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Elliptic Flow and Shear Viscosity of Color Glass Condensate

Mainly based on collaboration with:
F. Scardina, S. Plumari and V. Greco

St Goar, 2013 March 19

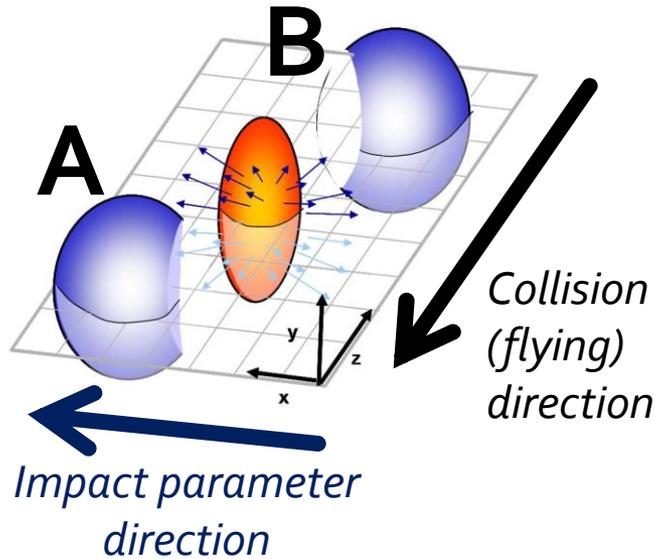


Plan of the talk

- Elliptic flow in heavy ion collisions (HICs)
- Initial conditions: Glauber and Color Glass Condensate (CGC)
- *Elliptic flow and viscosity of CGC*
- Conclusions and Outlook

Elliptic flow in HICs, 1

3D view of Heavy Ion Collision (HIC)

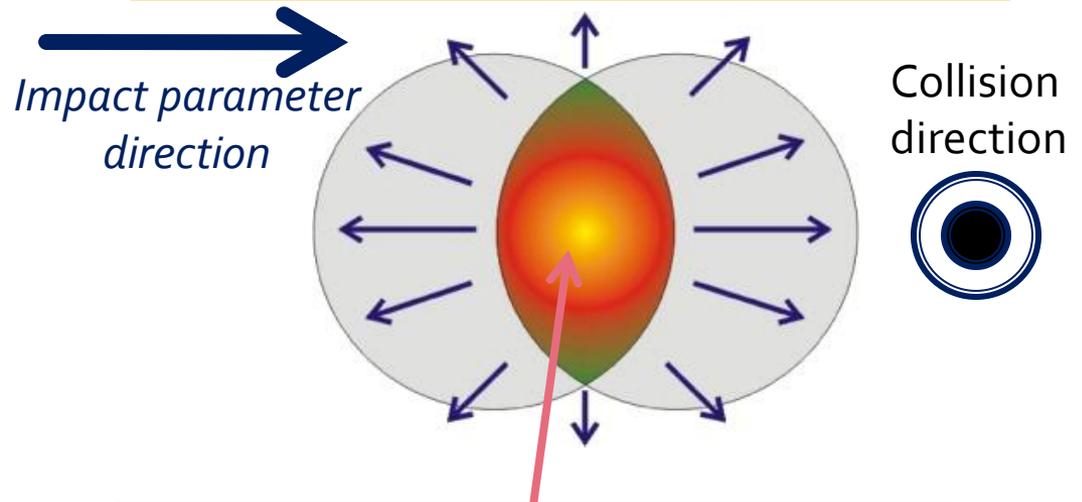


A, B: Cu, Au (RHIC)
Pb (LHC).

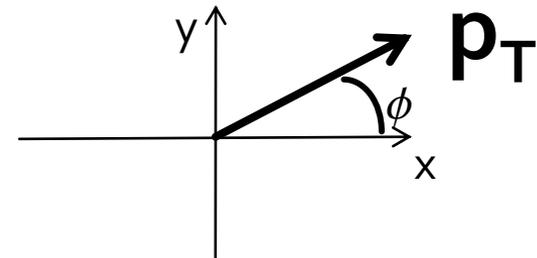
\sqrt{s} up to $200 \times A$ GeV, RHIC

\sqrt{s} up to $2.76 \times A$ TeV, LHC

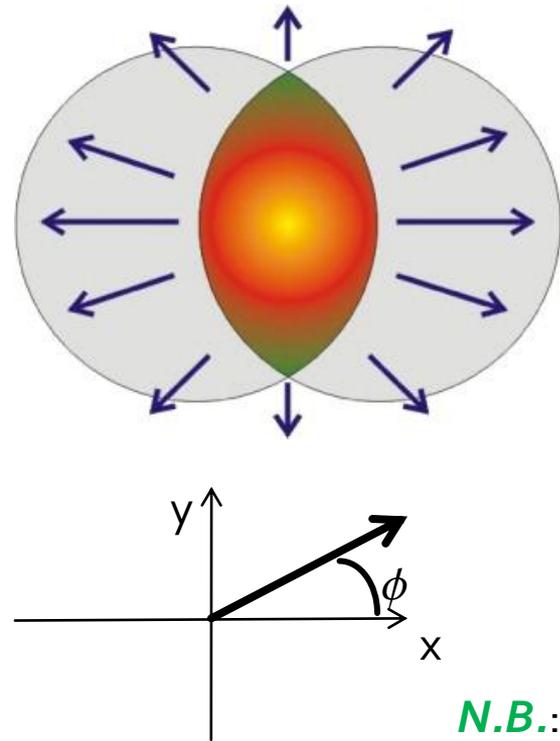
Projection of HIC on the transverse plane



FIREBALL:
Hot and dense expanding parton mixture
QUARK-GLUON-PLASMA (QGP)



Elliptic flow in HICs, 1



Particle multiplicity

$$\frac{d^3 N}{dy dp_T dp_T d\phi} = \frac{1}{2\pi} \frac{d^2 N}{dy dp_T dp_T} [1 + 2v_2(y, p_T) \cos 2\phi]$$

v_2 measures how efficiently the anisotropy of configuration in the initial state is transmitted to momenta in final states.

N.B.: STAR, PHENIX and ALICE collaborations report nonvanishing measured **odd flows**, which are however due to *initial state fluctuations* which we neglect, see e.g.

G. Eyyubova, Acta Phys. Pol. B Proceedings Supplement 5 (2012)

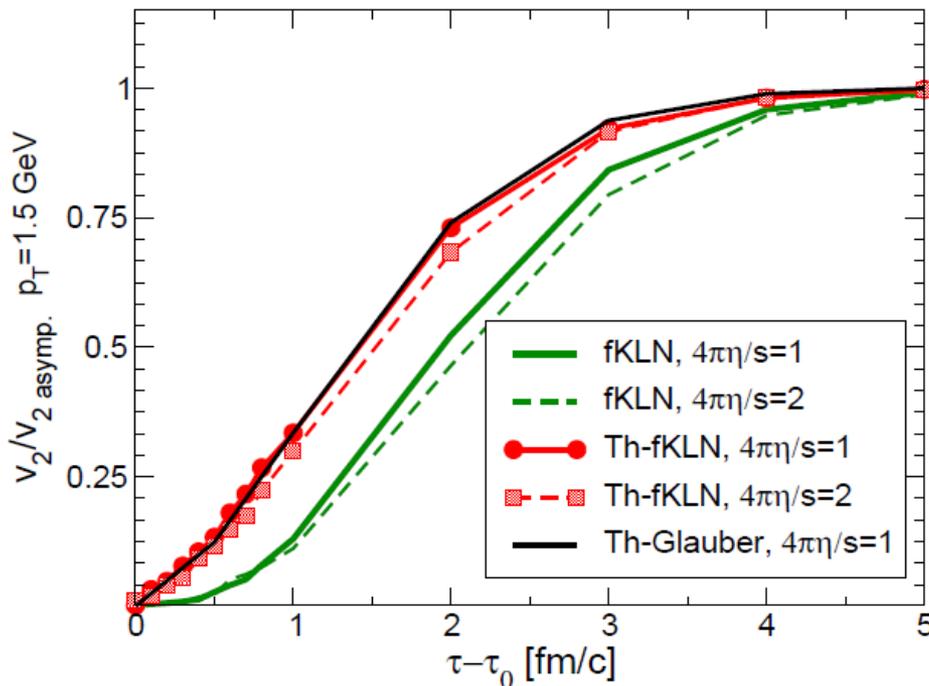
W. Cassing et al., 1202.3272 [nucl-th]

Thanks to Dr. Nan Su for pointing me out few details about this slide.

