

Description of Nuclear Collectivity Using Correlated Realistic Interactions

P. Papakonstantinou

Institute of Nuclear Physics, T.U.Darmstadt, Germany



Overview

- Introduction
 - The Unitary Correlation Operator Method (UCOM)
- Ground-state properties
 - Hartree-Fock and Perturbation Theory
- Collective excitations
 - RPA and beyond: Extended RPA and Second RPA
 - The UCOM Hamiltonian as an effective interaction
- Summary

Introduction

Correlated realistic interactions V_{UCOM}

- **Short-range** central and tensor correlations (**SRC**) described by a **unitary correlation operator** $C = C_{\Omega}C_r$
- Introduce SRC to uncorrelated A -body state or an operator of interest

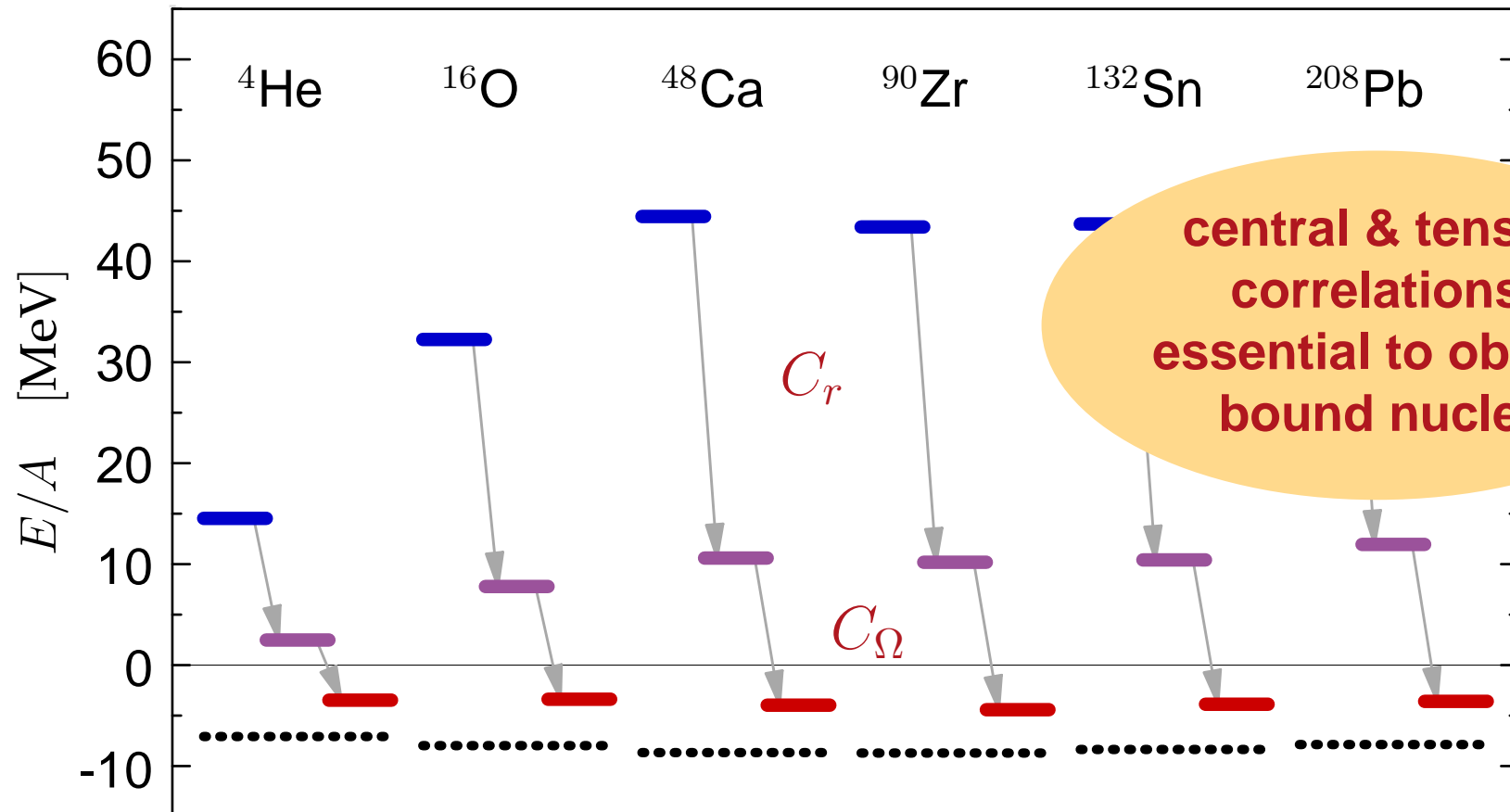
$$\langle \tilde{\Psi} | O | \tilde{\Psi} \rangle = \langle \Psi | C^{\dagger} O C | \Psi \rangle = \langle \Psi | \tilde{O} | \Psi \rangle$$

realistic NN interaction \rightarrow correlated interaction

- Same for **all nuclei**
- **Phase-shift equivalent** to the original NN interaction
- Suitable for use within **simple Hilbert spaces**

Introduction

- expectation value of Hamiltonian (with AV18) for Slater determinant of harmonic oscillator states

