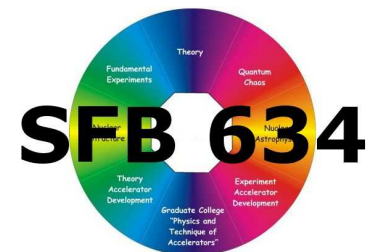


Pairing with Correlated Realistic NN Interactions

Heiko Hergert

Institut für Kernphysik, TU Darmstadt



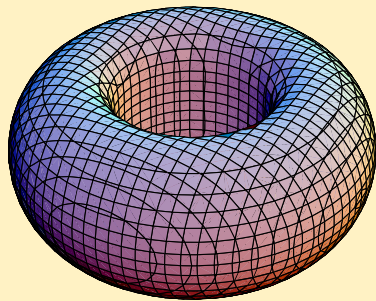
Overview

- The Unitary Correlation Operator Method
- Hartree-Fock-Bogoliubov
 - Particle-Number Projection
 - Inclusion of $3N$ -Forces
- Summary & Outlook

Motivation

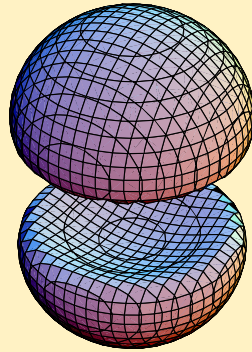
Argonne V18 Deuteron Solution

$$M_S = 0$$
$$\frac{1}{\sqrt{2}}(|\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle)$$



$$\rho_{1,M_S}^{(2)}(\vec{r})$$

$$M_S = \pm 1$$
$$|\uparrow\uparrow\rangle, |\downarrow\downarrow\rangle$$



- **central correlations:**
two-body density is suppressed at low distances
- **tensor correlations:**
angular distribution depends on the relative spin alignments

Central Correlator C_r

radial distance-dependent shift in the relative coordinate of a nucleon pair

$$C_r = \exp\left(-i \sum_{i,j}^A g_{r,ij}[s(r_{ij})]\right)$$

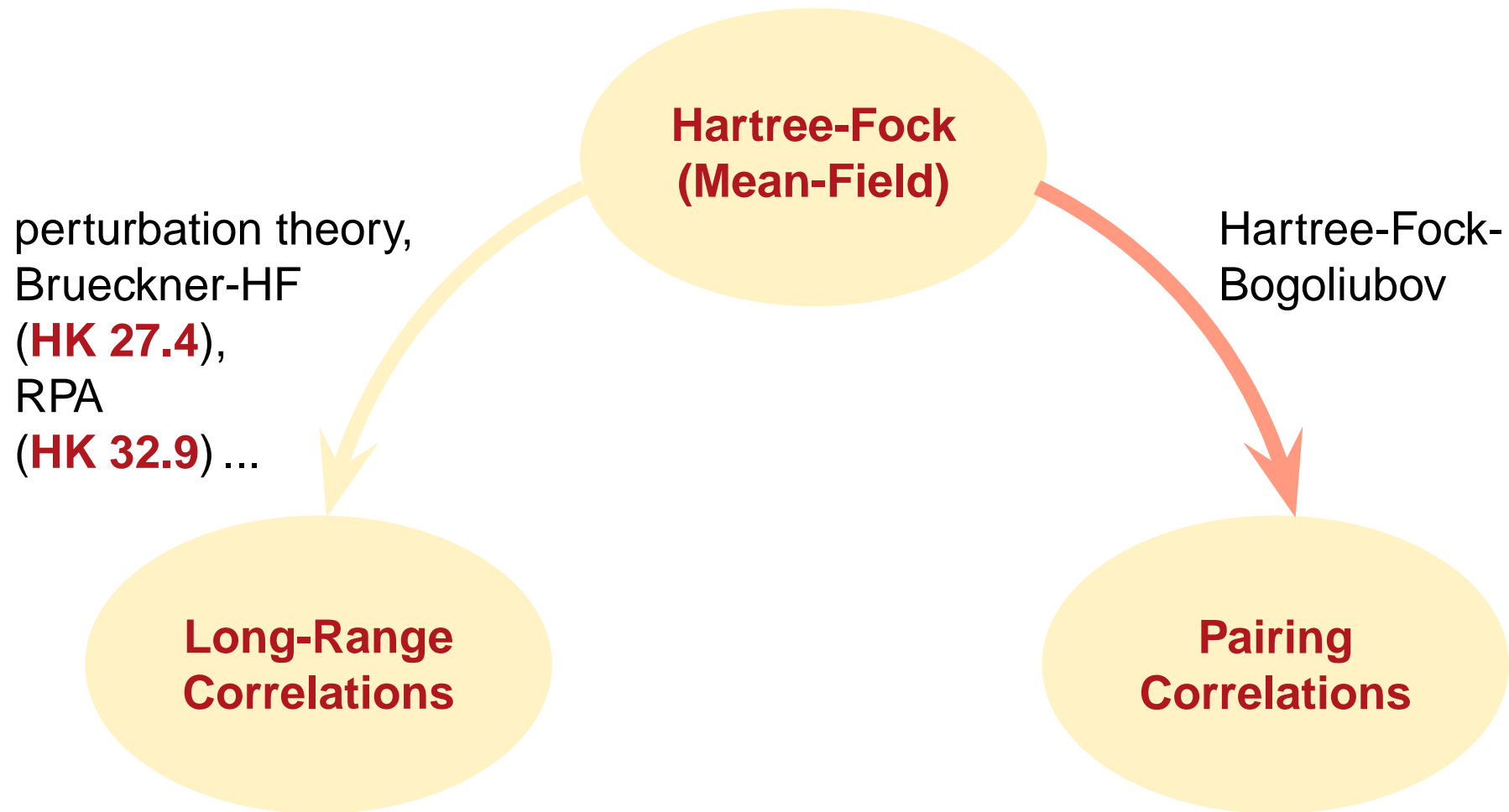
$s(r)$ and $\vartheta(r)$ encapsulate the physics of short-range correlations

