

# Quark and gluon propagators in dense 2-colour matter

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QHQCD, St Goar, 18 March 2008

# Outline

## Background

- Global symmetries of  $QC_2D$
- $QC_2D$  vs QCD

## Formalism

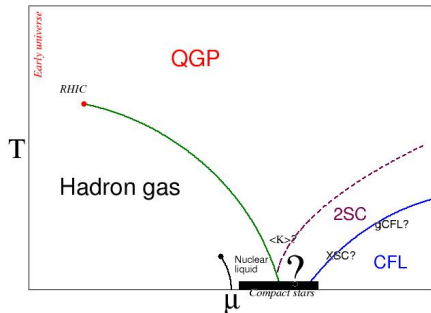
- Tensor structures
- Lattice formulation

## Results

- Bulk thermodynamics
- Gluon propagator results
- Quark propagator results

## Summary

# Background



- ▶ A plethora of phases at high  $\mu$ , low  $T$
- ▶ Based on models and perturbation theory
- ▶ Details depend on diquark gaps and strange quark mass

## Diquark condensation

- ▶ One-gluon exchange is **attractive**
- ▶ Energetically favourable to pair two quarks on opposite sides of Fermi surface  $\implies$  **BCS instability**
- ▶ In **QCD** this breaks the **gauge symmetry**  
 $\implies$  **colour superconductivity**
- ▶ Ground state at ultra-high densities has
$$SU(3)_L \otimes SU(3)_R \otimes SU(3)_c \rightarrow SU(3)_{L+R+c}$$
- ▶ Mismatch between strange and light quark Fermi momenta at intermediate densities  $\implies$  less symmetric pairing

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So standard Monte Carlo importance sampling can not be used!

We can still gain some insight:

- ▶ Approach from high(er) temperature
- ▶ Effective theories where problem is absent or reduced:  
HDET, NJL, ...
- ▶ **QCD-like theories without a sign problem**
- ▶ QC<sub>2</sub>D studies by Hands&Morrison; Muroya, Nakamura, Nonaka; Kogut&Sinclair; Allés, d'Elia, Lombardo; Chandrasekharan&Jiang, ...

## Global symmetries of QC<sub>2</sub>D

**Quarks** and **antiquarks** are in the same representation

Anti-unitary symmetry:  $\mathcal{K}\mathcal{M}\mathcal{K}^{-1} = \mathcal{M}^*$  with  $\mathcal{K} \equiv C\gamma_5\tau_2$

$$\begin{aligned}\mathcal{L} &= \bar{\psi}(\gamma_\nu D_\nu - \mu\gamma_0 + m)\psi \\ &= i\Psi^\dagger\sigma_\nu(D_\nu - \mu B_\nu)\Psi + \frac{1}{2}m\Psi^T\sigma_2\tau_2\hat{M}\Psi \\ \Psi &= \begin{pmatrix} \psi_L \\ \sigma_2\tau_2\psi_R^* \end{pmatrix}, \quad B_\nu = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}\delta_{\nu 0}, \quad \hat{M} = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}\end{aligned}$$

$m = \mu = 0$ : **global SU(2N<sub>f</sub>)** symmetry



# Chiral symmetry breaking

## Chiral condensate

$$\bar{\psi}\psi = -\frac{1}{2}\Psi^T \sigma_2 \tau_2 \hat{M} \Psi + h.c.$$

$\langle \bar{\psi}\psi \rangle \neq 0$  breaks  $SU(2N_f)$   $\longrightarrow$   $Sp(2N_f)$

$\Rightarrow N_f(2N_f - 1) - 1$  Goldstone modes

$N_f = 2$ : 5 modes

$\bar{\psi} \vec{\sigma} \gamma_5 \psi$  pion     $\psi^T \epsilon \tau_2 C \gamma_5 \psi$ ,  $\bar{\psi} \epsilon \tau_2 C \gamma_5 \bar{\psi}^T$  scalar diquark

**Note:** Staggered fermions have different symmetry breaking pattern!

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Diquarks are colour singlets in  $QC_2D$

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Condensation of tightly bound diquarks (Goldstone baryons)

↔ **Chiral perturbation theory**

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**Bardeen–Cooper–Schrieffer:**

Pairing of quarks near the **Fermi surface**

$$\langle \psi\psi \rangle \propto \Delta\mu^2$$

## $QC_2D$ vs $QCD$

### What we cannot learn

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# $QC_2D$ vs QCD

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## What we might learn

- ▶ **Gluodynamics**
- ▶ Checks on model studies
- ▶ Qualitative features of deconfinement
- ▶ Generic features of gauge theories at high densities
- ▶ **Medium effects on quarks, gluons and non-Goldstone hadrons**

## Issues of interest

### Gluodynamics — $SU(2)$ and $SU(3)$ very similar?

- ▶ Deconfinement at high density — effects on gluon propagator?
- ▶ Gap equation with effective or one-gluon interaction used to determine superconducting gap → more realistic input?
- ▶ Static magnetic gluon: unscreened at all orders in perturbation theory!