Elliptic flow in HICs, 1

Elliptic flow is mostly generated early in the nucleus-nucleus collision, and is present at the partonic level before partons hadronize.

Generation of elliptic flow (simulation)



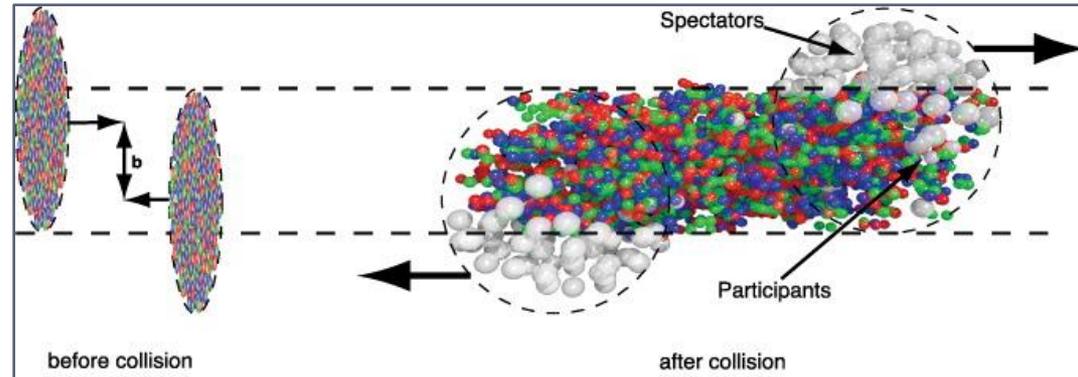
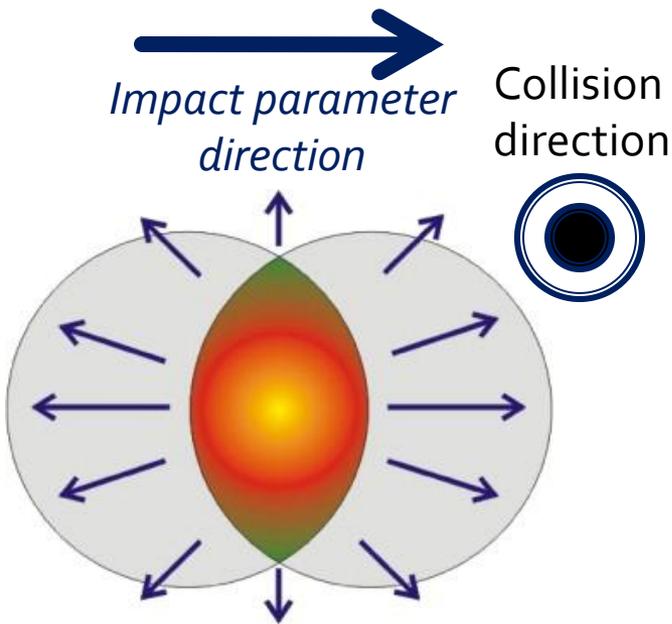
Elliptic flow acts as a **messenger**: it brings us the properties of the hot and dense state of partonic matter created after the collision.

Scenario in agreement with several independent computations, e.g.:

- Csernai *et al.*, PRL **97**, 152303 (2006)
- Molnar *et al.*, PRL **91**, 092301 (2003)
- Greco *et al.*, PRC **68**, 034904 (2003)
- Fries *et al.*, PRC **68**, 044902 (2003)
- Peschanski and Saridakis, PRC **80** (2009)

Initial condition: Glauber (briefly)

Geometrical description of the fireball:



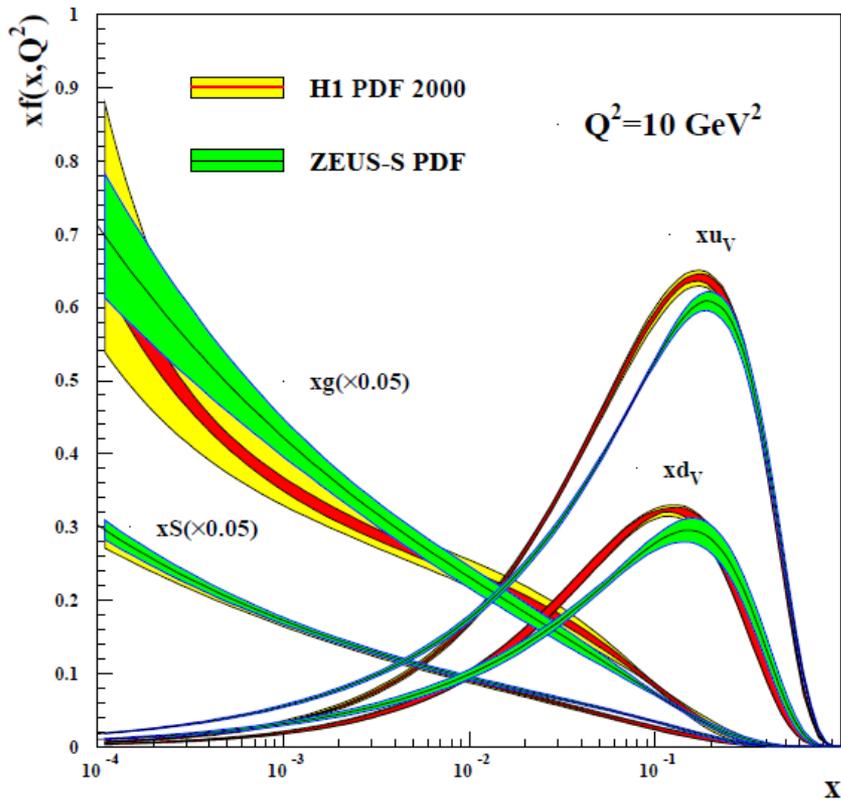
Assuming a nucleon distribution in the parents nuclei (typically a Woods-Saxon), one counts how many particles from each nucleus are present in the overlap region; among them, the **participants** are the nucleons that effectively can have an interaction (in fact, the particles that *are in the overlap region* but *do not interact*, are not taken into account).

For a review see:

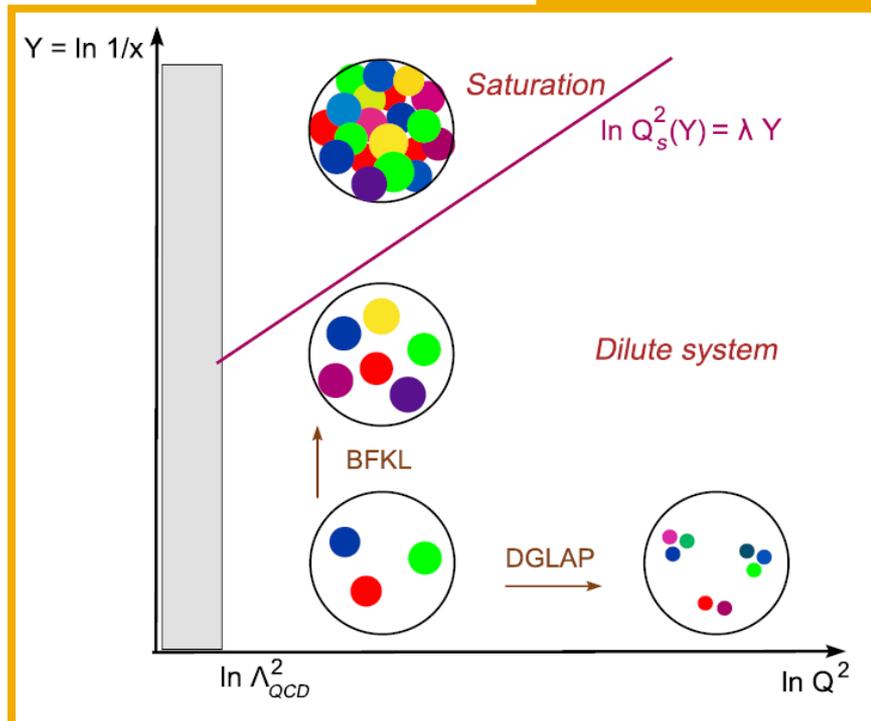
Miller *et al.*, Ann.Rev.Nucl.Part.Sci. **57**, 205 (2007)

