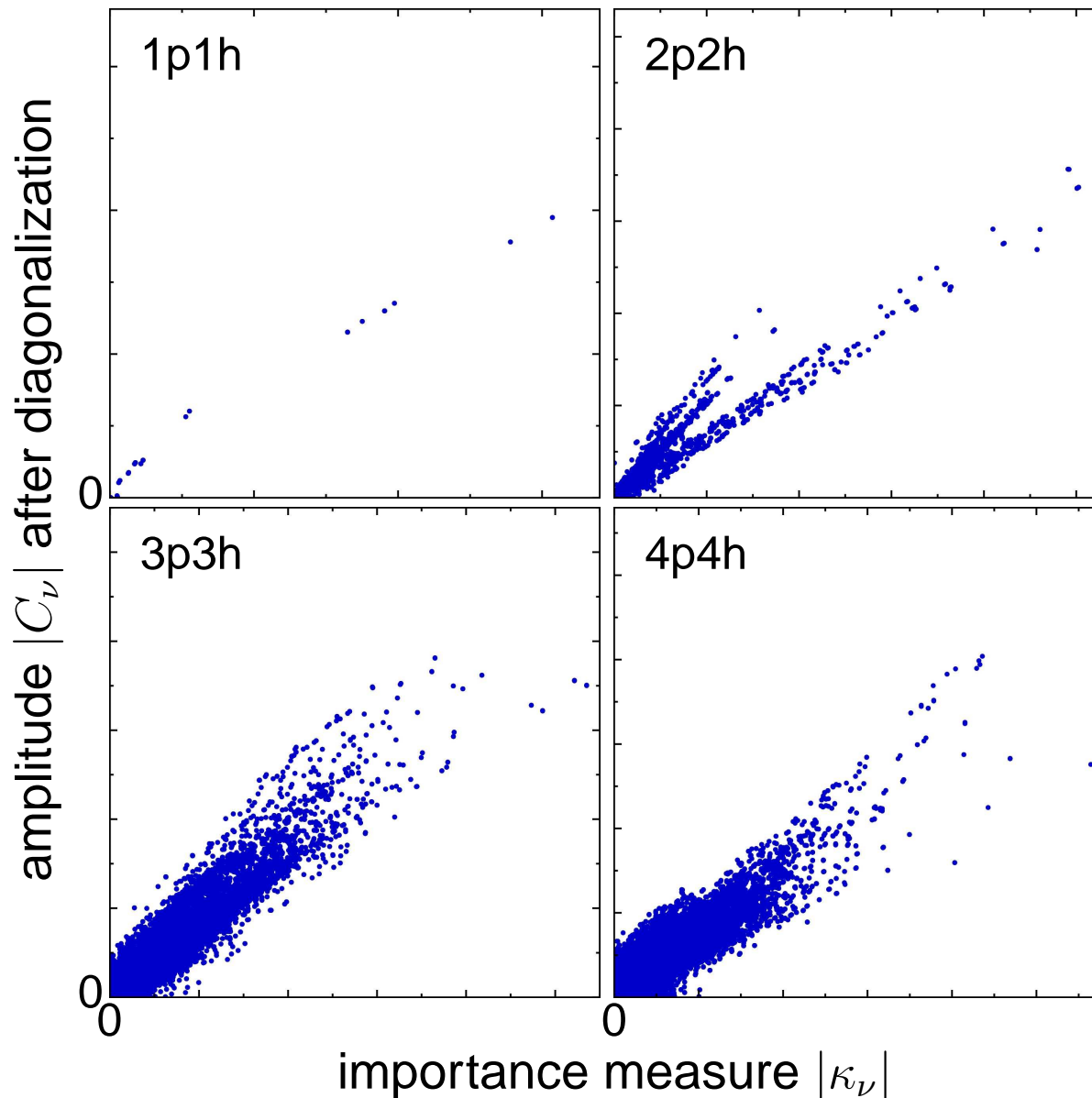
→ contains 2p2h excitations with respect to $|\Psi_{\text{ref}}\rangle$ at most

→ perturbative measure entails $N_p N_h$ hierarchy, i.e., higher-order $N_p N_h$ states only enter in higher orders of PT