Tensor Correlator C_Ω

angular shift, depending on orientation of spin and relative coordinate

$$C_\Omega = \exp\left(-i \sum_{i,j}^A g_{\Omega,ij}[\vartheta(r_{ij})]\right)$$

Beyond Hartree-Fock



HFB Theory Overview

Bogoliubov Transformation

$$\beta_k^\dagger = \sum_q U_{qk} c_q^\dagger + V_{qk} c_q$$

$$\beta_k = \sum_q U_{qk}^* c_q + V_{qk}^* c_q^\dagger$$

where

$$\{\beta_k, \beta_{k'}\} \stackrel{!}{=} \{\beta_k^\dagger, \beta_{k'}^\dagger\} \stackrel{!}{=} 0$$

$$\{\beta_k, \beta_{k'}^\dagger\} \stackrel{!}{=} \delta_{kk'}$$

HFB Densities & Fields

$$\rho_{kk'} \equiv \langle \Psi | c_{k'}^\dagger c_k | \Psi \rangle = (V^* V^T)_{kk'}$$

$$\kappa_{kk'} \equiv \langle \Psi | c_{k'} c_k | \Psi \rangle = (V^* U^T)_{kk'}$$

$$\Gamma_{kk'} = \sum_{qq'} \left(\frac{2}{A} \bar{t}_{\text{rel}} + \bar{v} \right)_{kq', k'q} \rho_{qq'}$$

$$\Delta_{kk'} = \sum_{qq'} \left(\frac{2}{A} \bar{t}_{\text{rel}} + \bar{v} \right)_{kk', qq'} \kappa_{qq'}$$

Energy

$$E[\rho, \kappa, \kappa^*] = \frac{\langle \Psi | \mathbf{H} | \Psi \rangle}{\langle \Psi | \Psi \rangle} \equiv \frac{1}{2} (\text{tr } \Gamma \rho - \text{tr } \Delta \kappa^*)$$

HFB Equations

$$(\mathcal{H} - \lambda \mathcal{N}) \begin{pmatrix} U \\ V \end{pmatrix} \equiv \begin{pmatrix} \Gamma - \lambda & \Delta \\ -\Delta^* & -\Gamma^* + \lambda \end{pmatrix} \begin{pmatrix} U \\ V \end{pmatrix} = E \begin{pmatrix} U \\ V \end{pmatrix}$$

Particle Number Projection

Variation of Projected Energy

$$\delta E(N_0) = \frac{1}{2\pi \langle \mathbf{P}_{N_0} \rangle} \int_0^{2\pi} d\phi \langle e^{i\phi(N-N_0)} \rangle \left\{ \delta \langle \mathbf{H} \rangle_\phi - \left(E(N_0) - \langle \mathbf{H} \rangle_\phi \right) \delta \log \langle e^{i\phi N} \rangle \right\}$$

$$\langle \mathbf{H} \rangle_\phi \equiv \langle \mathbf{H} e^{i\phi N} \rangle / \langle e^{i\phi N} \rangle$$

Lipkin-Nogami + PAV

- power series expansion
- expansion coefficients **not varied**
- indeterminate / numerically unstable at shell closures
- exact PNP after variation

VAP

- higher (but manageable) computational effort
- implement with care: **subtle cancellations between divergences of direct, exchange, and pairing terms**

☞ Structure of **HFB equations is preserved** by both methods!