Gluon saturation

Parton distribution functions

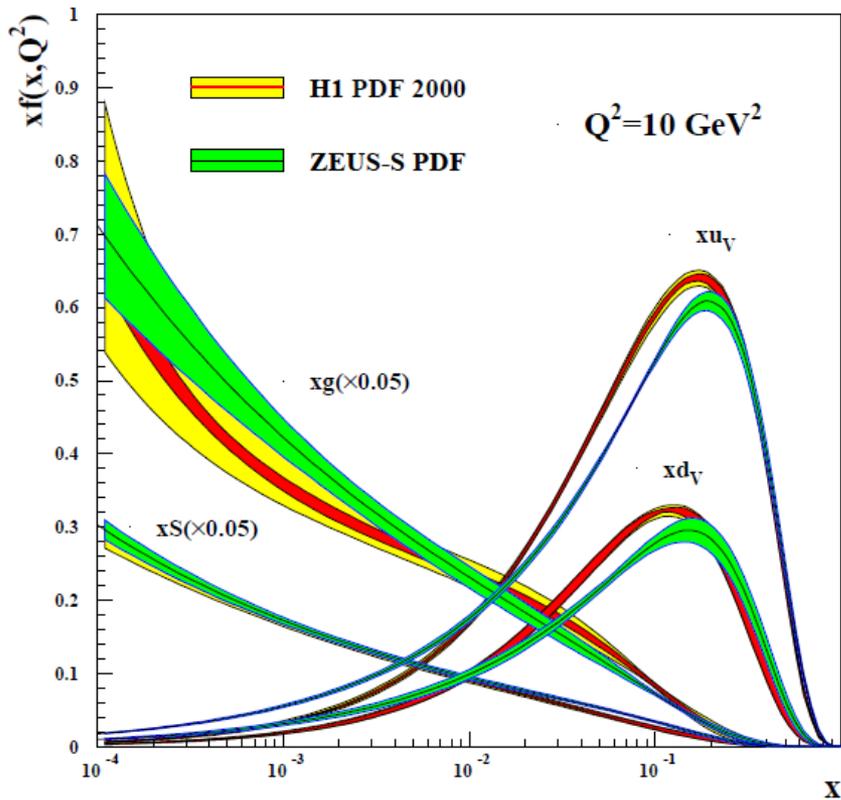


Phase Diagram



Gluon saturation

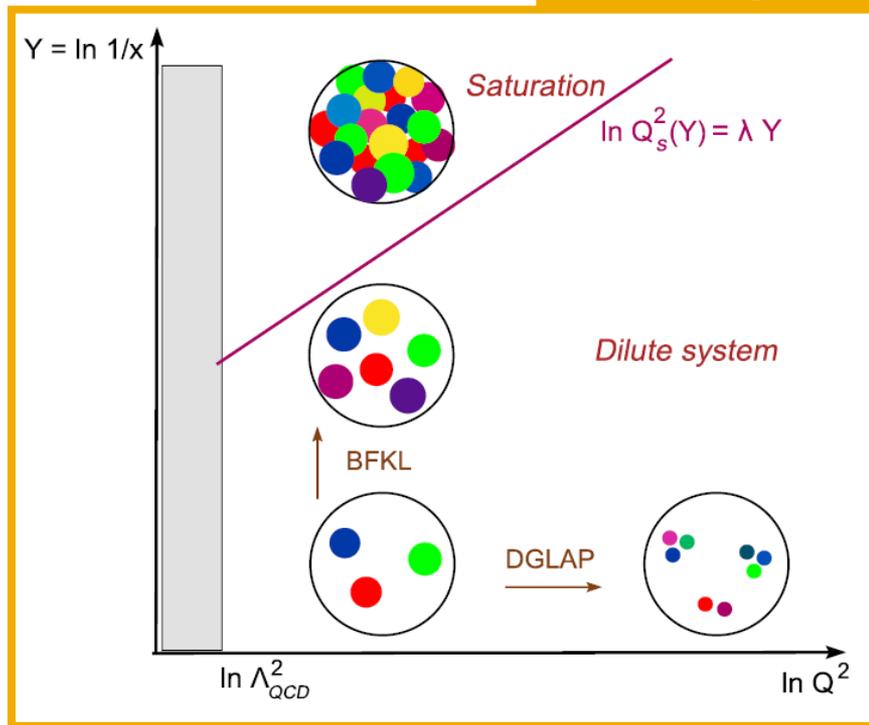
Parton distribution functions



Direction of decreasing transverse partons area ($1/Q^2$)

