

# Nuclear Structure with Correlated Realistic Interactions

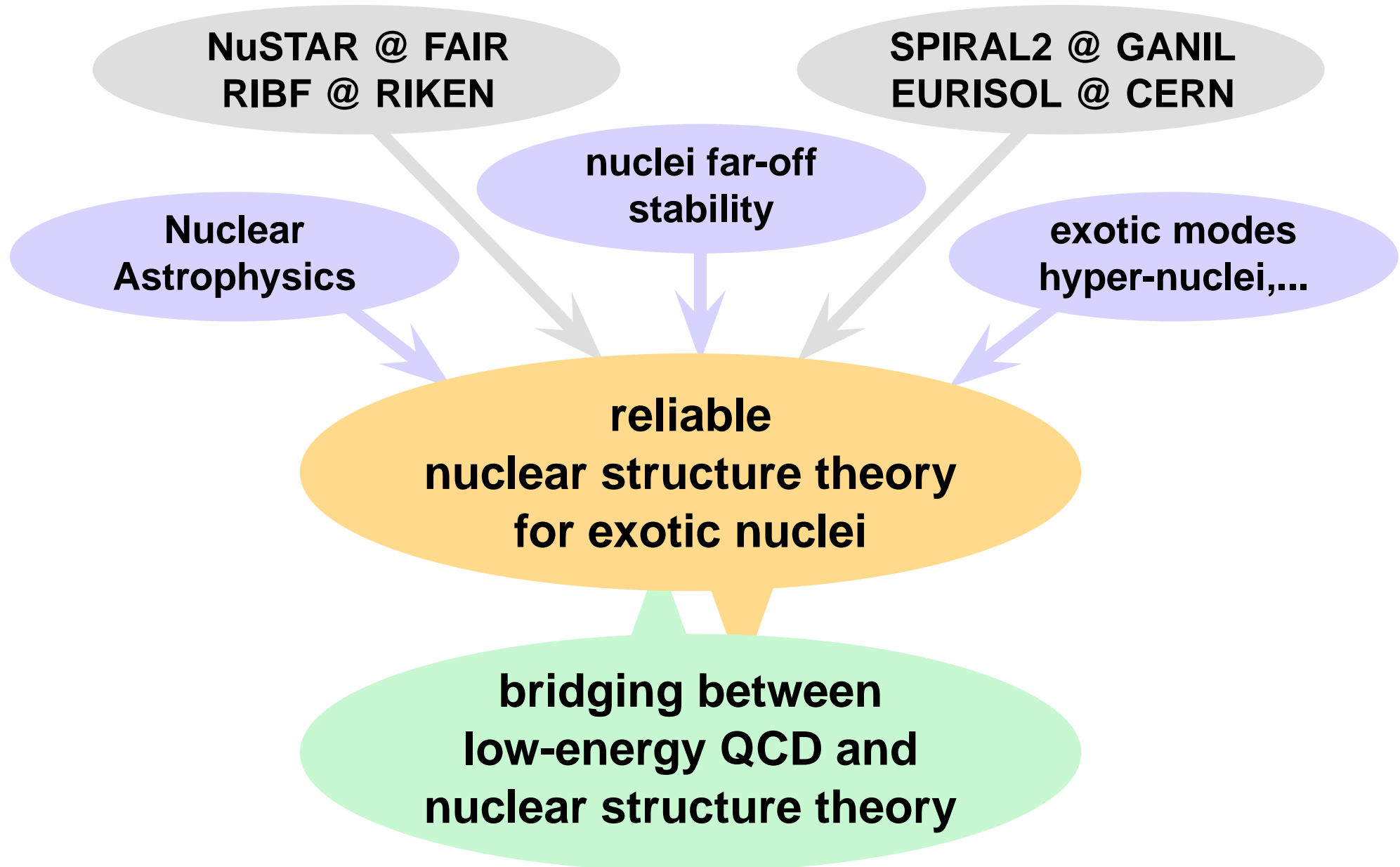


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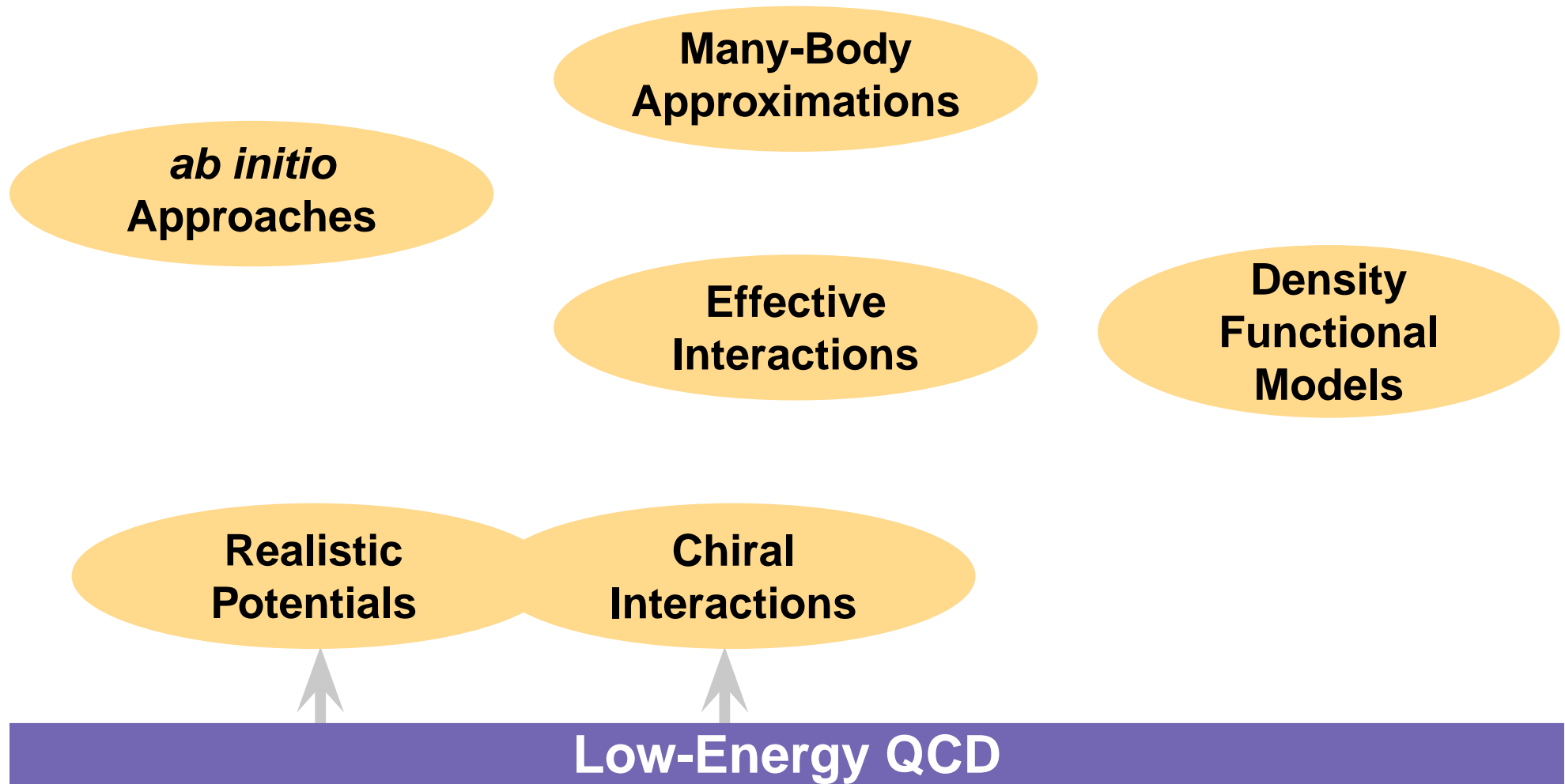
- Motivation
- Modern Effective Interactions
  - Unitary Correlation Operator Method
  - Similarity Renormalization Group
- Many-Body Methods
  - No Core Shell Model & Beyond
  - Hartree-Fock & Beyond
  - Random Phase Approximation & Beyond

