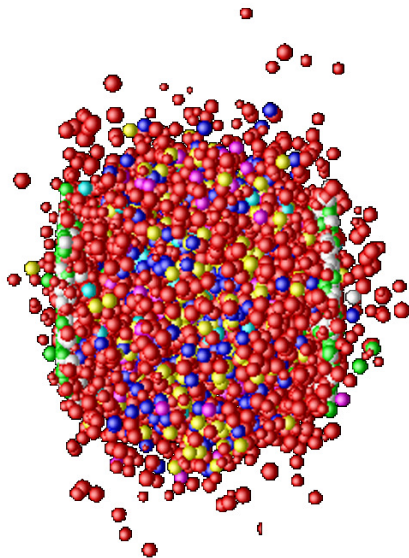


Institut für
Theoretische Physik I



Collective and electromagnetic observables from relativistic HIC and their relation to IQCD

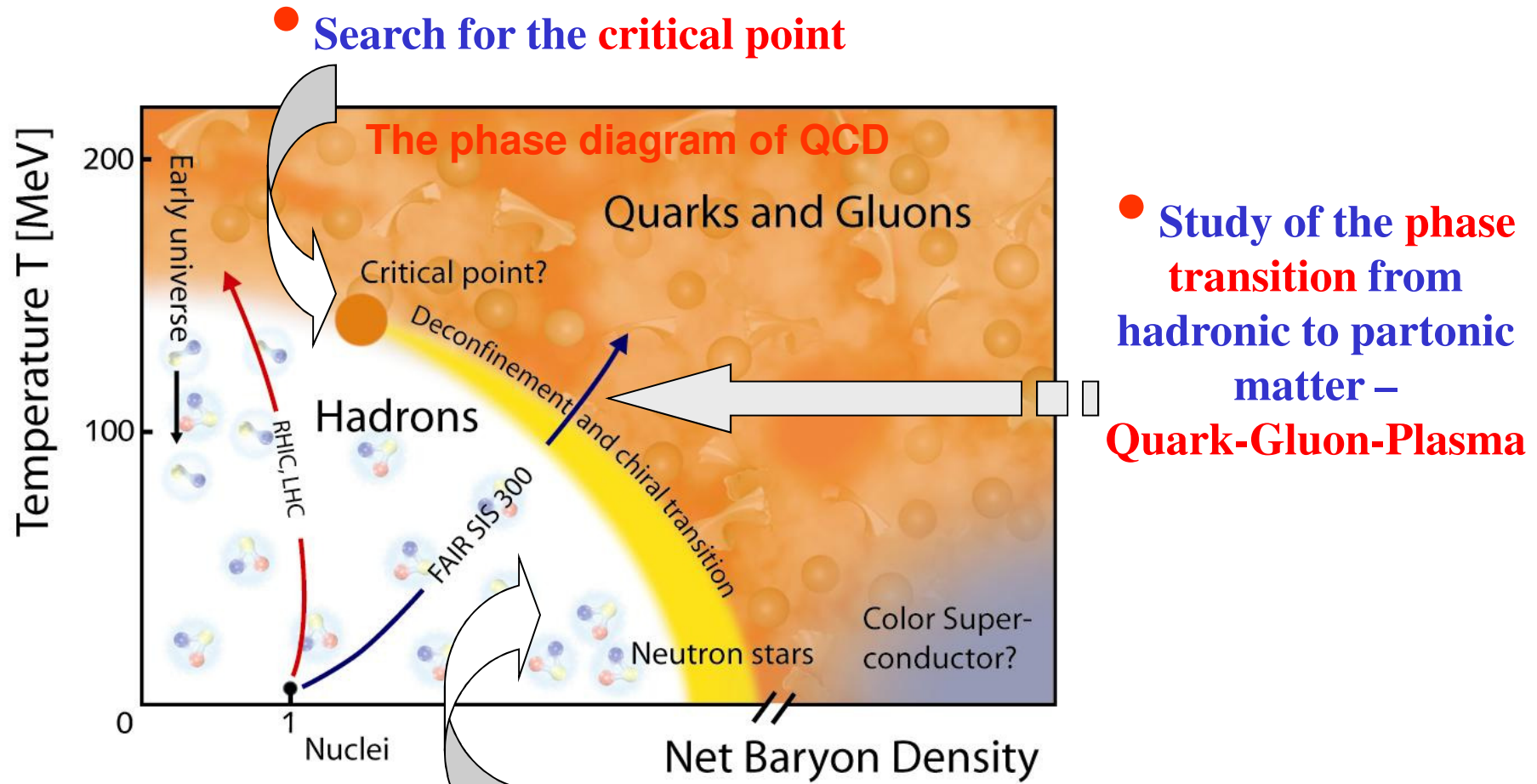


Wolfgang Cassing

St. Goar, 21.03.2013

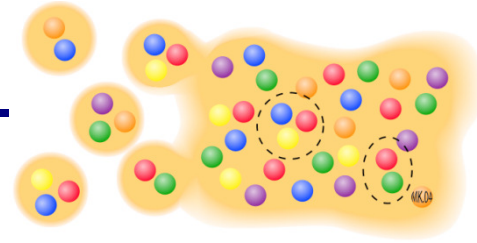


The holy grail of HIC



- Study of the **in-medium** properties of hadrons at high baryon density and temperature
- Study of the partonic medium beyond the phase boundary

From hadrons to partons



In order to study the **phase transition** from hadronic to partonic matter – **Quark-Gluon-Plasma** –

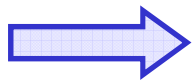
we need a **consistent non-equilibrium (transport) model with**

➤ **explicit parton-parton interactions** (i.e. between quarks and gluons) beyond strings!

➤ **explicit phase transition** from hadronic to partonic degrees of freedom

➤ **IQCD EoS** for partonic phase

Transport theory: off-shell Kadanoff-Baym equations for the Green-functions $S_h^<(x,p)$ in phase-space representation for the **partonic and hadronic phase**



Parton-Hadron-String-Dynamics (PHSD)



W. Cassing, E. Bratkovskaya, PRC 78 (2008) 034919;
NPA831 (2009) 215;
W. Cassing, EPJ ST 168 (2009) 3

QGP phase described by

Dynamical QuasiParticle Model (DQPM)

A. Peshier, W. Cassing, PRL 94 (2005) 172301;
Cassing, NPA 791 (2007) 365; NPA 793 (2007)

The Dynamical QuasiParticle Model (DQPM)

Basic idea: Interacting quasi-particles

- massive quarks and gluons (g, q, q_{bar}) with spectral functions :

$$\rho_i(\omega, T) = \frac{4\omega\Gamma_i(T)}{(\omega^2 - \vec{p}^2 - M_i^2(T))^2 + 4\omega^2\Gamma_i^2(T)} \quad (i = q, \bar{q}, g)$$

■ quarks

mass: $M_{q(\bar{q})}^2(T) = \frac{N_c^2 - 1}{8N_c} g^2 \left(T^2 + \frac{\mu_q^2}{\pi^2} \right)$

width: $\Gamma_{q(\bar{q})}(T) = \frac{1}{3} \frac{N_c^2 - 1}{2N_c} \frac{g^2 T}{8\pi} \ln\left(\frac{2c}{g^2} + 1\right)$

■ gluons:

$$M_g^2(T) = \frac{g^2}{6} \left(\left(N_c + \frac{N_f}{2} \right) T^2 + \frac{N_c}{2} \sum_q \frac{\mu_q^2}{\pi^2} \right)$$

$$\Gamma_g(T) = \frac{1}{3} N_c \frac{g^2 T}{8\pi} \ln\left(\frac{2c}{g^2} + 1\right)$$

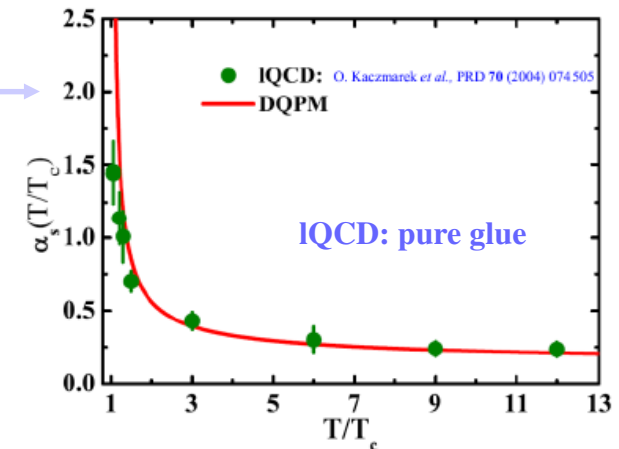
$$N_c = 3, N_f = 3$$

■ running coupling (pure glue):

$$\alpha_s(T) = \frac{g^2(T)}{4\pi} = \frac{12\pi}{(11N_c - 2N_f) \ln[\lambda^2(T/T_c - T_s/T_c)^2]}$$

□ fit to lattice (IQCD) results (e.g. entropy density)

with 3 parameters: $T_s/T_c=0.46$; $c=28.8$; $\lambda=2.42$
(for pure glue $N_f=0$)



➔ quasiparticle properties (mass, width)

DQPM: Peshier, Cassing, PRL 94 (2005) 172301;
Cassing, NPA 791 (2007) 365; NPA 793 (2007)

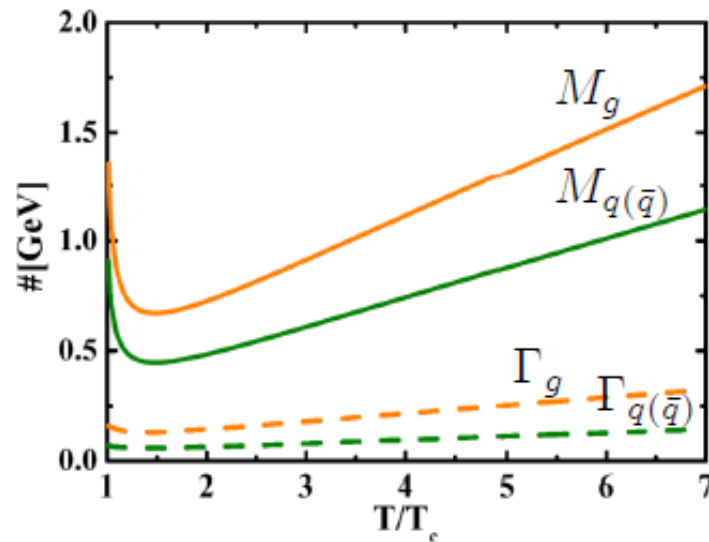
The Dynamical QuasiParticle Model (DQPM)

➤ **fit to lattice (IQCD) results** (e.g. entropy density)

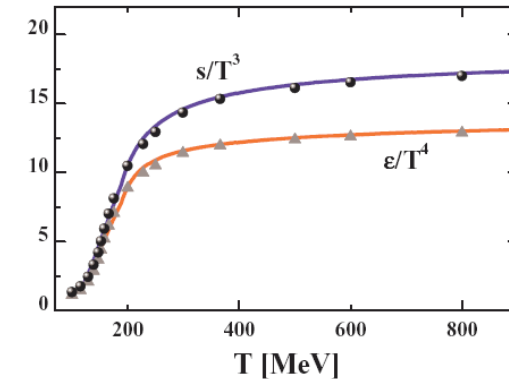
* fit to BMW IQCD data S. Borsanyi et al., JHEP 1009 (2010) 073

➔ **Quasiparticle properties:**

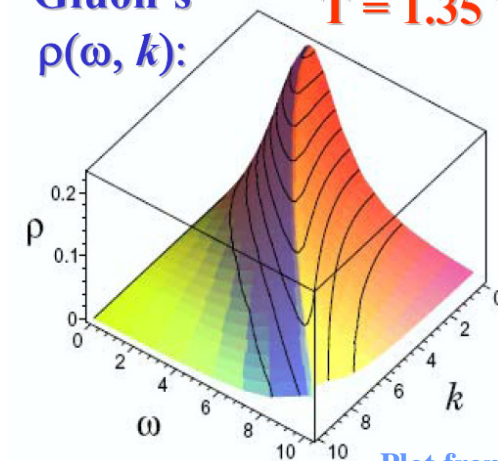
■ **large width and mass for gluons and quarks**