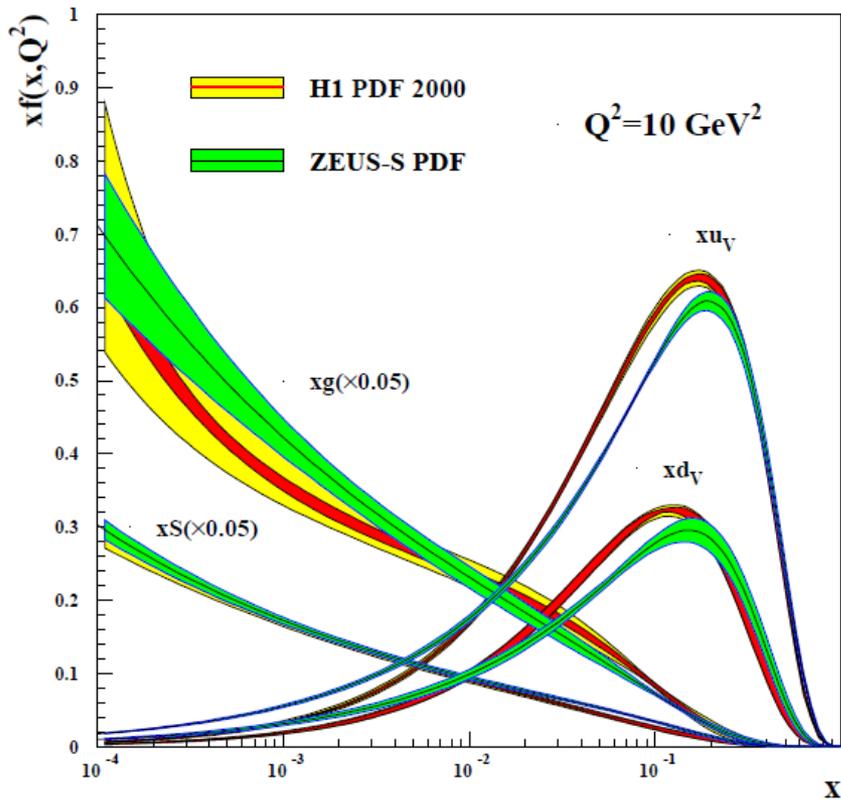


Phase Diagram



Gluon saturation

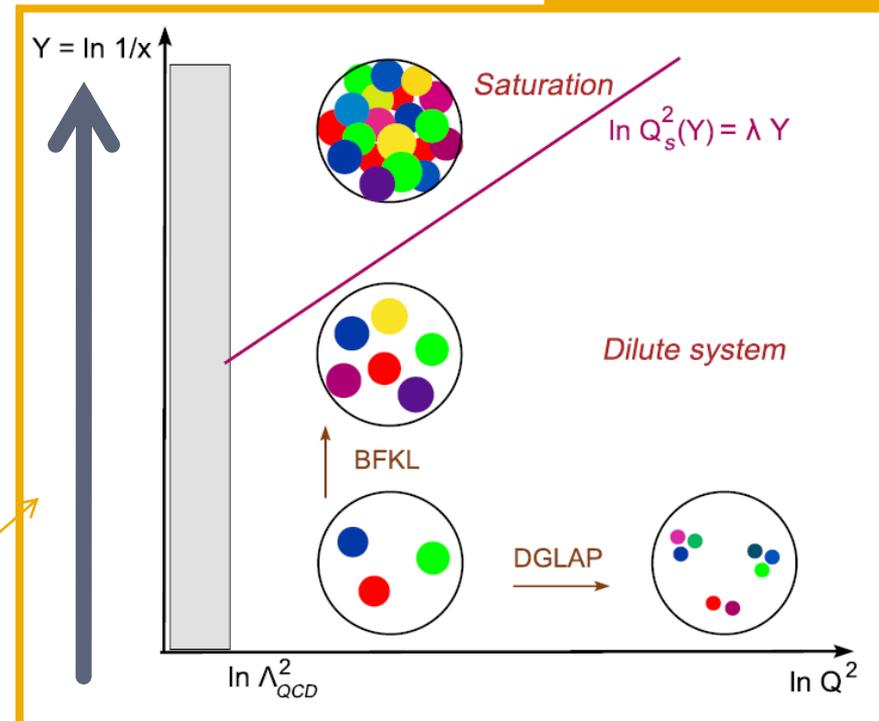
Parton distribution functions



Direction of decreasing transverse partons area ($1/Q^2$)



Phase Diagram

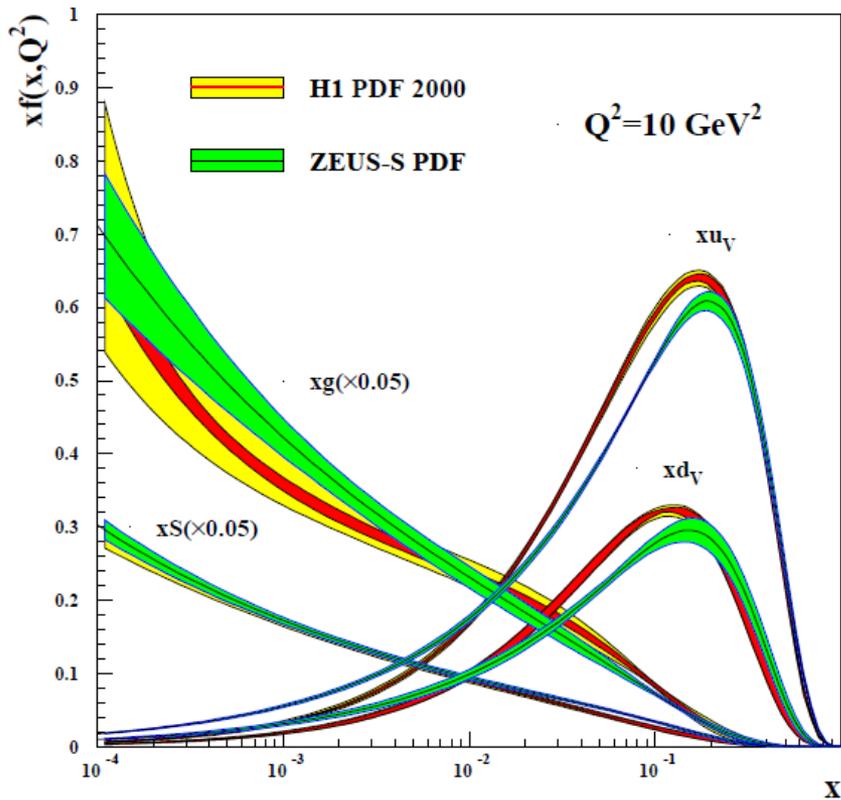


Direction of increasing gluon number



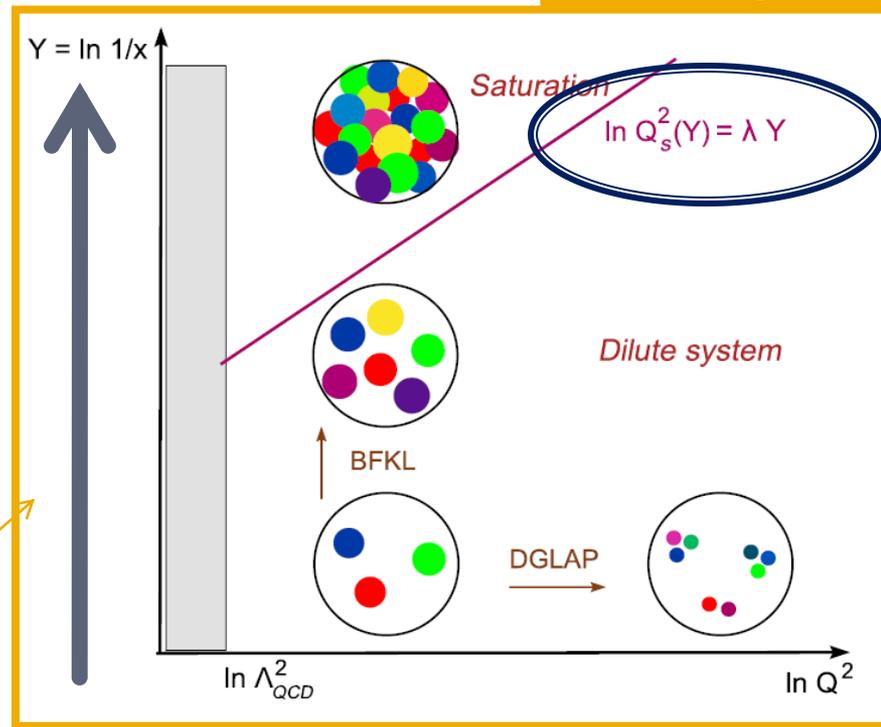
Gluon saturation

Parton distribution functions



Direction of decreasing transverse partons area ($1/Q^2$)

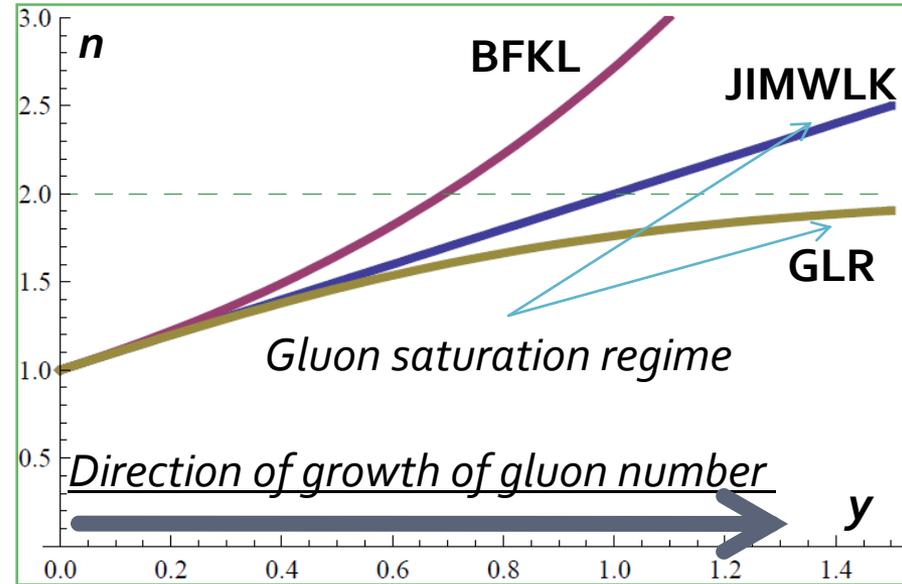
Phase Diagram



Direction of increasing gluon number

Gluon saturation

Small x (i.e. slow) gluon dynamics:
high density boson system,
living in the background of *large x (i.e. fast)*
color sources, the latter being *randomly distributed* in the nucleus.