# Nuclear Structure in the 21<sup>st</sup> Century



# Modern Nuclear Structure Theory

## Nuclear Structure



# Realistic Interactions

## ■ QCD motivated

- symmetries, meson-exchange picture
- chiral effective field theory

## ■ short-range phenomenology

- short-range parametrization or contact terms

## ■ experimental two-body data

- scattering phase-shifts & deuteron properties reproduced with high precision

## ■ supplementary three-nucleon force

- adjusted to spectra of light nuclei

Argonne V18

CD Bonn

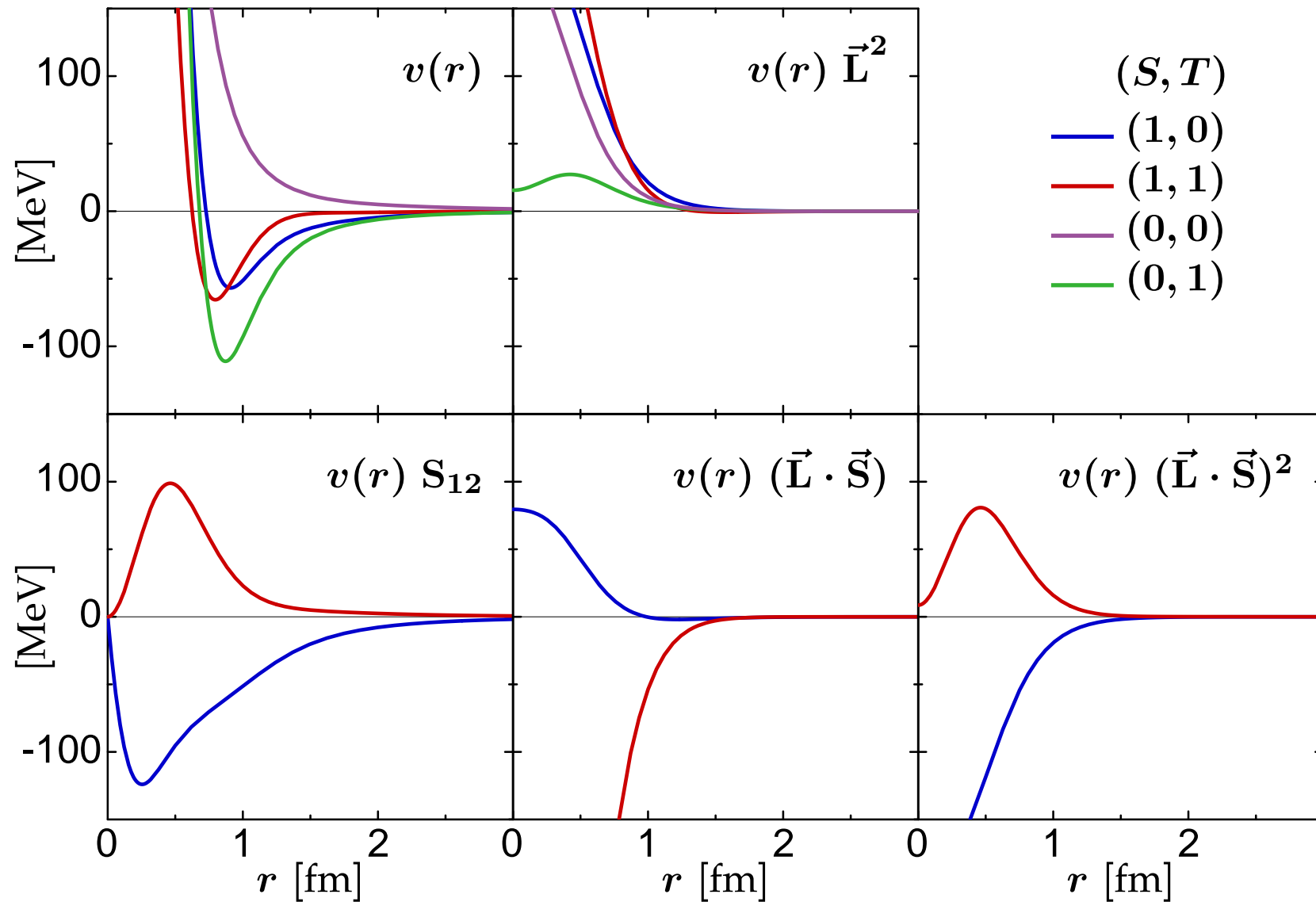
Nijmegen I/II

Chiral N3LO

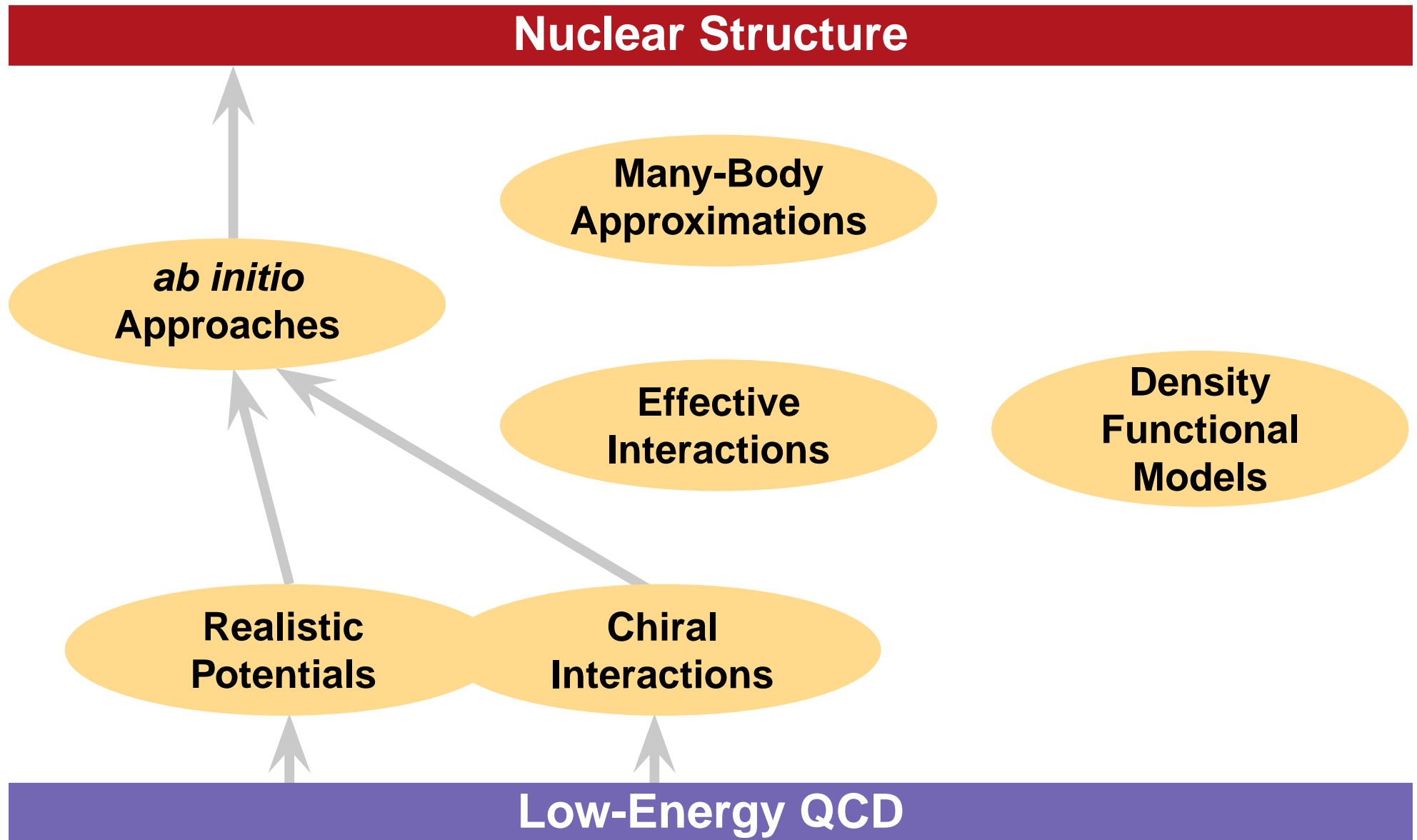
Argonne V18 +  
Illinois 2

Chiral N3LO +  
N2LO

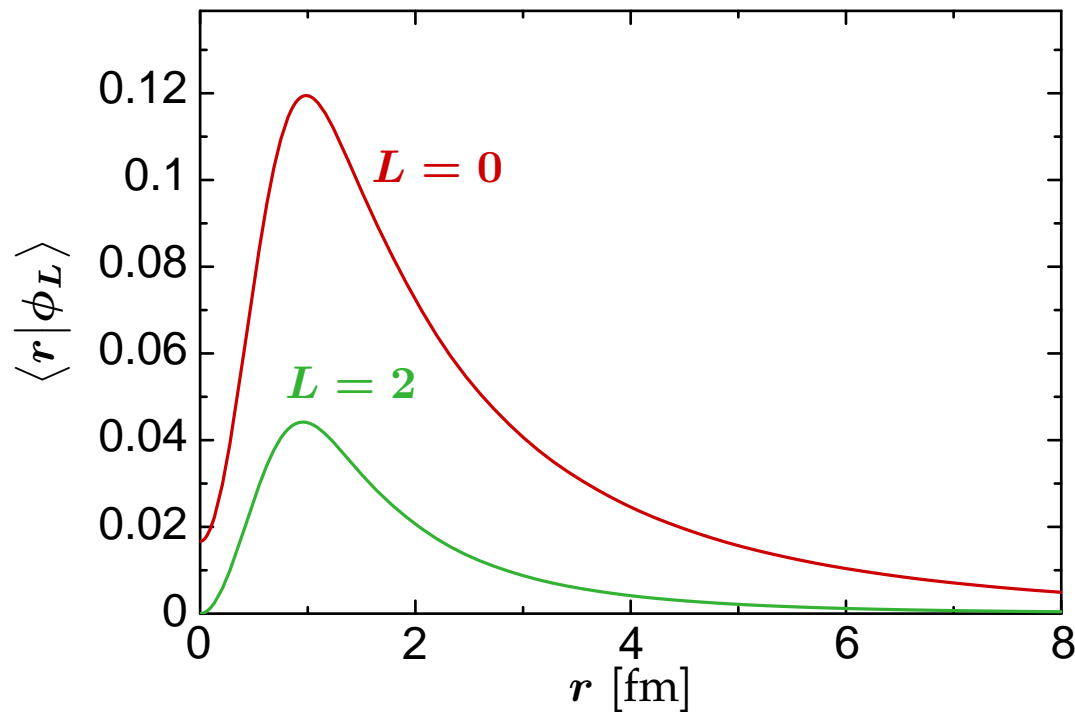
# Argonne V18 Potential



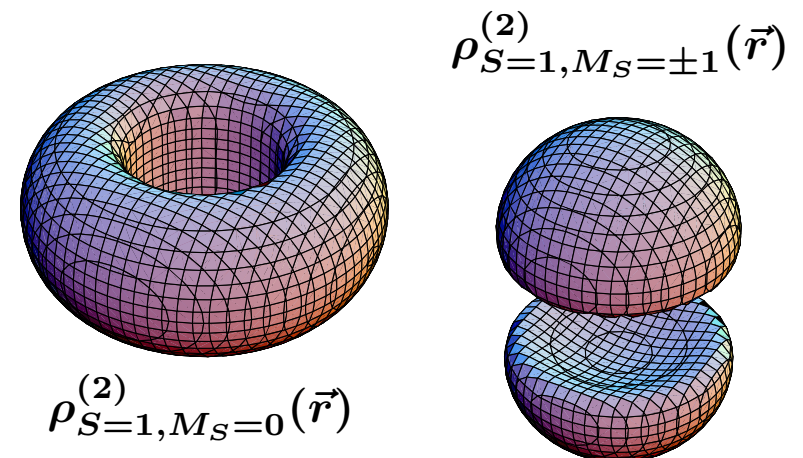
# Modern Nuclear Structure Theory



# Deuteron: Manifestation of Correlations



■ **exact deuteron solution**  
for Argonne V18 potential



short-range repulsion  
suppresses wavefunction at  
small distances  $r$

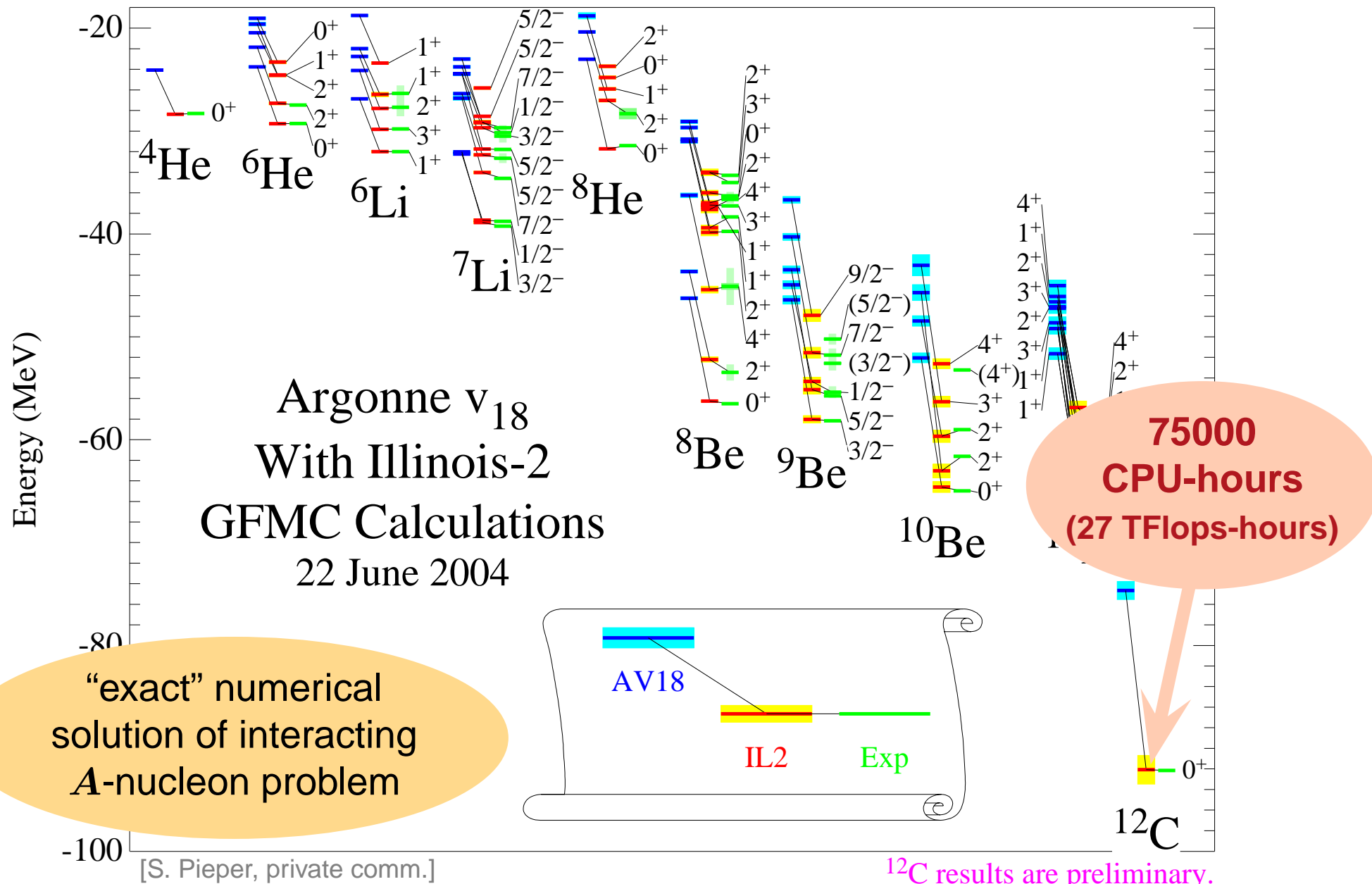
**central correlations**

tensor interaction  
generates D-wave admixture  
in the ground state

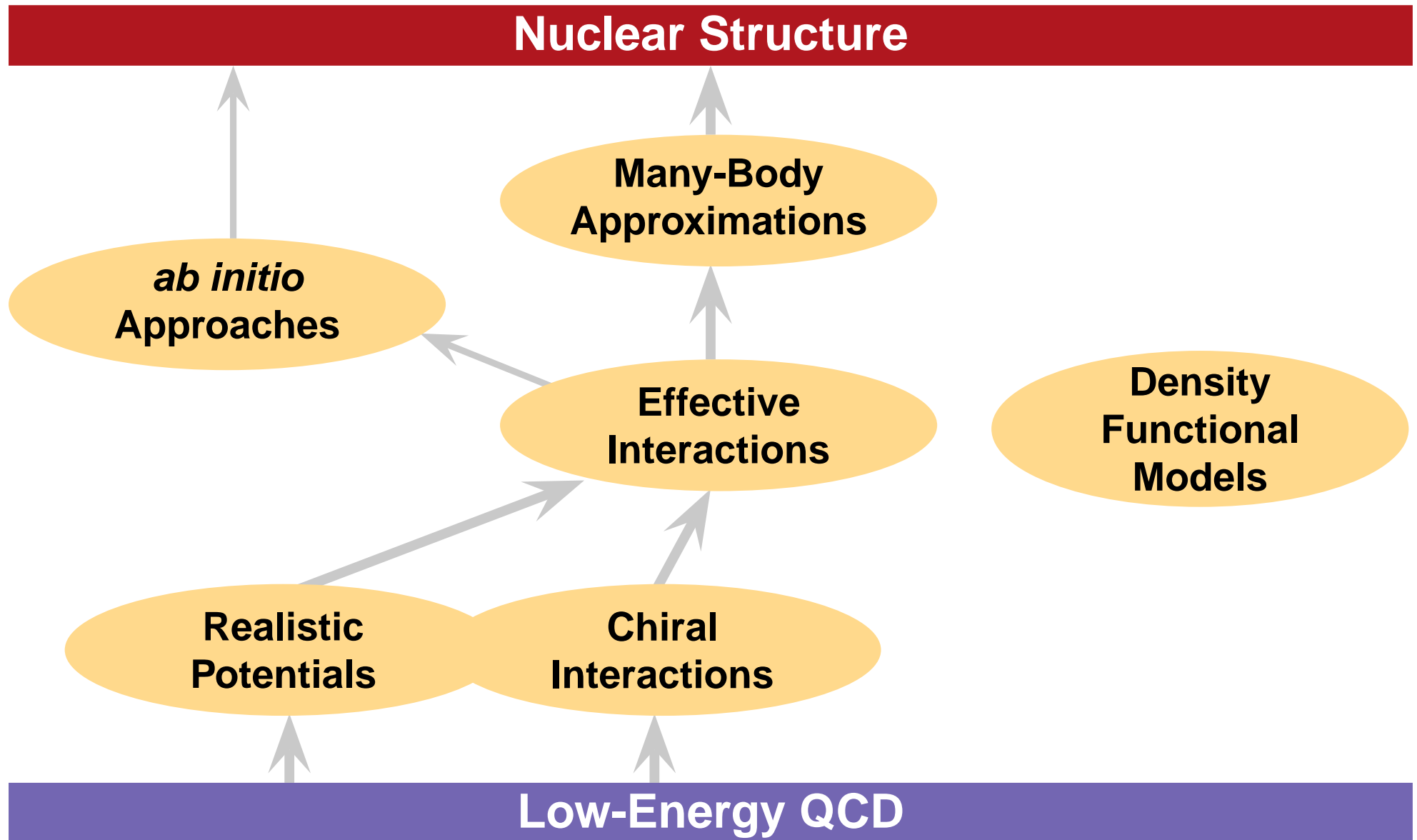
**tensor correlations**



# Ab initio Approaches: GFMC



# Modern Nuclear Structure Theory



# Why Effective Interactions?

## Realistic Potentials

- generate strong correlations in many-body states
- short-range central & tensor correlations most important

## Many-Body Approximations

- rely on truncated many-nucleon Hilbert spaces (model space)
- not capable of describing short-range correlations
- extreme: Hartree-Fock based on single Slater determinant

## Modern Effective Interactions

- adapt realistic potential to the available model space
- conserve experimentally constrained properties (phase shifts)

**can be viewed  
as realistic  
interactions**

Modern Effective Interactions I

# Unitary Correlation Operator Method (UCOM)

# Unitary Correlation Operator Method

## Correlation Operator

define an unitary operator  $\mathbf{C}$  to describe the effect of short-range correlations

$$\mathbf{C} = \exp[-i \mathbf{G}] = \exp\left[-i \sum_{i < j} g_{ij}\right]$$

## Correlated States

imprint short-range correlations onto uncorrelated many-body states

$$|\tilde{\psi}\rangle = \mathbf{C} |\psi\rangle$$

## Correlated Operators

adapt Hamiltonian and all other observables to uncorrelated many-body space

$$\tilde{\mathbf{O}} = \mathbf{C}^\dagger \mathbf{O} \mathbf{C}$$

$$\langle \tilde{\psi} | \mathbf{O} | \tilde{\psi}' \rangle = \langle \psi | \mathbf{C}^\dagger \mathbf{O} \mathbf{C} | \psi' \rangle = \langle \psi | \tilde{\mathbf{O}} | \psi' \rangle$$

# Unitary Correlation Operator Method

explicit ansatz for the correlation operator  
motivated by the **physics of short-range  
central and tensor correlations**