$T_C=158$ MeV
 $\epsilon_C=0.5$ GeV/fm³

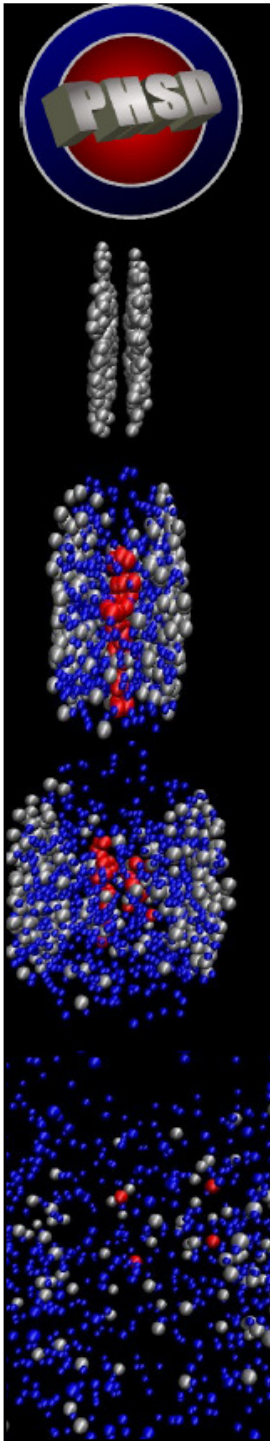


Gluon's
 $\rho(\omega, k):$ $T = 1.35 T_c$



Plot from Peshier,
 PRD 70 (2004)
 034016

- **DQPM matches well lattice QCD**
- **DQPM provides mean-fields (1PI) for gluons and quarks as well as effective 2-body interactions (2PI)**
- **DQPM gives transition rates for the formation of hadrons → PHSD**



PHSD - basic concept

Initial A+A collisions – HSD: string formation and decay to pre-hadrons

Fragmentation of pre-hadrons into quarks :

using the quark spectral functions from the **Dynamical QuasiParticle Model (DQPM)** - approximation to QCD

Partonic phase ($\epsilon > \epsilon_C$): quarks and gluons (= ,dynamical quasiparticles‘) with **off-shell spectral functions** (width, mass) defined by the DQPM

elastic and inelastic parton-parton interactions:

using the effective cross sections from the DQPM

- ✓ **q+qbar (flavor neutral) \Leftrightarrow gluon (colored)**
- ✓ **gluon + gluon \Leftrightarrow gluon** (possible due to large spectral width)
- ✓ **q+qbar (color neutral) \Leftrightarrow hadron resonances**

self-generated mean-field potential for quarks and gluons !

Hadronization: based on DQPM - massive, off-shell quarks and gluons with broad spectral functions hadronize to **off-shell mesons and baryons:**

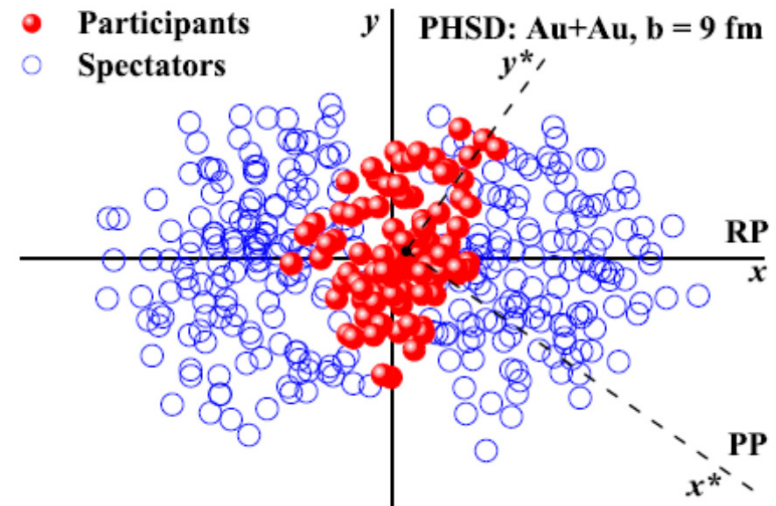
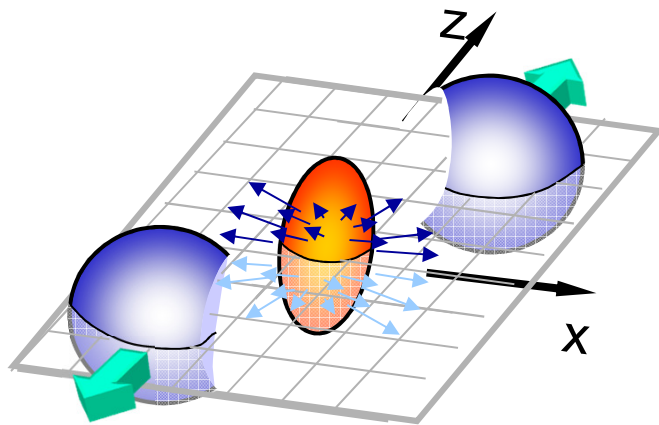
gluons \rightarrow q+qbar; q+qbar \rightarrow meson (or string);

q+q+q \rightarrow baryon (or string) (strings act as ,doorway states‘ for hadrons)

Hadronic phase: hadron-string interactions – **off-shell HSD**

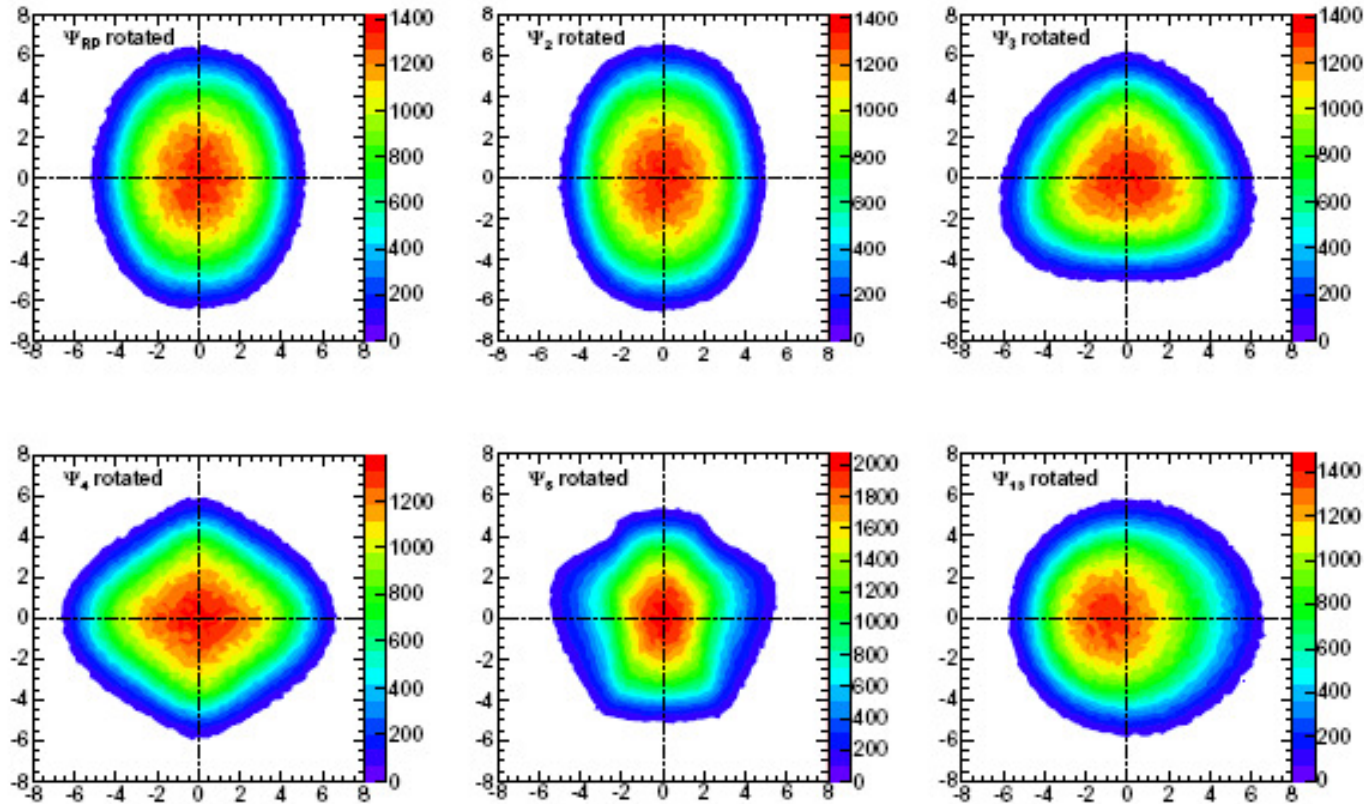
W. Cassing, E. Bratkovskaya, PRC 78 (2008) 034919;
NPA831 (2009) 215; EPJ ST 168 (2009) 3; NPA856 (2011) 162.

Collective flow: anisotropy coefficients (v_1, v_2, v_3, v_4) in A+A



Final angular distributions of hadrons

10k Au+Au collision events at b = 8 fm rotated to different event planes:



$$E \frac{d^3 N}{d^3 p} =$$

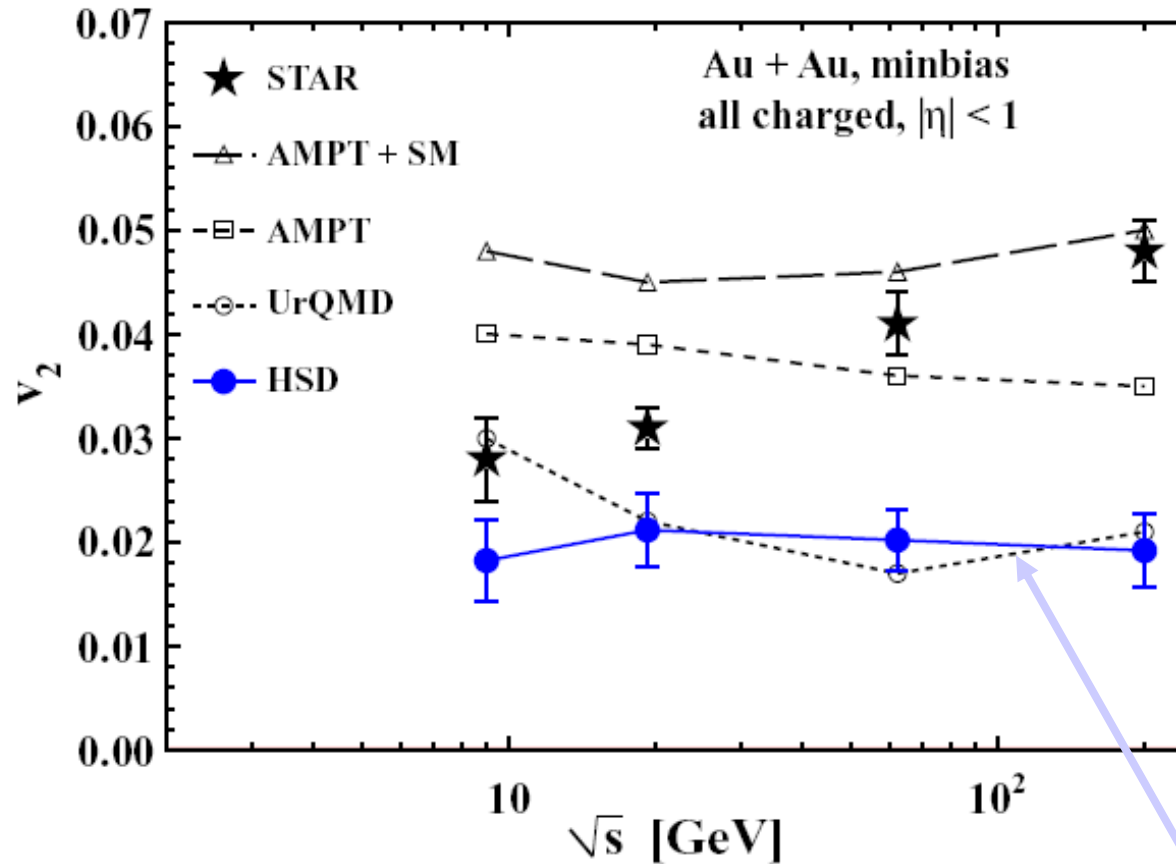
$$\frac{d^2 N}{2\pi p_T dp_T dy} \left(1 + \sum_{n=1}^{\infty} 2v_n(p_T) \cos(n(\psi - \Psi_n)) \right)$$