- L. V. Gribov et al., Phys. Rept. **100** (1983) 1.
A. H. Mueller and J. Qiu, Nucl. Phys. **B268** (1986) 427.
J. Jalilian-Marian et al., Nucl. Phys. **B504** (1997) 415.
E. Iancu et al., Nucl. Phys. **A692** (2001) 583.
A. Leonidov et al., Nucl. Phys. **A703** (2002) 489.

Color: gluons carry color charge

Glass: fast partons randomly distributed in the hadron,
with frozen dynamics (by Lorentz time dilatation).

Condensate: small x gluons are correlated for $\Delta x_{\perp} \approx 1/Q_s$

Initial condition: CGC

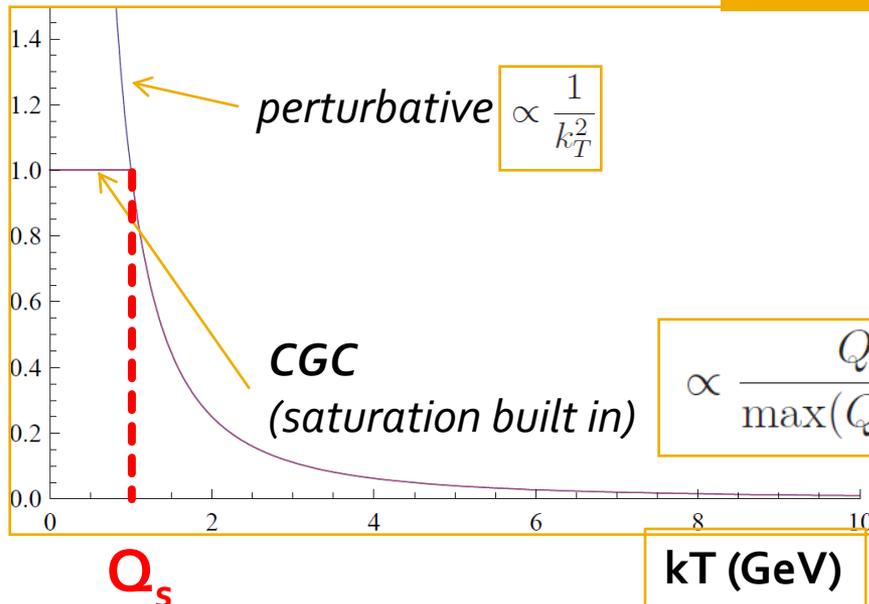
(f)KLN model

Factorization hypothesis:
convolution of distribution functions, assuming each one equal to the one the partons had in the parent nucleus.

Melting of the CGC leads to gluon production in the fireball:

$$\frac{dN_g}{d^2x_\perp dy} \propto \int \frac{d^2p_T}{p_T^2} \int_0^{p_T} d^2k_T \alpha_s(Q^2) \times \phi_A \left(x_A, \frac{(p_T + k_T)^2}{4}; \mathbf{x}_\perp \right) \times \phi_B \left(x_B, \frac{(p_T - k_T)^2}{4}; \mathbf{x}_\perp \right)$$

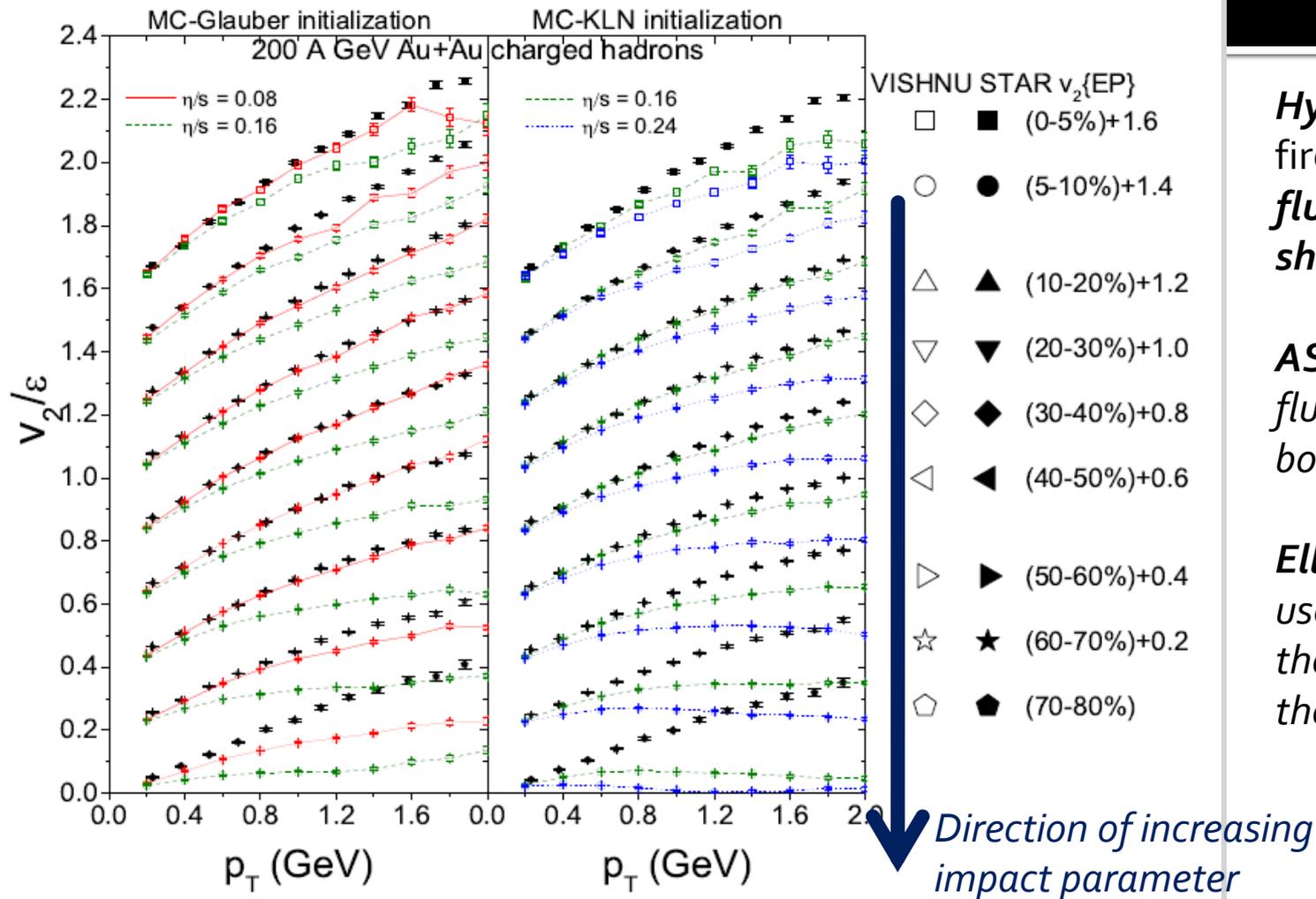
Unintegrated distribution functions (uGDFs)



Saturation scale **Q_s** depends on:
1.) **position in transverse plane;**
2.) **gluon rapidity.**

- Nardi *et al.*, Nucl. Phys. A**747**, 609 (2005)
- Kharzeev *et al.*, Phys. Lett. B**561**, 93 (2003)
- Nardi *et al.*, Phys. Lett. B**507**, 121 (2001)
- Drescher and Nara, PRC**75**, 034905 (2007)
- Hirano and Nara, PRC**79**, 064904 (2009)
- Hirano and Nara, Nucl. Phys. A**743**, 305 (2004)