Importance Truncation: General Idea

- solve **eigenvalue problem** in importance truncated space
 - rigorous variational upper bound
- **iterative scheme**: repeat construction of importance truncated model space using eigenstate as improved reference $|\Psi_{\text{ref}}\rangle$
 - recovers full $\mathcal{N}_{\text{max}}\hbar\omega$ space after $A/2$ iterations in the limit $\kappa_{\text{min}} \rightarrow 0$
 - typically 2 or 3 iterations to convergence
- **threshold extrapolation**: constrained extrapolation $\kappa_{\text{min}} \rightarrow 0$ of energies recovers contribution of excluded configurations
- **perturbative estimates** and **Davidson corrections** for the contribution of the next iteration also possible

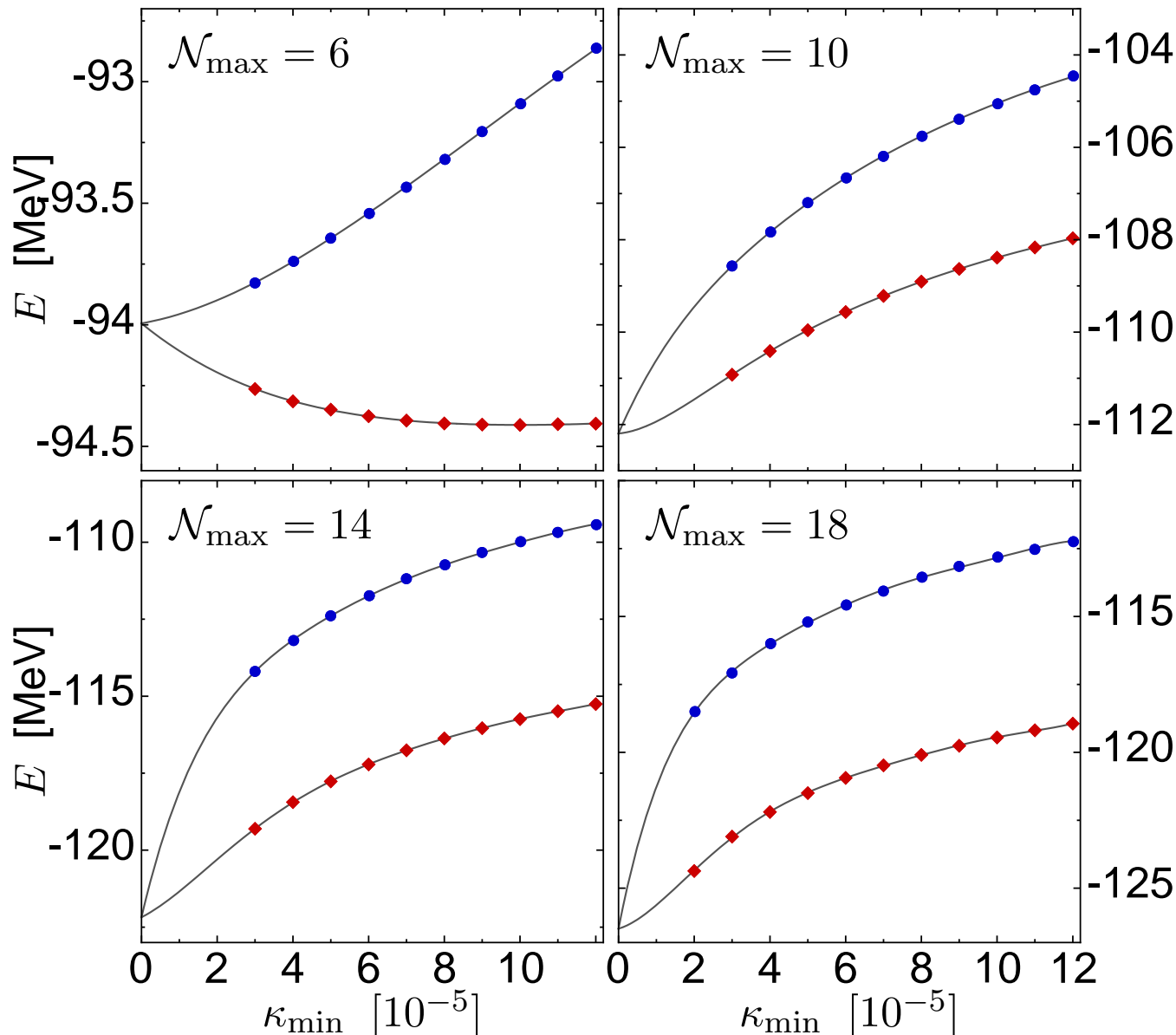
Technicalities: Importance Measure



- importance measure κ_ν provides **reliable a priori estimate** of the a posteriori amplitude C_ν obtained from diagonalization