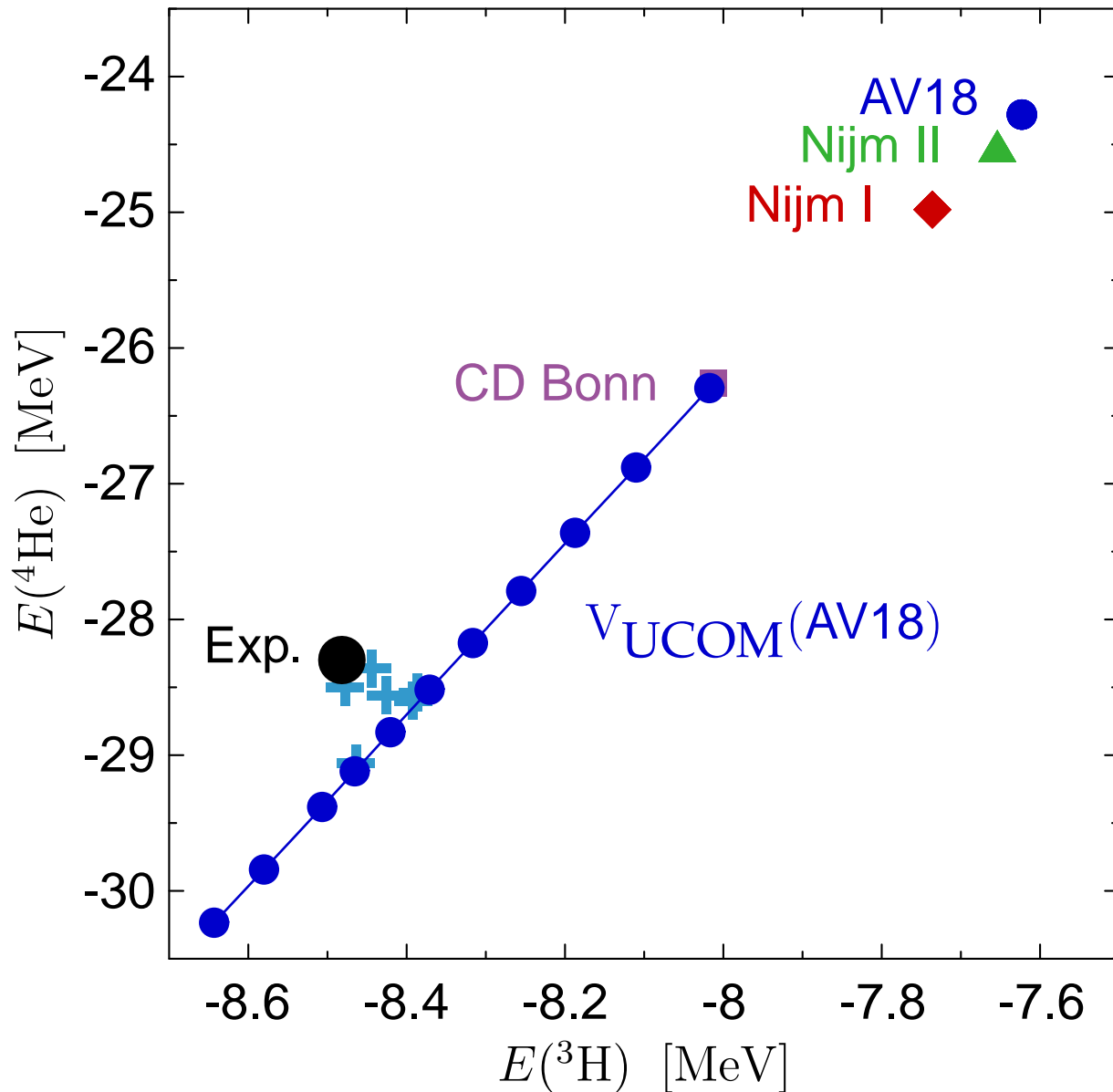


Introduction

Use of the V_{UCOM} in many-body calculations across the nuclear chart:

- **Ground state** properties and **excited states** of closed-shell nuclei:
 - Hartree-Fock calculations
 - Second-order **perturbation theory**
 - Versions of the **RPA**: Standard, Extended, Second RPA
- ...and open-shell ones:
 - **Hartree-Fock-Bogolyubov, Quasi-particle RPA...**
- In what follows, a UCOM Hamiltonian based on the **Argonne V18** NN interaction is used

Tjon Line and Correlator Range



- Tjon line: $E(^4\text{He})$ vs $E(^3\text{H})$ for phase-shift equivalent NN interactions
- Change of tensor-correlator range results in shift along the Tjon line

minimize net
three-body force
by choosing correlator
giving energies close to
the experimental point

Ground-State Properties

Standard Hartree-Fock

- Ground state approximated by a single Slater determinant

$$|\text{HF}\rangle = \mathcal{A}\{|\phi_1\rangle \otimes |\phi_2\rangle \otimes \cdots |\phi_A\rangle\} \longrightarrow \text{no correlations}$$

- Single-particle states are expanded in a H.O. basis

$$|\phi_i\rangle = \sum_{\alpha} D_{i\alpha} |\alpha\rangle \quad ; \quad |\alpha\rangle = |n, (\ell \frac{1}{2}) j m, \frac{1}{2} m_t\rangle$$

- Expansion coeff's $D_{i\alpha}$ determined by minimizing the energy

$$E_{\text{HF}} = \langle \text{HF} | \hat{H}_{\text{int}} | \text{HF} \rangle = \frac{1}{2} \sum_{i,j=1}^A \langle \phi_i \phi_j | T_{\text{rel}} + V_{\text{UCOM}} | \phi_i \phi_j \rangle$$

inclusion of SRC

LRC: extending the model space

Second-order perturbation theory

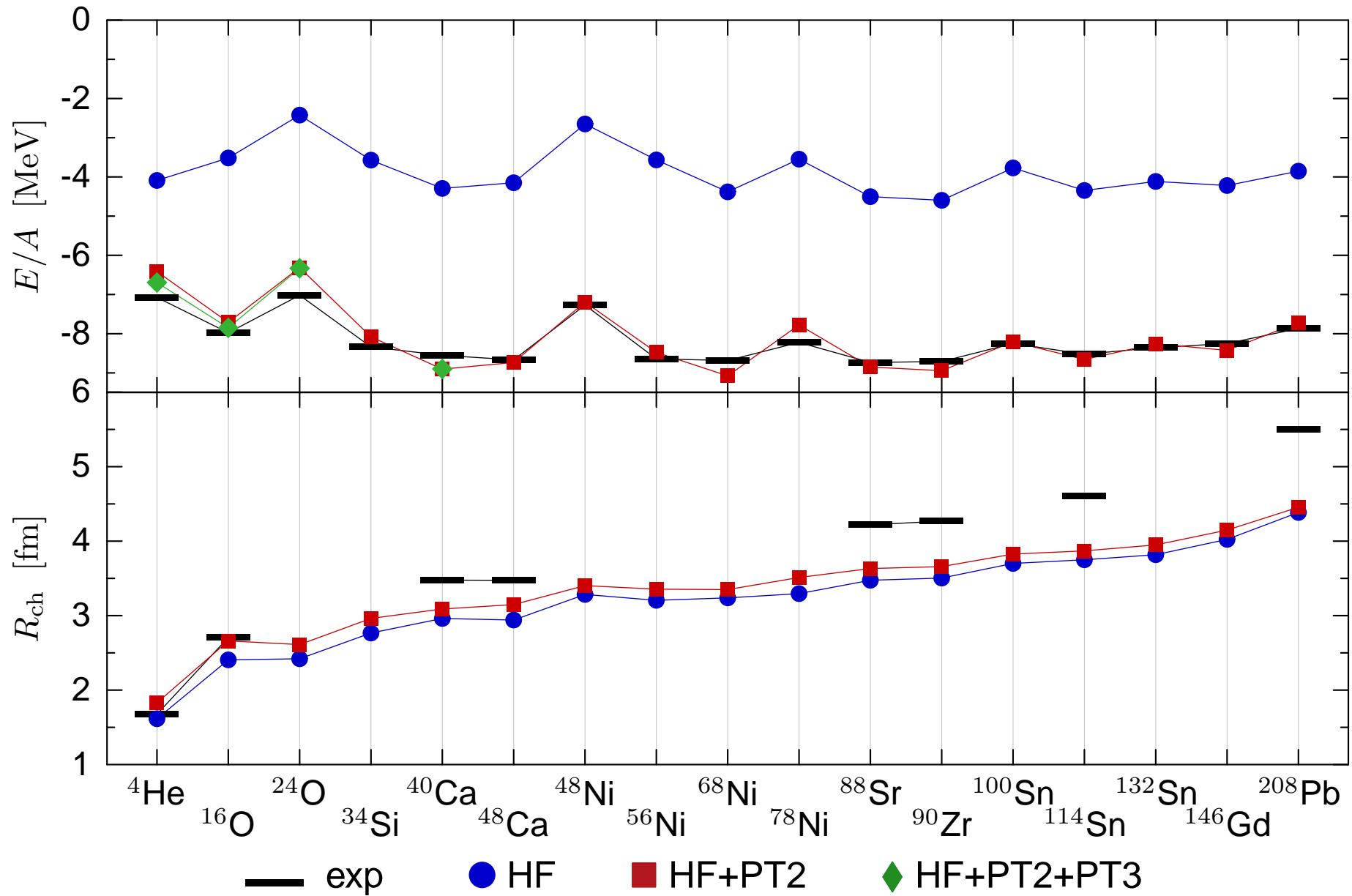
- Binding-energy correction:

$$E^{(2)} = -\frac{1}{4} \sum_{i,j}^{\text{occ}} \sum_{a,b}^{\text{unocc}} \frac{|\langle ij|H_{\text{int}}|ab\rangle|^2}{e_a + e_b - e_i - e_j} ; \quad H_{\text{int}} = T_{\text{rel}} + V_{\text{UCOM}}$$

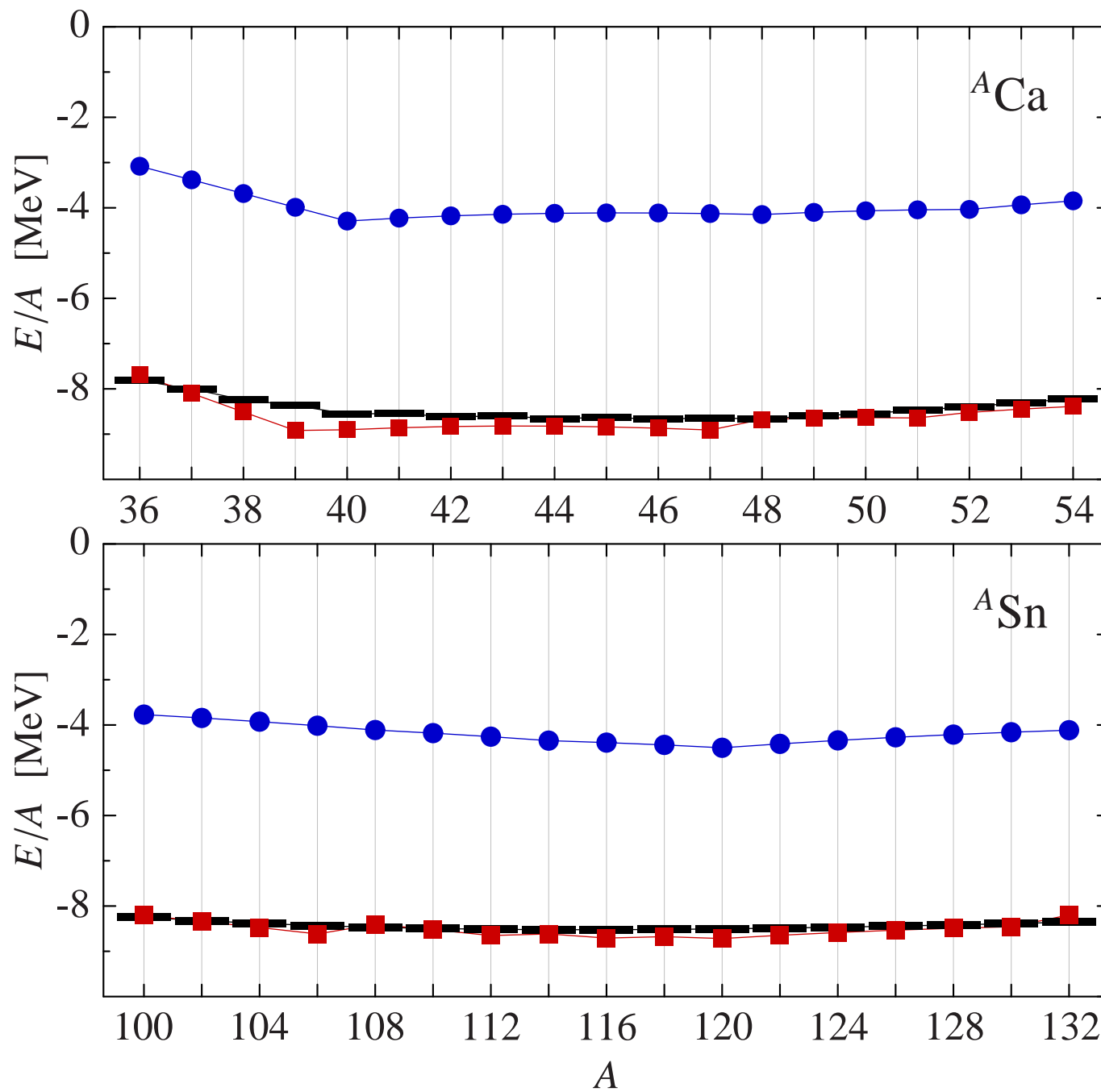
- Modified density matrix and occupation numbers