### Quark propagator

- ▶ Details of phase diagram depend critically on the effective quark mass in the medium.
- ▶ Dynamical quark masses → effective **strange** quark mass?
- ▶ Direct determination of diquark gap?

## Tensor structure in medium

The medium breaks Lorentz (Euclidean) symmetry to  $O(3)$

$\Rightarrow$  1  $\rightarrow$  2 scalar functions in gluon, 2  $\rightarrow$  4 in quark:

$$D_{\mu\nu}(\vec{q}, q_t) = P_{\mu\nu}^T D_M(\vec{q}^2, q_t^2) + P_{\mu\nu}^E D_E(\vec{q}^2, q_t^2) + \xi \frac{q_\mu q_\nu}{q^4}$$

$$S^{-1}(\vec{p}, \tilde{\omega}) = i\vec{p}A(\vec{p}^2, \tilde{\omega}^2) + i\gamma_4\tilde{\omega}C(\vec{p}^2, \tilde{\omega}^2) + B(\vec{p}^2, \tilde{\omega}^2) \\ + i\gamma_4\vec{p}D(\vec{p}^2, \tilde{\omega}^2)$$

$$S(\vec{p}, \tilde{\omega}) = i\vec{p}S_a + i\gamma_4\tilde{\omega}S_c + S_b + i\gamma_4\vec{p}S_d$$

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where  $\tilde{\omega} \equiv p_t - i\mu$ .

In general the form factors are **complex**!

## Gor'kov formalism

Quarks and antiquarks are in the same representation.

Construct Gor'kov spinor  $\Psi = \begin{pmatrix} \psi \\ \bar{\psi}^T \end{pmatrix}$

$$\implies \langle \Psi(x) \bar{\Psi}(y) \rangle \equiv \mathcal{G}(x, y) = \begin{pmatrix} S_N & -S_A \\ \bar{S}_A & \bar{S}_N \end{pmatrix}$$

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$S_A$  contains information about anomalous propagation

The corresponding self-energies are diquark gaps  $\Delta$   
(superfluid/superconducting)

## Lattice formulation

We use **Wilson fermions**:

- ▶ Correct symmetry breaking pattern, Goldstone spectrum
- ▶  $N_f < 4$  needed to guarantee continuum limit
- ▶ No problems with locality, fourth root trick
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$$S = \bar{\psi}_1 M(\mu) \psi_1 + \bar{\psi}_2 M(\mu) \psi_2 - J \bar{\psi}_1 (C \gamma_5) \tau_2 \bar{\psi}_2^T + \bar{J} \psi_2^T (C \gamma_5) \tau_2 \psi_1$$
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**Diquark source  $J$**  introduced to

- ▶ lift low-lying eigenmodes in the superfluid phase
- ▶ study diquark condensation without uncontrolled approximations

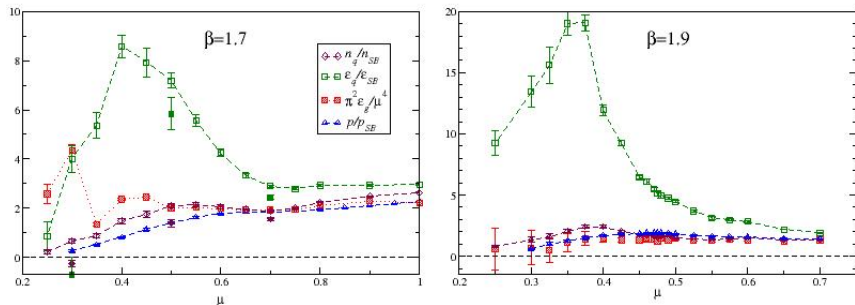
## Simulation Parameters

We work on two lattices, 'coarse' and 'fine':

Name	$\beta$	$\kappa$	Volume	$a$	$am_\pi$	$m_\pi/m_\rho$
coarse	1.7	0.178	$8^3 \times 16$	0.26fm	0.79	0.80
fine	1.9	0.168	$12^3 \times 24$	0.20fm	0.65	0.80

- ▶ Simulations performed with  $j = J/\kappa = 0.04$  for  $\mu = 0.3 - 1.0$
- ▶ 300–500 trajectories for each  $\mu$ .
- ▶ Simulations with  $j = 0.02, 0.06$  for  $\mu = 0.3, 0.5, 0.7, 0.9$   
→ enable extrapolation to  $j = 0$ .