Elliptic flow in HICs, 2

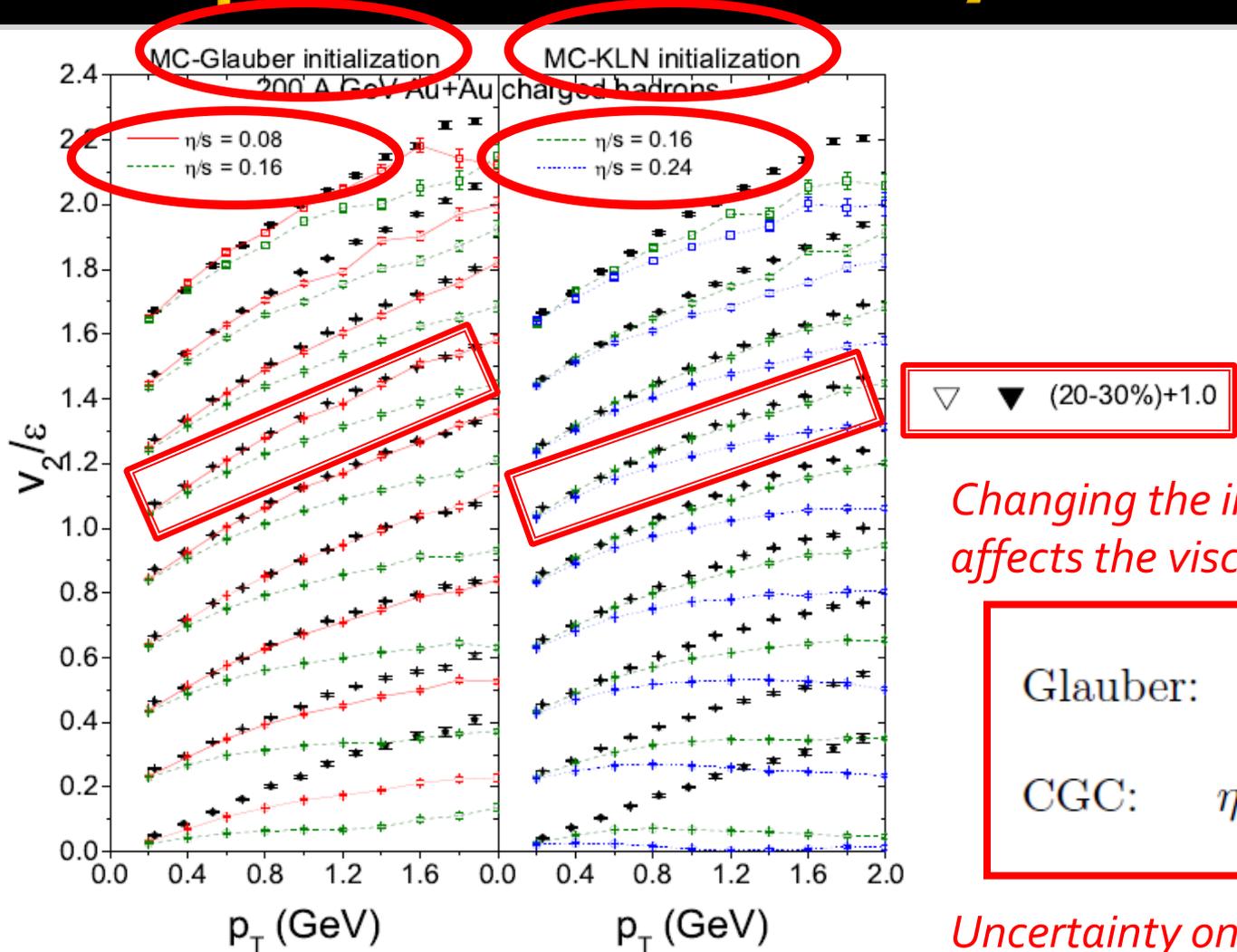


Hydro simulations:
fireball treated as a
fluid with a given
shear viscosity.

ASSUMPTION:
*fluid is thermalized in
both cases.*

**Elliptic flow data are
useful to estimate
the *shear viscosity* of
the QGP.**

Elliptic flow in HICs, 2



Hydro simulations:
fireball treated as a
fluid with a given
shear viscosity.

*Changing the initial condition
affects the viscosity of the fireball:*

$$\text{Glauber: } \eta/s \approx \frac{1}{4\pi}$$

$$\text{CGC: } \eta/s \approx \frac{2}{4\pi}$$

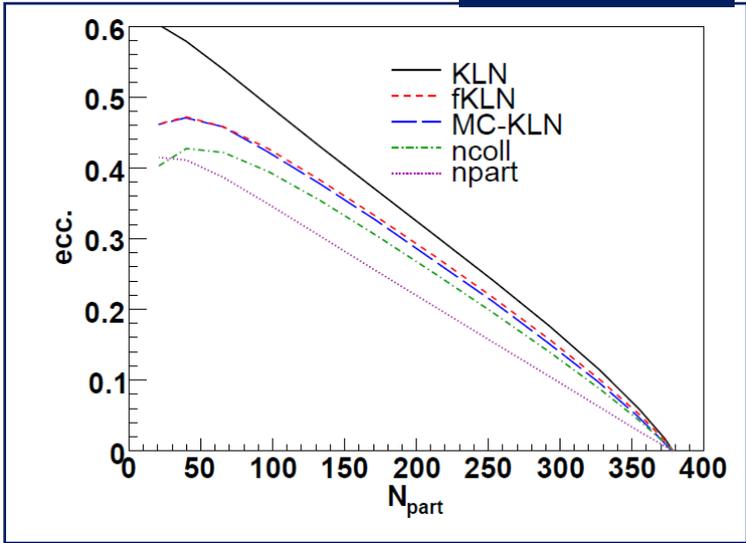
*Uncertainty on the initial condition
implies uncertainty on the ratio η/s
of the produced quark-gluon plasma.*

Elliptic flow in HICs, 2

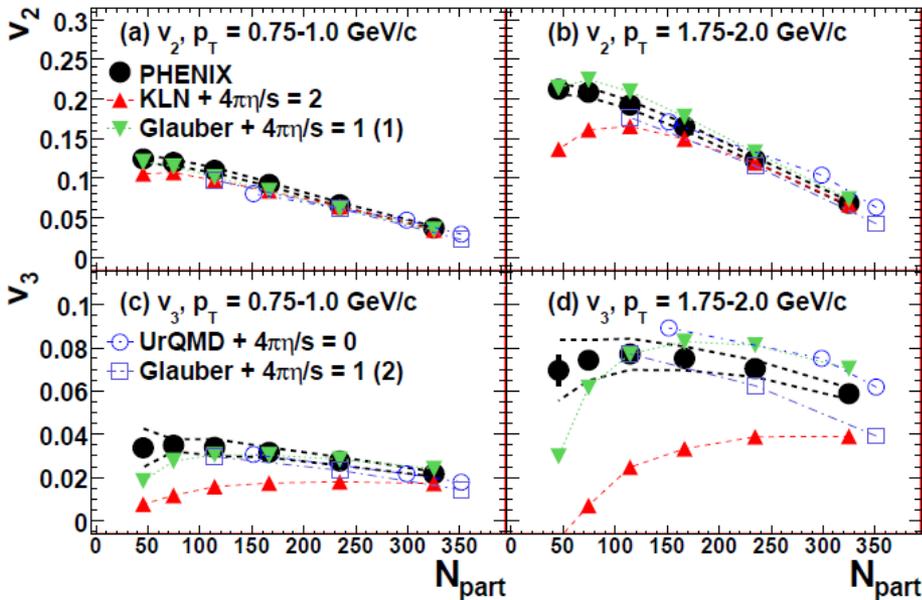
Glauber: $\eta/s \approx \frac{1}{4\pi}$

CGC: $\eta/s \approx \frac{2}{4\pi}$

Eccentricity



Adare *et al.*, [PHENIX Collaboration], PRL **107**, 252301 (2011).



The value of η/s affects the theoretical estimate of the higher harmonics.