## Central Correlator $C_r$

- radial distance-dependent shift in the relative coordinate of a nucleon pair

$$g_r = \frac{1}{2} [s(r) \mathbf{q}_r + \mathbf{q}_r s(r)]$$

$$\mathbf{q}_r = \frac{1}{2} \left[ \frac{\vec{r}}{r} \cdot \vec{q} + \vec{q} \cdot \frac{\vec{r}}{r} \right]$$

## Tensor Correlator $C_\Omega$

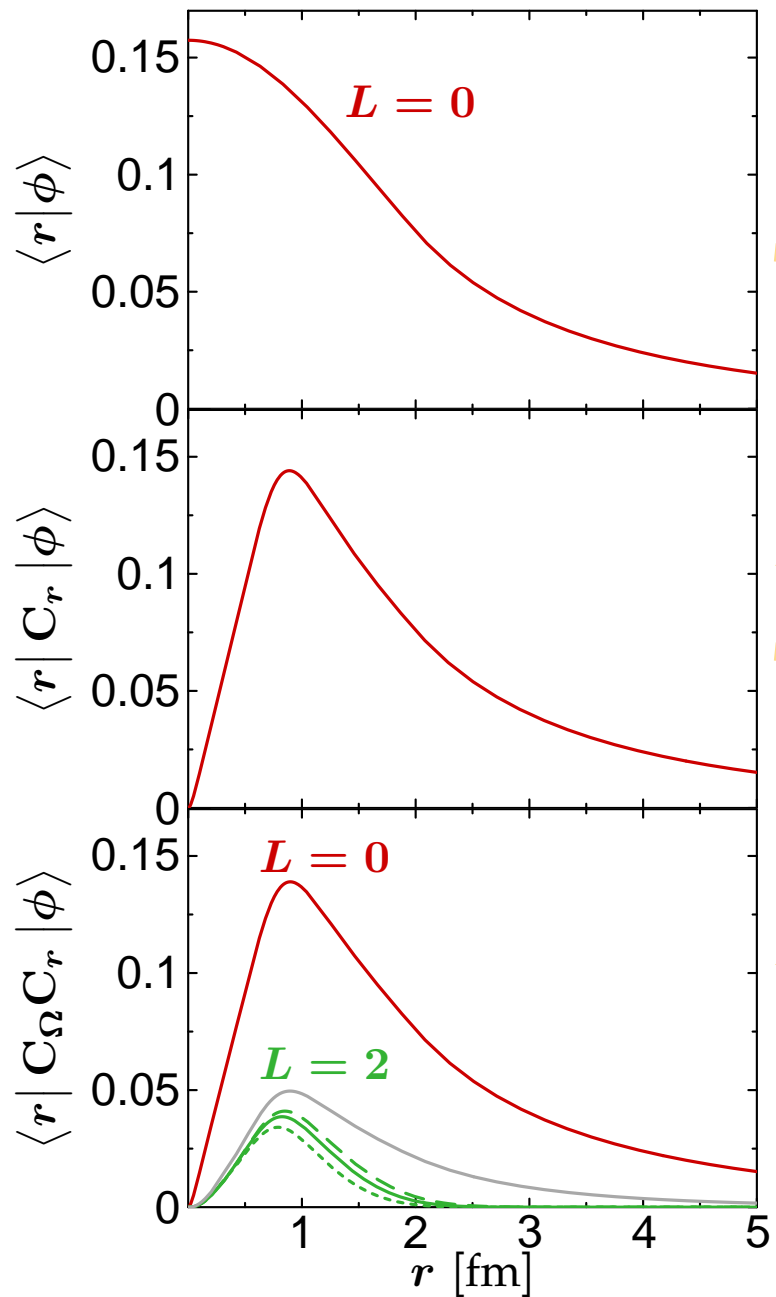
- angular shift depending on the orientation of spin and relative coordinate of a nucleon pair

$$g_\Omega = \frac{3}{2} \vartheta(r) [(\vec{\sigma}_1 \cdot \vec{q}_\Omega)(\vec{\sigma}_2 \cdot \vec{r}) + (\vec{r} \leftrightarrow \vec{q}_\Omega)]$$

$$\vec{q}_\Omega = \vec{q} - \frac{\vec{r}}{r} \mathbf{q}_r$$

- $s(r)$  and  $\vartheta(r)$  for given potential determined by energy minimization in the two-body system (for each  $S, T$ )

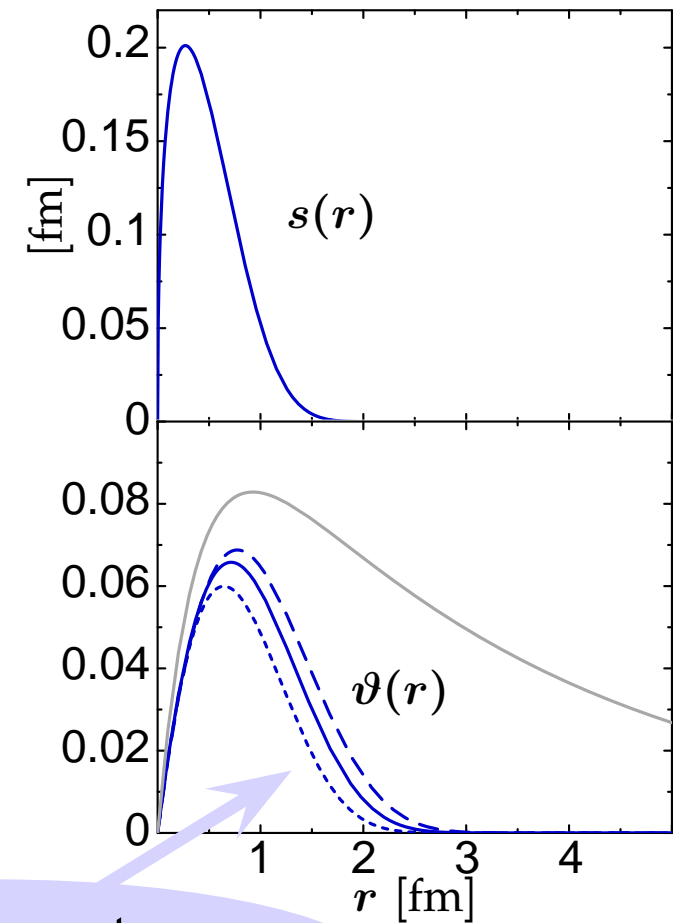
# Correlated States: The Deuteron



central correlations

tensor correlations

only short-range tensor correlations treated by  $C_\Omega$



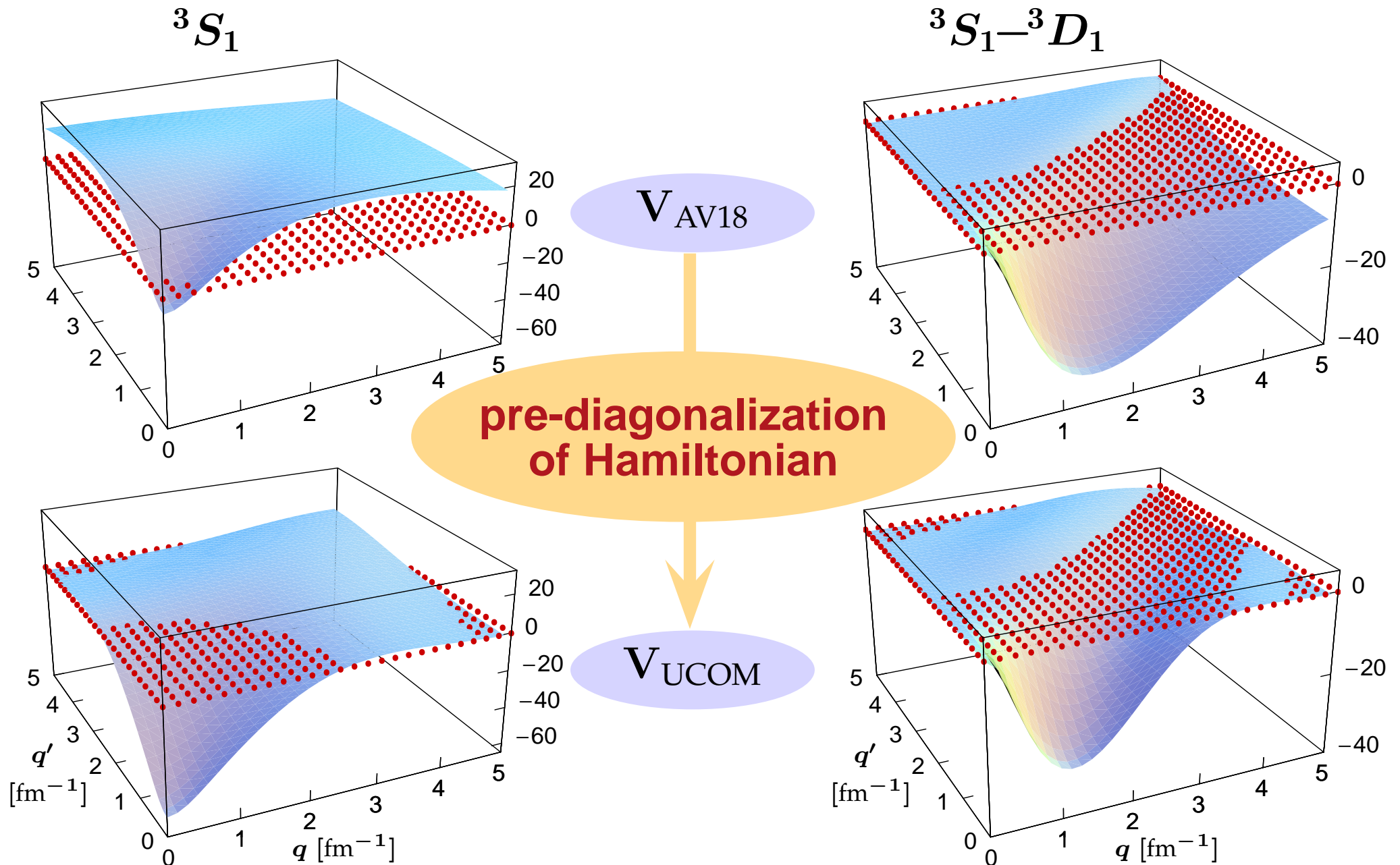
# Correlated Interaction: $V_{\text{UCOM}}$

$$\tilde{H} = T + V_{\text{UCOM}} + V_{\text{UCOM}}^{[3]} + \dots$$

- **closed operator expression** for the correlated interaction  $V_{\text{UCOM}}$  in two-body approximation
- correlated interaction and original NN-potential are **phase shift equivalent** by construction
- unitary transformation results in a **pre-diagonalization** of Hamiltonian (similar to renormalization group methods)
- operators of **all observables** (densities, transitions) have to be and can be **transformed consistently**



# Momentum-Space Matrix Elements



Modern Effective Interactions II

# Similarity Renormalization Group (SRG)

# Similarity Renormalization Group

unitary transformation of the **Hamiltonian to a band-diagonal form** with respect to a given uncorrelated many-body basis

## Flow Equation for Hamiltonian

- evolution equation for Hamiltonian

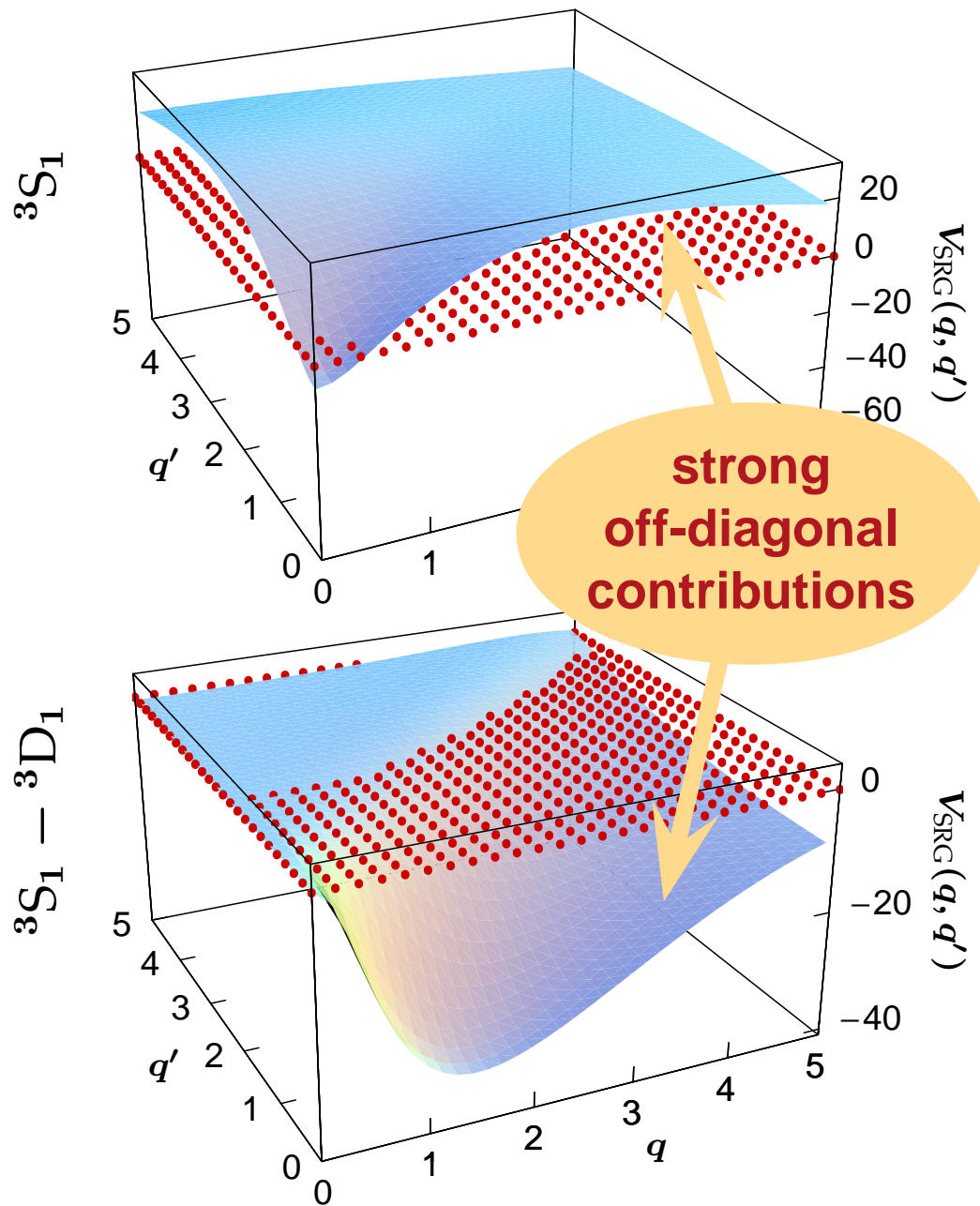
$$\tilde{H}(\alpha) = C^\dagger(\alpha) H C(\alpha) \quad \rightarrow \quad \frac{d}{d\alpha} \tilde{H}(\alpha) = [\eta(\alpha), \tilde{H}(\alpha)]$$