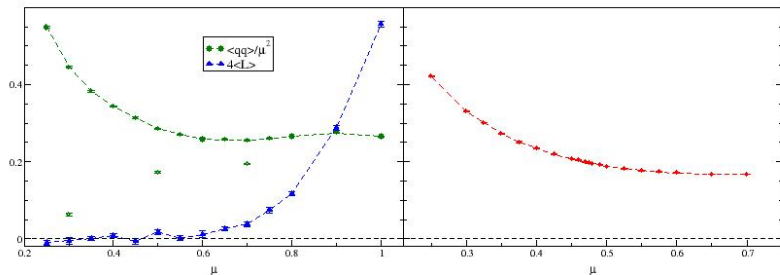
# Thermodynamics results



- ▶ Close to SB scaling for  $\mu > \mu_d$
- ▶  $\epsilon_q \sim 2\epsilon_{SB} \rightarrow k_F > E_F \implies$  binding energy?
- ▶ 30–40% of total energy from gluons!?

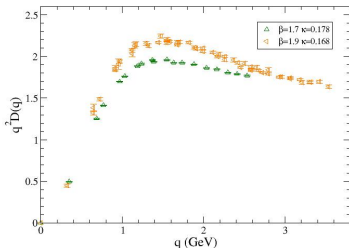


# Phase transitions



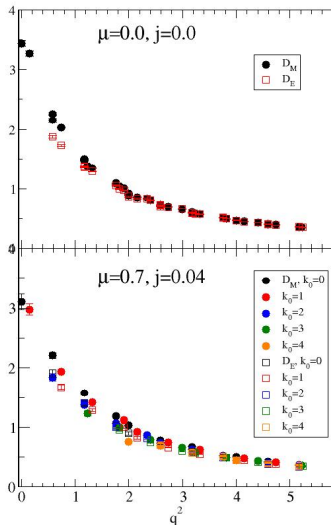
- ▶ Deconfining transition on coarse lattice  
— goes away on fine lattice?
- ▶ BEC  $\rightarrow$  BCS crossover becoming softer?

# Gluon propagator results

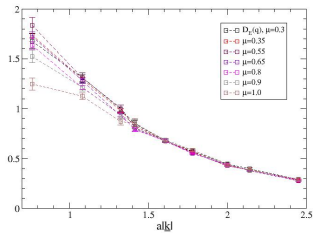
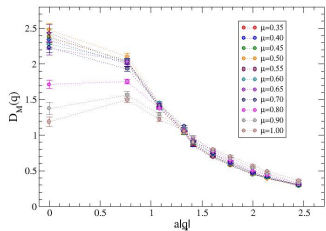
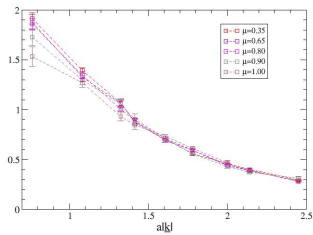
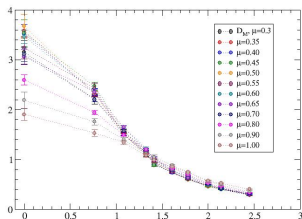