^{16}O
 V_{UCOM}
 $\hbar\Omega = 22 \text{ MeV}$
 $\mathcal{N}_{\text{max}} = 6$

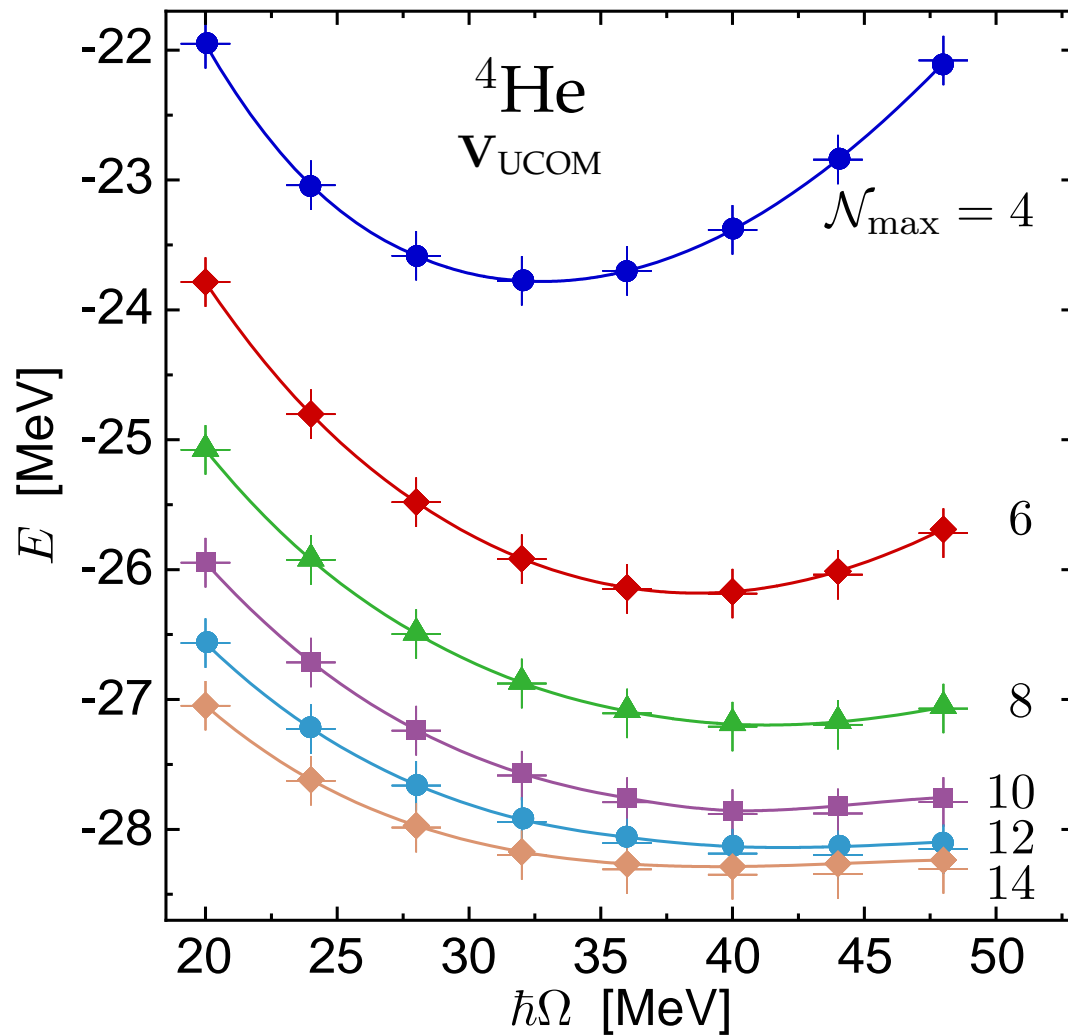
Technicalities: Threshold Extrapolation



- **smooth κ_{\min} -dependence** allows for robust extrapolation
- include perturbative estimate of excluded configurations for simultaneous **constrained extrapolation** $\kappa_{\min} \rightarrow 0$