- dynamical generator defined as commutator with the operator in whose eigenbasis  $H$  shall be diagonalized

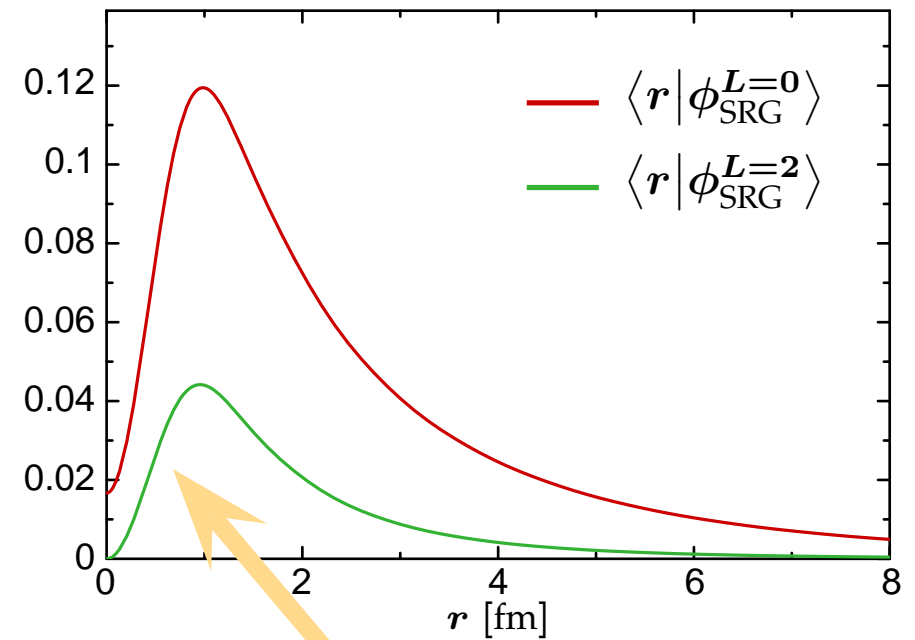
$$\eta(\alpha) = [T_{\text{int}}, \tilde{H}(\alpha)] \stackrel{2B}{=} \frac{1}{2\mu} [\vec{q}^2, \tilde{H}(\alpha)]$$

- $\eta(0)$  has the same structure as the UCOM generators  $g_r$  and  $g_\Omega$

# SRG Evolution: The Deuteron

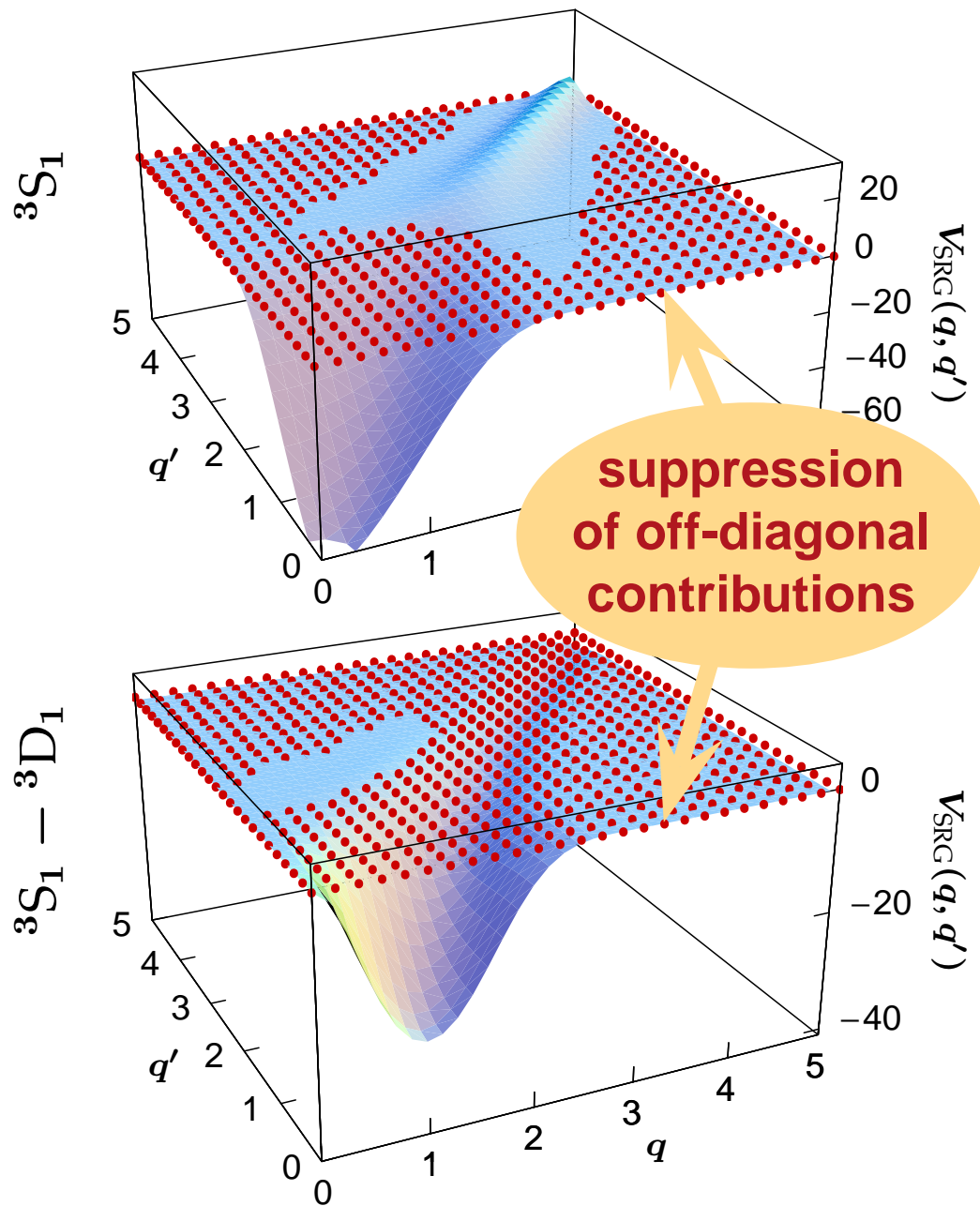


Argonne V18

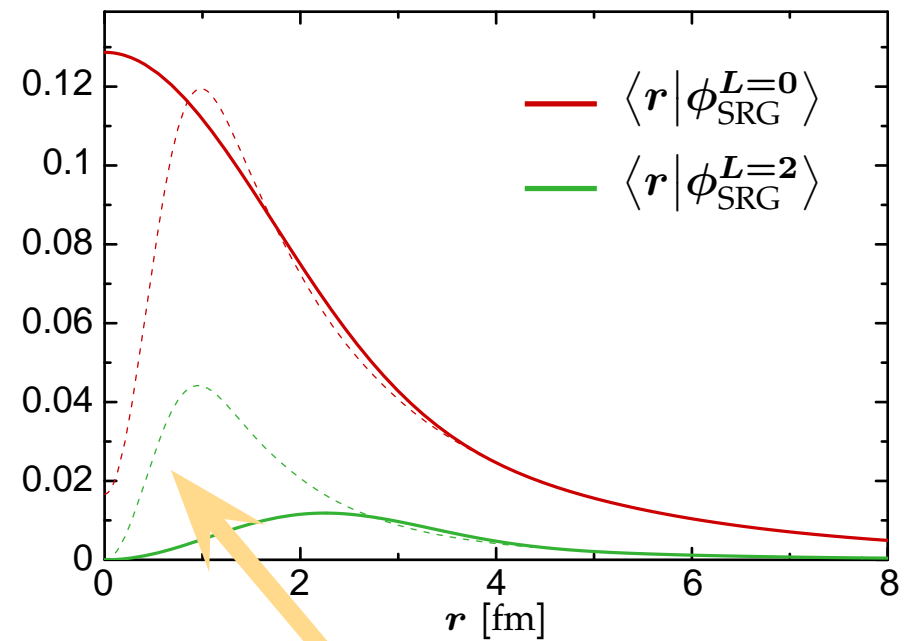


**short-range central & tensor correlations**

# SRG Evolution: The Deuteron



$$\alpha = 0.1000 \text{ fm}^4$$



Many-Body Methods I

# No-Core Shell Model

in collaboration with  
Petr Navrátil (LLNL)

# Reminder: No-Core Shell Model

- many-body state is **expanded in Slater determinants**  $|\text{SD}_i\rangle$  composed of harmonic oscillator single-particle states

$$|\Psi\rangle = \sum_i C_i |\text{SD}_i\rangle$$

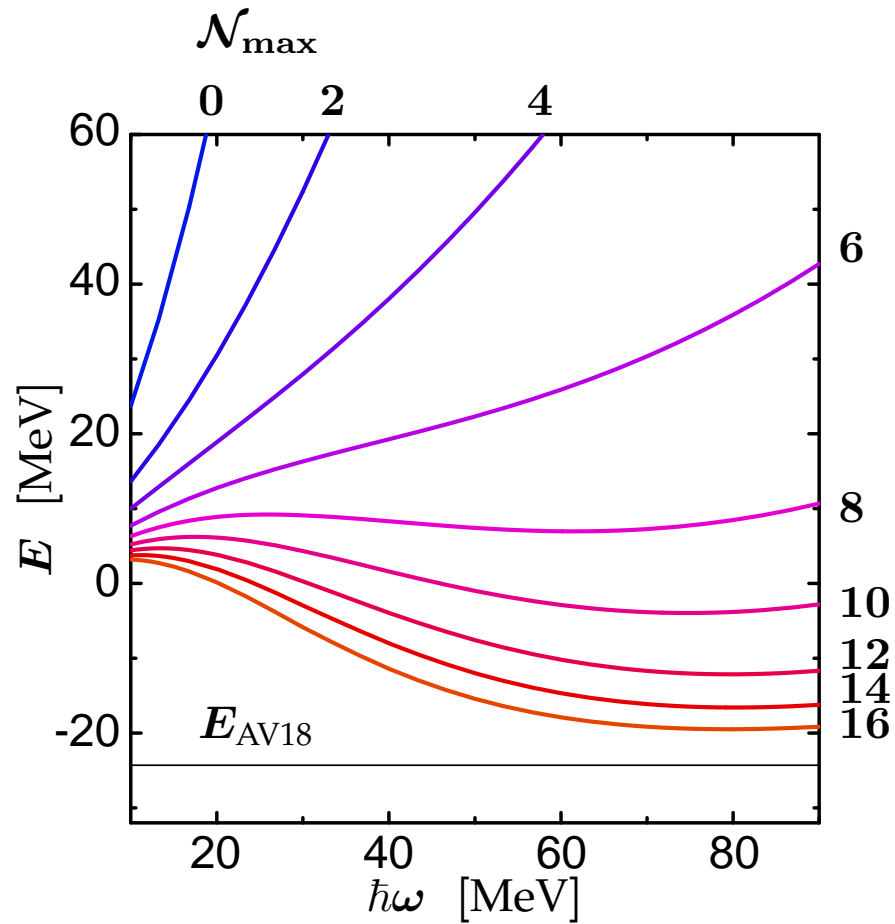
- $\mathcal{N}_{\max} \hbar\omega$  **model space**: truncate basis of Slater determinants with respect to number of oscillator quanta (unperturbed excitation energy)

with increasing model space size more and more **correlations can be described** by the shell model states

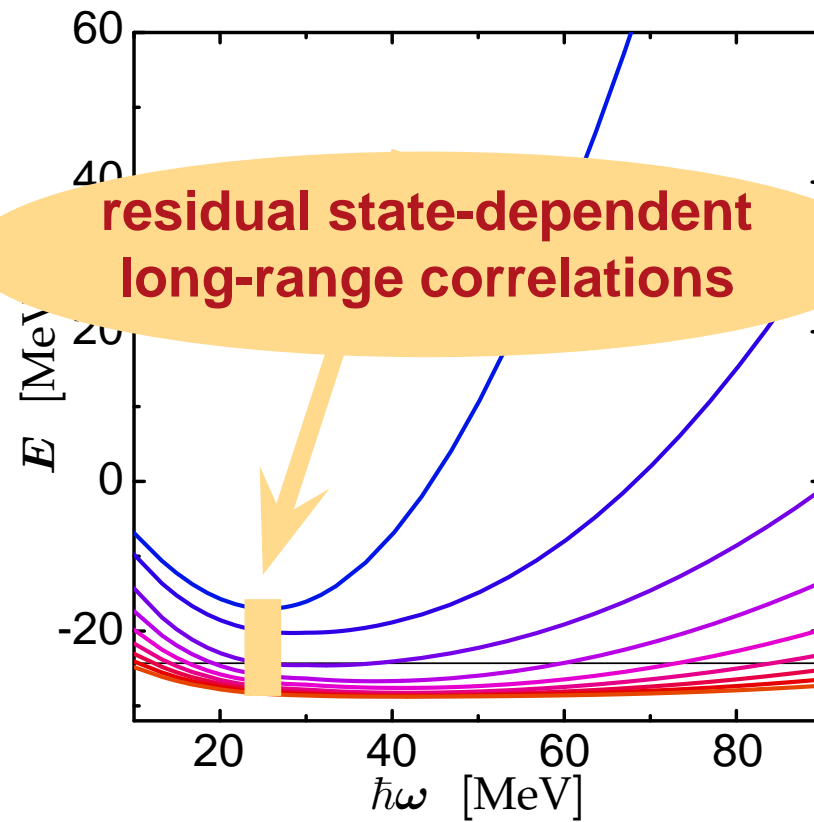
facilitates systematic study of short- and long-range correlations