show higher order harmonics v_n

S. A. Voloshin, arXiv:1111.7241



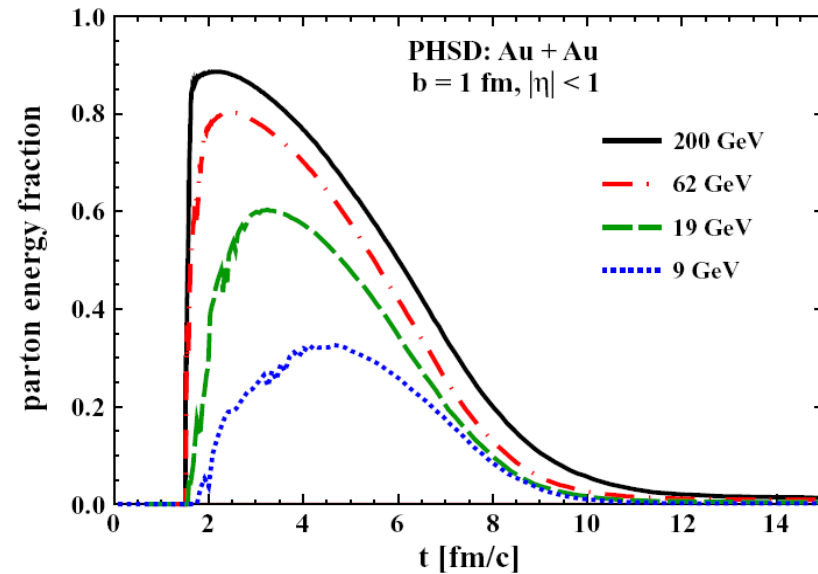
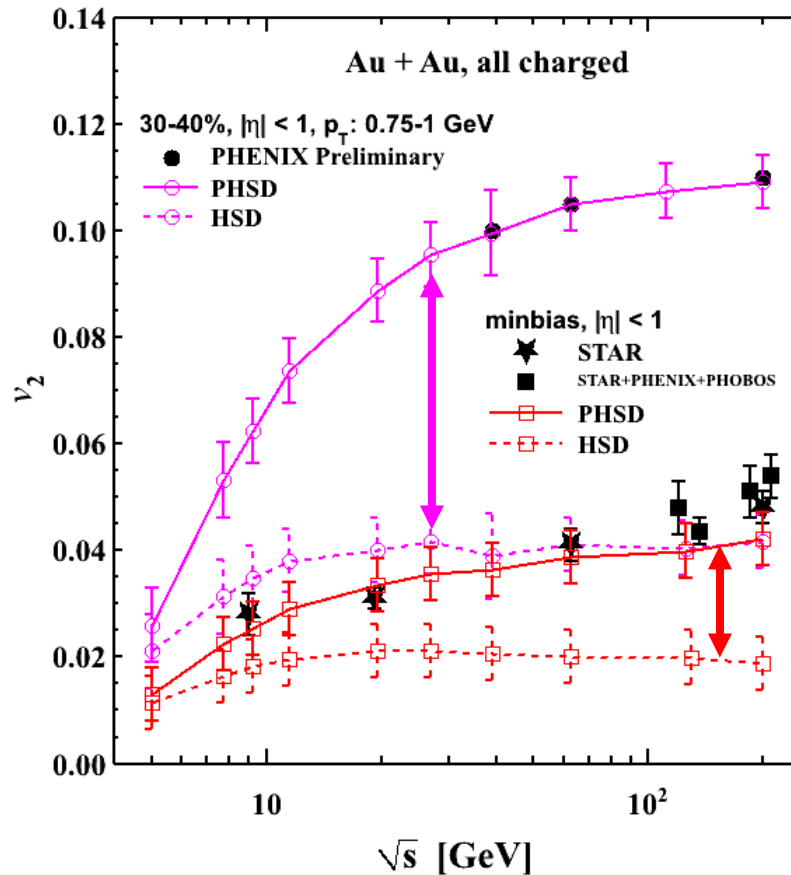
The problem: excitation function of elliptic flow



Excitation function of elliptic flow is not described by **hadron-string** or **purely partonic** models !



Elliptic flow v_2 vs. collision energy for Au+Au



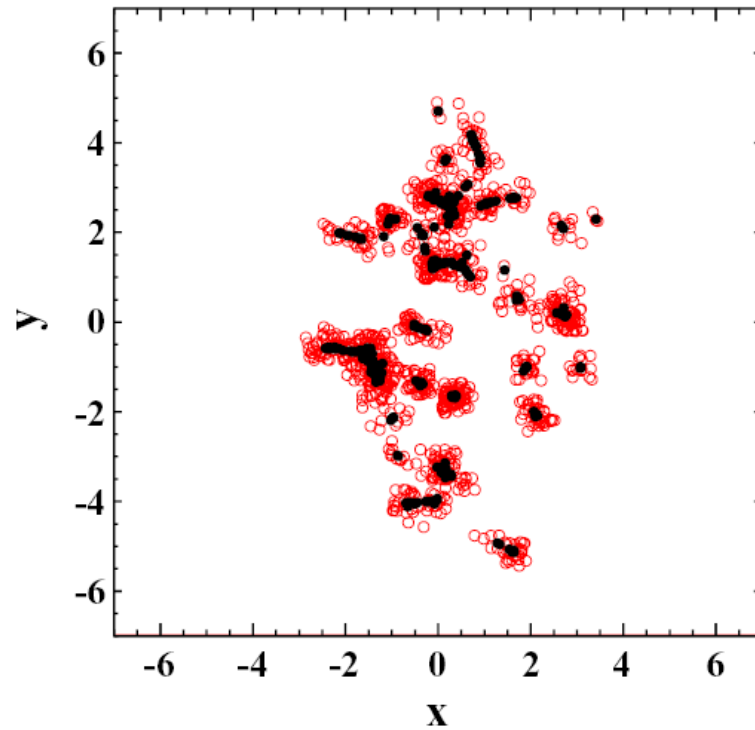
- v_2 in PHSD is larger than in HSD due to the repulsive scalar mean-field potential $U_s(\rho)$ for partons
- v_2 grows with bombarding energy due to the increase of the parton fraction !

V. Konchakovski, E. Bratkovskaya, W. Cassing, V. Toneev, V. Voronyuk,
 Phys. Rev. C 85 (2012) 011902

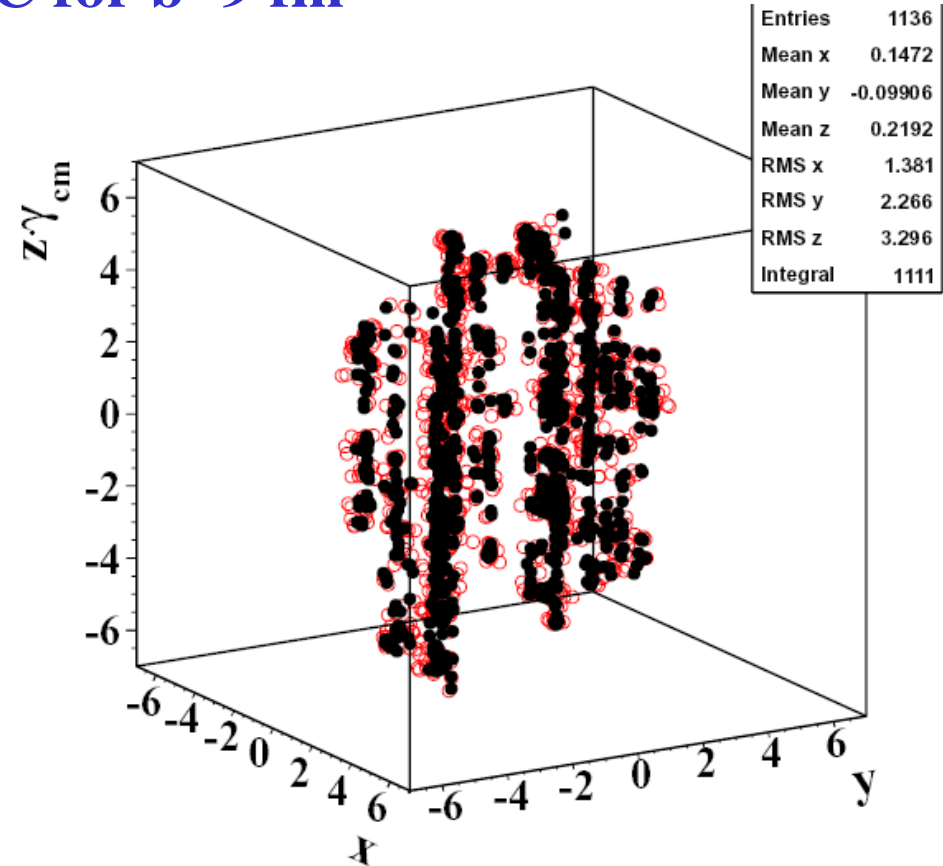


Initial state fluctuations in energy density

Au + Au @ RHIC for $b=9$ fm



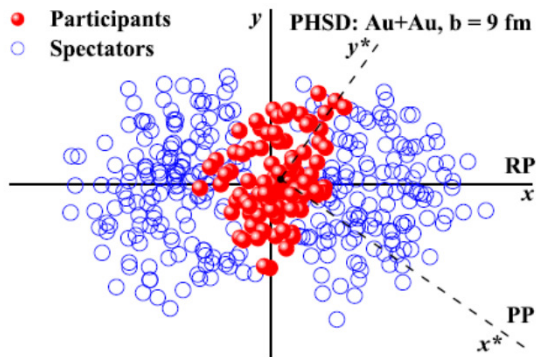
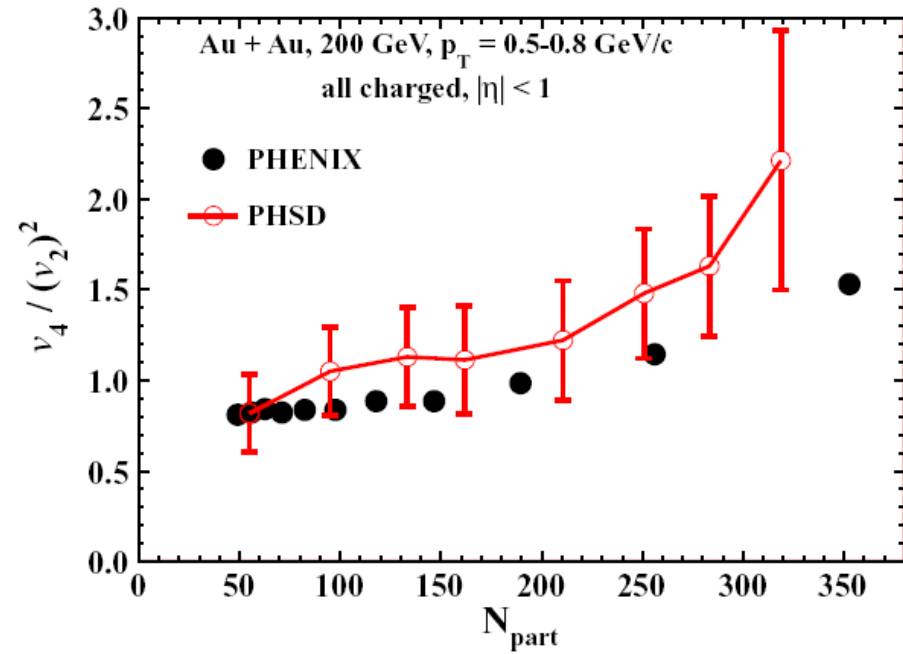
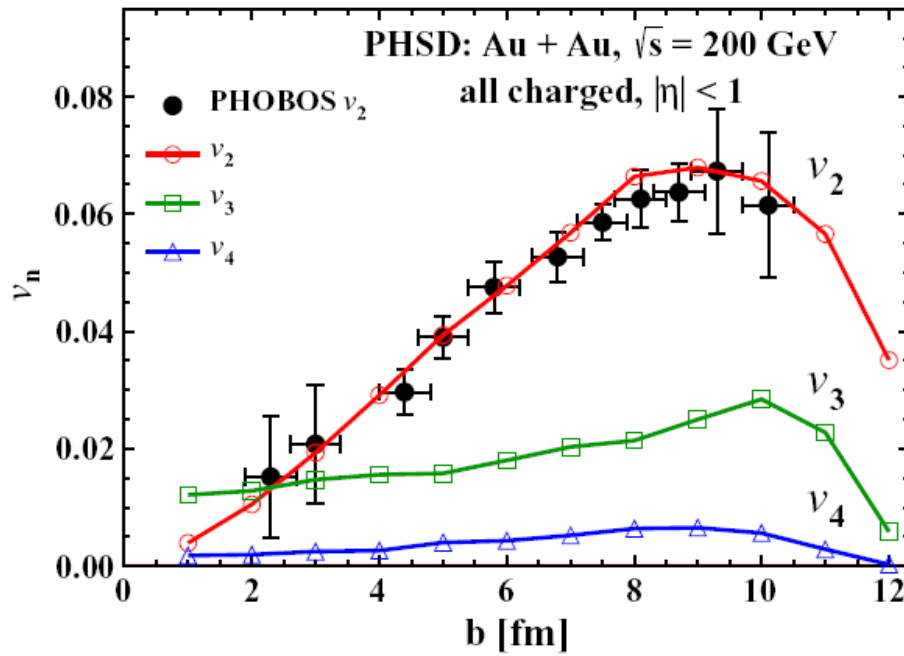
Transverse plane