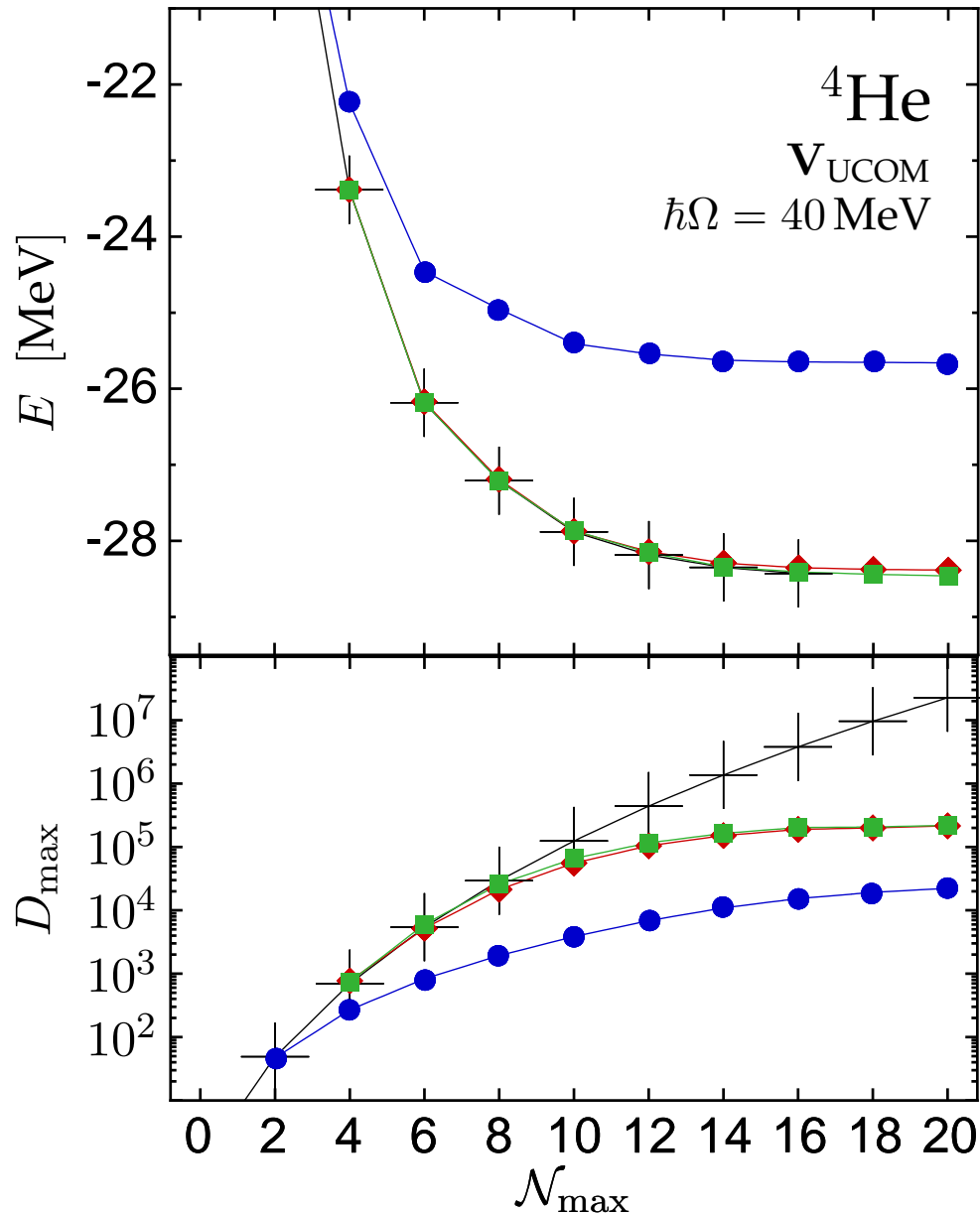
^{16}O
 \mathbf{V}_{UCOM}
 $\hbar\Omega = 22 \text{ MeV}$

^4He : Importance Truncated NCSM



- reproduces exact NCSM result for all $\hbar\omega$ and \mathcal{N}_{max}
 - importance measure and threshold extrapolation are reliable
 - no center-of-mass contamination of states
- + + ... full NCSM
 ● ◆ ... IT-NCSM(2 iter, 4p4h)