☞ Modified charge radii

UCOM-HF + PT



UCOM-HF + PT



Missing Pieces

long-range
correlations

genuine
three-body forces

three-body cluster
contributions

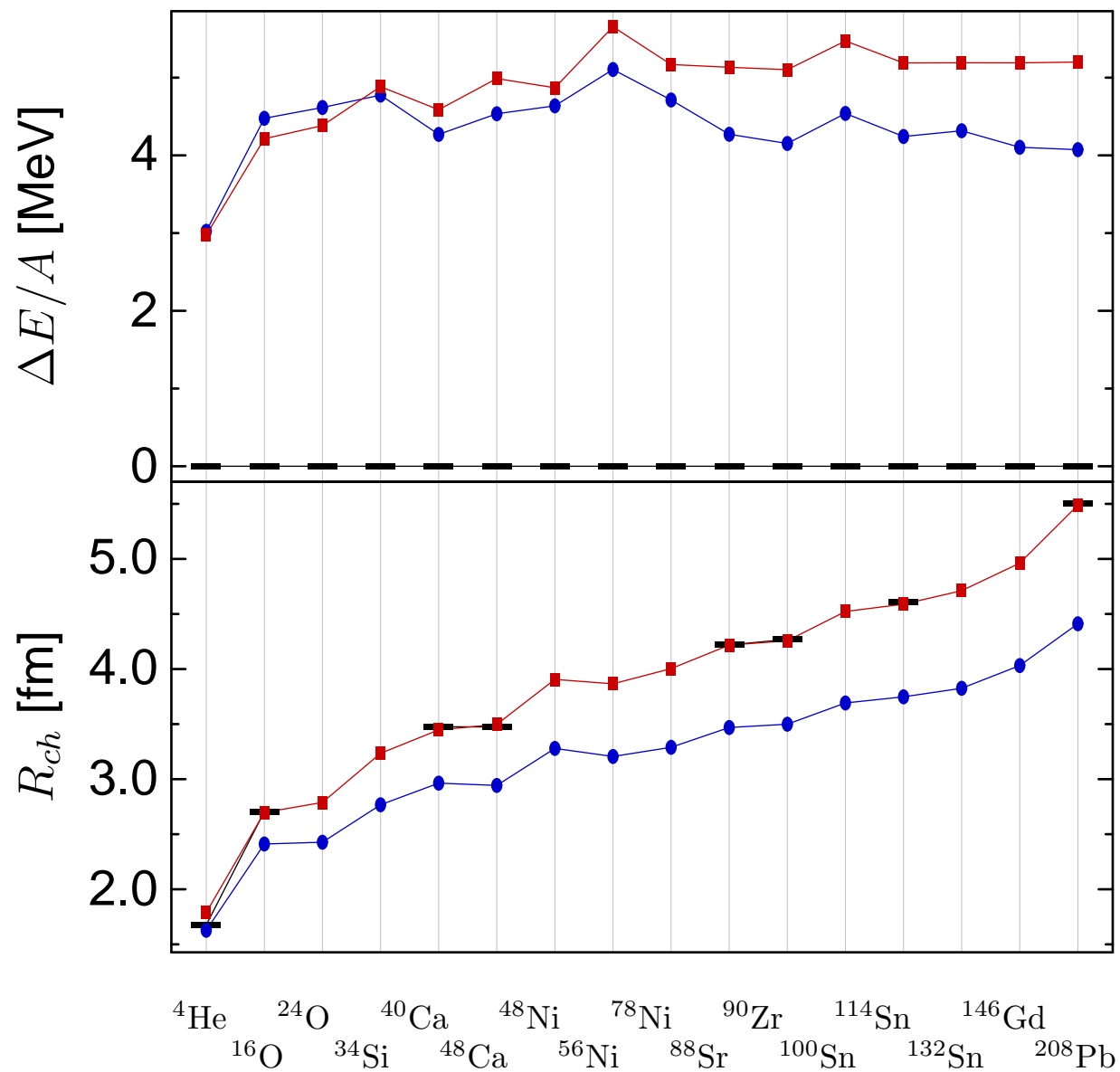
Beyond Hartree-Fock

- residual long-range correlations are **perturbative**
- mostly long-range **tensor correlations**
- easily tractable within MBPT, CI, CC,...

Net Three-Body Force

- small effect on binding energies for all masses
- cancellation does not work for all observables
- construct simple effective three-body force

Three-body force - In progress



Collective Excitations: RPA and beyond

Standard RPA

- Vibration creation operator:

$$Q_\nu^\dagger = \sum_{ph} X_{ph}^\nu O_{ph}^\dagger - \sum_{ph} Y_{ph}^\nu O_{ph} \quad ; \quad Q_\nu |\text{RPA}\rangle = 0 \quad ; \quad Q_\nu^\dagger |\text{RPA}\rangle = |\nu\rangle$$

- Standard RPA - the RPA vacuum is approximated by the HF ground state:

$$\langle \text{RPA} | \dots | \text{RPA} \rangle \rightarrow \langle \text{HF} | \dots | \text{HF} \rangle \quad ; \quad O_{ph} \rightarrow a_p^\dagger a_h$$

- RPA equations in ph -space:

$$\begin{pmatrix} A & B \\ -B^* & -A^* \end{pmatrix} \begin{pmatrix} X^\nu \\ Y^\nu \end{pmatrix} = \hbar\omega_\nu \begin{pmatrix} X^\nu \\ Y^\nu \end{pmatrix}$$