Longitudinal plane



Flow coefficients versus centrality at RHIC

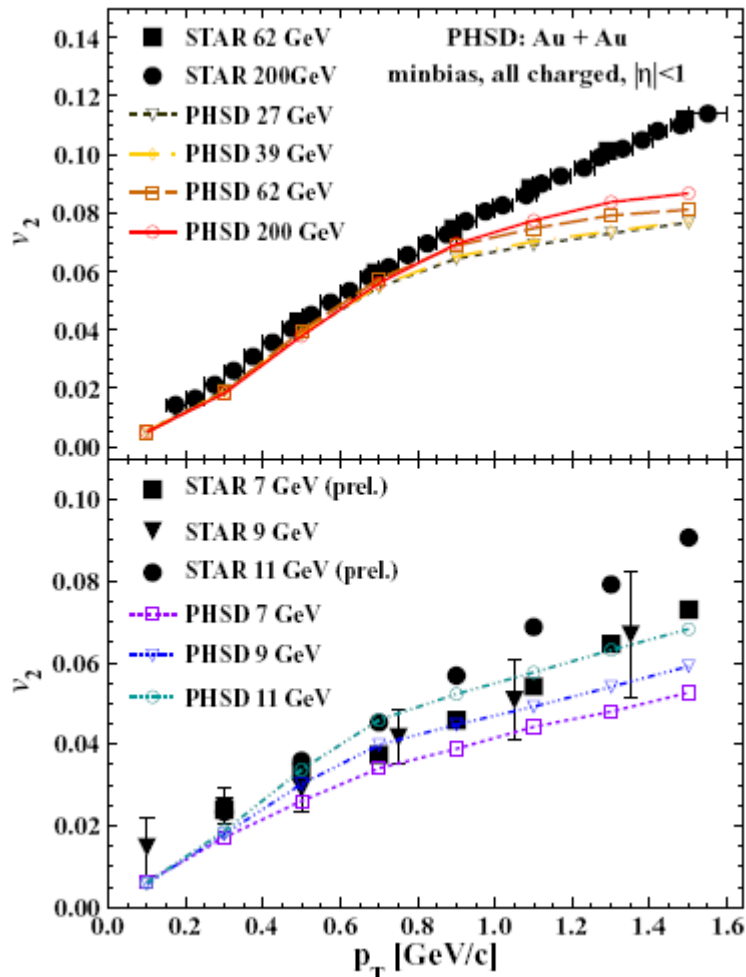


□ increase of v_2 with impact parameter but flat v_3 and v_4

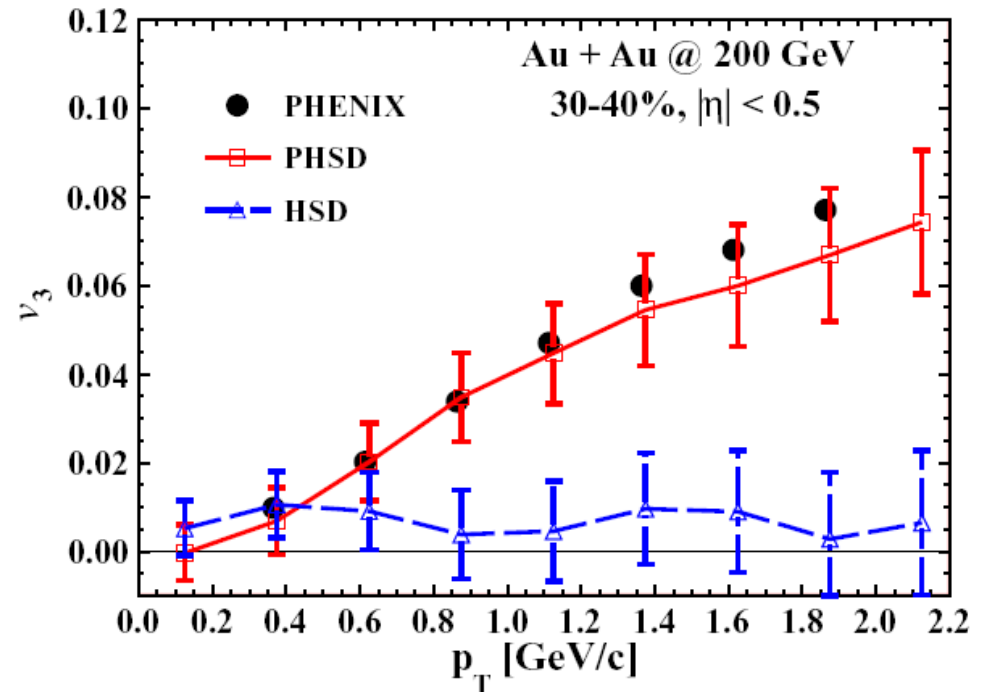


Transverse momentum dependence at RHIC

elliptic flow



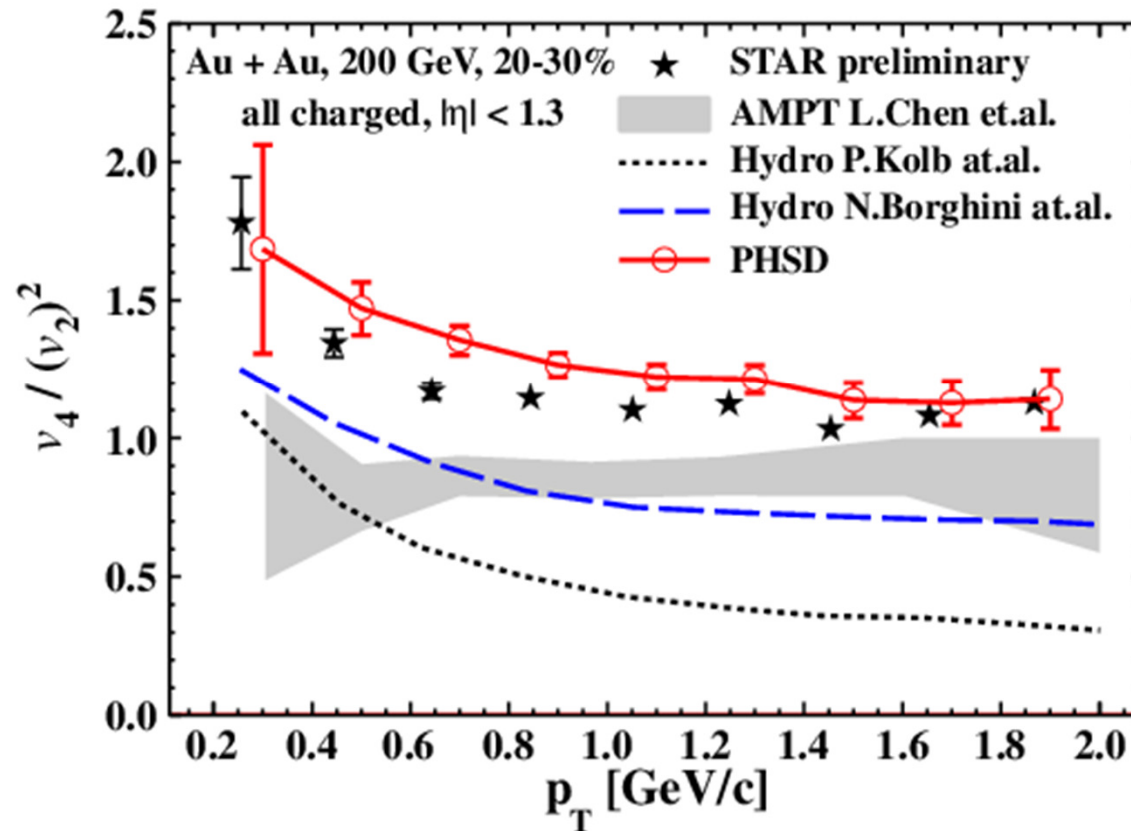
triangular flow



- v_2 vs. p_T follows an approximate scaling for high invariant energies $s^{1/2}=27, 39, 62, 200$ GeV
- v_3 : needs partonic degrees-of-freedom !



Ratio $v_4/(v_2)^2$ vs. p_T at RHIC

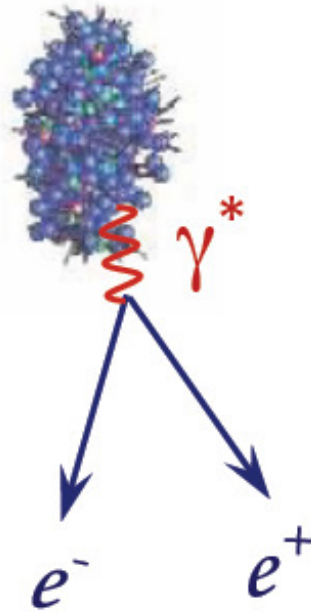


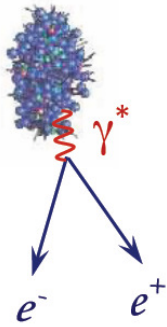
The ratio $v_4/(v_2)^2$:

- is very sensitive to the microscopic dynamics
- PHSD: ratio grows at low p_T - in line with exp. data

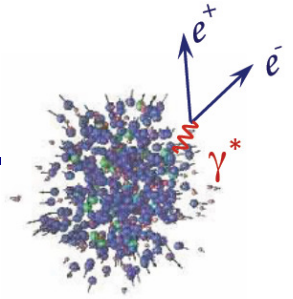
V. Konchakovski, E. Bratkovskaya, W. Cassing, V. Toneev, V. Voronyuk,
Phys. Rev. C 85 (2012) 044922

Dileptons





Electromagnetic probes: dileptons and photons



- Dileptons are emitted from different stages of the reaction and not much effected by final-state interactions

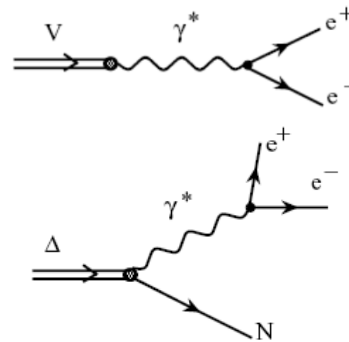
Dilepton sources:

- from the QGP via partonic (q,qbar, g) interactions:

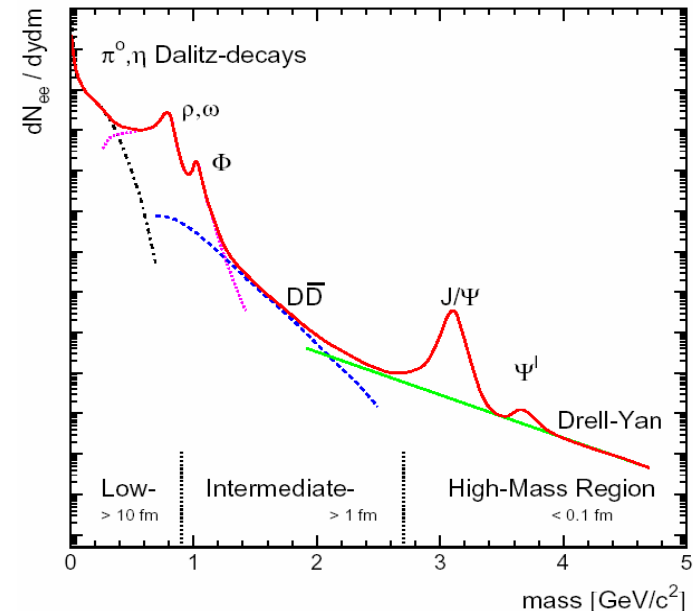


- from hadronic sources:

- direct decay of vector mesons ($\rho, \omega, \phi, J/\Psi, \Psi'$)
- Dalitz decay of mesons and baryons ($\pi^0, \eta, \Delta, \dots$)
- correlated $D+Dbar$ pairs
- radiation from multi-meson reactions ($\pi+\pi, \pi+\rho, \pi+\omega, \rho+\rho, \pi+a_1$) - , 4π '



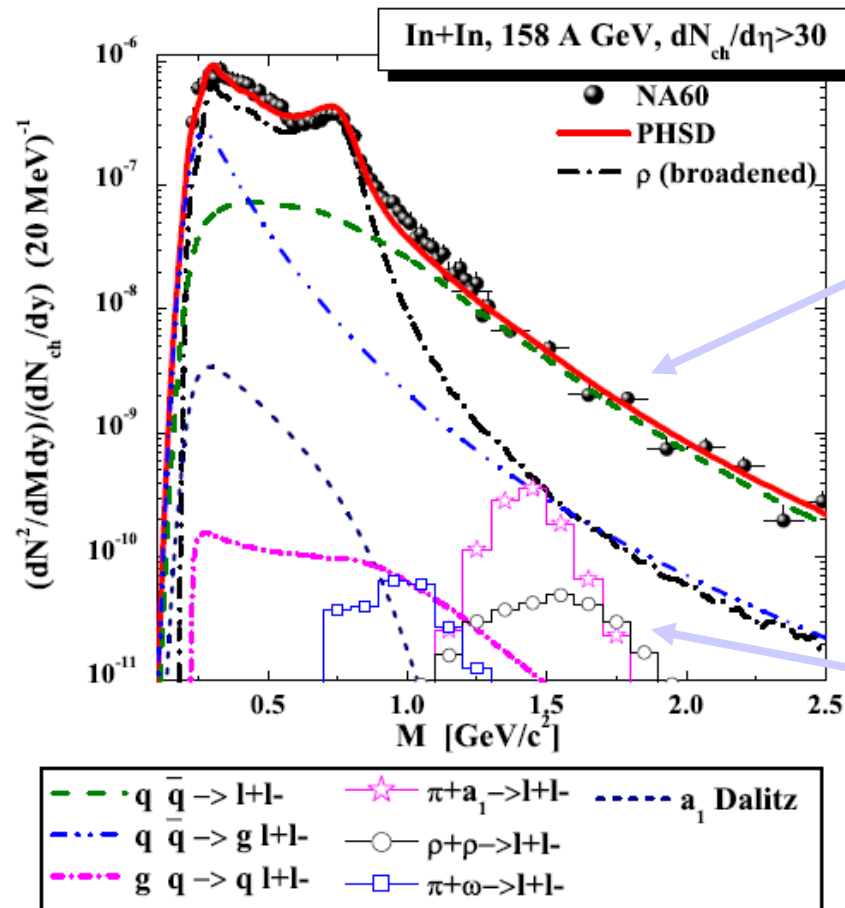
➔ Dileptons are an ideal probe to study the properties of the hot and dense medium



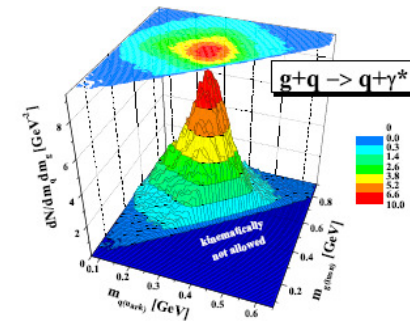
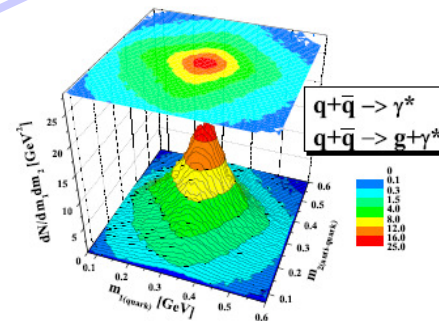


Dileptons at SPS: NA60

Acceptance corrected NA60 data



Mass region above 1 GeV is dominated by **partonic radiation** !

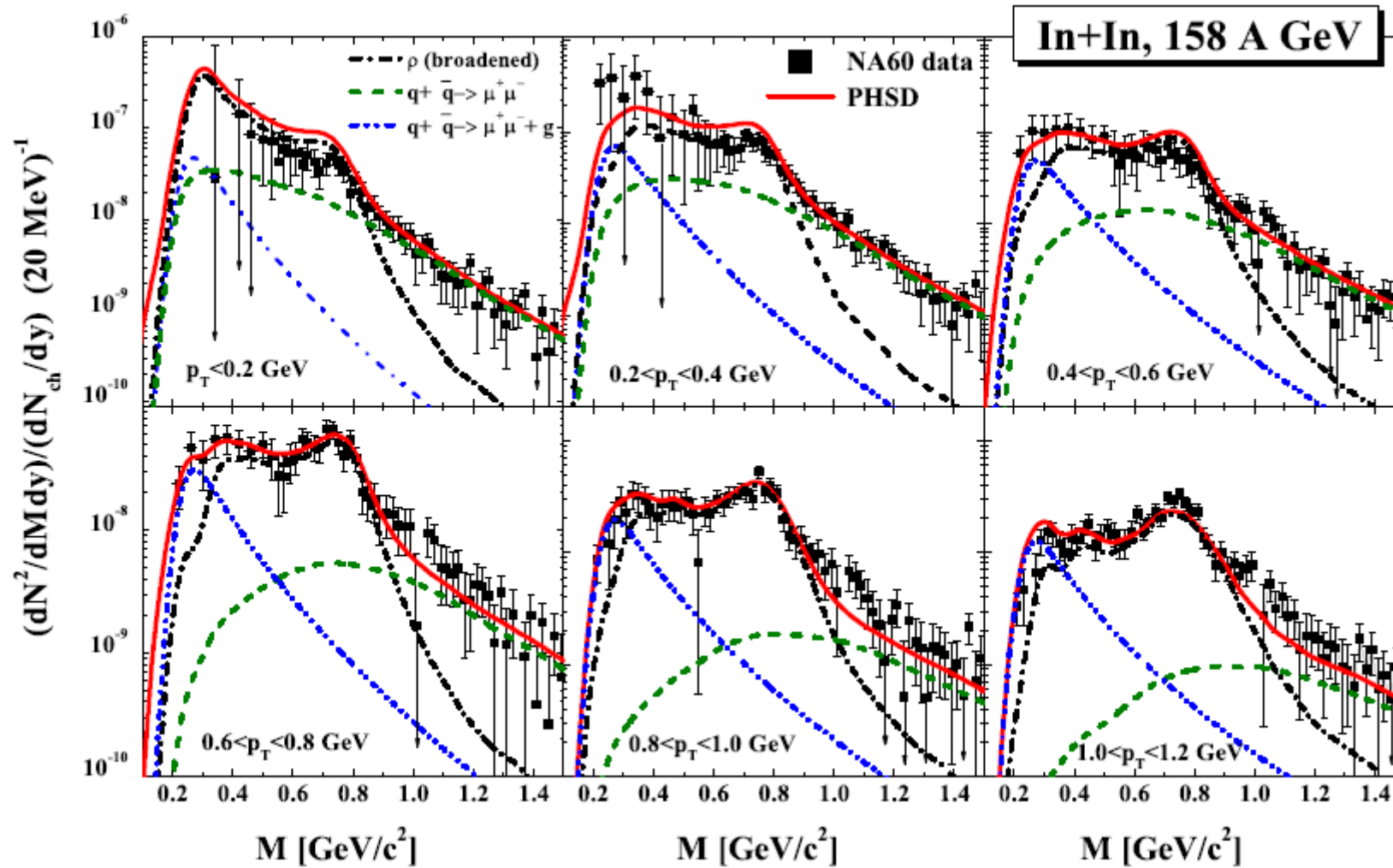


Contributions of “**4π**” channels (radiation from multi-meson reactions) are **small**

O. Linnyk, E.B., V. Ozvenchuk, W. Cassing and C.-M. Ko, PRC 84 (2011) 054917



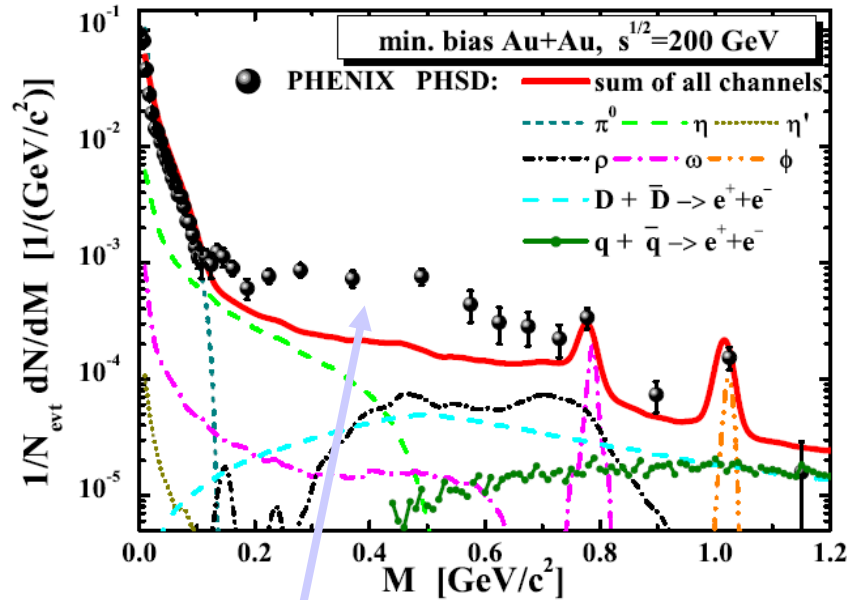
Dileptons at SPS: NA60



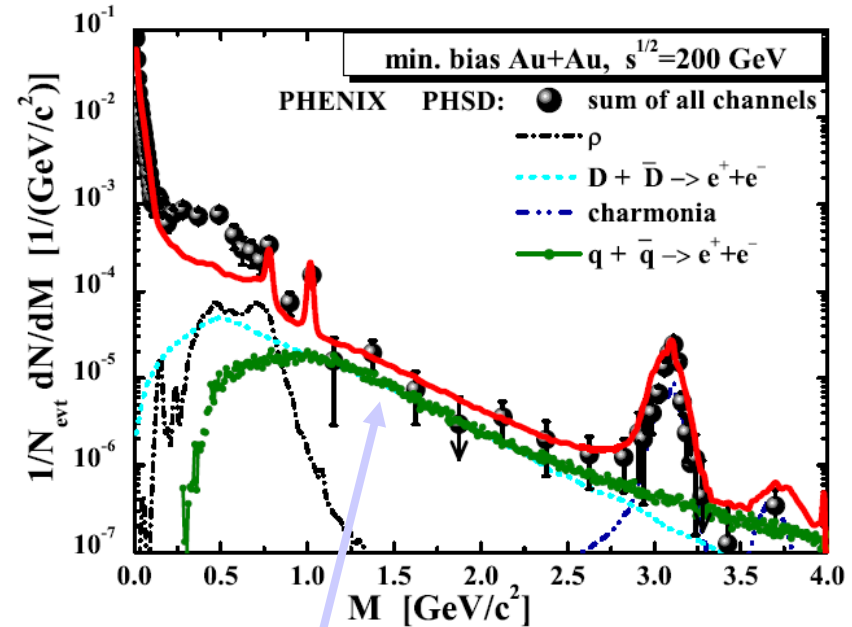
O. Linnyk, E.B., V. Ozvenchuk, W. Cassing
and C.-M. Ko, PRC 84 (2011) 054917