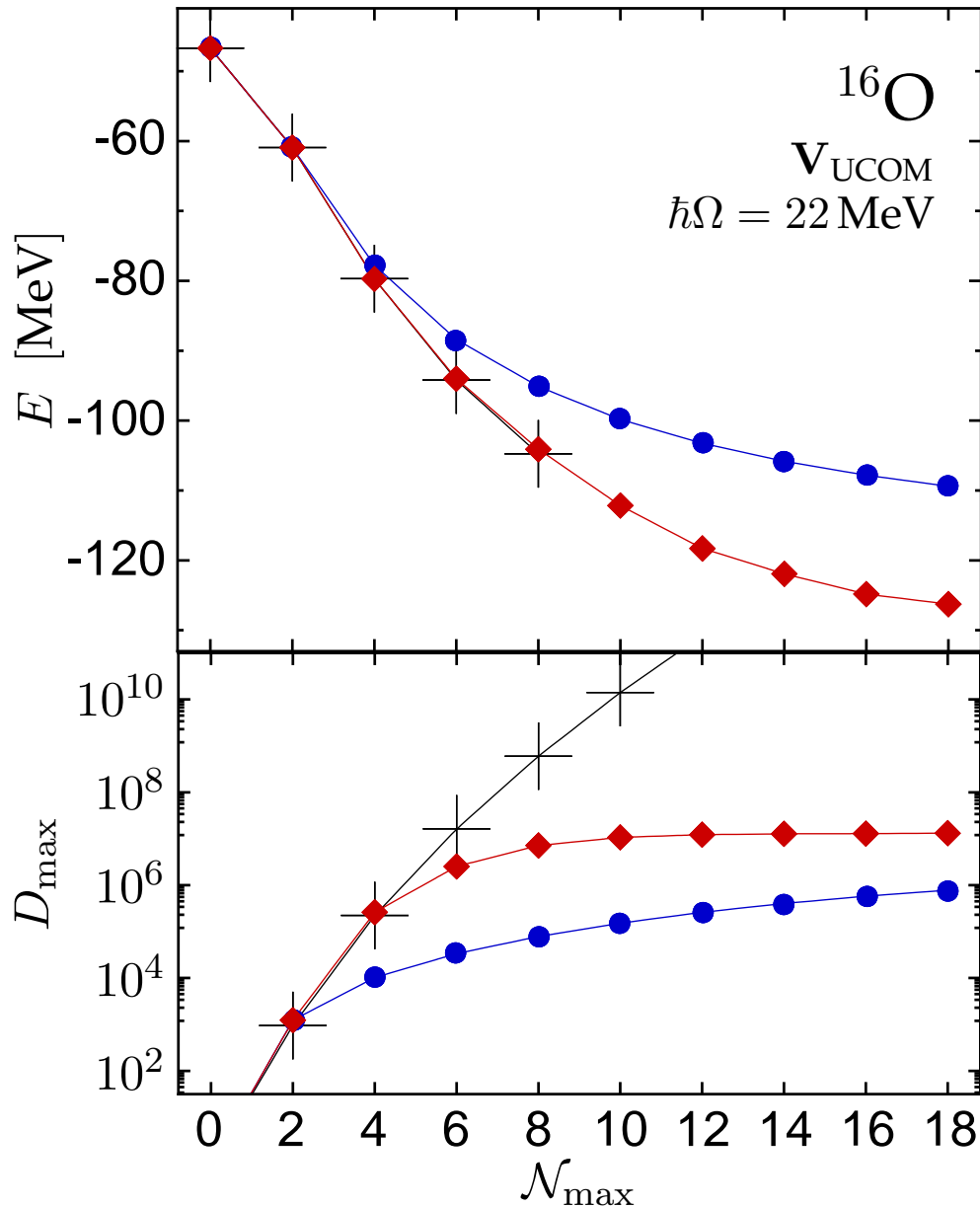
^4He : Importance Truncated NCSM



- reproduces exact NCSM result for all $\hbar\omega$ and \mathcal{N}_{max}
- iterations converge very fast
- reduction of basis by more than two orders of magnitude w/o loss of precision
- saturation of IT-NCSM dimension indicates convergence

- + full NCSM
- IT-NCSM(1 iter, 2p2h)
- ◆ IT-NCSM(2 iter, 4p4h)
- IT-NCSM(3 iter, 4p4h)