Some finite volume and lattice spacing effects at  $\mu = 0$

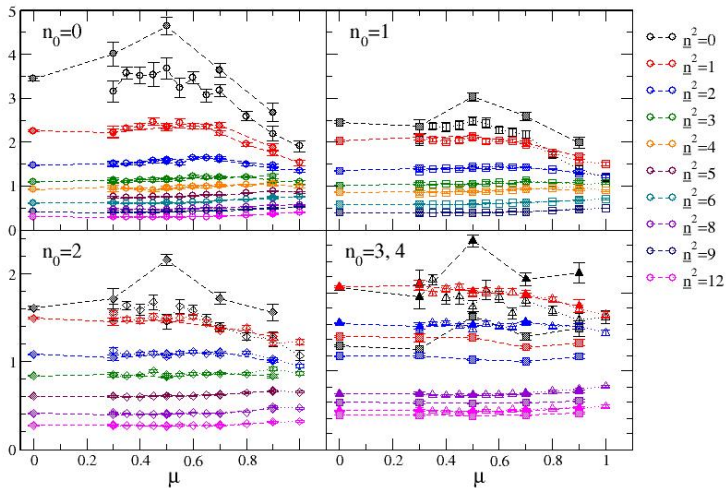
In-medium modifications, incl. violations of Lorentz symmetry, visible in magnetic gluon at  $\mu = 0.7$



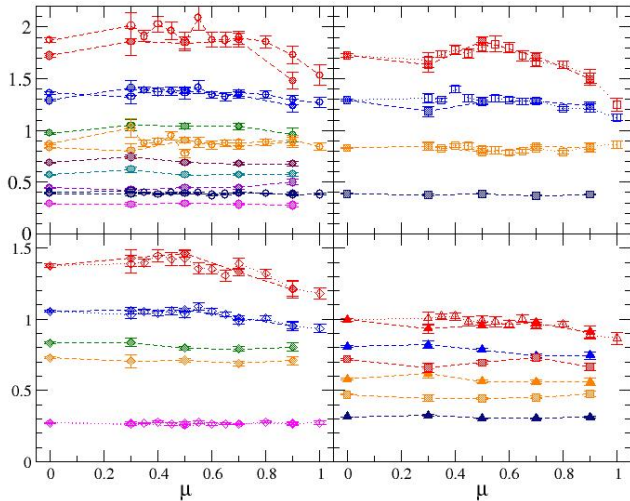
# Coarse lattice results



## Magnetic gluon (coarse lattice)

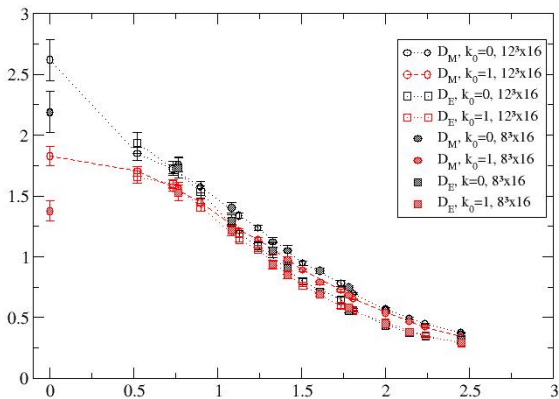


## Electric gluon (coarse lattice)

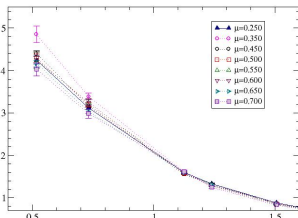
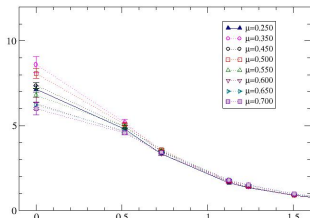
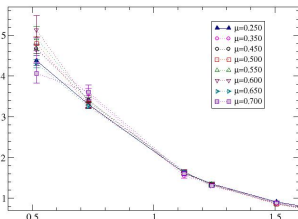
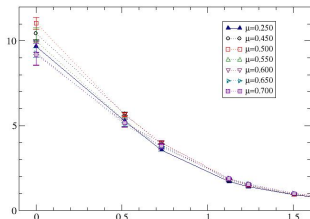