PHENIX: dileptons from partonic channels



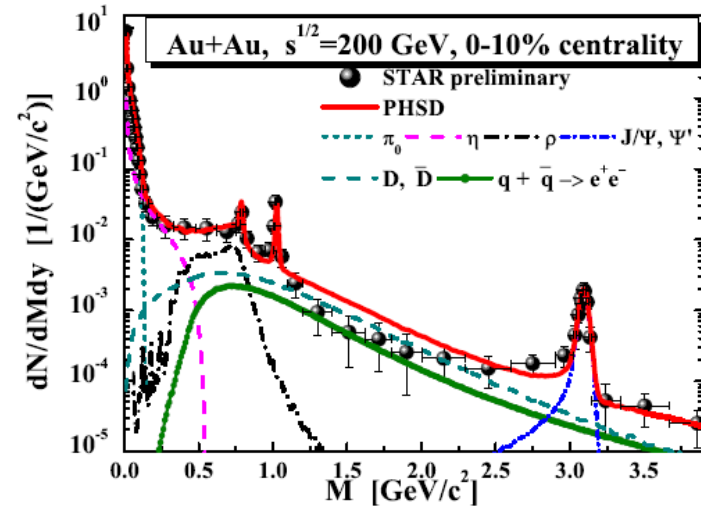
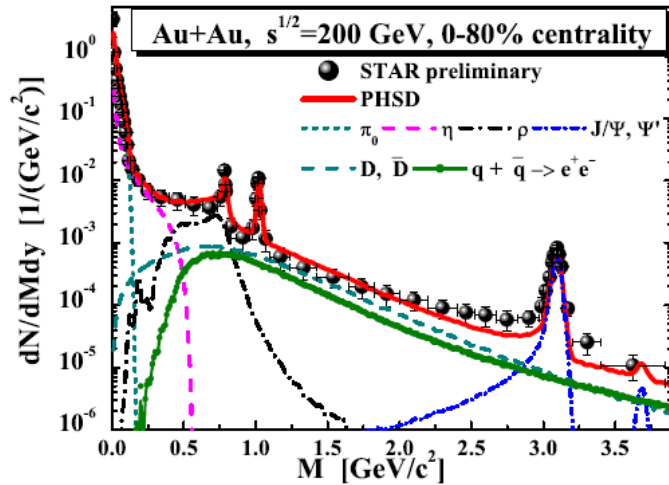
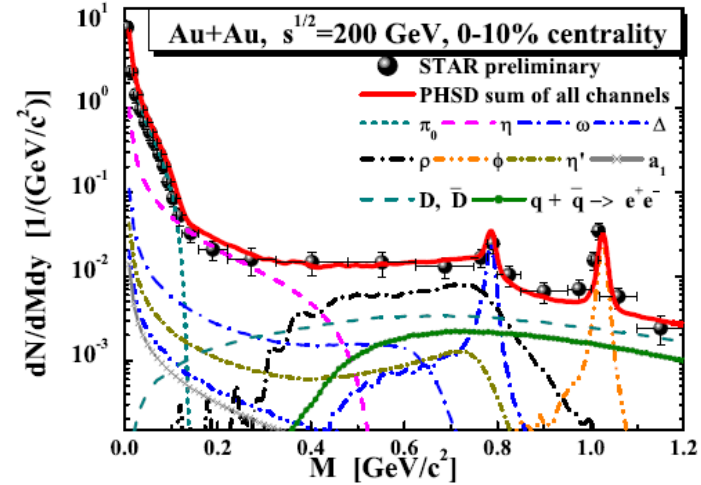
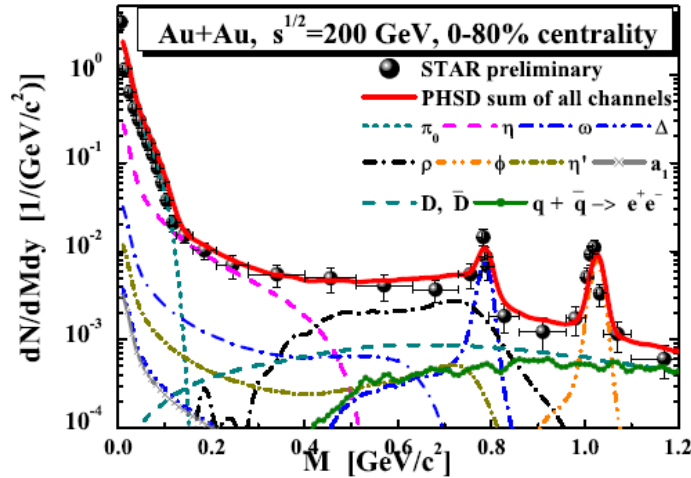
- The **excess** over the considered mesonic sources for $M=0.15-0.6$ GeV is not explained by the QGP radiation as incorporated in PHSD



- The **partonic channels** fill up the discrepancy between the hadronic contributions and the data for $M > 1$ GeV



STAR: mass spectra

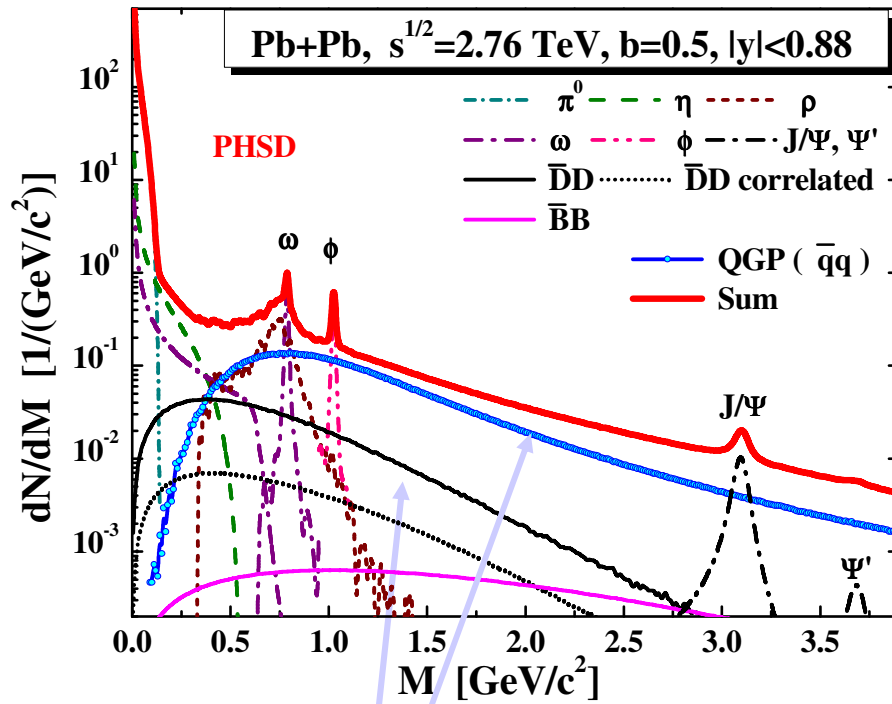


▪ STAR data are well described!

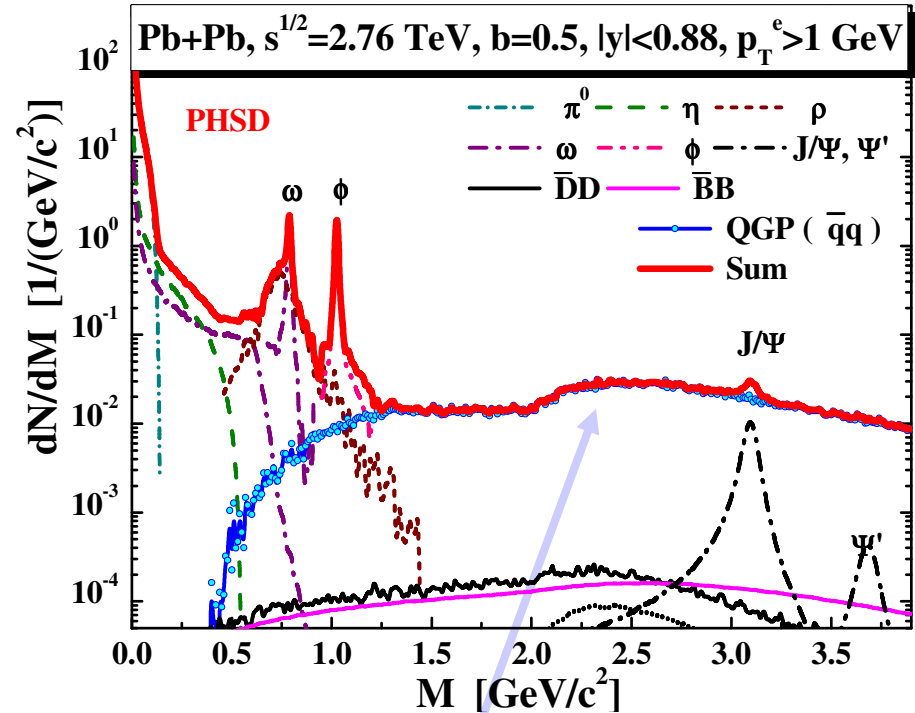
O. Linnyk, W. Cassing, J. Manninen, E.B. and C.-M. Ko,
 PRC 85 (2012) 024910