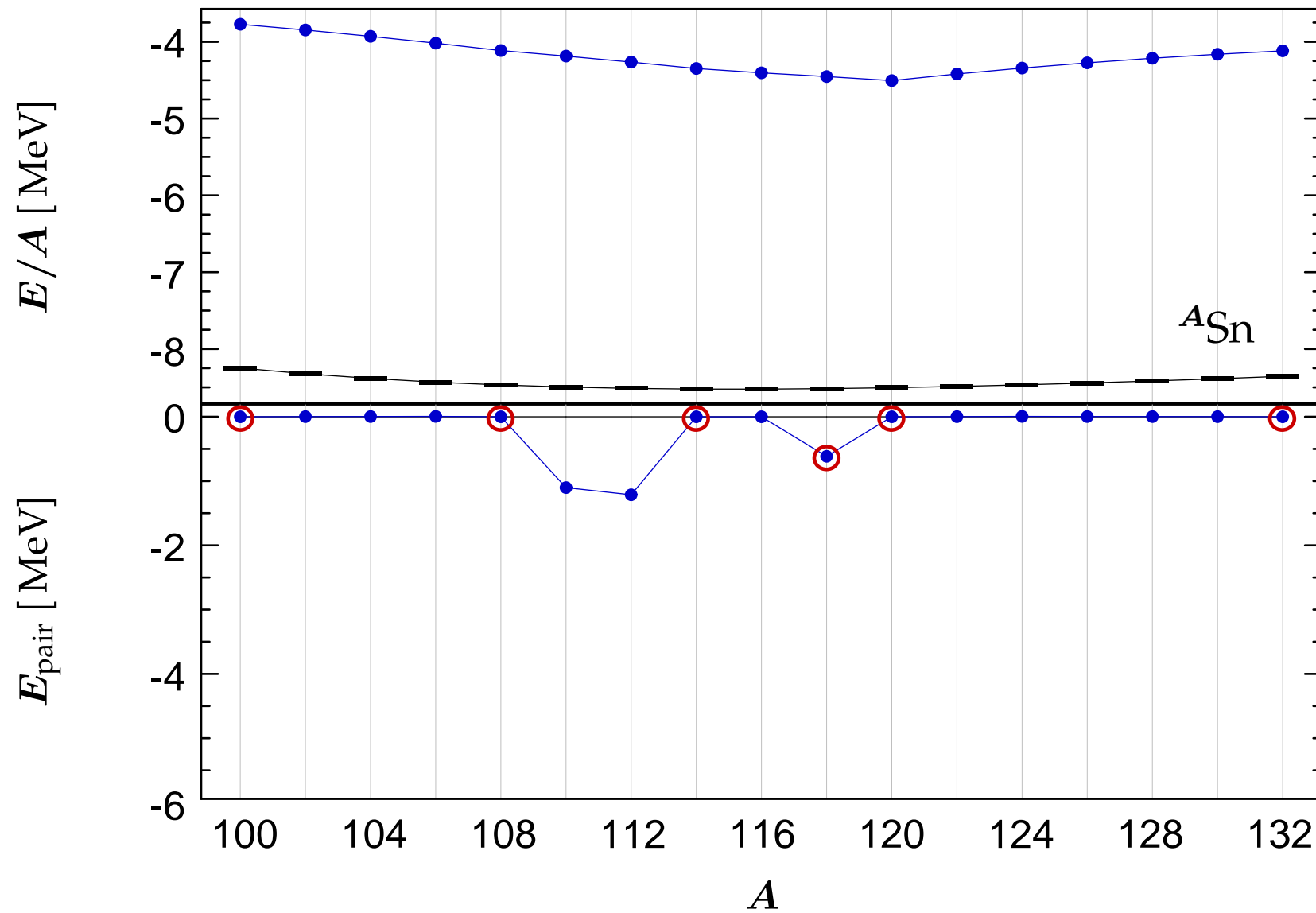
Implementation

- vary **intrinsic energy** $H_{\text{int}} = H - T_{\text{cm}}$
- project on proton and neutron numbers simultaneously:
 $P_{N_0 Z_0} = P_{N_0} P_{Z_0}$

consistent treatment of direct, exchange & pairing terms

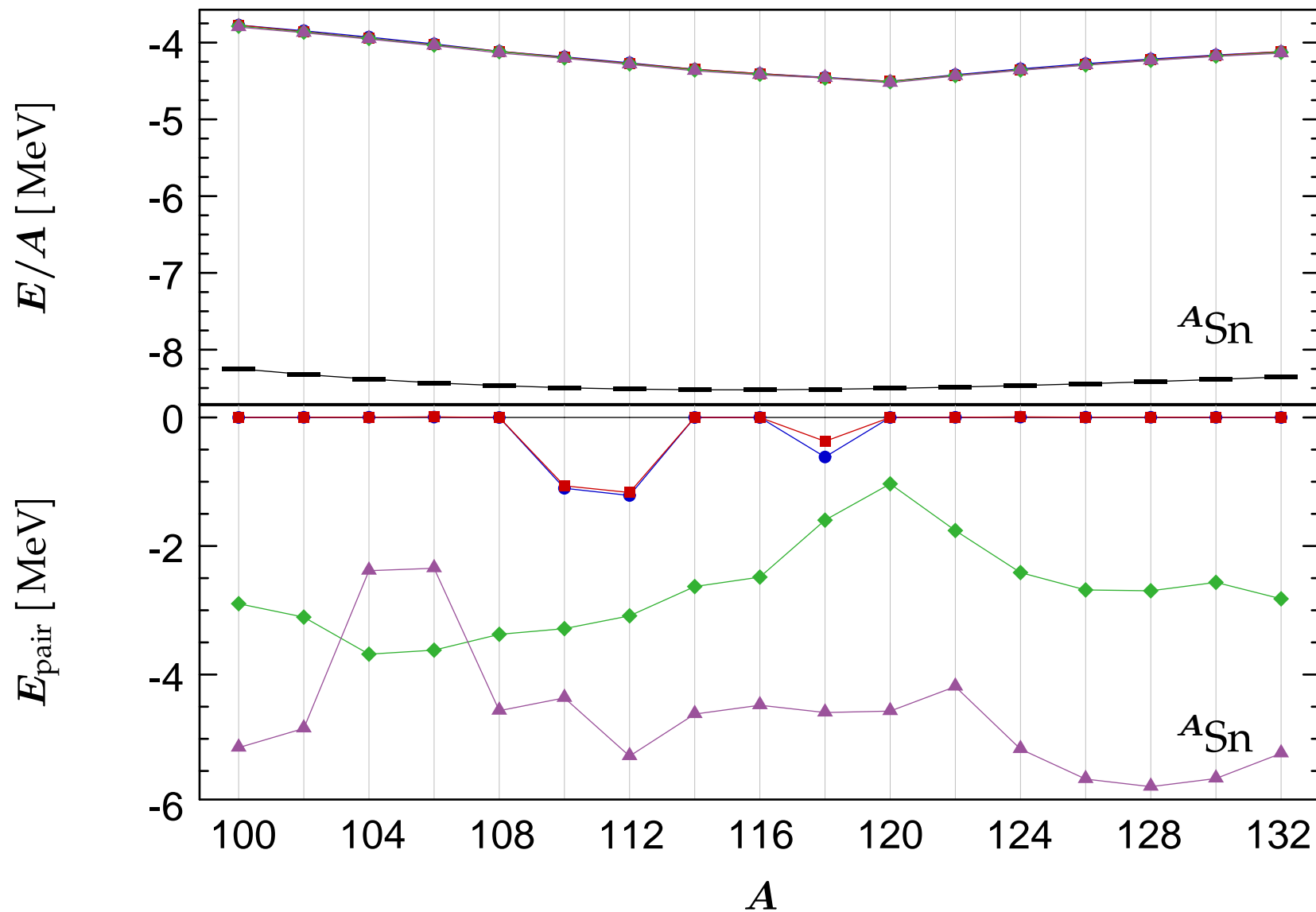
- V_{UCOM} in **ph- and pp-channel** !
- include (anti-)pairing effects from intrinsic kinetic energy and Coulomb interaction
- this level of consistency is **crucial for particle-number projection**