# Volume dependence

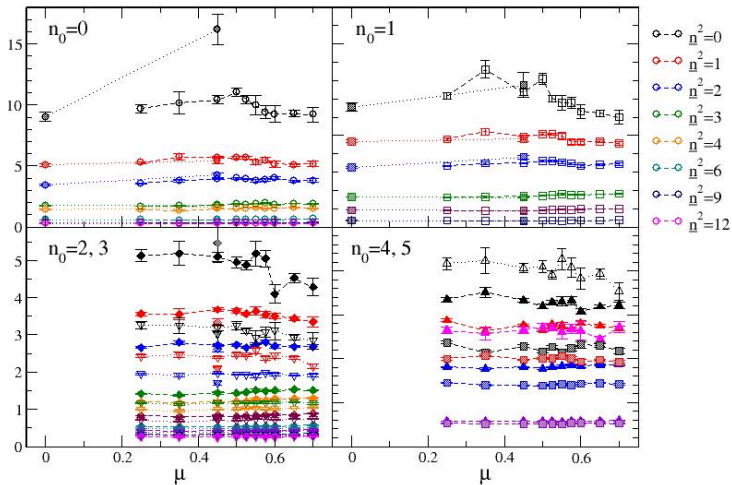
$[\mu = 0.9, j = 0.04]$



# Fine lattice results

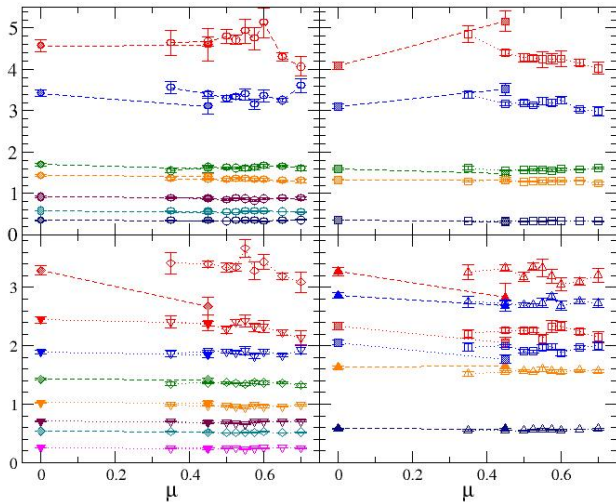


## Magnetic gluon (fine lattice)





## Electric gluon (fine lattice)



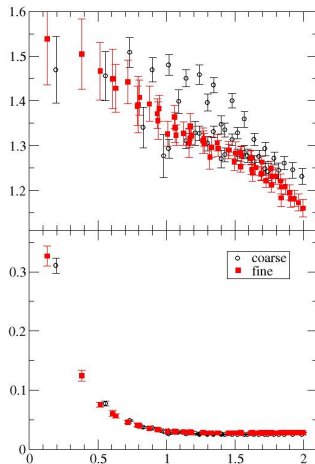
# Quark propagator results

Quark propagator in vacuum

Raw data!

Large lattice artefacts on coarse  
lattice

Unusual momentum behaviour?



## Tensor structure

Extracting form factors with the most general Ansatz for the tensor structure is complicated!

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- ▶ Computed **all** parts of the quark propagator at  $\mu = 0.5, j = 0.04$  on **coarse** lattice
  - ▶ 4 Dirac tensors
  - ▶ Normal and anomalous propagator
  - ▶ Real and imaginary part

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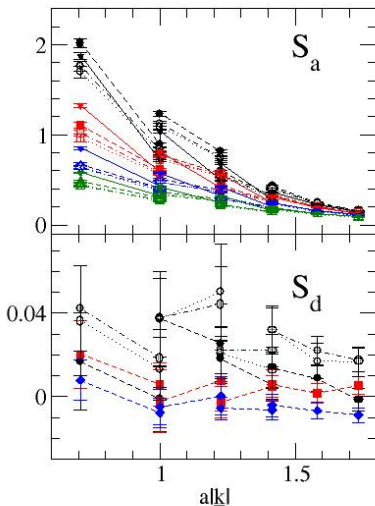
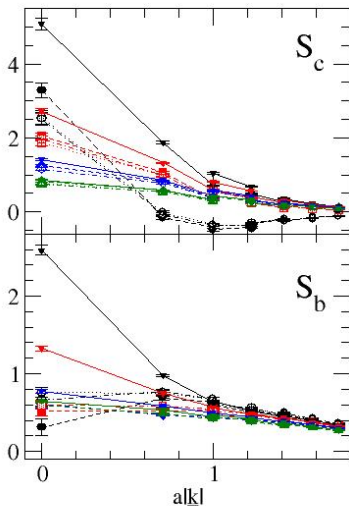
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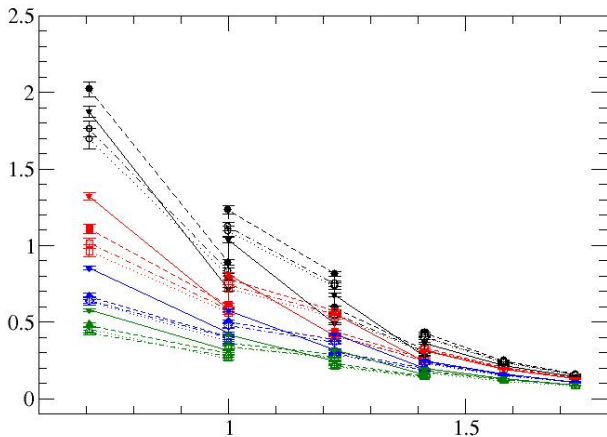
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- ▶ Anomalous parts are **complex**
- ▶ All are consistent with zero??