# $^4\text{He}$ : Convergence

$V_{\text{AV18}}$



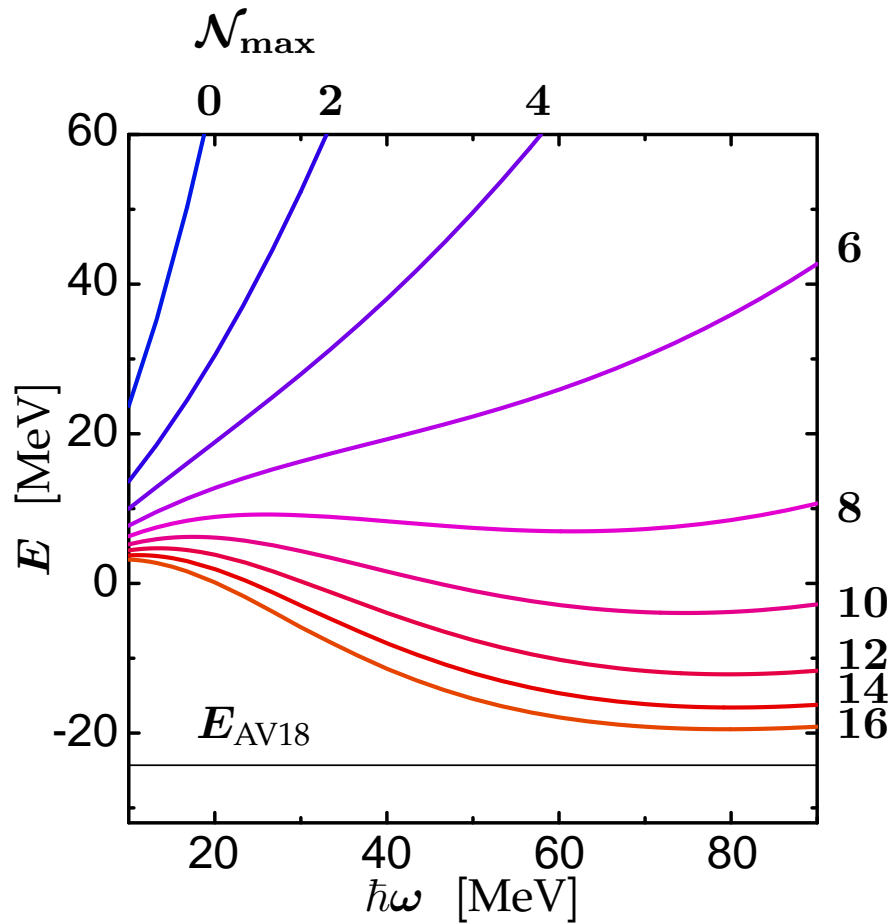
$V_{\text{UCOM}}$



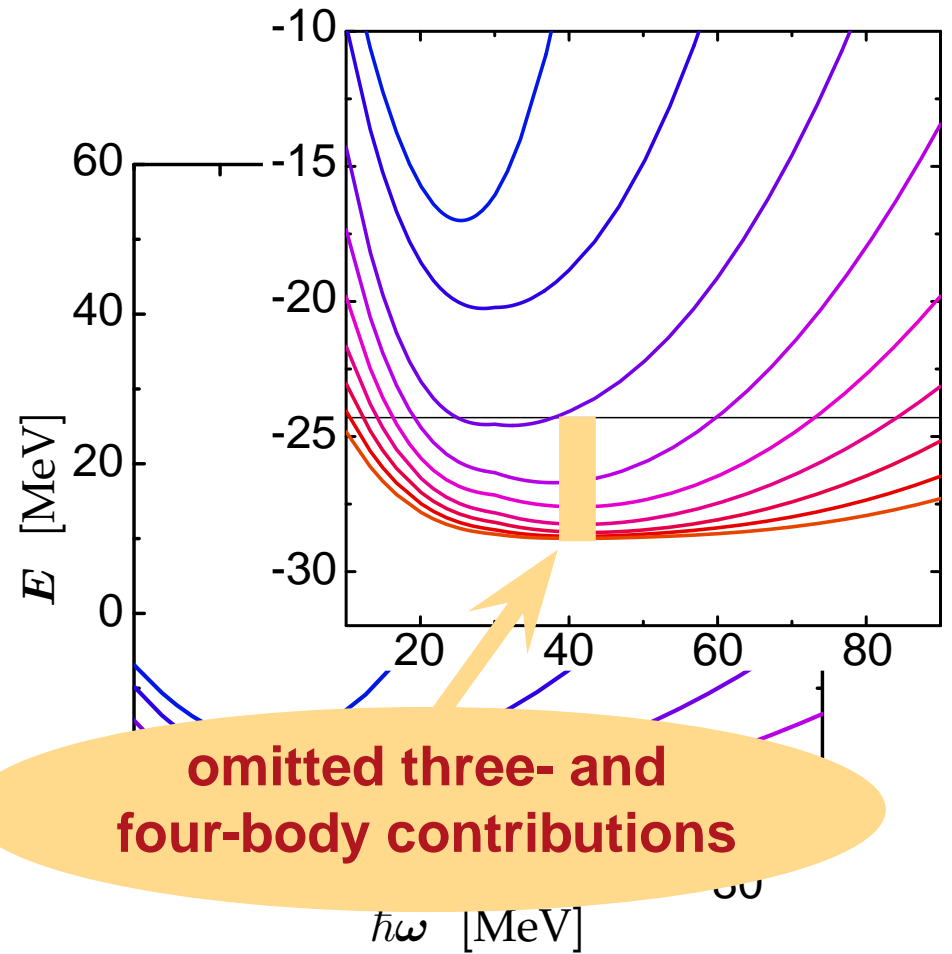


# $^4\text{He}$ : Convergence

$V_{\text{AV18}}$



$V_{\text{UCOM}}$



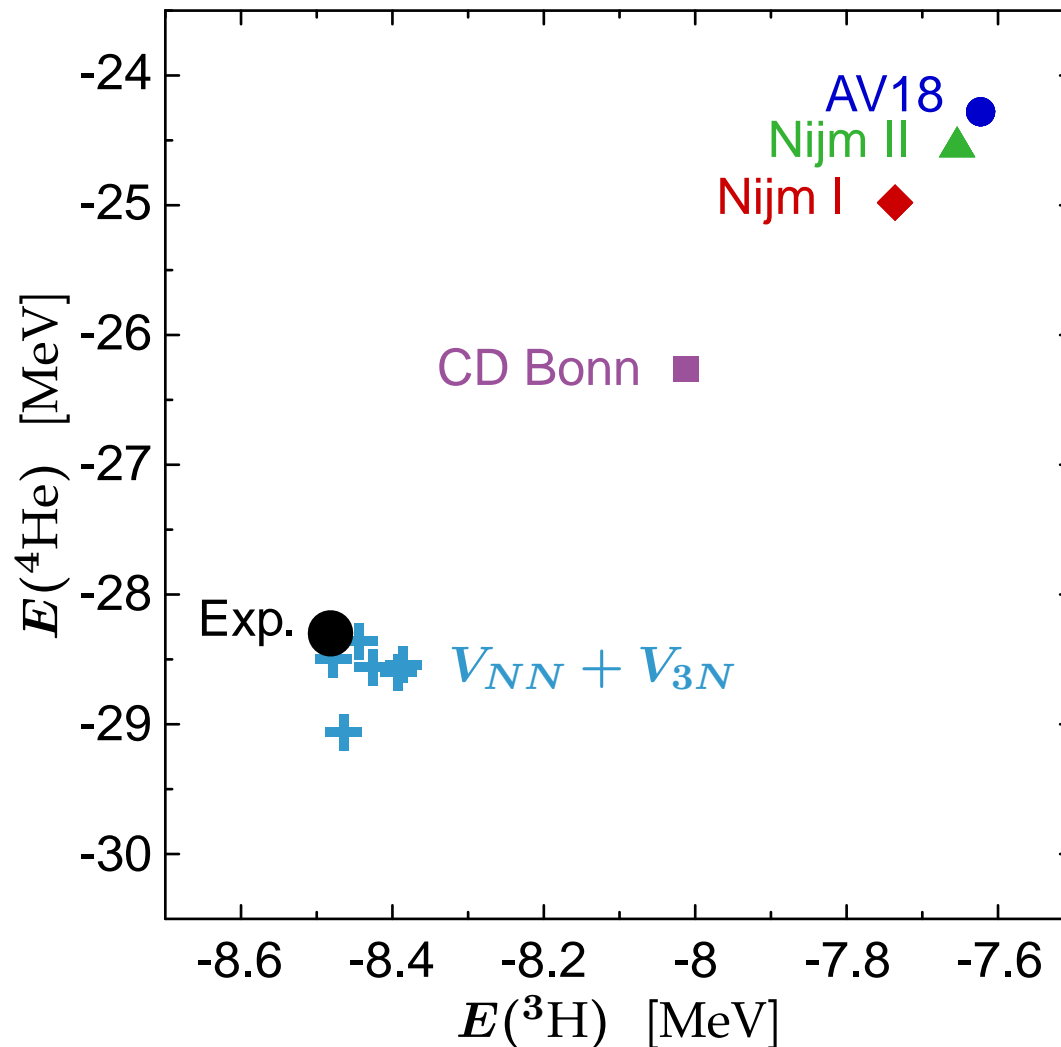
# Three-Body Interactions — Strategies

## Correlated Hamiltonian in Many-Body Space

$$\begin{aligned}\tilde{H} &= C^\dagger (\mathbf{T} + \mathbf{V}_{NN} + \mathbf{V}_{3N}) C \\ &= \tilde{\mathbf{T}}^{[1]} + (\tilde{\mathbf{T}}^{[2]} + \tilde{\mathbf{V}}_{NN}^{[2]}) + (\tilde{\mathbf{T}}^{[3]} + \tilde{\mathbf{V}}_{NN}^{[3]} + \tilde{\mathbf{V}}_{3N}^{[3]}) + \dots \\ &= \mathbf{T} + \mathbf{V}_{UCOM} + \mathbf{V}_{UCOM}^{[3]} + \dots\end{aligned}$$

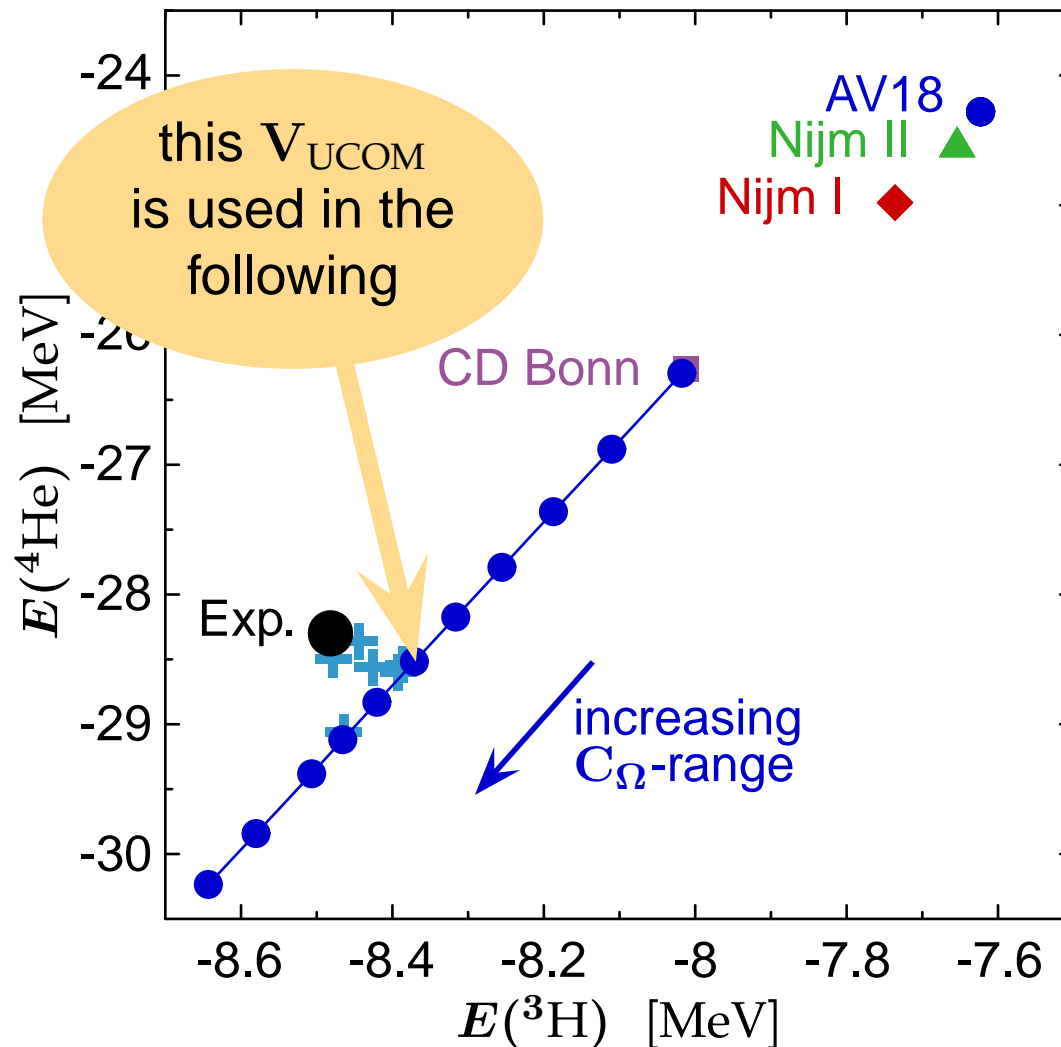
- **include full**  $\mathbf{V}_{UCOM}^{[3]}$  consisting of genuine and induced 3N terms  
(not really feasible beyond lightest isotopes)
- **replace**  $\mathbf{V}_{UCOM}^{[3]}$  by phenomenological three-body force  
(tractable also for heavier nuclei)
- **minimize**  $\mathbf{V}_{UCOM}^{[3]}$  by proper choice of unitary transformation  
(calculation with a pure two-body interaction)

# Three-Body Interactions — Tjon Line



- **Tjon-line:**  $E({}^4\text{He})$  vs.  $E({}^3\text{H})$  for phase-shift equivalent NN-interactions