Sn Isotopes: Binding & Pairing Energies



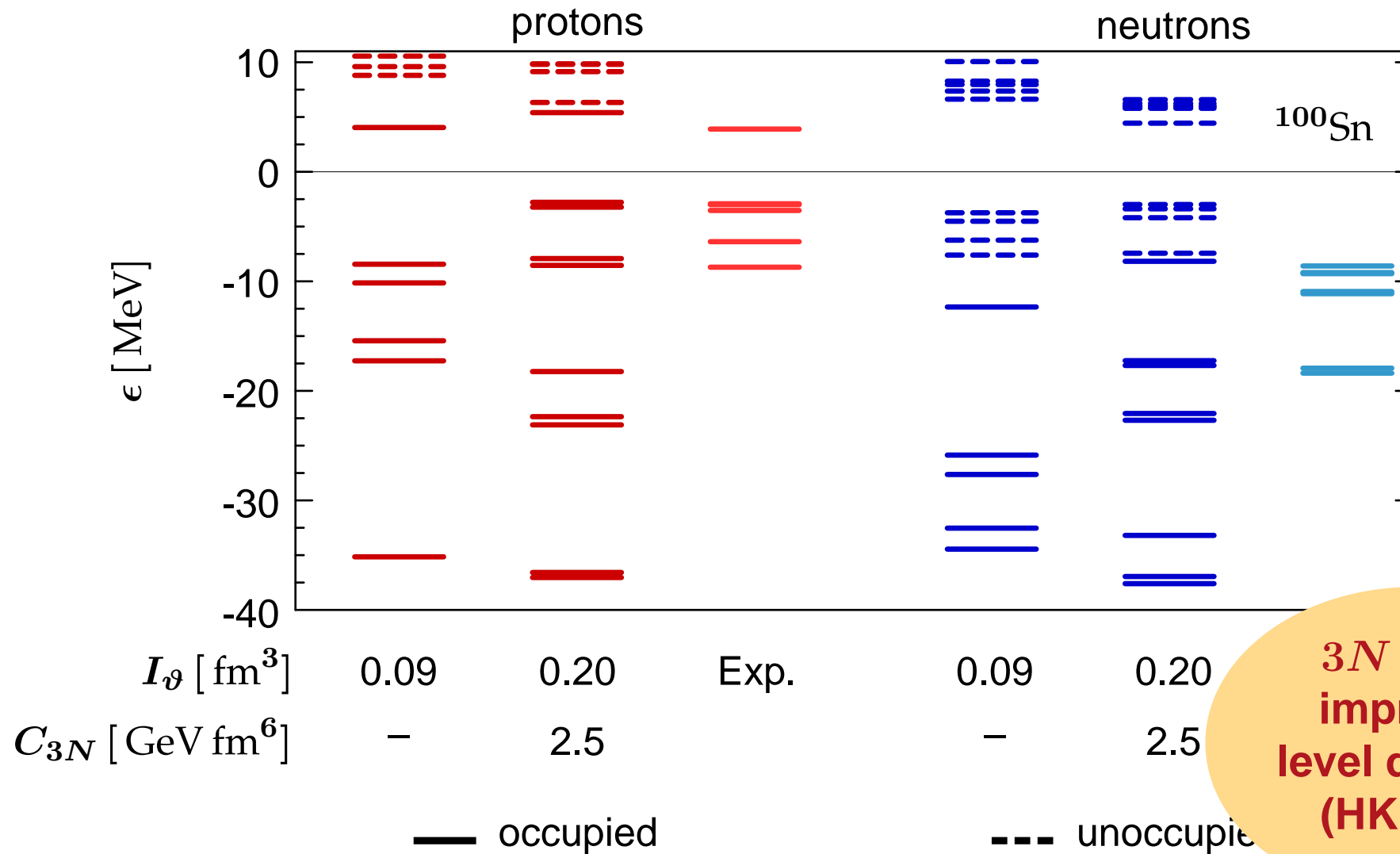
$e_{\text{max}} = 14, l_{\text{max}} = 10$ — exp. ● HFB

