Time Evolution of Bosonic Gases in 1D Two-Colour Superlattices

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Introduction

- Framework : Bose-Hubbard Model
- Two-Colour Superlattice
- Zero-Temperature Phase Diagramm
- Approach
 - Eigenproblem / Time Evolution
 - Adaptive Basis Truncation
 - Excitation by Temporal Modulation
- Results

Summary

Motivation

- bosonic gases in optical lattices are well suited for studying quantum phase transitions
- all relevant parameters can be tuned
- effect of an irregular lattice potential ?
- dynamical signatures as a probe for quantum phases ?

2D-Random Lattice Potential



Framework



Basis Representation
$$|\Psi
angle = \sum_{lpha}^{D} c_{lpha} |\{n_{1}n_{2}\dots n_{I}\}_{lpha}
angle$$

- states are defined by coefficients c_{α}
- coefficients $c_{\alpha}^{(0)}$ of the groundstate are obtained by diagonalisation of the Hamilton matrix

Two-Colour Superlattice



- superposition of two standing waves
- detuning of the wavelengths leads to a sinusoidal modulation of the optical lattice

 Bose-Hubbard model: corresponding on-site energies ε_i with modulation amplitude Δ



Zero-Temperature Phase Diagram 10 Bosons / 10 Lattice Sites



(Quasi-) Bose Glass

rearrangements driven by the competition between interaction strength U and spatial modulation amplitude Δ

Superfluid Phase

in the region of weak interaction strengths / ext. potential all number states contribute to the groundstate

Homogenous Mott Insulator

the interaction strength exceeds the modulation amplitude $(U > \Delta)$ \implies groundstate is dominated by a single number state

Adaptive Basis Truncation

Problem

- solving eigenproblem / time evolution
- basis dimension increases rapidly with number of atoms & lattice sites

\mathbb{K}^{\leq} Basis Truncation

- few number states contribute to low-lying eigenstates
- diagonal elements of Hamiltonian provide estimate for importance of basis states
- relevant number states $|\{n_1 n_2 \cdots n_l\}_{\alpha}\rangle$ satisfy the inequality

$$E_{trunc} \ge \sum_{i=1}^{l} \epsilon_{i} n_{i} + \frac{U}{2} \sum_{i=1}^{l} n_{i} (n_{i} - 1)$$

with the cut-off energy Etrunc

precise description in the vicinity of the Mott insulating phase

poster Q30.2 : Tuesday, 16:30-18:30

Temporal Lattice Modulation

Probing the excitation spectrum...

... by lattice oscillation with amplitude ${\mathcal F}$ and frequency $\omega {:}$

$$V_{ ext{lattice}}(x,t) = V_{ ext{0,lattice}}(x) ig(ext{1} {+} \mathcal{F} \sin \left(\omega t
ight) ig)$$

Time-Dependent Bose-Hubbard Parameters

tunneling strength $J(t) \approx J_0 \exp(-\mathcal{F}\sin(\omega t))$ interaction strength $U(t) \approx U_0 (1 + \mathcal{F}\sin(\omega t))^{1/4}$ on-site energy $\epsilon_i(t) \approx \epsilon_{i,0} (1 + \mathcal{F}\sin(\omega t))$

Roadmap



- 1) Superfluid Phase
- 2) Mott Insulator
- 3) Homogenous Mott Phase
- 4, 5) Bose Glass Phase

(1) Superfluid Phase



(2) Strongly Interacting Regime



- N = 10 bosons, I = 10 sites
- interaction strength $\frac{U_0}{J_0} = 30$

• lattice modulation amplitude $\frac{\Delta}{de} = 0$

- immediate response
- strong resonance at $\omega \approx U_0$ \implies particle-hole excitations of the groundstate



(3) Homogenous Mott Insulator & Superlattice

Setup

- N = 10 bosons, I = 10 sites
- interaction strength $\frac{U_0}{J_0} = 30$
- lattice modulation amplitude $\frac{\Delta}{d_0} = 20$

- resonance at frequency $\omega \approx U_0$ \implies particle-hole excitations of the groundstate
- broadening due to differences in the on-site energies



(4) Quasi-Bose Glass Phase

Setup

- N = 10 bosons, I = 10 sites
- interaction strength $\frac{U_0}{J_0} = 30$
- lattice modulation amplitude $\frac{\Delta}{J_0} = 35$

- rich resonance structure / resonances at low energies appear
- peaks agree with particle-hole excitations of the most propable number state



(5) Quasi-Bose Glass Phase

Setup

- N = 10 bosons, I = 10 sites
- interaction strength $\frac{U_0}{J_0} = 30$

• lattice modulation amplitude $\frac{\Delta}{J_0} = 50$

- resonances are "pushed" towards lower energies
- peaks agree with particle-hole excitations of the most propable number state



- strong resonance in the Mott insulator phase
- broadened by spatial lattice modulation
- larger modulation amplitude "pushes" the resonances towards lower energies
- dynamical signatures seem to be a promising tool to probe quantum phases