# Three-Body Interactions — Tjon Line

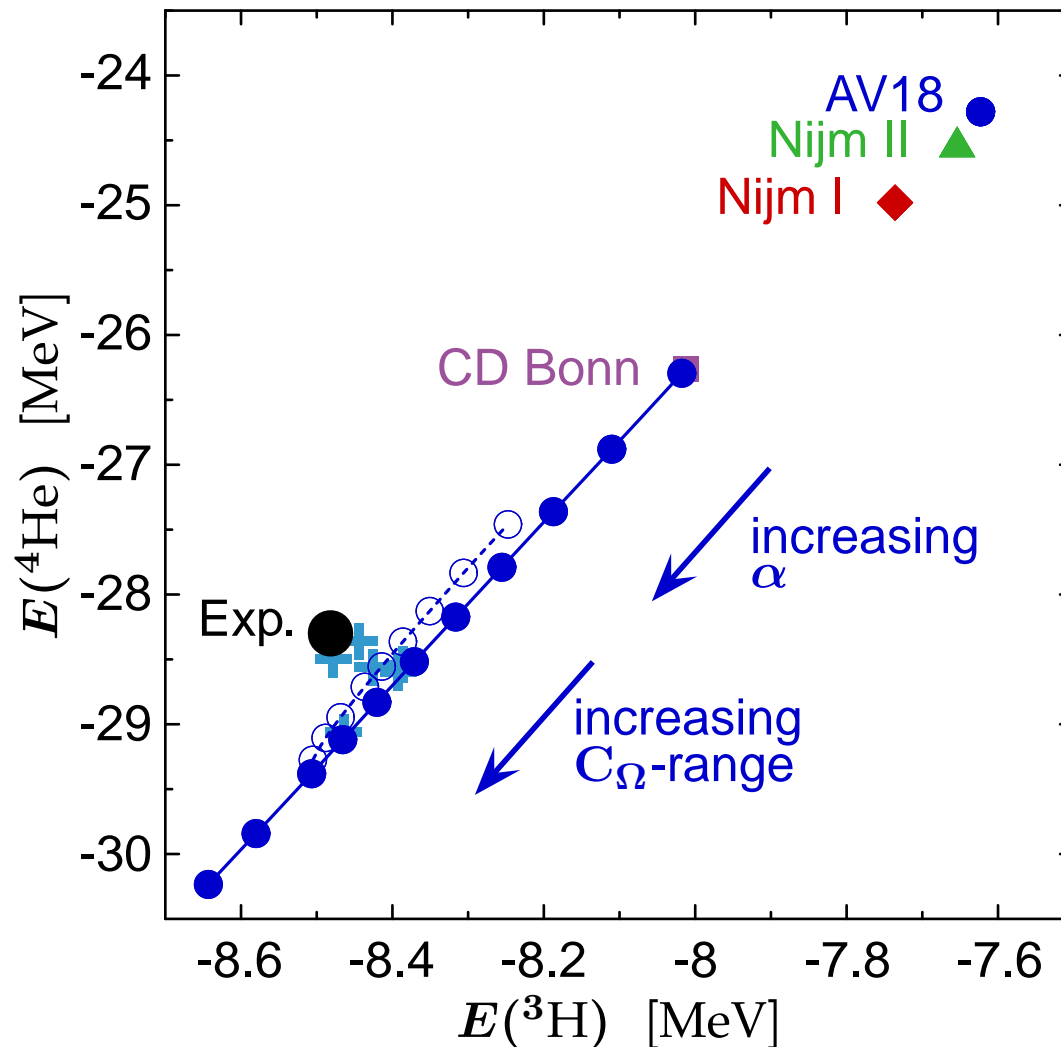


- **Tjon-line:**  $E(^4\text{He})$  vs.  $E(^3\text{H})$  for phase-shift equivalent NN-interactions

- change of  $C_\Omega$ -correlator range results in shift along Tjon-line

**minimize net three-body force** by choosing correlator with energies close to experimental value

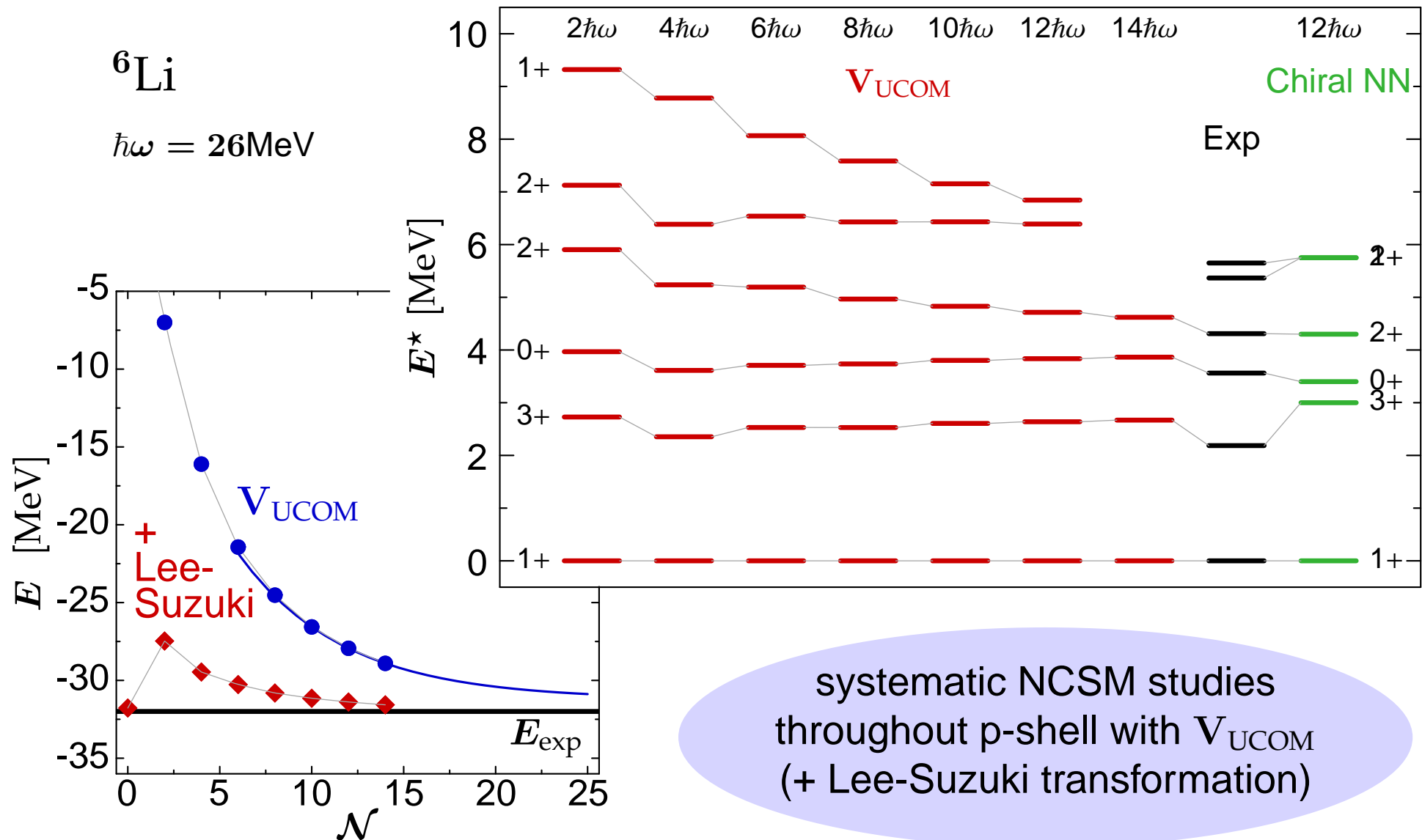
# Three-Body Interactions — Tjon Line



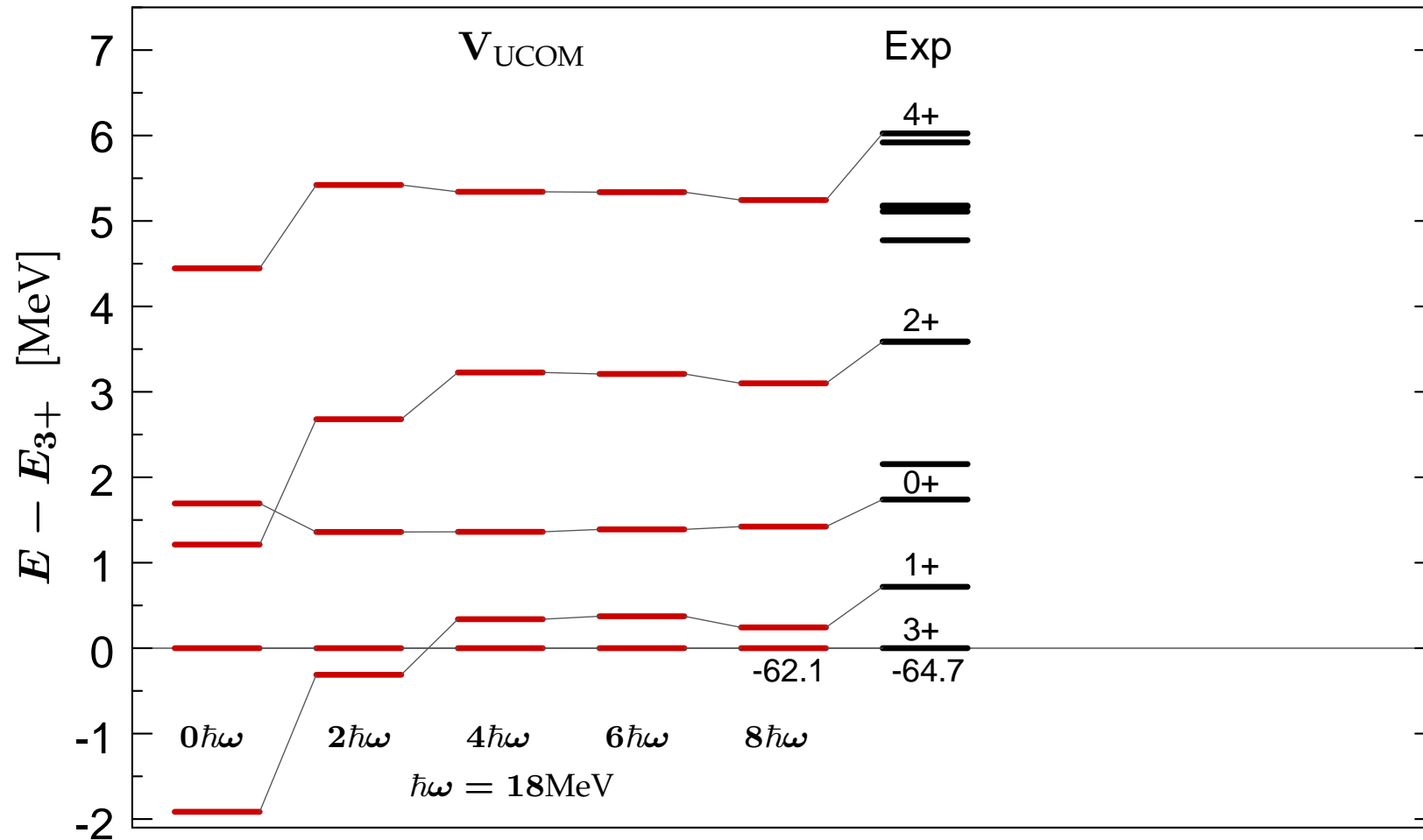
- **Tjon-line:**  $E(^4\text{He})$  vs.  $E(^3\text{H})$  for phase-shift equivalent NN-interactions
- same behavior for the SRG interaction as function of  $\alpha$

**minimize net  
three-body force**  
by choosing correlator  
with energies close to  
experimental value

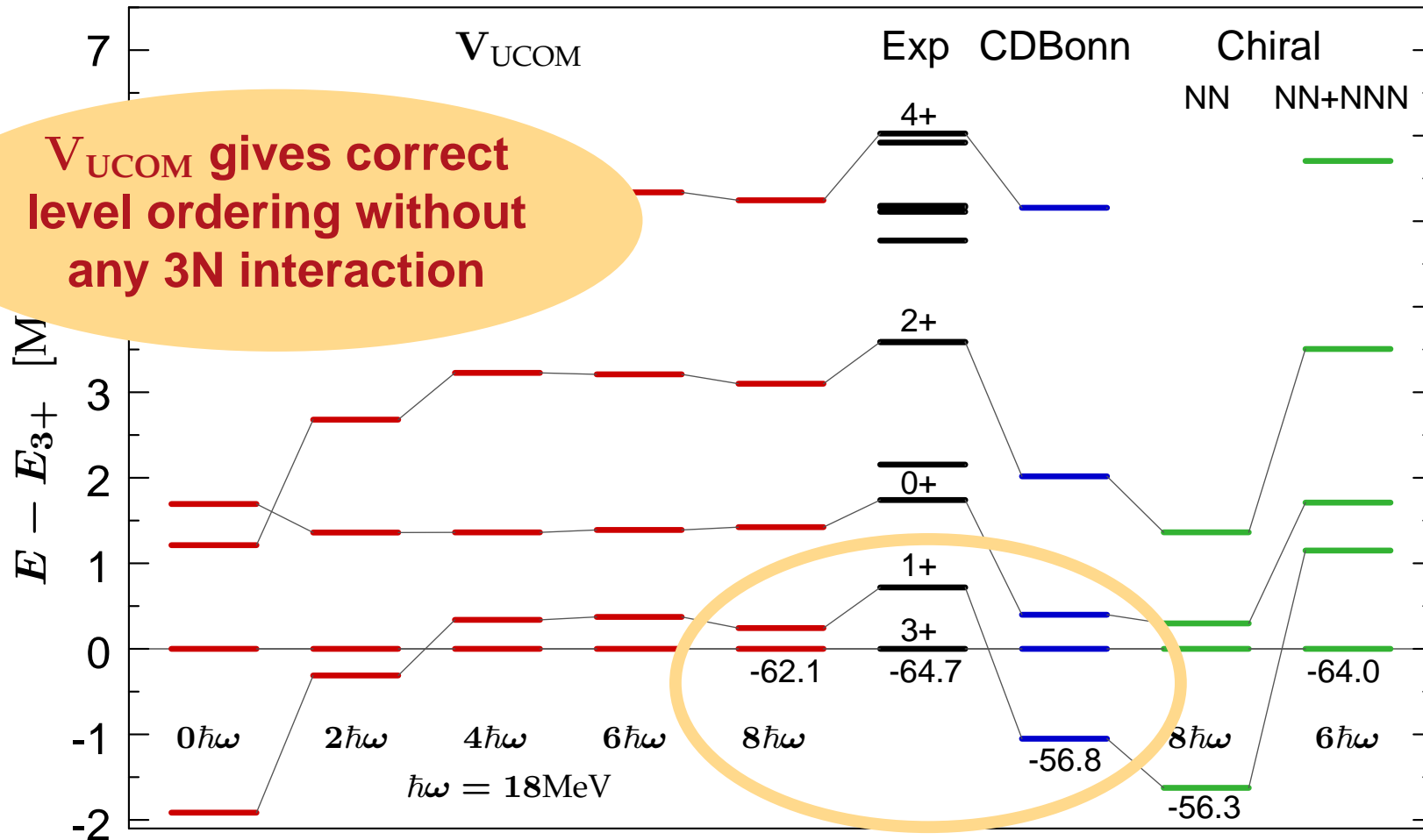
# ${}^6\text{Li}$ : NCSM throughout the p-Shell