Sn Isotopes: Binding & Pairing Energies

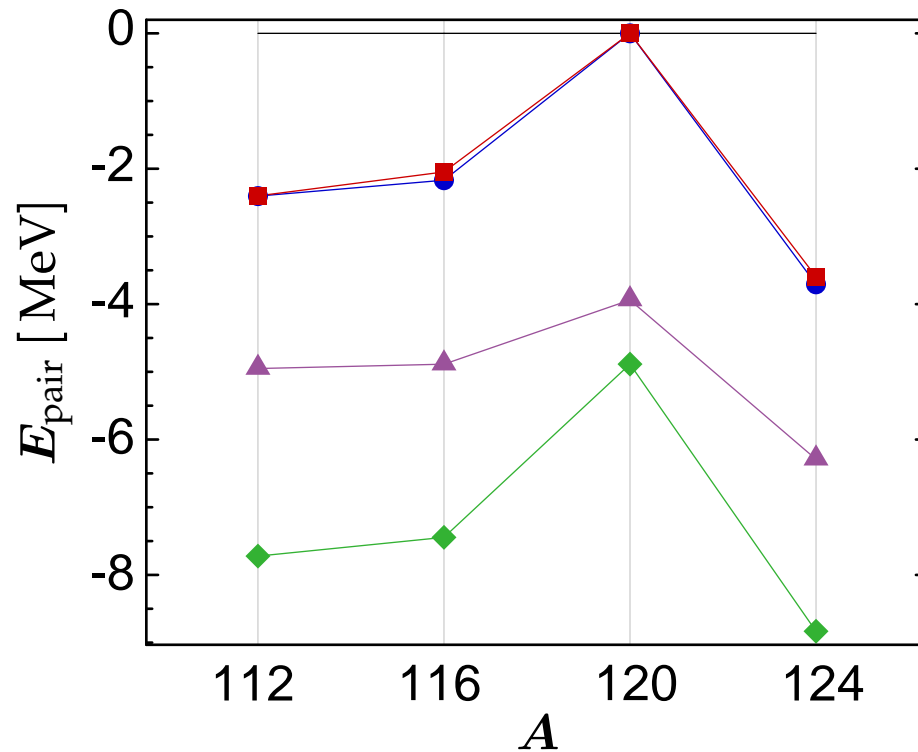


$e_{\text{max}} = 14, l_{\text{max}} = 10$ — exp. ● HFB ■ PAV ◆ LN+PAV ▲ VAP

3N Forces: HF Single-Particle Energies



3N Forces: Pairing



$$e_{\max} = 10, l_{\max} = 8$$

- HFB
- ◆ LN+PAV
- PAV
- ▲ VAP

$$\bar{v}_{kk',qq'}^{[2]} \rightarrow \bar{v}_{kk',qq'}^{[2]} + f \cdot \sum_{rr'} \bar{v}_{kk'r,qq'r'}^{[3]} \rho_{r'r}^{HF}, \quad f = \begin{cases} \frac{1}{3} & \text{expect. values} \\ \frac{1}{2} & \text{fields} \end{cases}$$

phenomenological VAP calculations: $E_{\text{pair}} \simeq 10 - 20 \text{ MeV}$

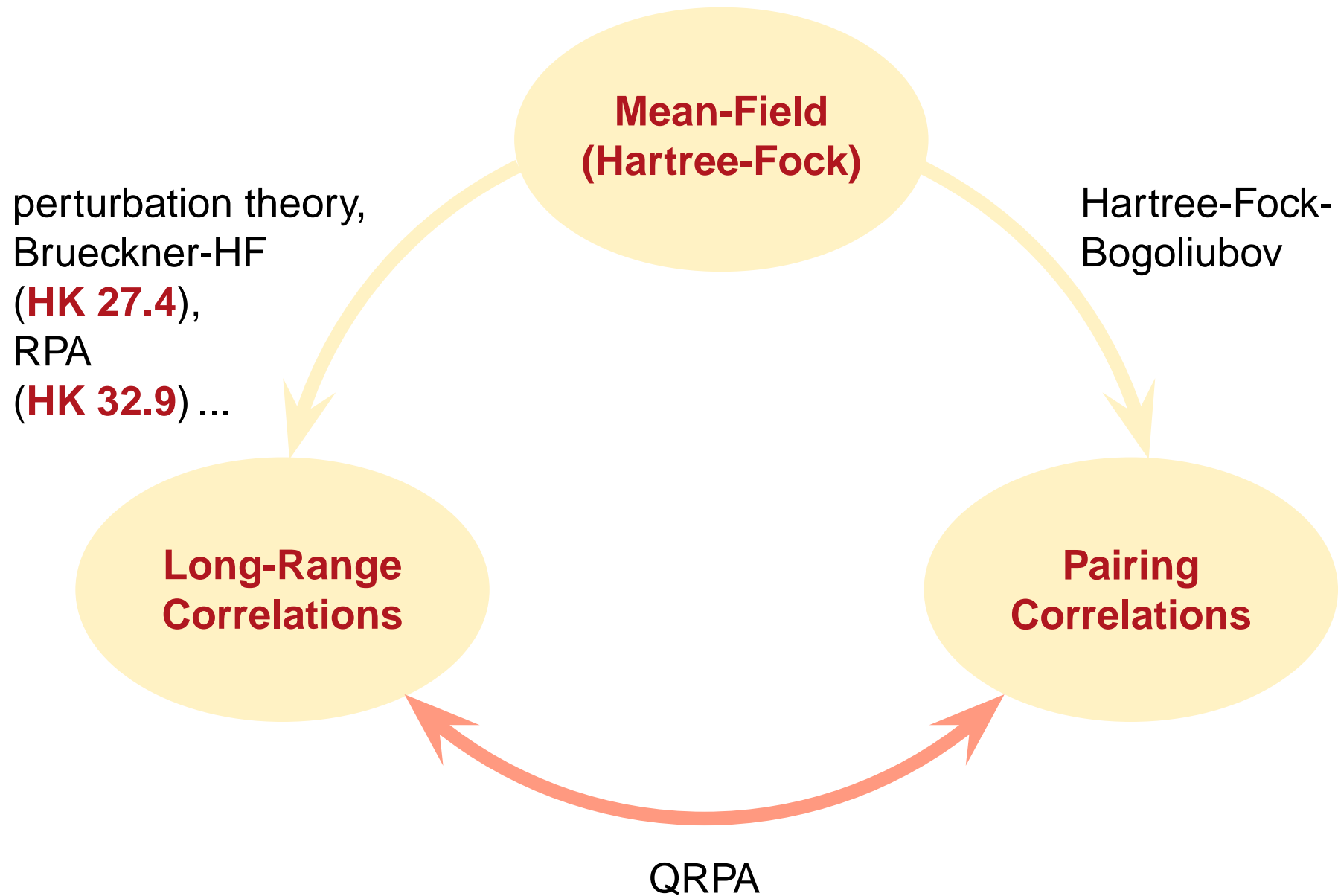
(Stoitsov et al., nucl-th/0610061; Anguiano et al., Phys. Lett. **B545** (2002), 62)

“correct” order
of magnitude
**with realistic
NN int.**

Summary

- **fully consistent** HFB calculations with particle number projection, **based on a Hamiltonian**
 - using V_{UCOM} — a “universal” phase-shift equivalent NN interaction
 - **inclusion of $3N$ -forces** fixes the two-body V_{UCOM} 's problems with single-particle level density (\rightarrow A. Zapp, **HK 32.8**)
 - HFB results are encouraging
- \Rightarrow contribute to a **consistent description of different aspects** of nuclear structure based on V_{UCOM} : HF(B), (Q)RPA, SRPA, MBPT, FMD, NCSM...