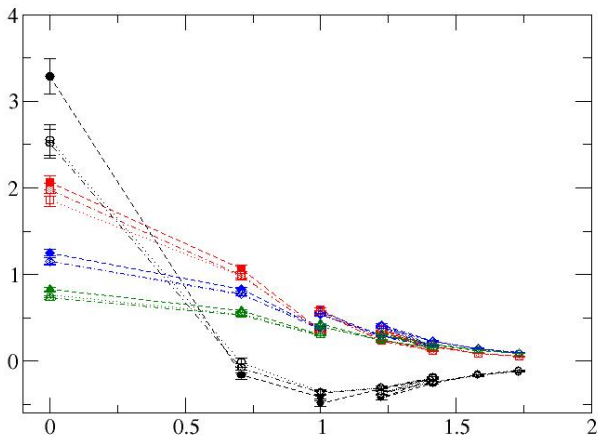
Quark propagator  $\mu = 0.5$  (Preliminary!)

## Quark propagator: spatial vector part

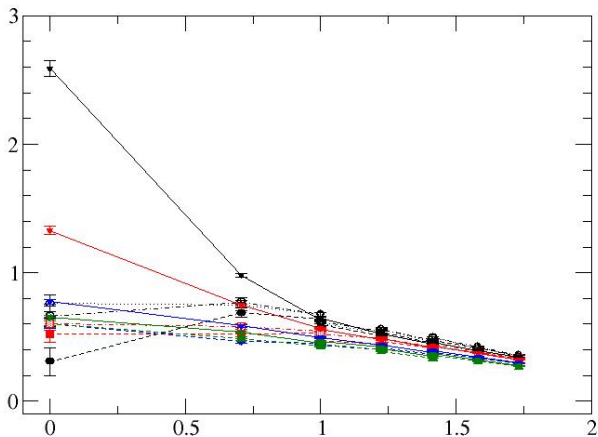




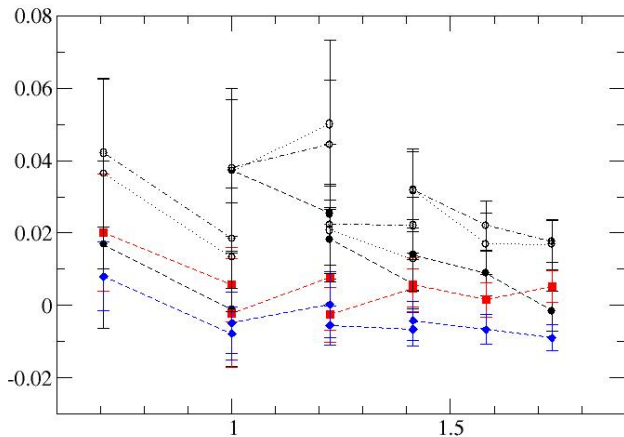
# Quark propagator: temporal vector part



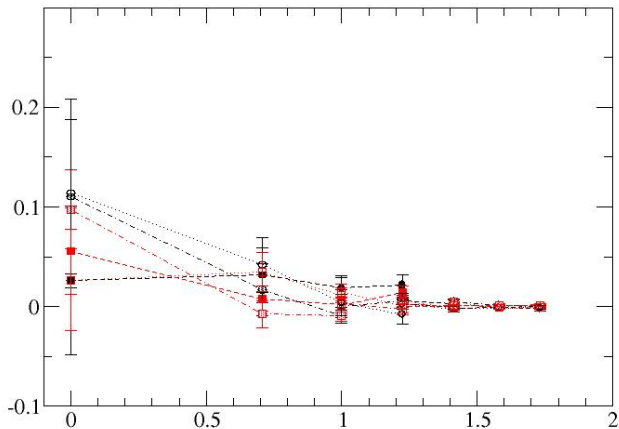
## Quark propagator: scalar part



# Quark propagator: tensor part



# Anomalous propagation



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- ▶ Need to understand **diquark source** dependence
- ▶ Need to understand **lattice artefacts**