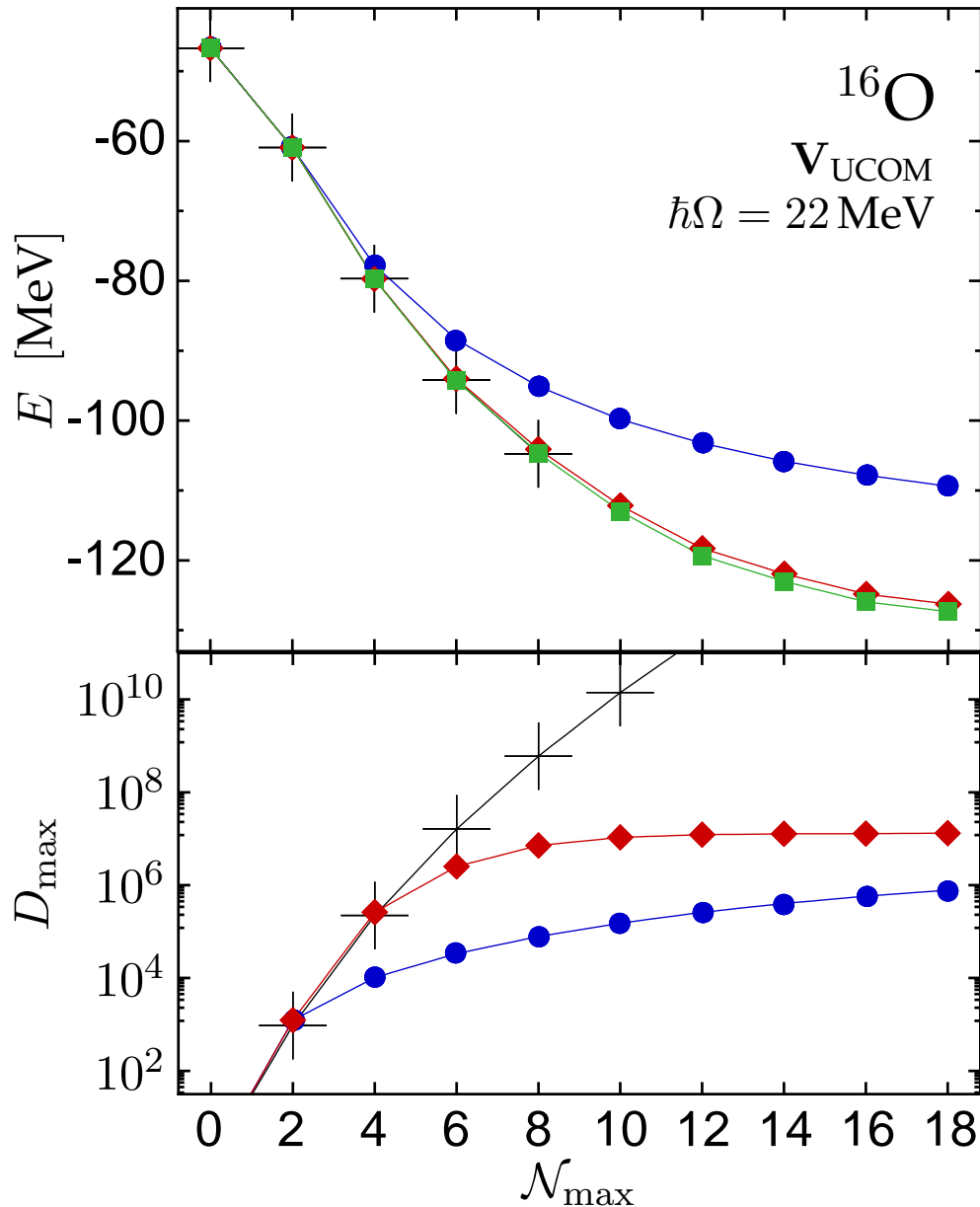
^{16}O : Importance Truncated NCSM



- **excellent agreement with full NCSM** calculation although configurations beyond 4p4h are not included
- dimension reduced by **several orders of magnitude**; possibility to go way beyond the domain of the full NCSM

- + full NCSM
- IT-NCSM(1 iter, 2p2h)
- ◆ IT-NCSM(2 iter, 4p4h)