Outlook



Outlook

- **Quasiparticle RPA** (benchmark calculations in progress)
- investigation of ***pn* pairing** & ***pn*-QRPA**
⇒ Isobaric Analog & Gamow-Teller Resonances
- **“deformed” UCOM-HFB** (in progress)
- **symmetry restoration** by projection (isospin, parity, angular momentum)

Epilogue...

My Collaborators

- R. Roth, P. Papakonstantinou, A. Zapp, P. Hedfeld
Institut für Kernphysik, TU Darmstadt
- T. Neff, H. Feldmeier
Gesellschaft für Schwerionenforschung (GSI)
- N. Paar
Department of Physics — Faculty of Science, University of Zagreb, Croatia

References

- H. Hergert, R. Roth, nucl-th/0703006, submitted to Phys. Rev. C
- N. Paar, P. Papakonstantinou, H. Hergert, and R. Roth, Phys. Rev. **C74**, 014318 (2006)
- R. Roth, P. Papakonstantinou, N. Paar, H. Hergert, T. Neff, and H. Feldmeier, Phys. Rev. **C73**, 044312 (2006)
- <http://crunch.ikp.physik.tu-darmstadt.de/tnp/>

Optional

Correlated Interaction

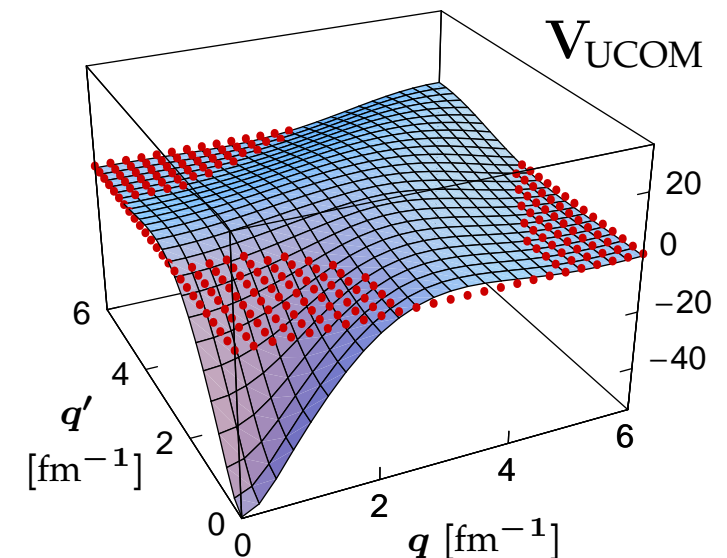
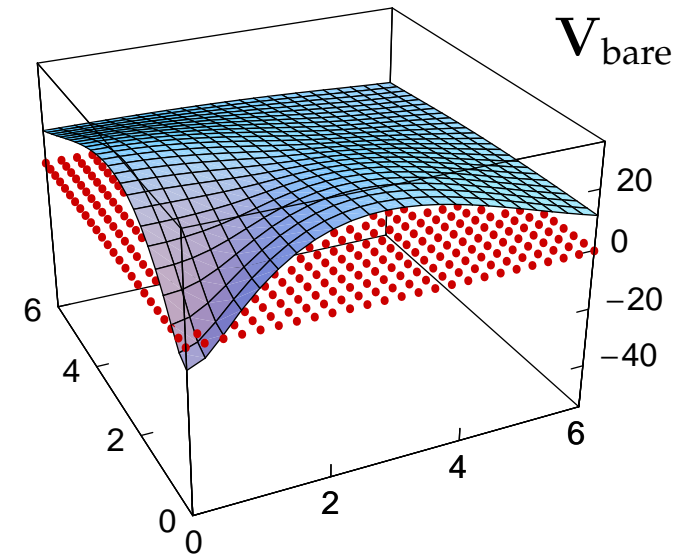
Correlated Hamiltonian