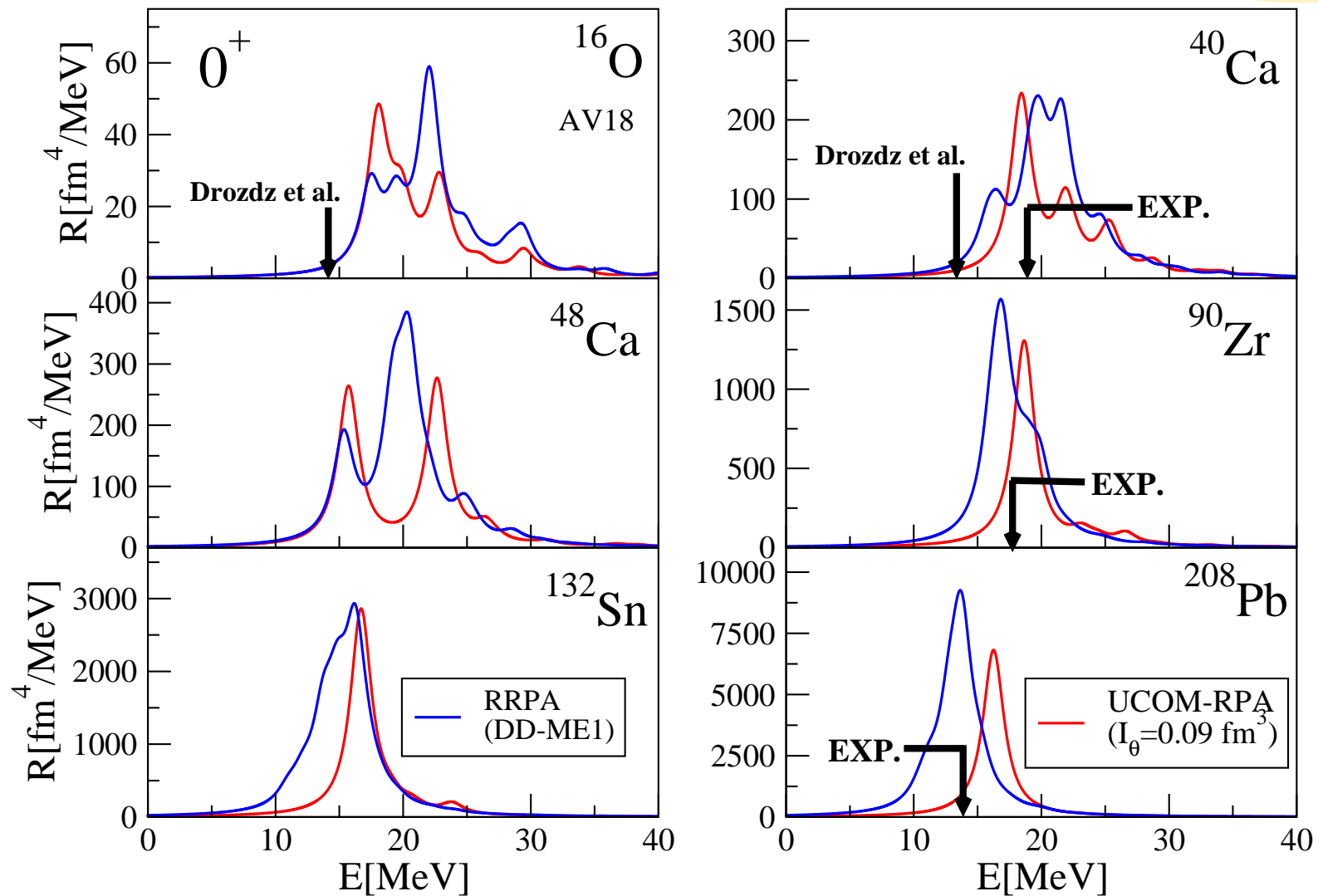
$$A_{ph,p'h'} = \delta_{pp'} \delta_{hh'} (e_p - e_h) + H_{hp',ph'} \quad ; \quad B_{ph,p'h'} = H_{hh',pp'} \quad ; \quad H = H_{\text{int}} = T_{\text{rel}} + V_{\text{UCOM}}$$