^{16}O : Importance Truncated NCSM



■ extrapolation to $\mathcal{N}_{\text{max}} \rightarrow \infty$

$$E_{\text{IT-NCSM}(2 \text{ iter})} \approx -129 \pm 1 \text{ MeV}$$

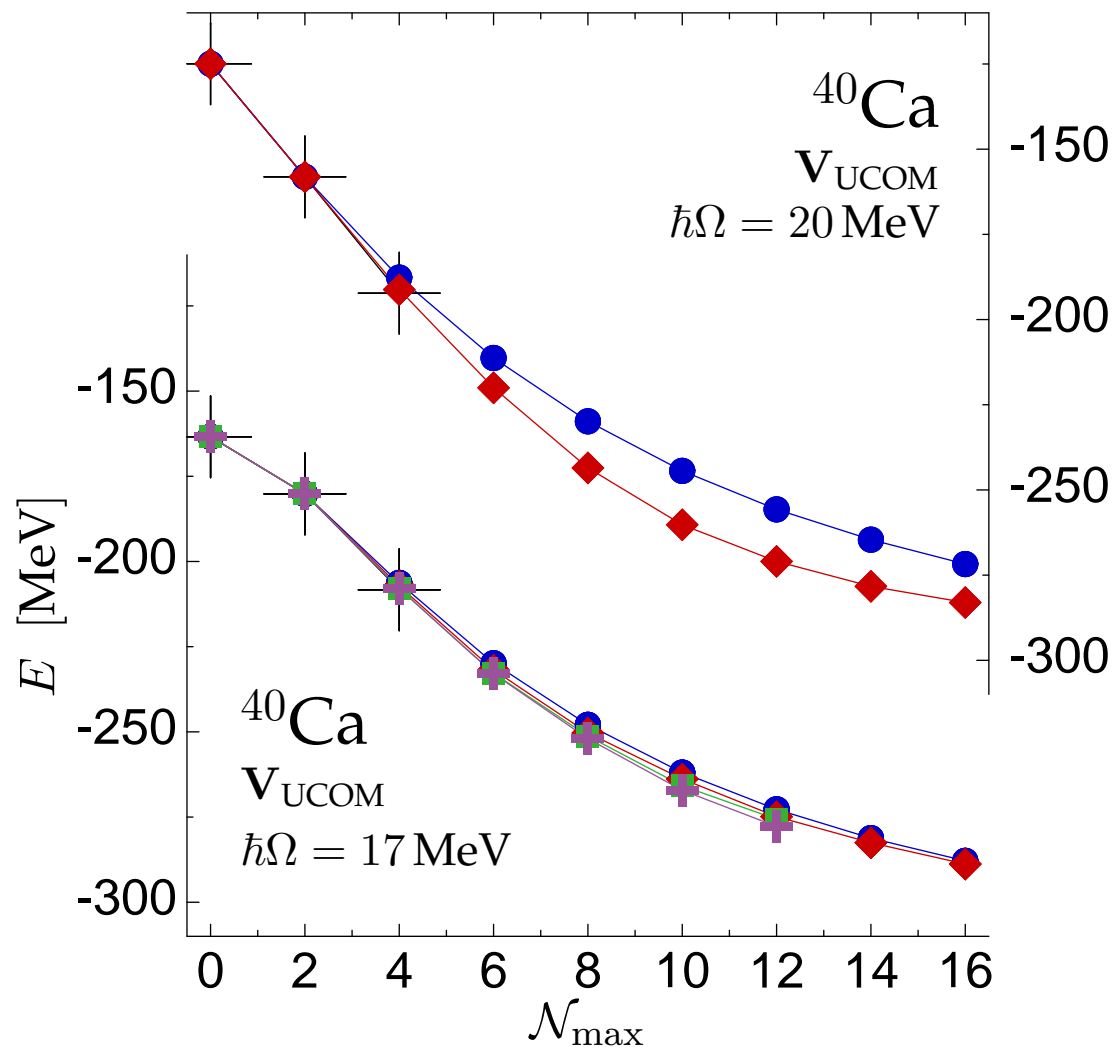
$$E_{\text{IT-NCSM}(2 \text{ iter})+\text{MRD}} \approx -130 \pm 1 \text{ MeV}$$