Predictions for LHC



QGP($\bar{q}q$) dominates at $M > 1.2$ GeV



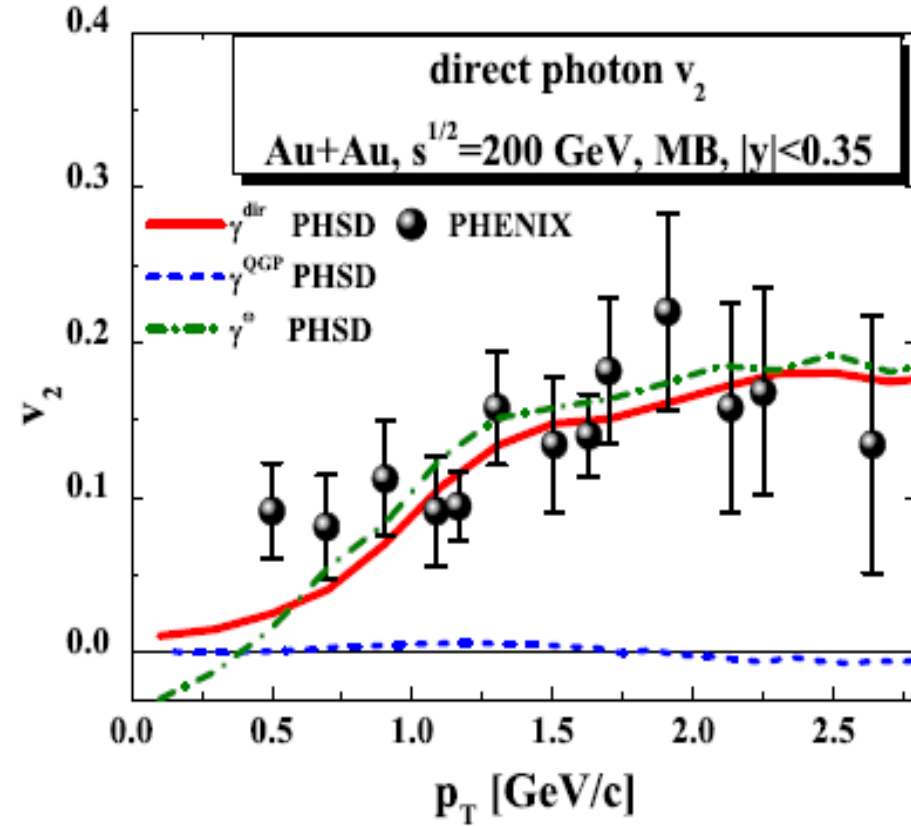
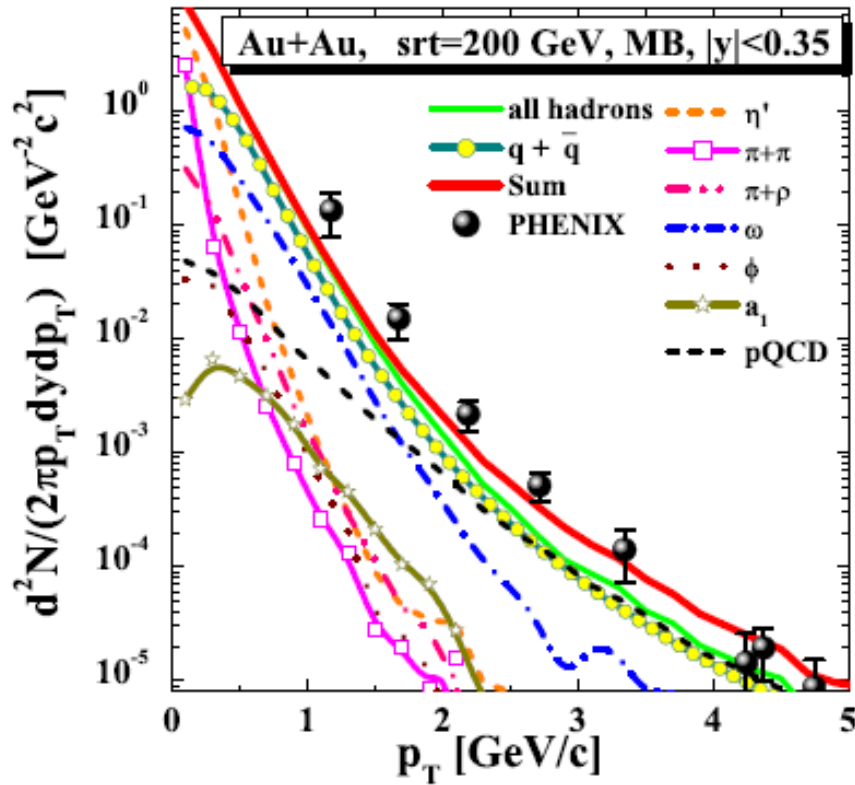
p_T cut enhances the signal of QGP($\bar{q}q$)

- ❑ D-, B-mesons energy loss from Pol-Bernard Gossiaux and Jörg Aichelin
- ❑ J/ Ψ and Ψ' nuclear modification from Che-Ming Ko and Taesoo Song

O. Linnyk, W. Cassing, J. Manninen, E. L. B., P. B. Gossiaux, J. Aichelin, T. Song, C. M. Ko, PRC 87 (2013) 014905



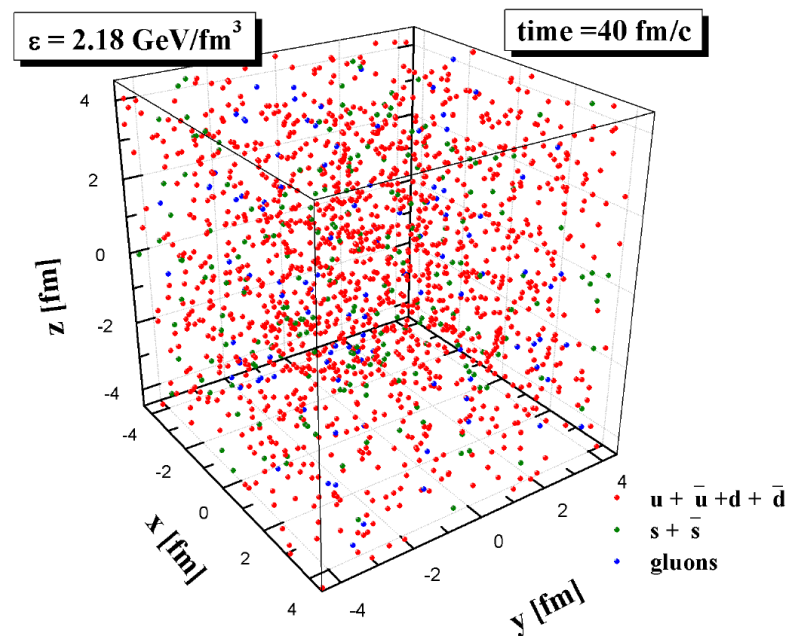
Direct photons and photon v_2 @ RHIC



appears to be compatible with PHENIX data !

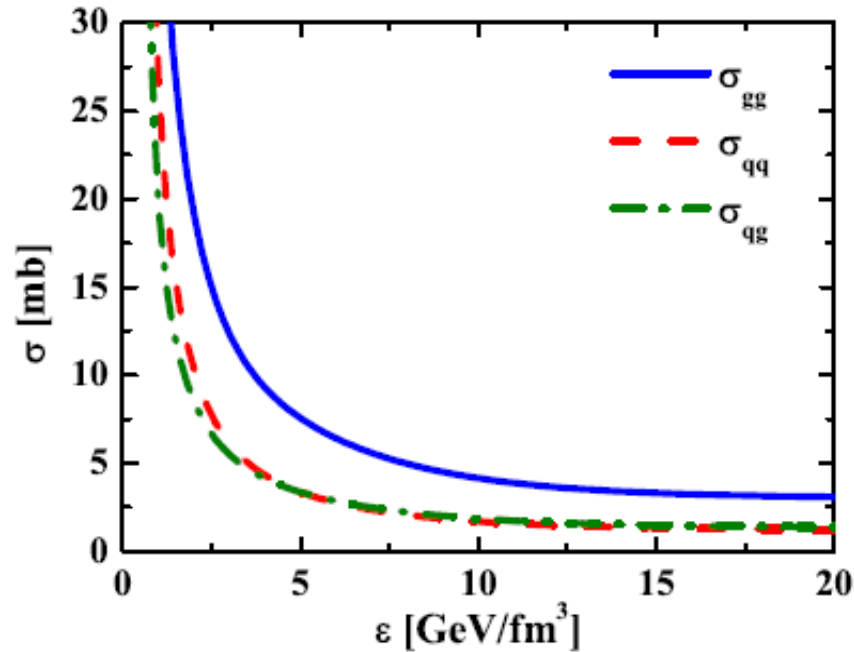
Transport coefficients from PHSD in a box

■ **Kubo:**
$$\eta = \frac{1}{T} \int d^3r \int_0^\infty dt \langle \pi^{xy}(\mathbf{0}, 0) \pi^{xy}(\mathbf{r}, t) \rangle_{\text{equil}}$$



■ V. Ozvenchuk et al., arXiv: 1212.5393

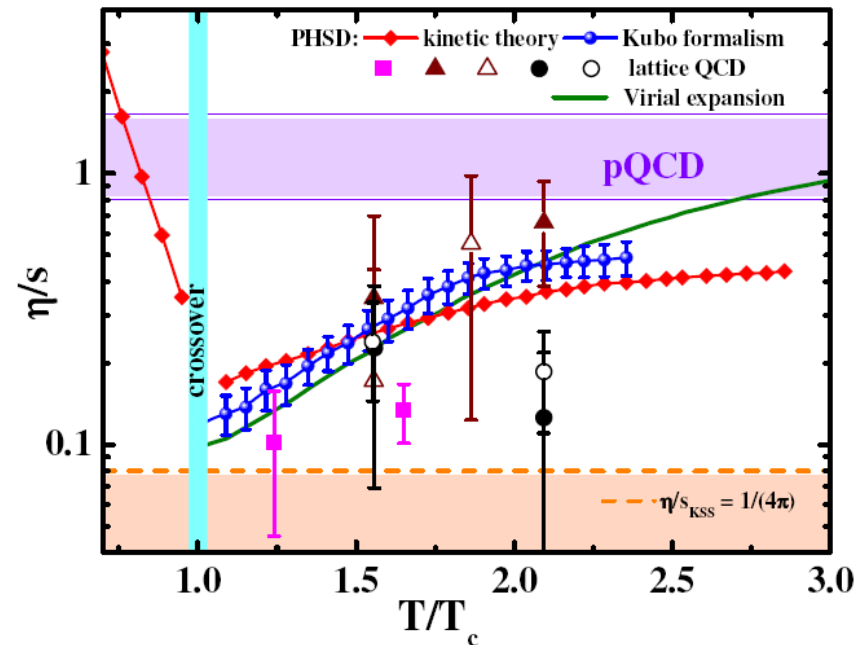
Transport coefficients



$$\eta = \frac{1}{T} \int d^3r \int_0^\infty dt \langle \pi^{xy}(\mathbf{0}, 0) \pi^{xy}(\mathbf{r}, t) \rangle,$$

Cross sections in PHSD

$$\pi^{xy}(\mathbf{r}, t) \equiv T^{xy}(\mathbf{r}, t) = \int \frac{d^3p}{(2\pi)^3} \frac{p^x p^y}{E} f(\mathbf{r}, \mathbf{p}; t).$$



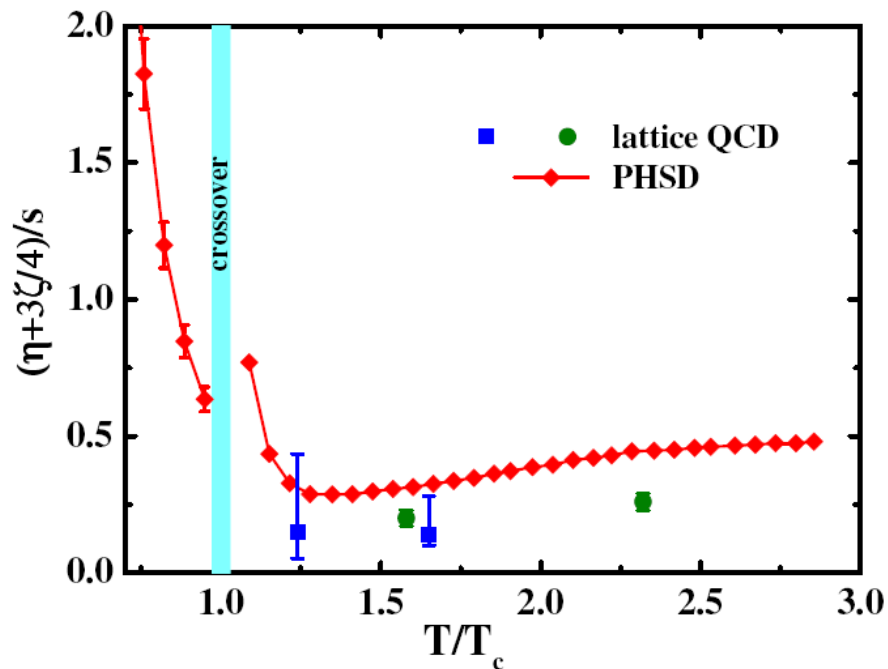
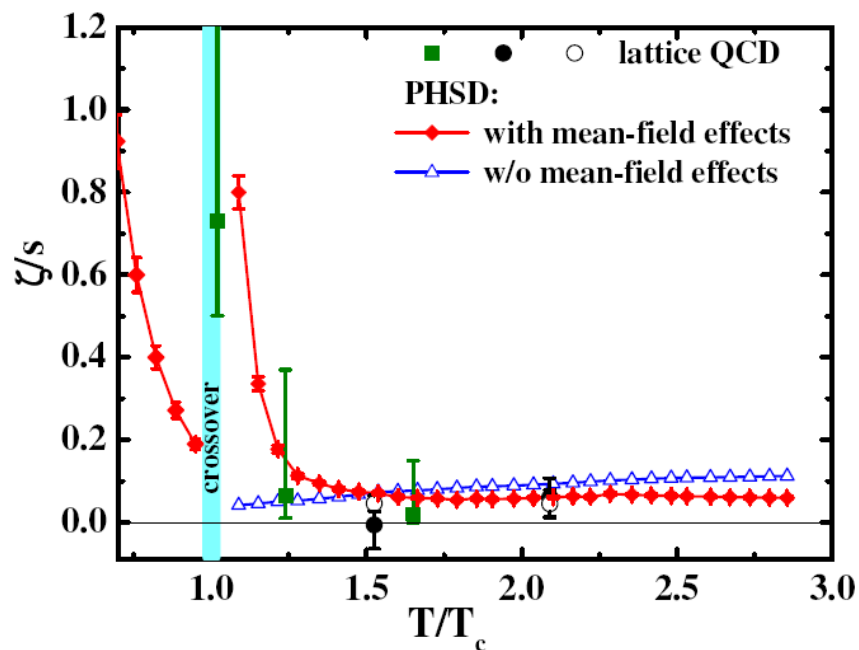
- Shear viscosity shows a minimum close to T_c !