➡ Self-consistent HF+RPA: spurious state and sum rules

Standard RPA

Isoscalar monopole response

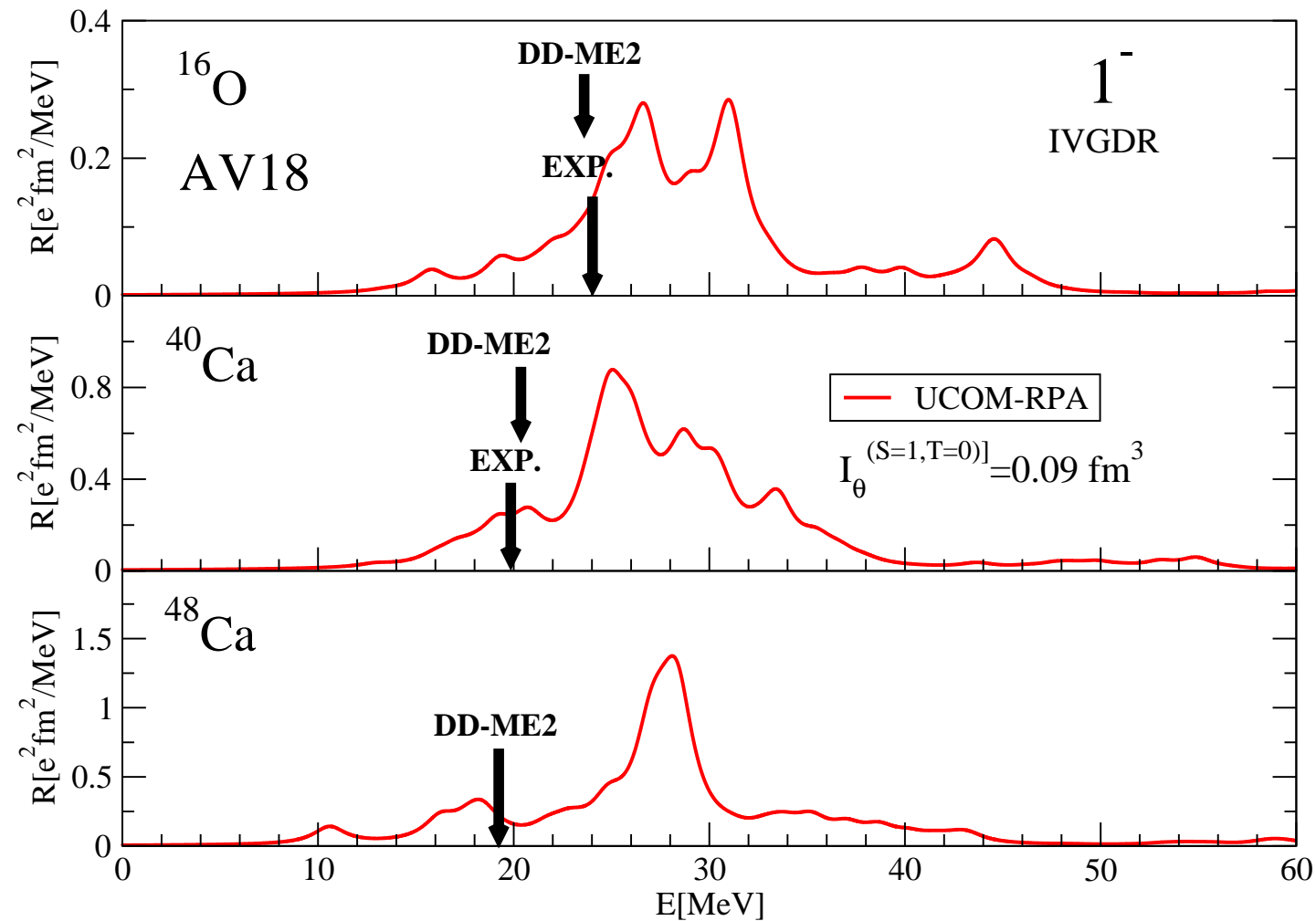
$$N_{\max} = 12$$



Standard RPA

Isovector dipole response

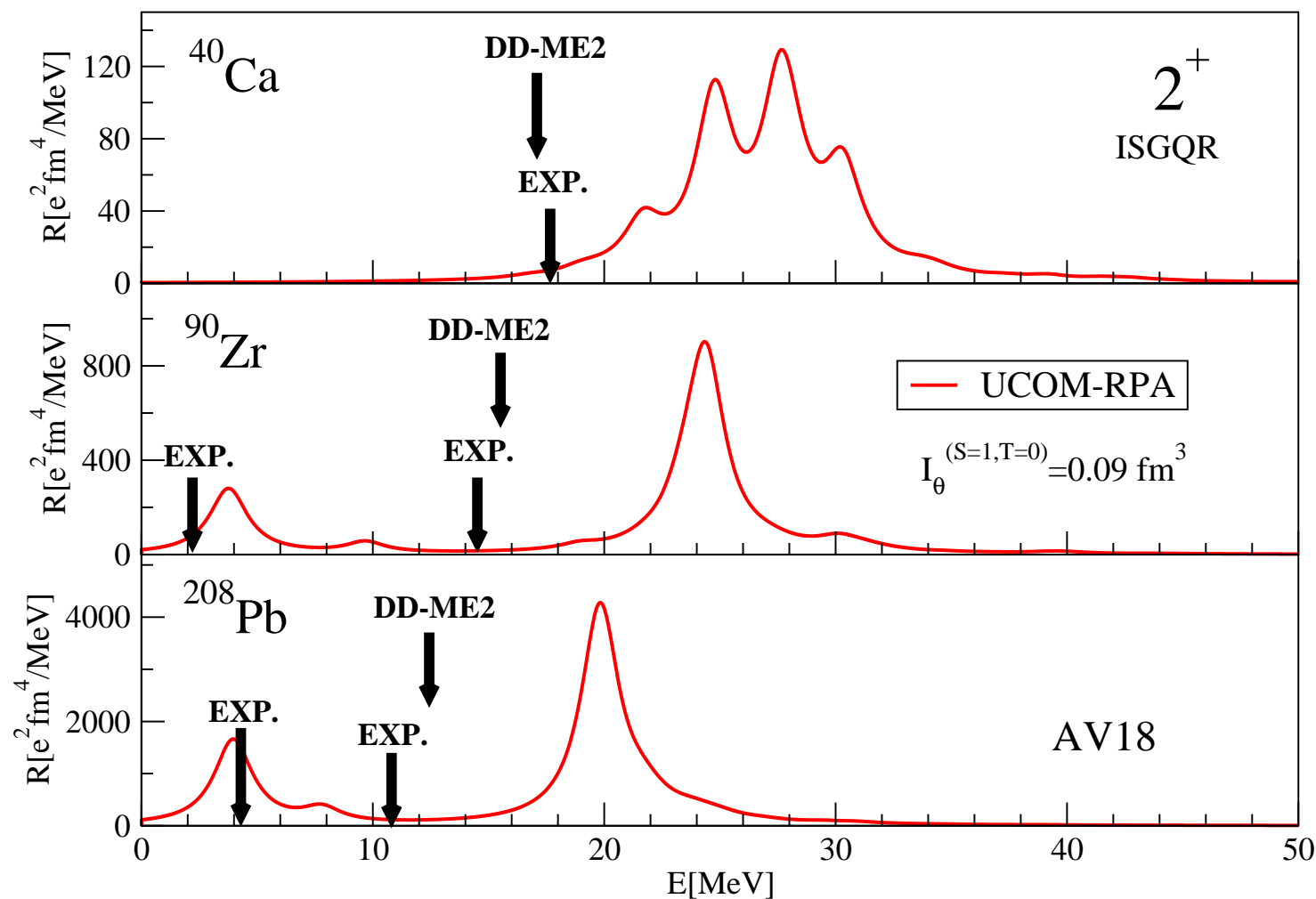
$$N_{\max} = 12$$



Standard RPA

Isoscalar quadrupole response

$$N_{\max} = 12$$



Beyond Standard RPA

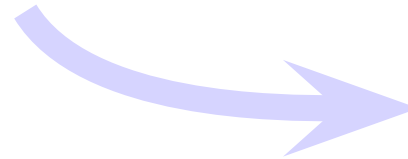
The HF+RPA method is based mainly on the following **approximations**:

☞ Coupling to higher order excitations
($np - nh$) is neglected



Second RPA

☞ The ground state does not deviate much
from the HF ground state



Renormalized RPA,
“Extended” RPA, ...

- Vibration creation operator:

$$Q_\nu^\dagger = \sum_{ph} X_{ph}^\nu O_{ph}^\dagger - \sum_{ph} Y_{ph}^\nu O_{ph} \quad ; \quad Q_\nu |\text{RPA}\rangle = 0 \quad ; \quad Q_\nu^\dagger |\text{RPA}\rangle = |\nu\rangle$$

- Excitations are built on the RPA vacuum. In general,

$$O_{ph} = \sum_{p'h'} N_{ph,p'h'} a_{p'}^\dagger a_{h'}$$

- ERPA is formulated in the natural-orbital basis:

$$O_{ph} \rightarrow D_{ph}^{-1/2} a_p^\dagger a_h \quad ; \quad D_{ph} \equiv n_h - n_p$$