$$\tilde{H} = T^{[1]} + V_{\text{UCOM}} + V_{\text{UCOM}}^{[3]} + \dots$$

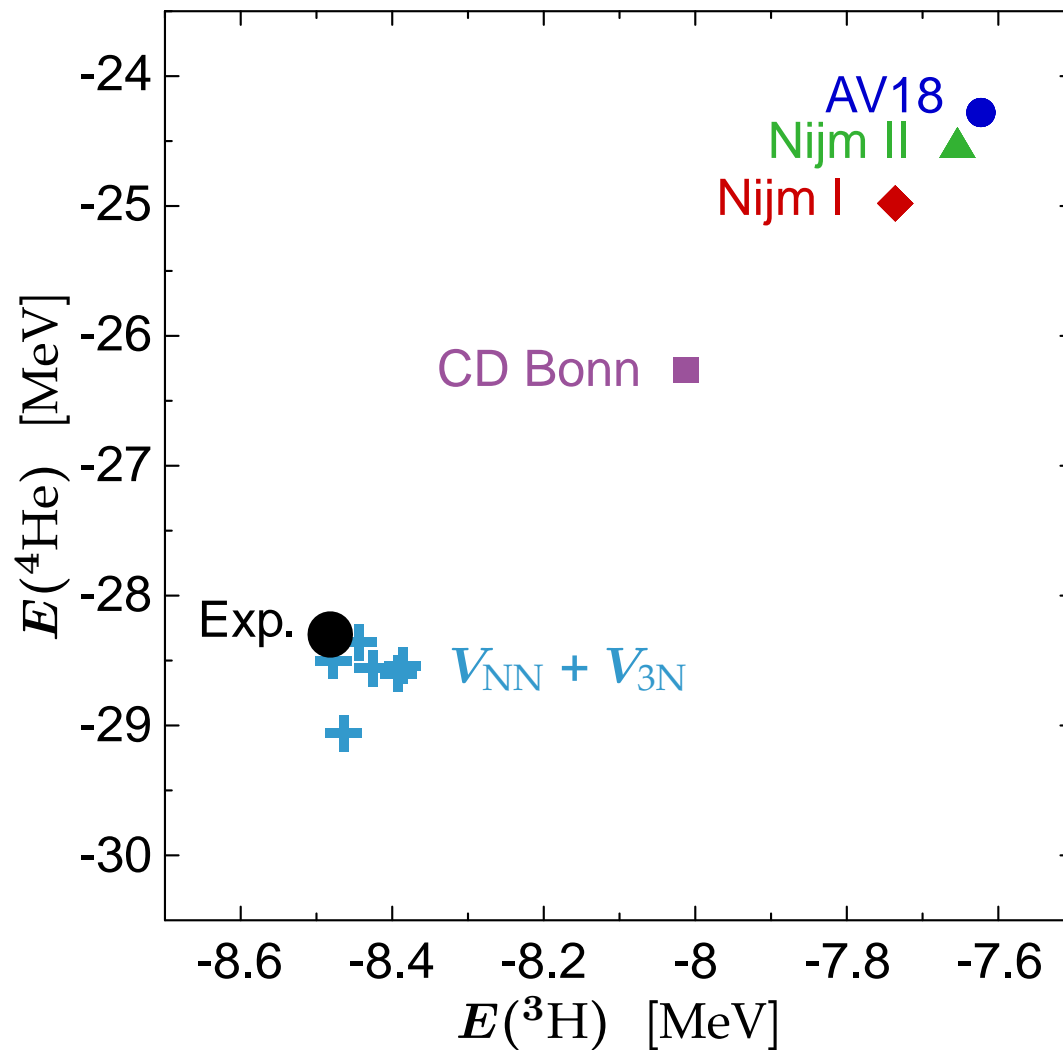
- **closed operator representation** of V_{UCOM} in two-body approximation
⇒ usable with **arbitrary many-body basis**
- V_{UCOM} is **phase-shift equivalent** to the underlying bare nucleon-nucleon interaction
- V_{UCOM} is pre-diagonalized in momentum space, i. e. **high-momentum components are decoupled** (similar to $V_{\text{low-k}}$)

AV18

3S_1



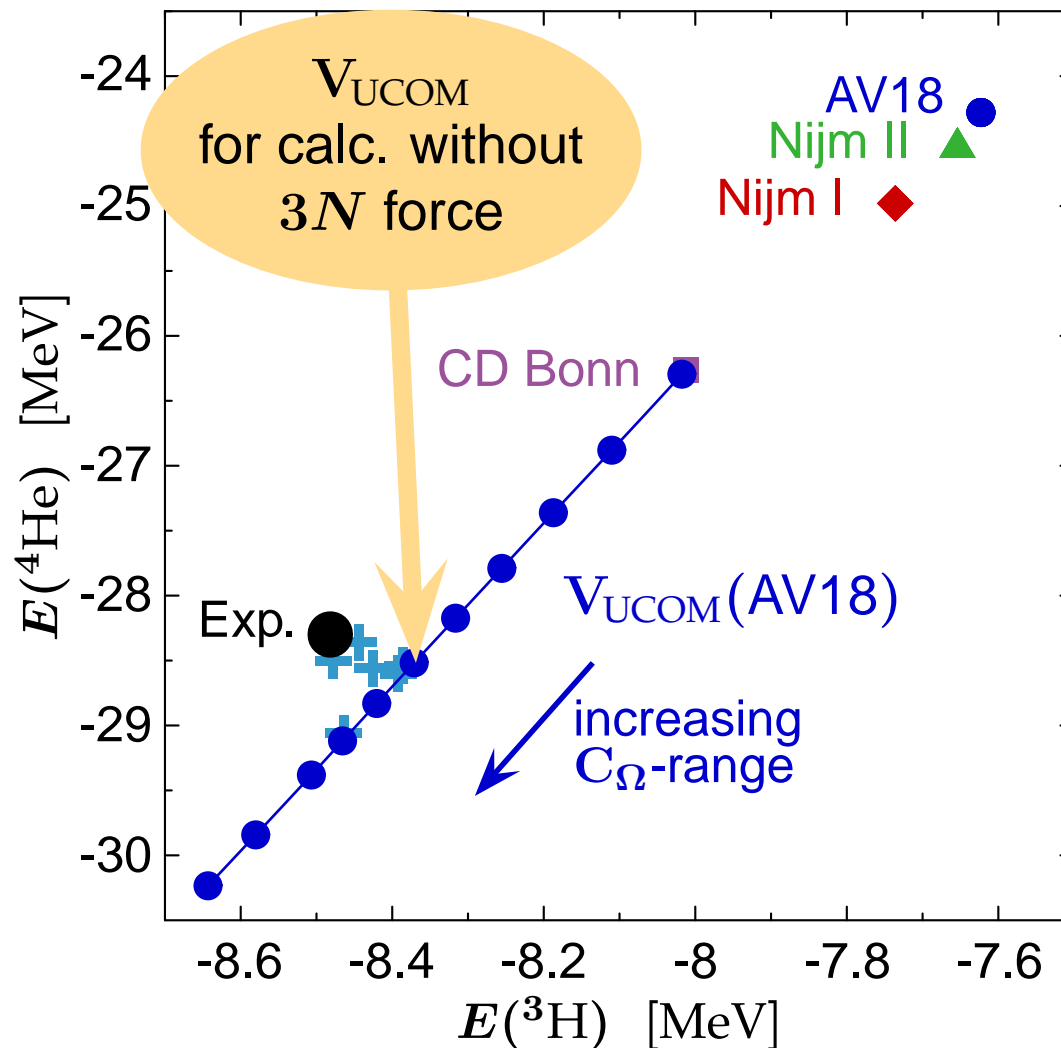
Tjon-Line and Correlator Range



- **Tjon-line:** $E({}^4\text{He})$ vs. $E({}^3\text{H})$ for phase-shift equivalent NN-interactions