# $^{10}\text{B}$ : Hallmark of a 3N Interaction?



# $^{10}\text{B}$ : Hallmark of a 3N Interaction?





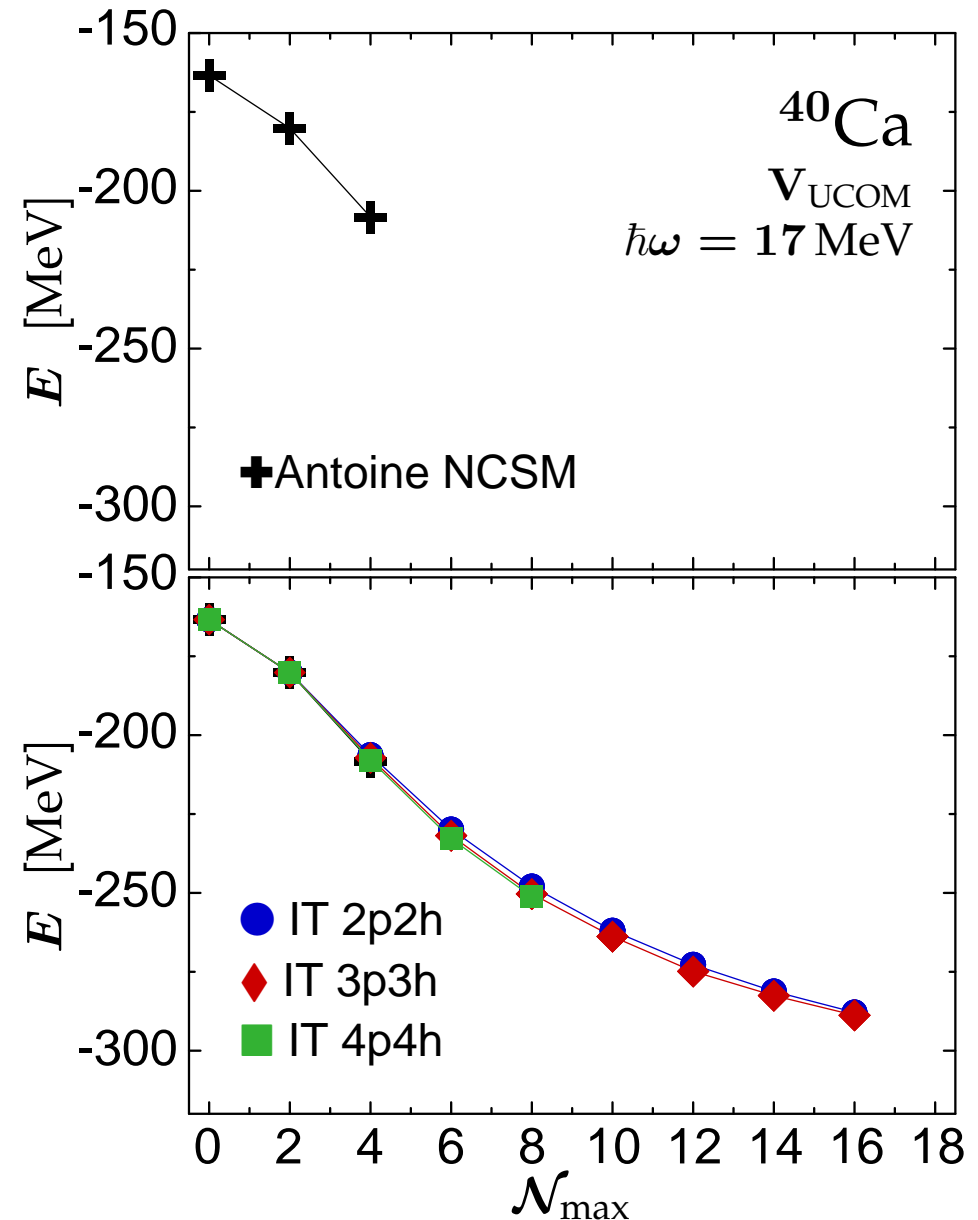
# Importance Truncated NCSM

## NCSM

- converged calculations essentially restricted to p-shell
- $6\hbar\omega$  for  $^{40}\text{Ca}$  presently not feasible ( $\sim 10^{10}$  states)

## Importance Truncation

- diagonalization in space of **important** configurations
- **a priori importance measure** given by perturbation theory



Many-Body Methods II:

# Hartree-Fock & Beyond

# Reminder: Hartree-Fock Approximation

- ground state approximated by a **single Slater determinant**

$$|\Psi\rangle \approx |\text{HF}\rangle = |\phi_1, \phi_2, \dots, \phi_A\rangle_a$$

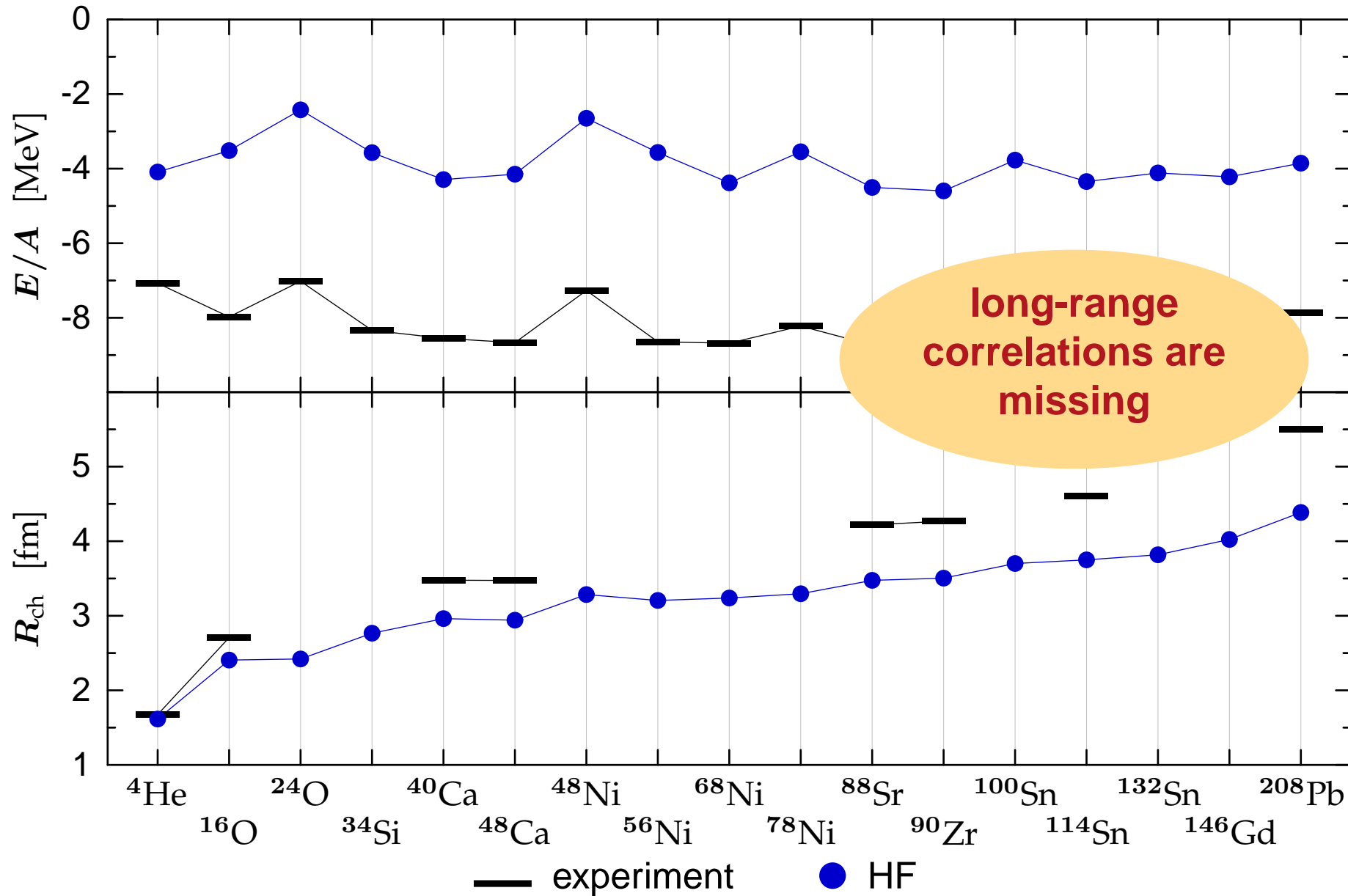
- **variational calculation**: single-particle states  $|\phi_i\rangle$  determined by minimizing the energy expectation value

$$E_{\text{HF}} = \langle \text{HF} | \mathbf{H}_{\text{int}} | \text{HF} \rangle = \frac{1}{2} \sum_{i,j=1}^A \langle \phi_i \phi_j | (\mathbf{T}_{\text{int}} + \mathbf{V}) | \phi_i \phi_j \rangle_a$$

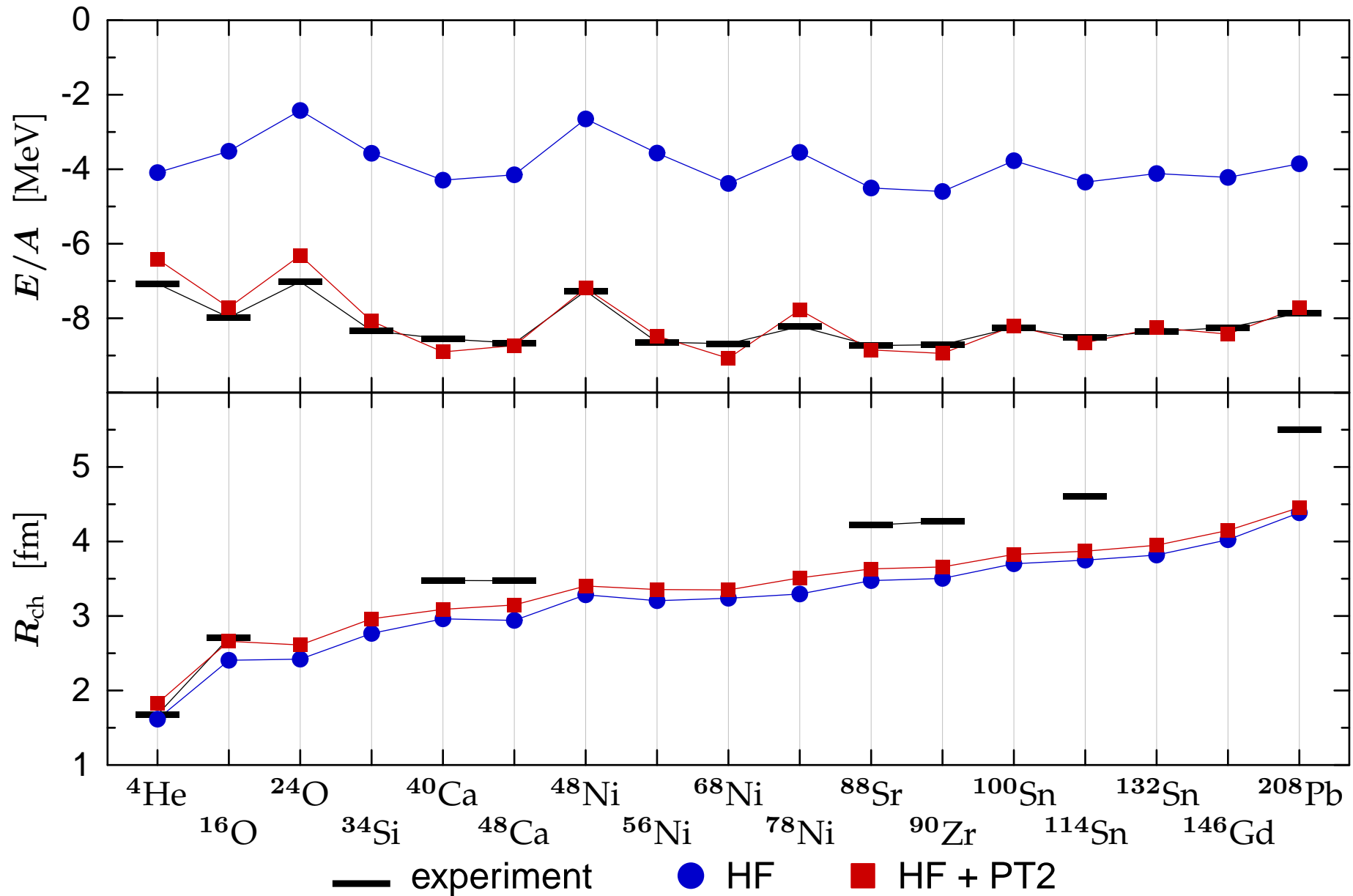
single Slater determinant by definition **cannot describe any correlations** (independent particle state)

Hartree-Fock solution is **starting point for improved calculations**

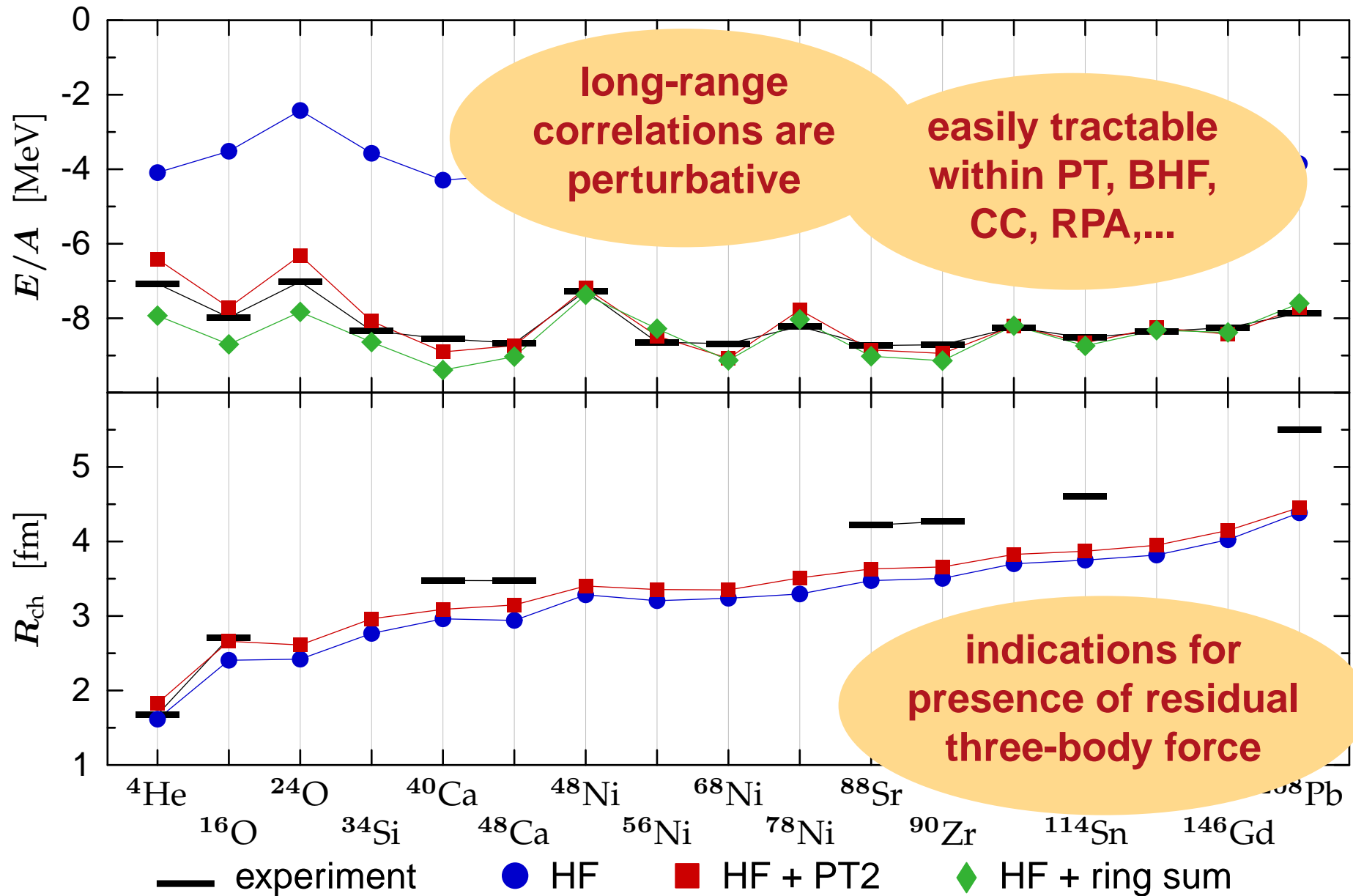
# Hartree-Fock with $V_{UCOM}$



# Perturbation Theory with $V_{\text{UCOM}}$



# RPA Ring Summation with $V_{UCOM}$



Many-Body Methods III

# RPA & Beyond

# Reminder: Random Phase Approximation

- describe **excited states** via vibration creation operator  $Q_\nu^\dagger$

$$Q_\nu |\text{RPA}\rangle = 0 \quad Q_\nu^\dagger |\text{RPA}\rangle = |\nu\rangle$$