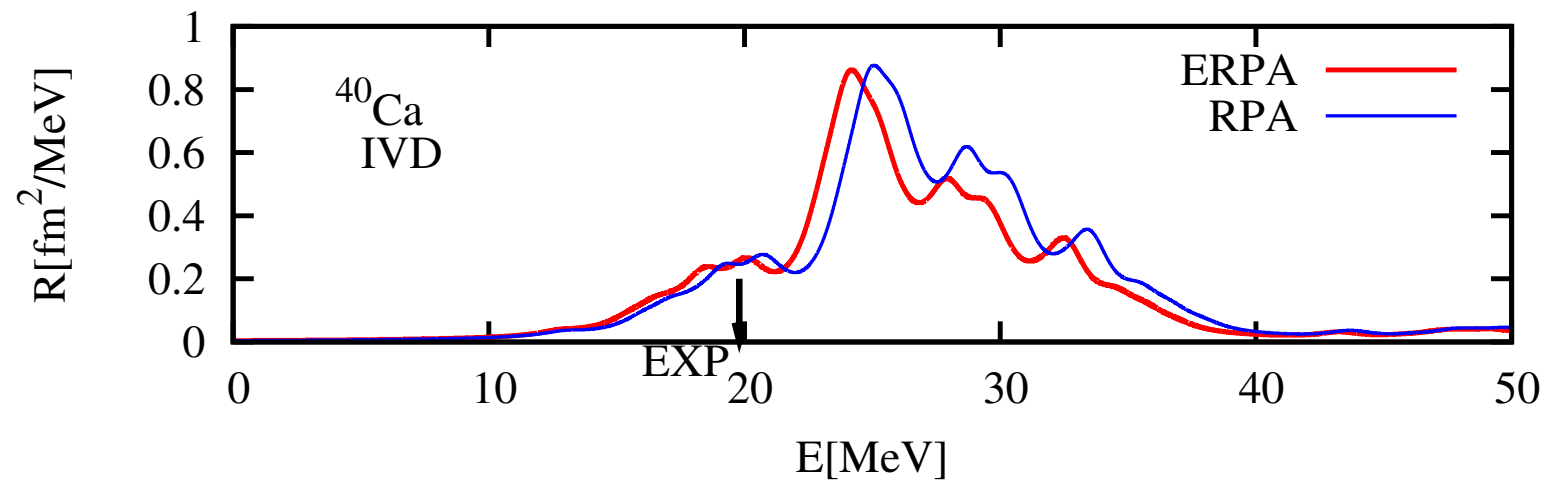
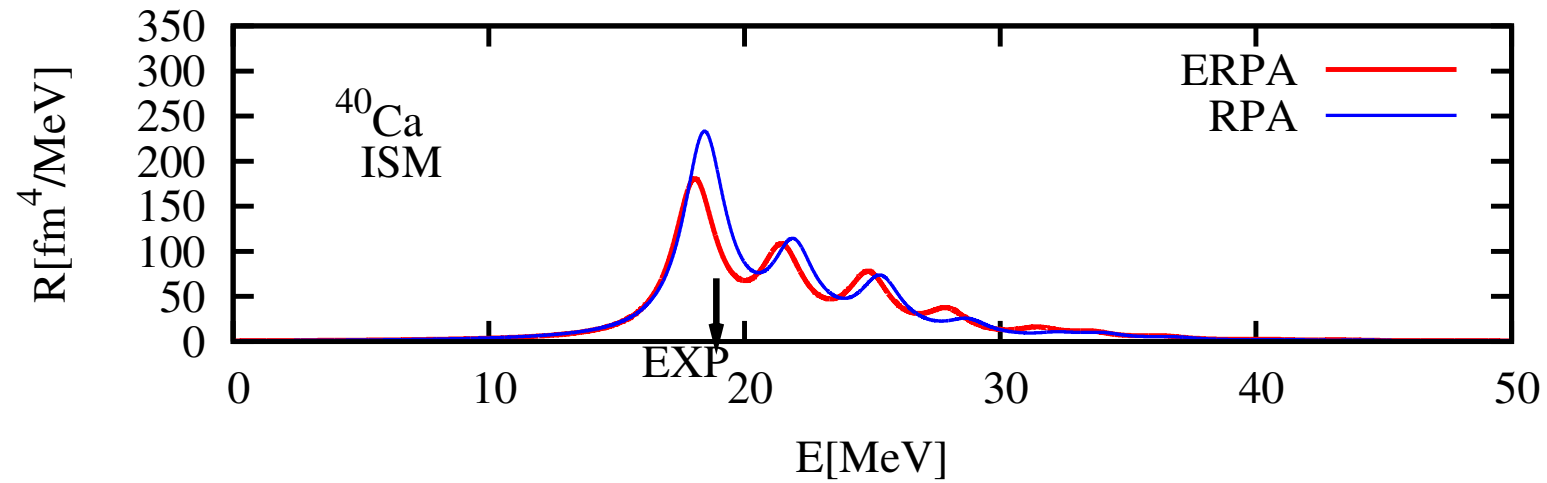
ERPA equations: solved iteratively

$$\begin{pmatrix} A & B \\ -B^* & -A^* \end{pmatrix} \begin{pmatrix} X^\nu \\ Y^\nu \end{pmatrix} = \hbar\omega_\nu \begin{pmatrix} X^\nu \\ Y^\nu \end{pmatrix}$$

$$A_{ph,p'h'} = \delta_{hh'} e_{pp'} - \delta_{pp'} e_{hh'} + D_{ph}^{1/2} D_{p'h'}^{1/2} H_{hp',ph'} \quad ; \quad B_{ph,p'h'} = D_{ph}^{1/2} D_{p'h'}^{1/2} H_{hh',pp'}$$

$$e_{ij} = \sum_k n_k H_{ik,jk}$$

Extended RPA



Second RPA

- **Vibration creation operator:** Includes $2p2h$ configurations

$$Q_{\nu}^{\dagger} = \sum_{ph} X_{ph}^{\nu} O_{ph}^{\dagger} - \sum_{ph} Y_{ph}^{\nu} O_{ph} + \sum_{p_1 h_1 p_2 h_2} \mathcal{X}_{p_1 h_1 p_2 h_2}^{\nu} O_{p_1 h_1 p_2 h_2}^{\dagger} - \sum_{p_1 h_1 p_2 h_2} \mathcal{Y}_{p_1 h_1 p_2 h_2}^{\nu} O_{p_1 h_1 p_2 h_2}$$

- The **SRPA vacuum** is approximated by the HF ground state:

$$\langle \text{SRPA} | \dots | \text{SRPA} \rangle \rightarrow \langle \text{HF} | \dots | \text{HF} \rangle$$

- **SRPA equations** in $ph \oplus 2p2h$ -space:

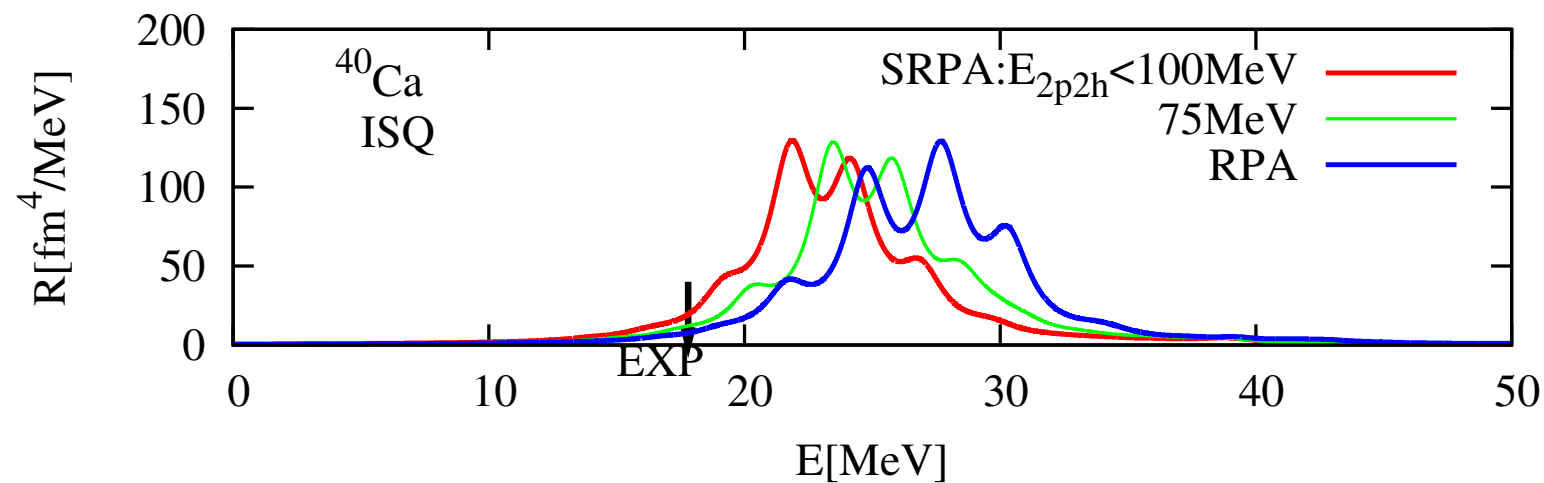
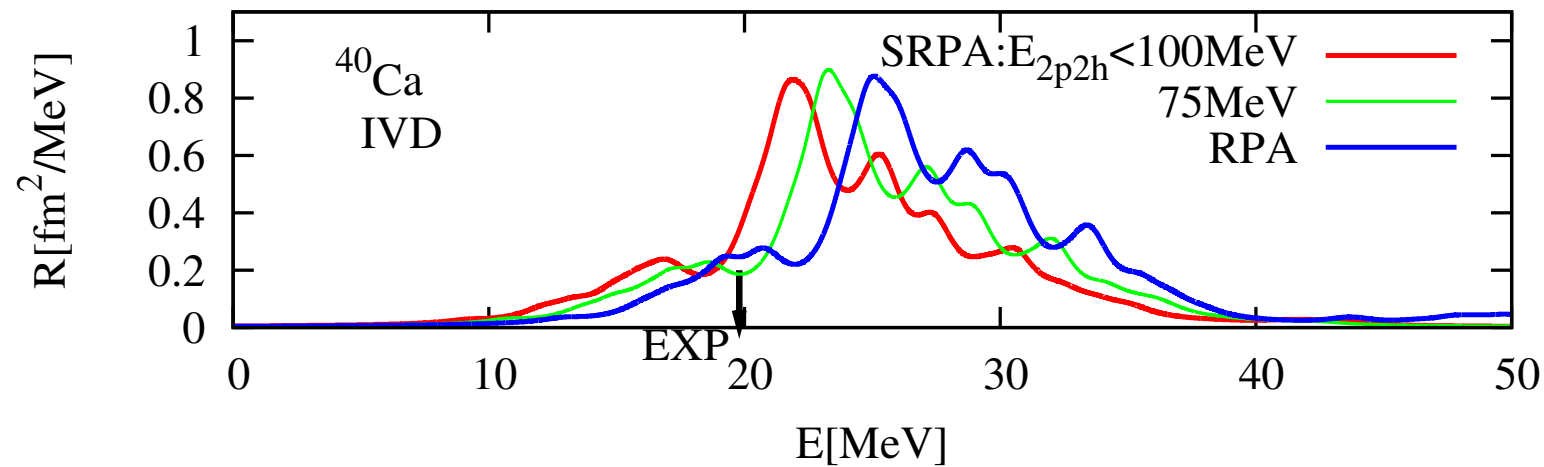
$$\left(\begin{array}{cc|cc} A & \mathcal{A}_{12} & B & 0 \\ \mathcal{A}_{21} & \mathcal{A}_{22} & 0 & 0 \\ \hline -B^* & 0 & -A^* & -\mathcal{A}_{12}^* \\ 0 & 0 & -\mathcal{A}^* & -\mathcal{A}_{22}^* \end{array} \right) \begin{pmatrix} X^{\nu} \\ \mathcal{X}^{\nu} \\ Y^{\nu} \\ \mathcal{Y}^{\nu} \end{pmatrix} = \hbar\omega_{\nu} \begin{pmatrix} X^{\nu} \\ \mathcal{X}^{\nu} \\ Y^{\nu} \\ \mathcal{Y}^{\nu} \end{pmatrix}$$

$$A_{ph,p'h'} = \delta_{pp'} \delta_{hh'} (e_p - e_h) + H_{hp',ph'} ; \quad B_{ph,p'h'} = H_{hh',pp'} ; \quad H = H_{\text{int}} = T_{\text{rel}} + V_{\text{UCOM}}$$

\mathcal{A}_{12} : interactions between ph and $2p2h$ states

\mathcal{A}_{22} : $\delta_{p_1 p'_1} \delta_{h_1 h'_1} \delta_{p_1 p'_1} \delta_{h_1 h'_1} (e_{p_1} + e_{p_2} - e_{h_1} - e_{h_2})$ + interactions among $2p2h$ states

Second RPA - Preliminary



Summary

Nuclear collective excitations using the V_{UCOM}

■ Standard RPA

- GMR properties well reproduced
- GDR, GQR centroids overestimated
- Properties of the V_{UCOM} as an “effective interaction”

■ “Extended” RPA

- Small effect of explicit ground-state correlations

■ Second RPA

- Preliminary results show sizable effect of extended space

👉 Role of residual long-range correlations and three-body terms

Thank you!

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References and further information

- <http://crunch.ikp.physik.tu-darmstadt.de/tnp/>