Data points: A. Nogga et al., Phys. Rev. Lett. **85**, 944 (2000)

Tjon-Line and Correlator Range



- **Tjon-line:** $E(^4\text{He})$ vs. $E(^3\text{H})$ for phase-shift equivalent NN-interactions
- change of C_{Ω} -correlator range results in shift along Tjon-line

minimize net three-body force by choosing correlator with energies close to experimental value

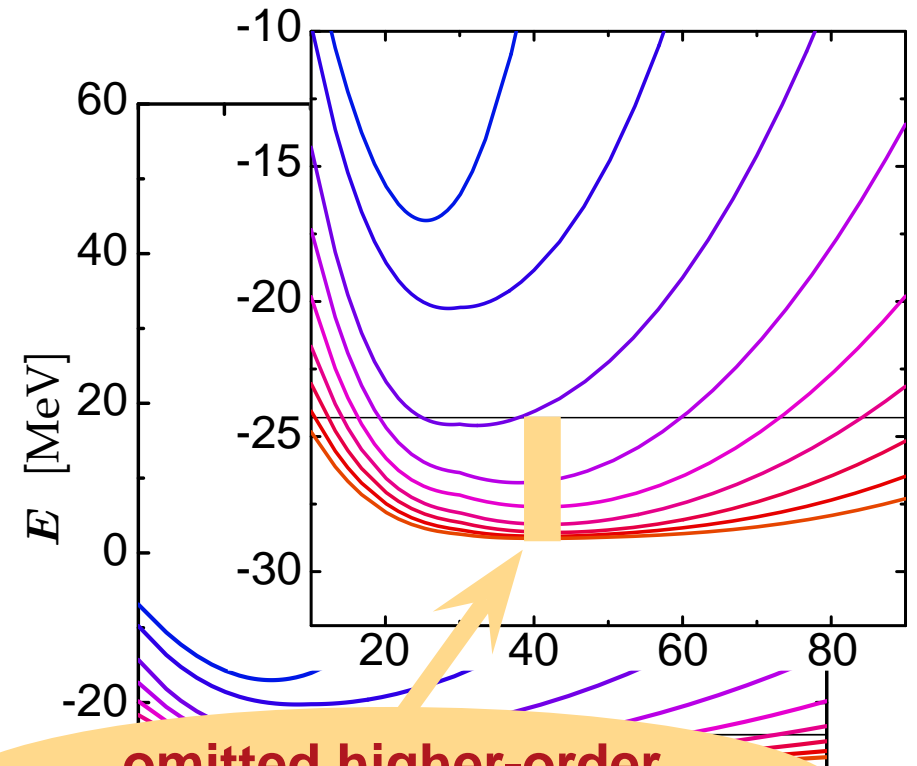
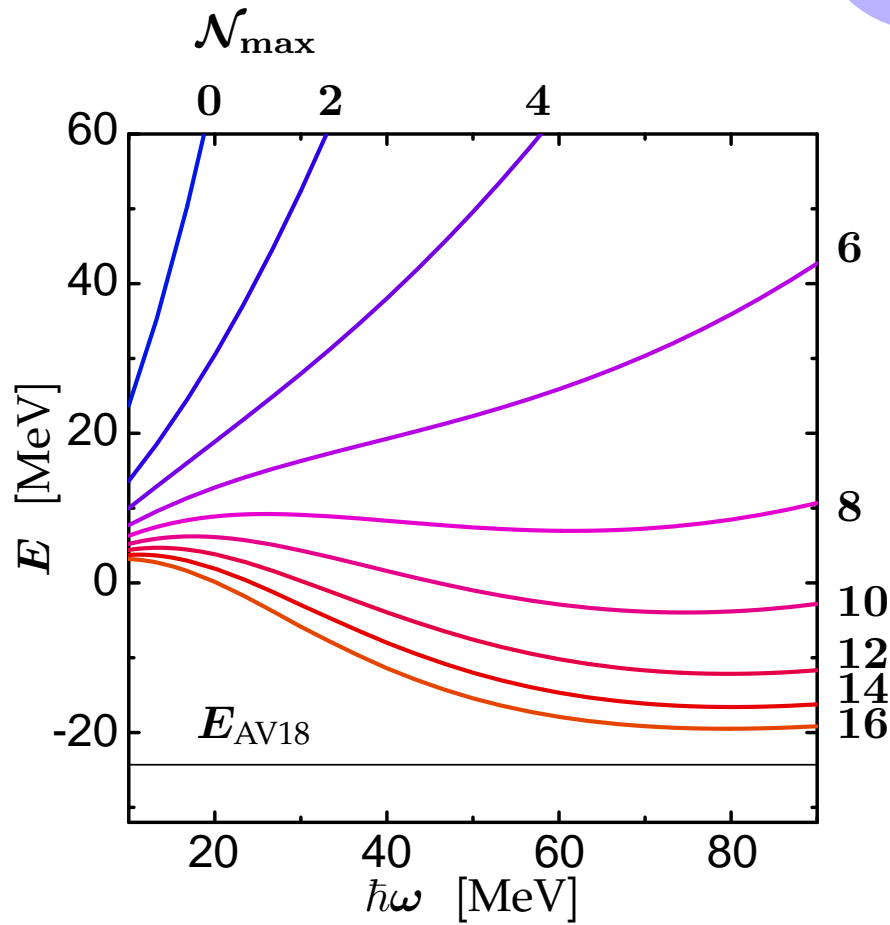
Data points: A. Nogga et al., Phys. Rev. Lett. **85**, 944 (2000)

^4He : Convergence

AV18

^4He

V_{UCOM}



omitted higher-order
cluster contributions

NCSM code by P. Navrátil [PRC 61, 044001 (2000)]