- ansatz for **vibration creation operator**  $Q_\nu^\dagger$  including 1p1h excitations with respect to HF single-particle basis

$$Q_\nu^\dagger = \sum_{ph} X_{ph}^\nu a_p^\dagger a_h - \sum_{ph} Y_{ph}^\nu a_h^\dagger a_p$$

- formal solution of eigenvalue problem via equations of motion method approximating vacuum state by  $|\text{HF}\rangle$  yields **RPA equations**

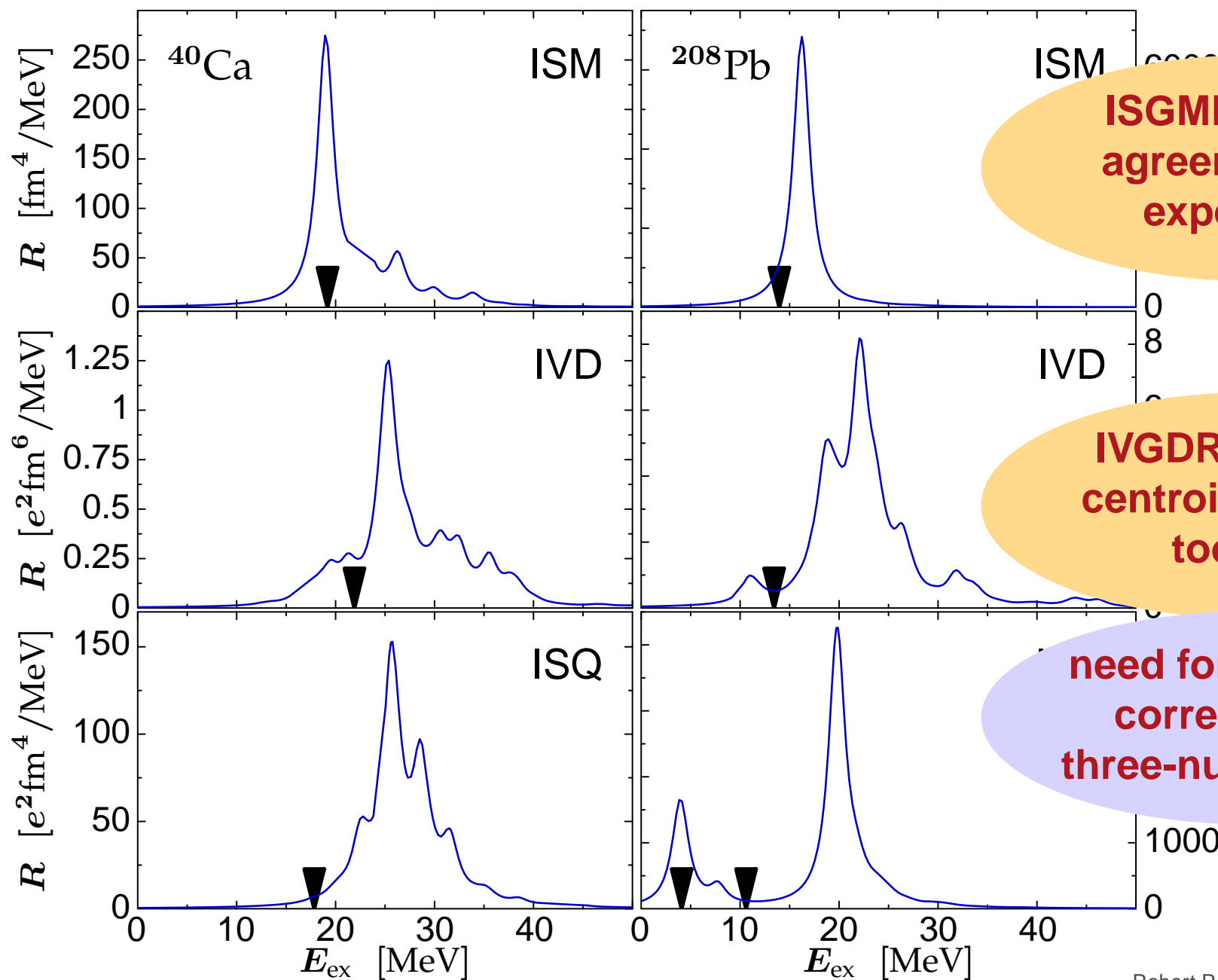
$$\begin{pmatrix} A & B \\ -B^* & -A^* \end{pmatrix} \begin{pmatrix} X^\nu \\ Y^\nu \end{pmatrix} = E_\nu \begin{pmatrix} X^\nu \\ Y^\nu \end{pmatrix}$$

$$A_{ph,p'h'} = \delta_{pp'} \delta_{hh'} (\epsilon_p - \epsilon_h) + \langle hp' | \mathbf{H}_{\text{int}} | ph' \rangle \quad B_{ph,p'h'} = \langle hh' | \mathbf{H}_{\text{int}} | pp' \rangle$$

- **self-consistent** solution using the same Hamiltonian  $\mathbf{H}_{\text{int}}$  as in HF



# RPA with $V_{UCOM}$

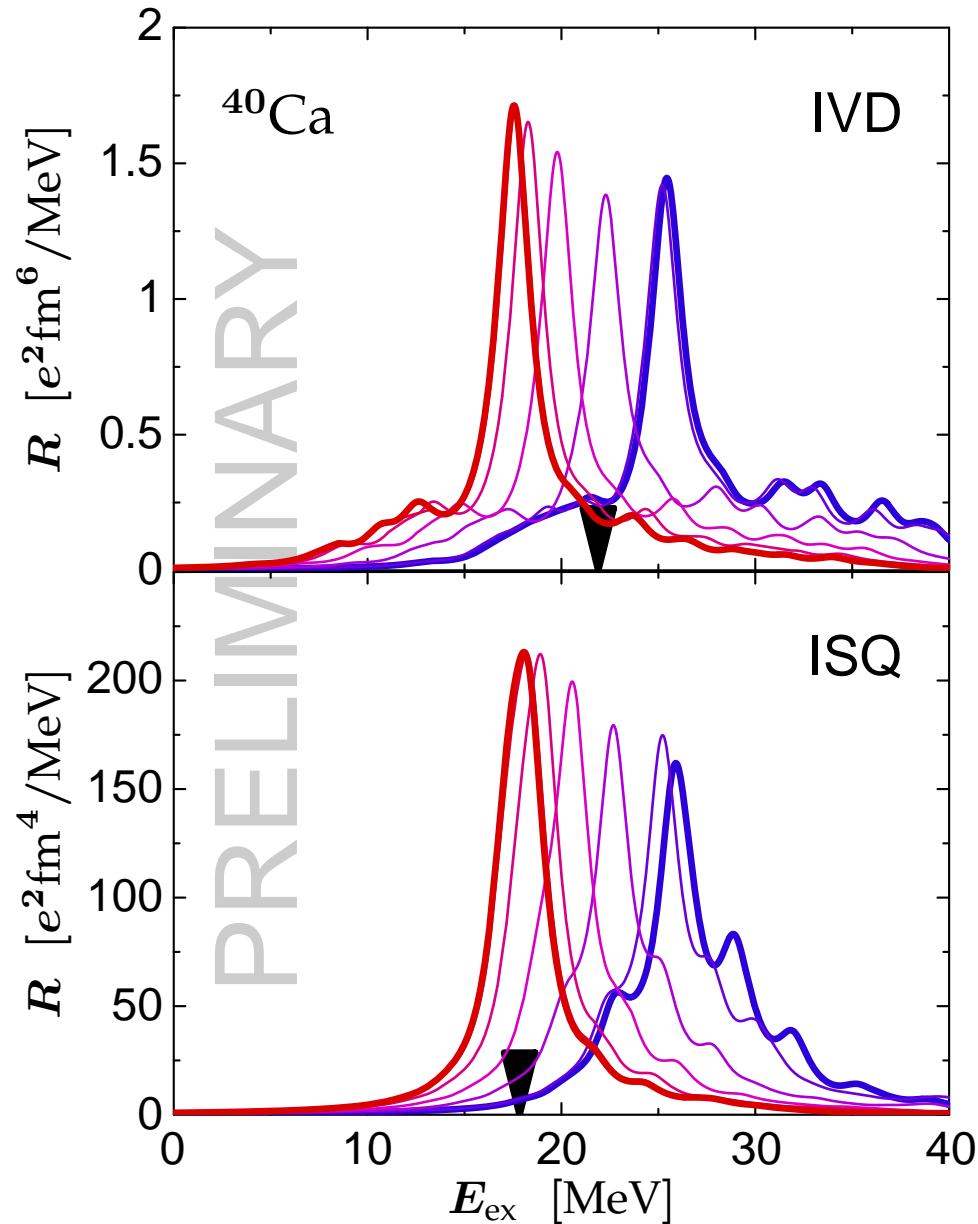


**ISGMR in good agreement with experiment**

**IVGDR & ISGQR centroid energies too high**

**need for additional correlations & three-nucleon force**

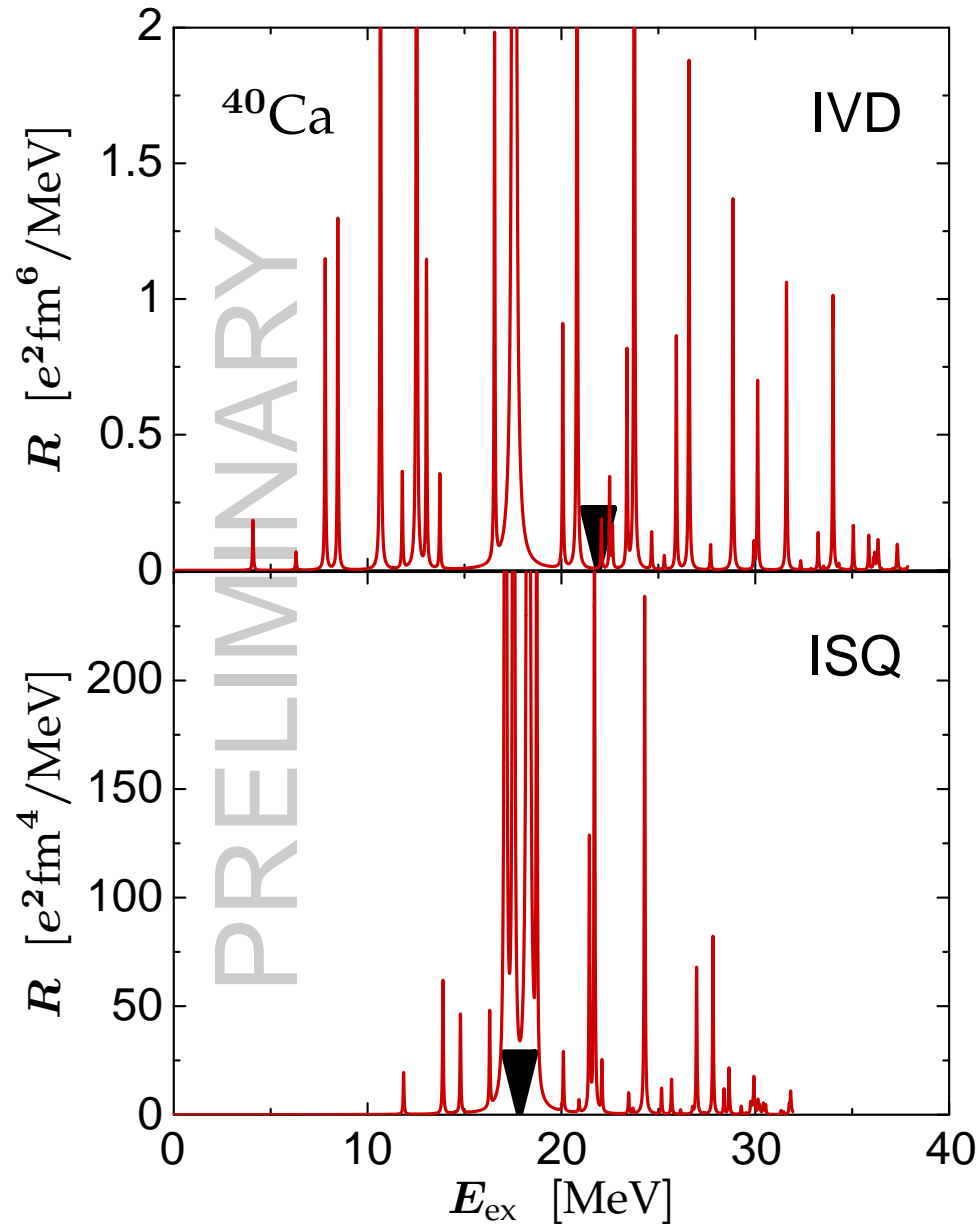
# SRPA: Complex Configurations



- RPA (full 1p1h space)
  - SRPA (full 1p1h+2p2h space)
- (11 major shells,  $\Gamma = 2 \text{ MeV}$ )

**complex  
configurations have  
significant impact on  
response**

# SRPA: Complex Configurations



- RPA (full 1p1h space)
  - SRPA (full 1p1h+2p2h space)
- (11 major shells,  $\Gamma = 50$  keV)

**complex configurations have significant impact on response**

**possibility to investigate fine structure**

# Conclusions

## ■ Modern Effective Interactions

- treatment of short-range central and tensor correlations by unitary transformations: UCOM, SRG, Lee-Suzuki,...
- universal phase-shift equivalent correlated interaction  $V_{\text{UCOM}}$

## ■ Innovative Many-Body Methods

- No-Core Shell Model, Importance Truncation
- Hartree-Fock, MBPT, CC, RPA, ERPA, SRPA,...
- Fermionic Molecular Dynamics

**unified description of nuclear  
structure across the whole  
nuclear chart is within reach**

## ■ thanks to my group & my collaborators

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