Boltzmann equation

In order to simulate the temporal evolution of the fireball we solve the **Boltzmann equation** for the parton distribution functions:

$$p^\mu \partial_\mu f(x, p) = \mathcal{C}[f]$$

f : distribution function of gluons
 $\mathcal{C}[f]$: collision integral, computed stochastically

In $\mathcal{C}[f]$ we consider only the **elastic massless gluon scattering** process:

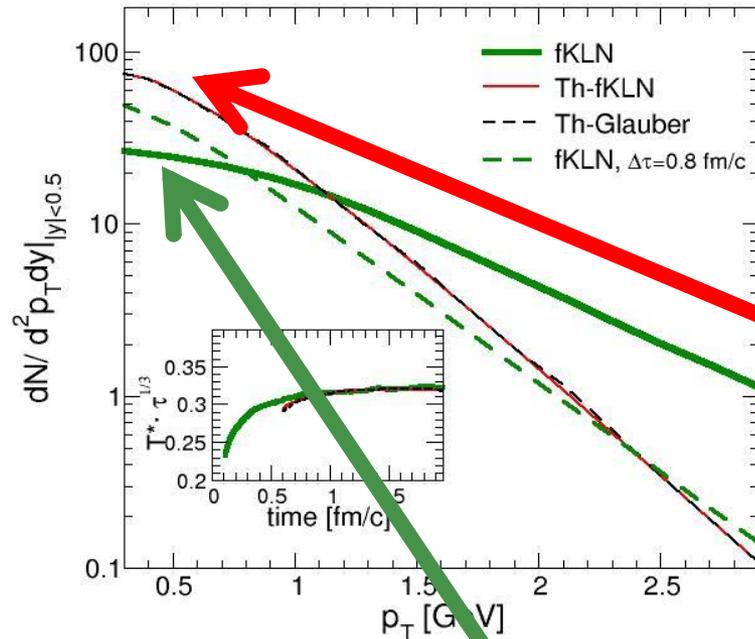
$$gg \rightarrow gg$$

Distribution function is sampled at each time step by means of **test particles** which evolve according to the classical Hamilton equations of motion:

$$\frac{d\mathbf{r}_i}{dt} = \frac{\mathbf{p}_i}{E_i}, \quad \frac{d\mathbf{p}_i}{dt} = \text{coll}$$

We use Boltzmann equation to simulate a fluid at **fixed** η/s ; the **cross section is computed in each phase space cell** to give the wished value of η/s .

Initial condition and thermalization

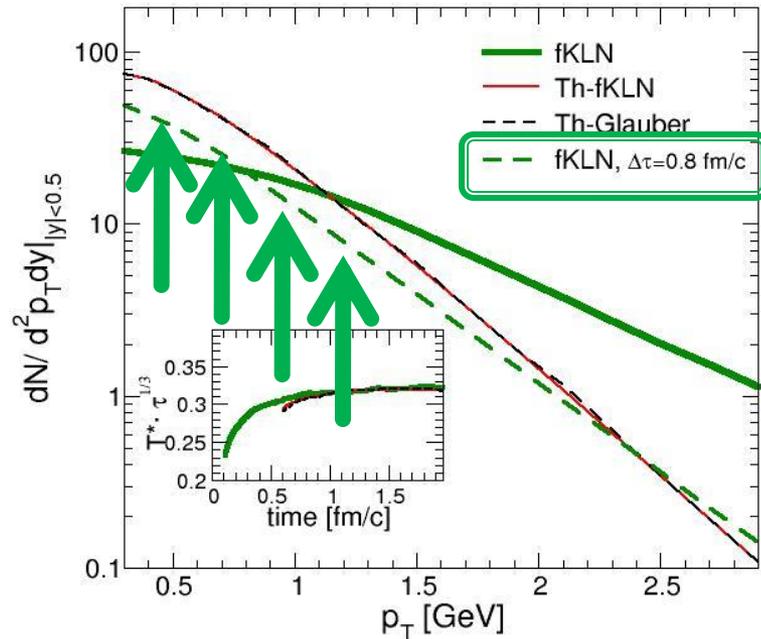


We use Transport to simulate a fluid at *fixed* η/s ; the *cross section* is *computed* in *each phase space cell* to give the wished value of η/s .

Initial conditions as in hydro simulations: thermalized distribution, $t_0=0.6$ fm/c, free streaming

CGC initial conditions, with saturation at low p_T built in, $t_0=0.1$ fm/c, no free streaming.

Initial condition and thermalization

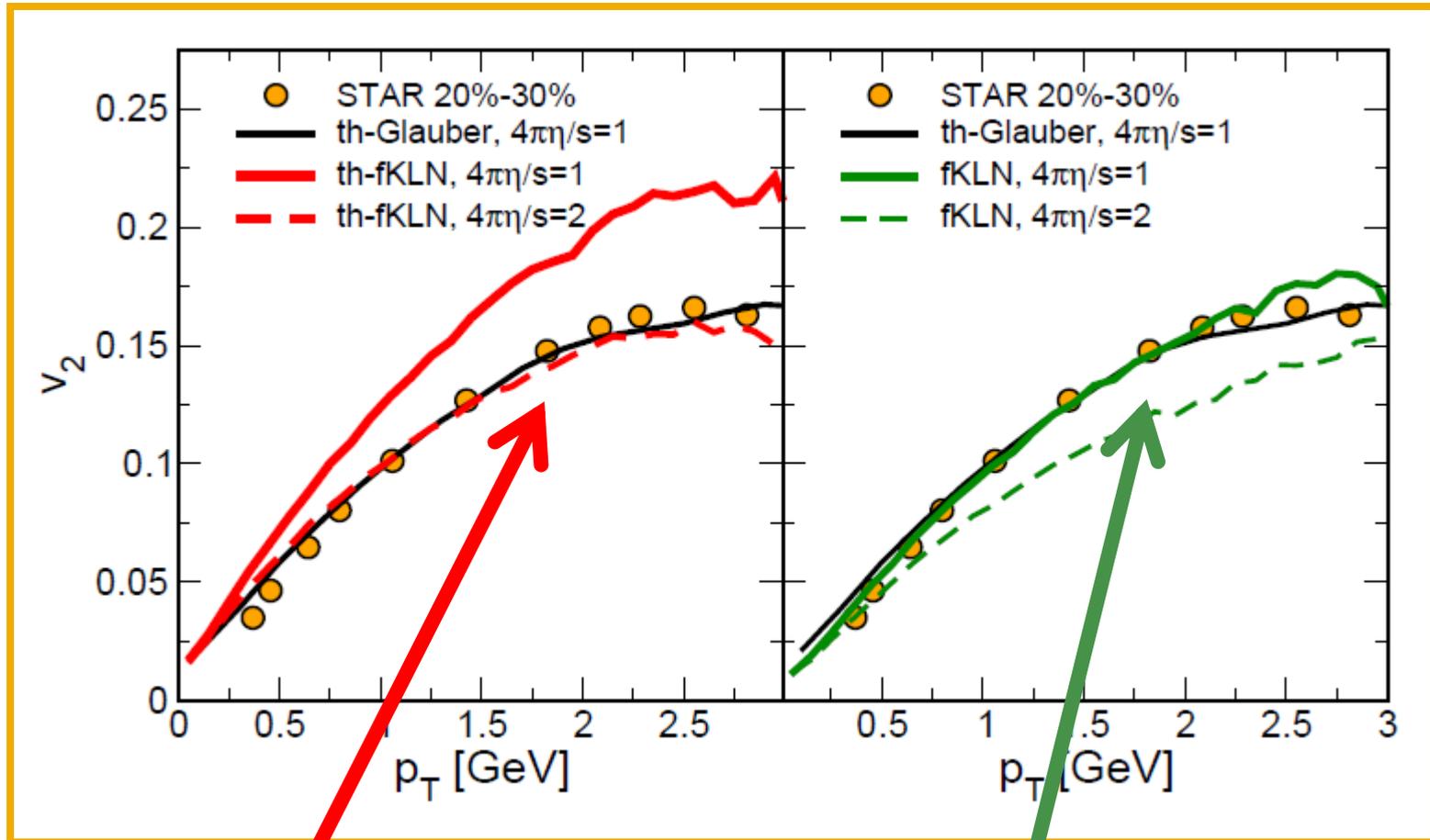


We use Transport to simulate a fluid at *fixed* η/s ; the *cross section* is *computed* in *each phase space cell* to give the wished value of η/s .

Thermalization in less than 1 fm/c,
in agreement with:
Greiner *et al.*, Nucl. Phys. A806, 287 (2008).