$$E_{\text{exp}} = -127.6 \text{ MeV}$$

■ V_{UCOM} predicts **reasonable binding energies** also for heavier nuclei

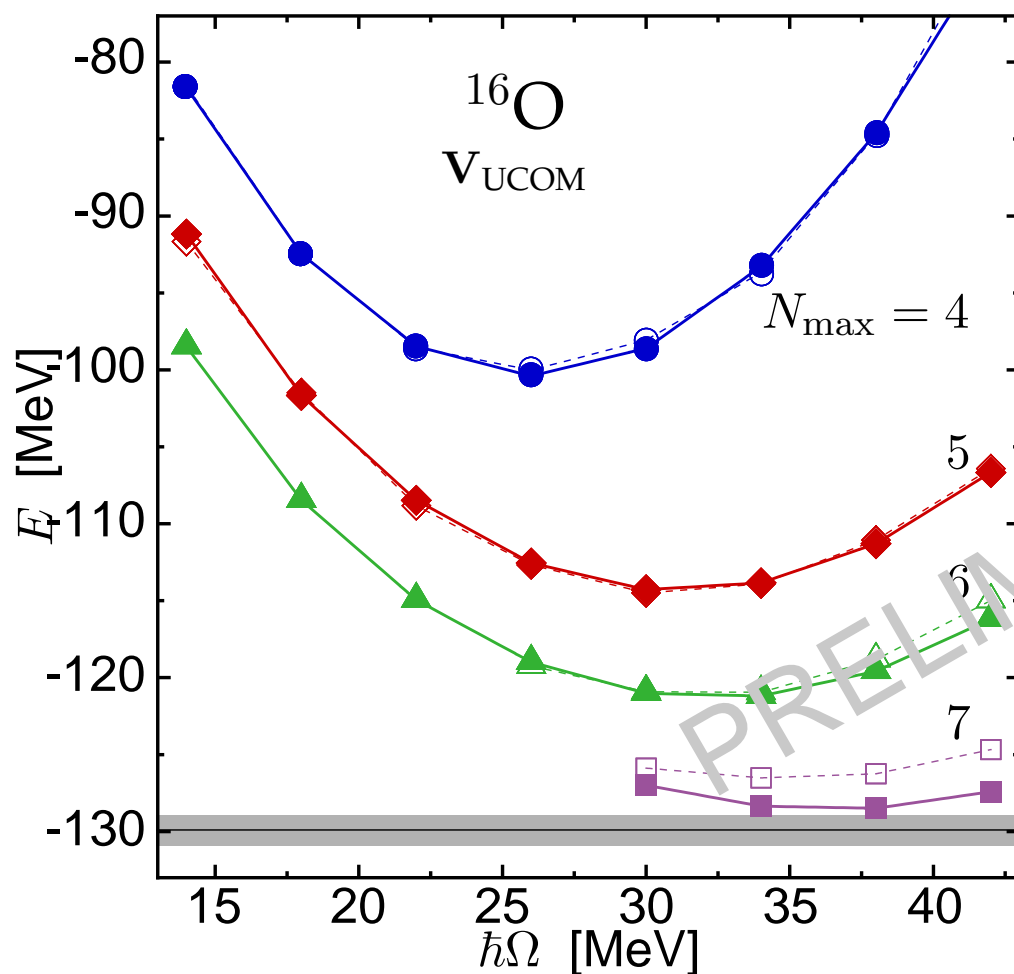
- + full NCSM
- IT-NCSM(1 iter, 2p2h)
- ◆ IT-NCSM(2 iter, 4p4h)
- IT-NCSM(2 iter, 4p4h) + MRD

^{40}Ca : Importance Truncated NCSM



- **$16\hbar\Omega$ and more are feasible** for ^{40}Ca in IT-NCSM(2 iter)
 - size of individual N_pN_h -contributions depends on $\hbar\Omega$
 - result consistent with experimental binding energy
-
- + full NCSM
 - IT-NCSM(2 iter, 2p2h)
 - ◆ IT-NCSM(2 iter, 3p3h)
 - IT-NCSM(2 iter, 4p4h)
 - ⊕ IT-NCSM(2 iter, 4p4h) + MRD

Direct Comparison: CC vs. IT-CI



● ◆ ... CR-CC(2,3)
○ ◇ ... IT-CI(2 iter,4p4h) + MRD

■ CR-CC vs. IT-CI:

- good agreement for all $\hbar\Omega$ and models spaces
- lack of strict size extensivity in the IT-CI is irrelevant

■ CR-CC/IT-CI vs. IT-NCSM:

- CC/CI seems to tend to a lower binding energy than IT-NCSM
- CC/CI suffer from center-of-mass contamination

IT-NCSM: Pros and Cons

- ✓ rigorously fulfills **variational principle** and Hylleraas-Undheim theorem
- ✓ **no sizable center-of-mass contamination** induced by IT in $\mathcal{N}_{\max} \hbar\Omega$ space
- ✓ constrained **threshold extrapolation** $\kappa_{\min} \rightarrow 0$ recovers contribution of excluded configurations efficiently and accurately
- ✓ **perturbative correction** and Davidson correction for perturbative energy correction can be used
- ✓ **compatible with shell-model**: excited states and angular-momentum projection via Lanczos, eigenstates in shell-model representation, computation of observables
- ✗ only **approximate size-extensivity** if working with few iterations
- ✗ computationally still demanding

■ Modern Effective Interactions

- treatment of short-range central and tensor correlations by unitary transformations: UCOM, SRG, Lee-Suzuki,...
- phase-shift equivalent correlated interaction V_{UCOM} which is soft and requires minimal three-body forces
- universal input for...

■ Innovative Many-Body Methods

- No-Core Shell Model,...
- Importance Truncated NCSM, Coupled Cluster Method,...
- Hartree-Fock plus MBPT, Padé Resummed MBPT, BHF, HFB, RPA,...
- Fermionic Molecular Dynamics,...

■ thanks to my group & my collaborators

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