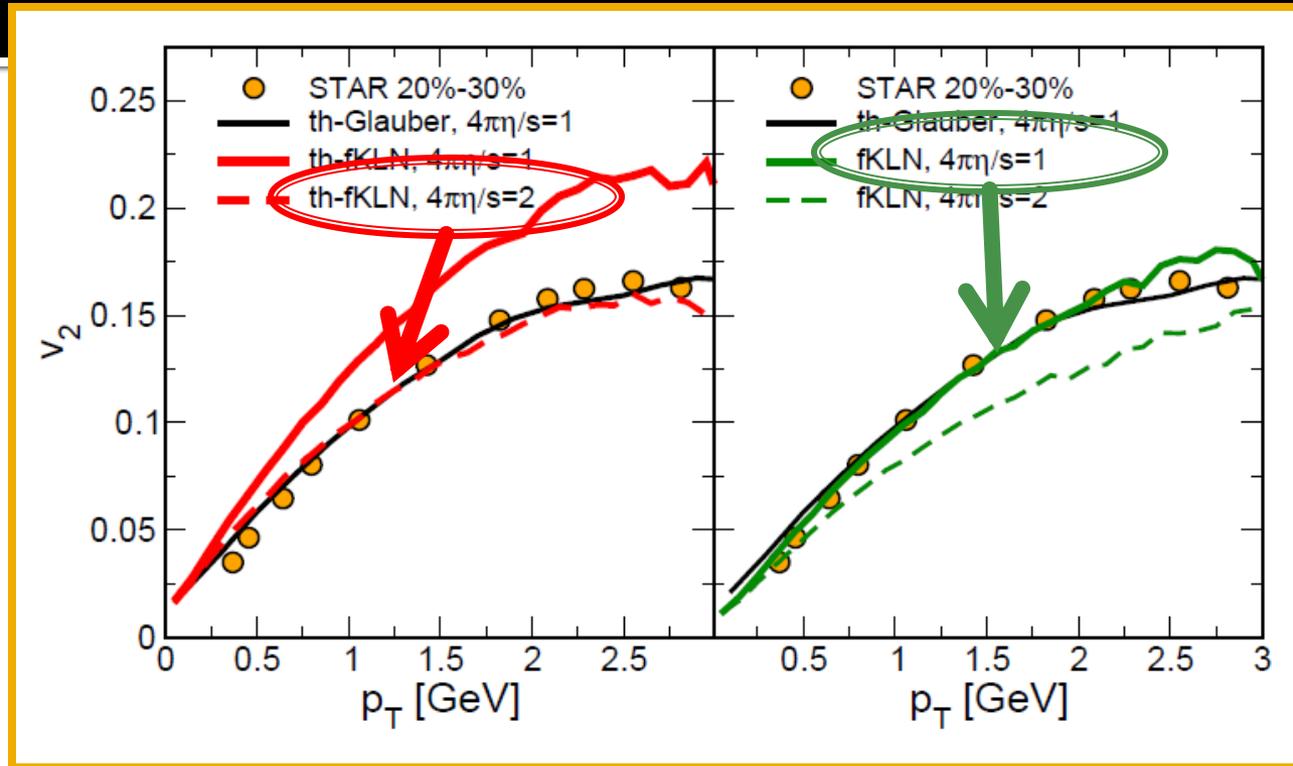
Elliptic flow from Transport



Hydro initial condition

CGC initial condition

Elliptic flow from Transport



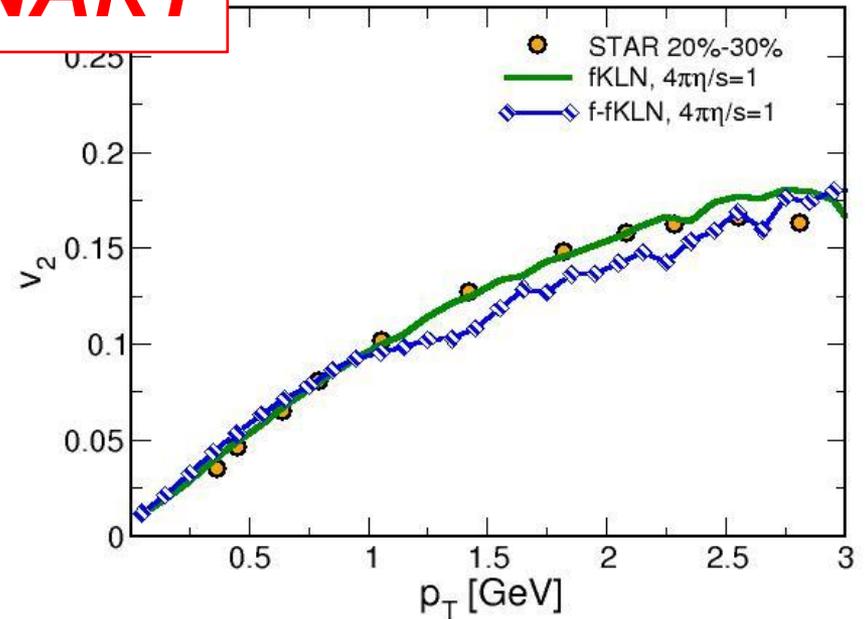
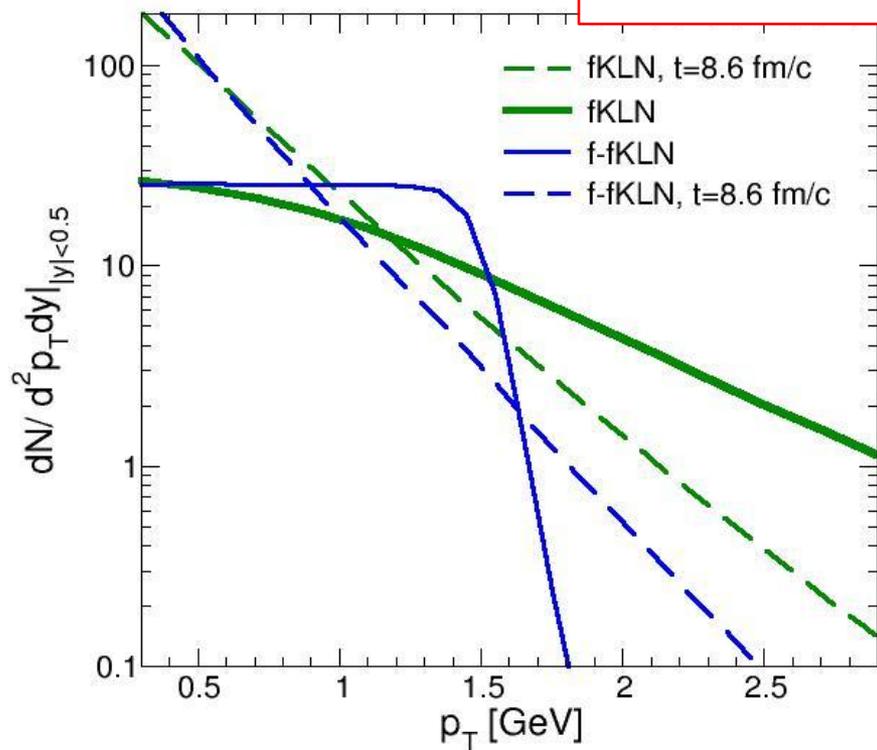
Implementing the proper initial condition in momentum space, which embeds saturation proper of a CGC, leads to the estimate of η/s in agreement with the Glauber initial condition.

Incidentally, η/s turns out to be of the order of the **universal lower limit** conjectured in: Kovtun *et al.* in PRL **94**, 111691 (2005).

Hence, the expanding fireball behaves as an **almost perfect fluid**.

Elliptic flow from Transport

PRELIMINARY



Taking the initial spectrum artificially flat leads to a further reduction of the elliptic flow.

Conclusions and Outlook

- Transport approach allows to study collective flow of CGC.
- Proper initial distribution in momentum space, which embeds saturation, has an important effect on the formation of the CGC v_2 .

Outlook

- Inclusion of inelastic processes and of thermal masses.
- *Detailed study of pre-equilibration times.*
- *Computation of higher order harmonics.*
- Extension to the LHC energies.



**THANK
YOU
FOR
YOUR
ATTENTION**

There are two breeds of fools: those who do not doubt anything, those who doubt everything.

(Charles-Joseph de Ligne)