Bulk viscosity

$$\zeta = \frac{1}{T} \sum_a \int \frac{d^3p}{(2\pi)^3} \frac{\tau_a(E_a)}{E_a^2} f_a^{eq}(E_a/T) \times \left[\left(\frac{1}{3} - v_s^2 \right) |\mathbf{p}|^2 - v_s^2 \left(m_a^2 - T^2 \frac{dm_a^2}{dT^2} \right) \right]^2$$

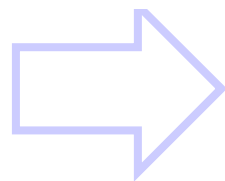
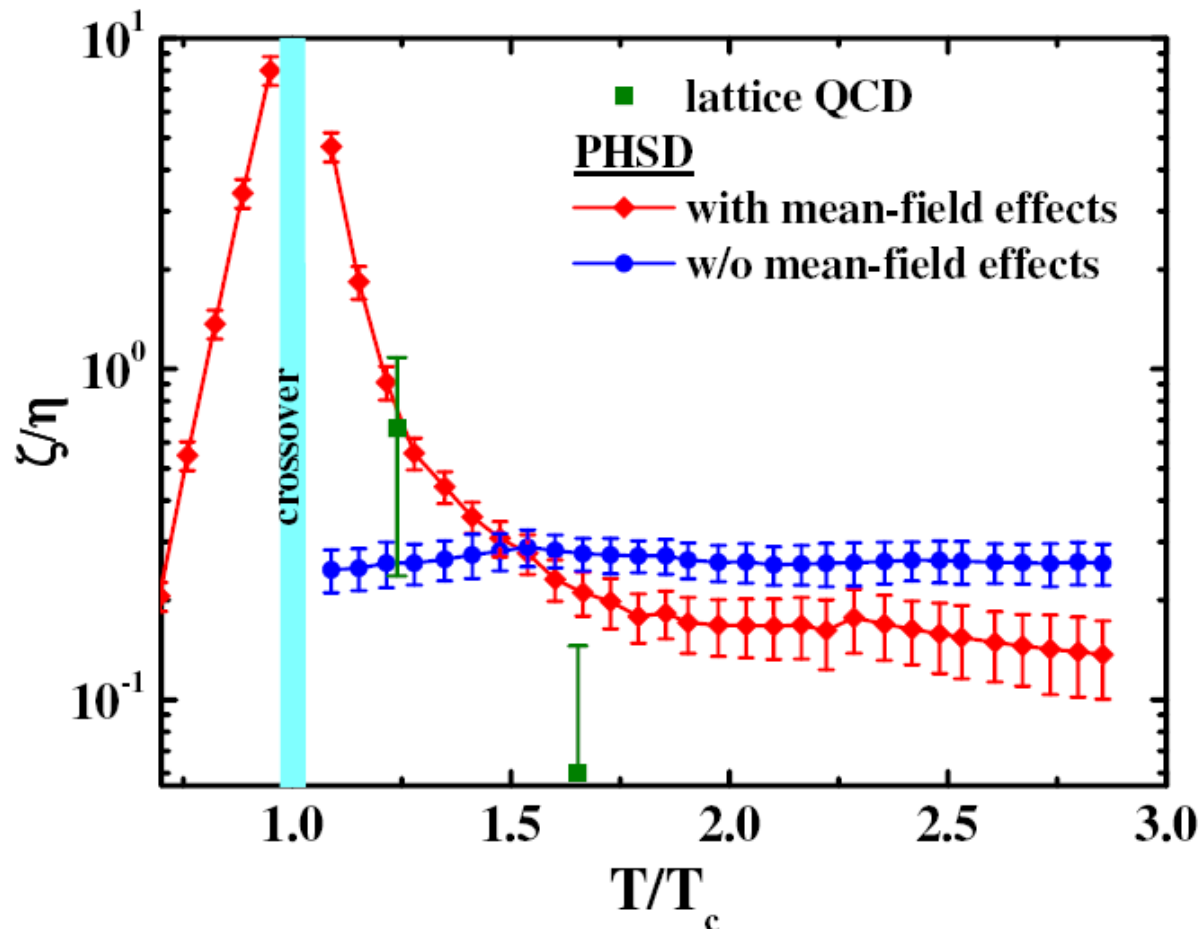
■ in comparison to IQCD

■ specific sound



■ shows a maximum close to T_c

bulk/shear versus temperature T

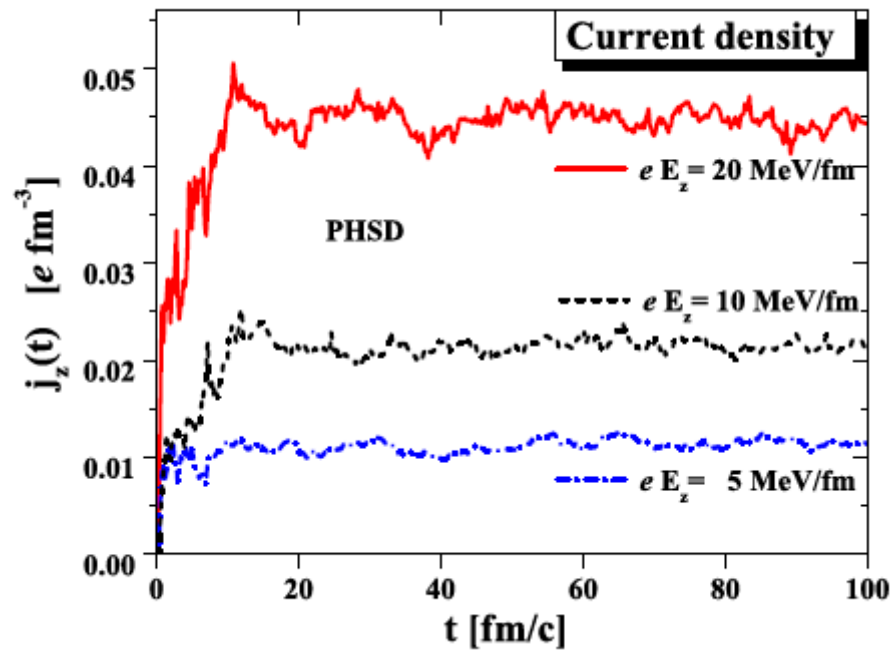


■ pronounced maximum at T_c !

electric conductivity I

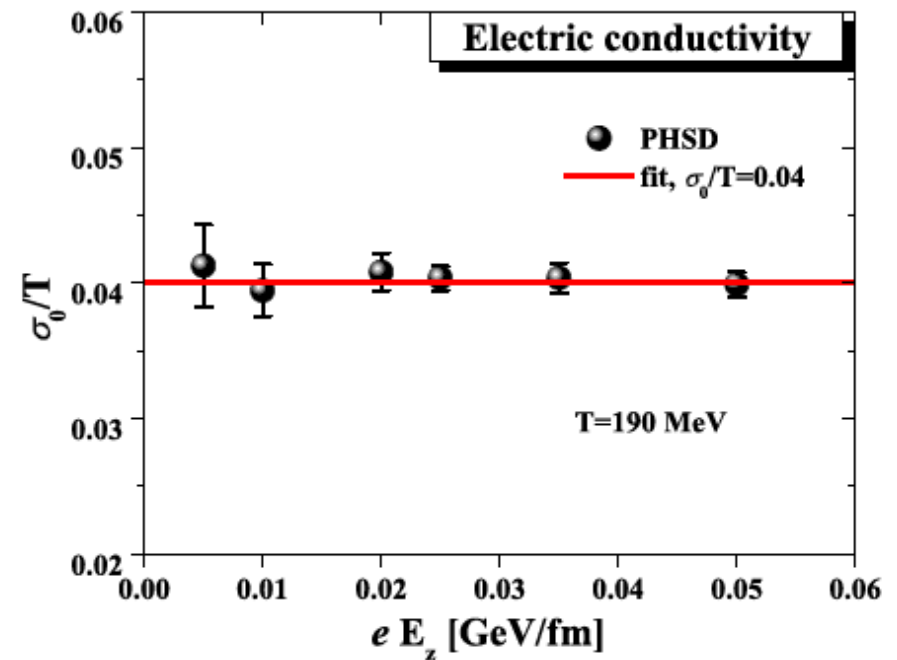
■ induced current j_z

$$j_z(t) = \frac{1}{V} \sum_j eq_j \frac{p_z^j(t)}{M_j(t)}$$

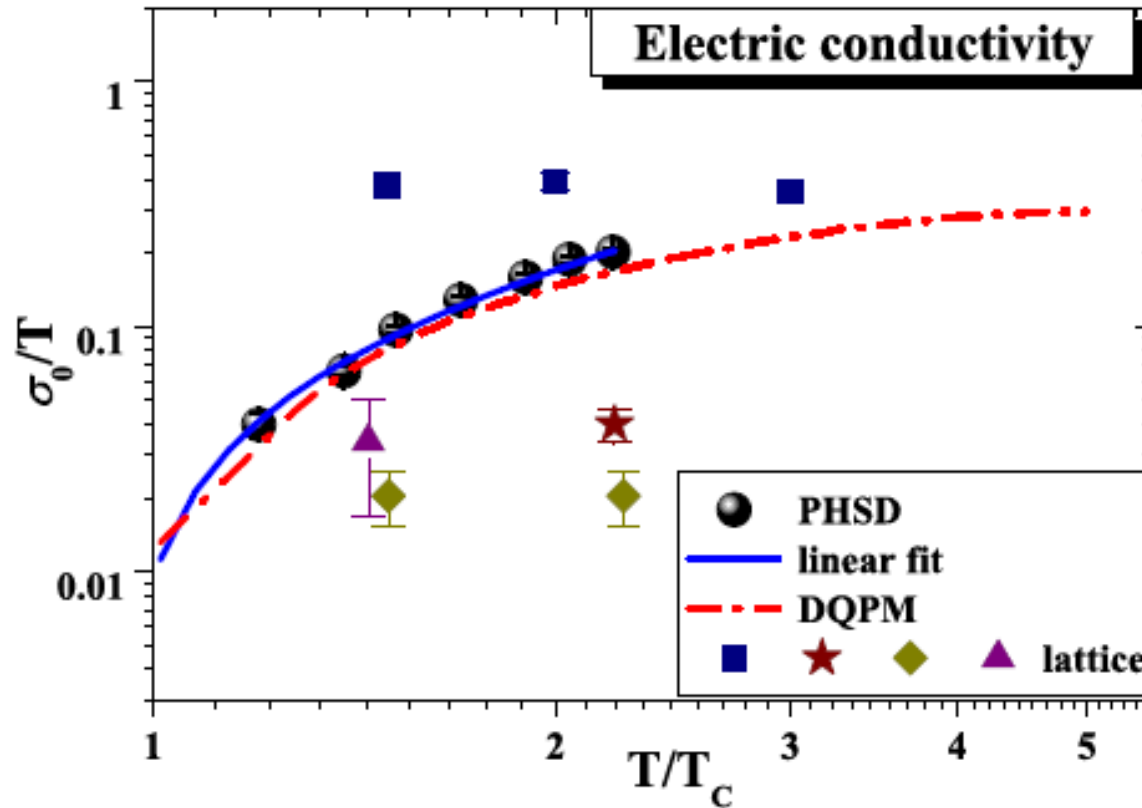


■ electric conductivity

$$\frac{\sigma_0}{T} = \frac{j_{eq}}{E_z T}$$



electric conductivity II



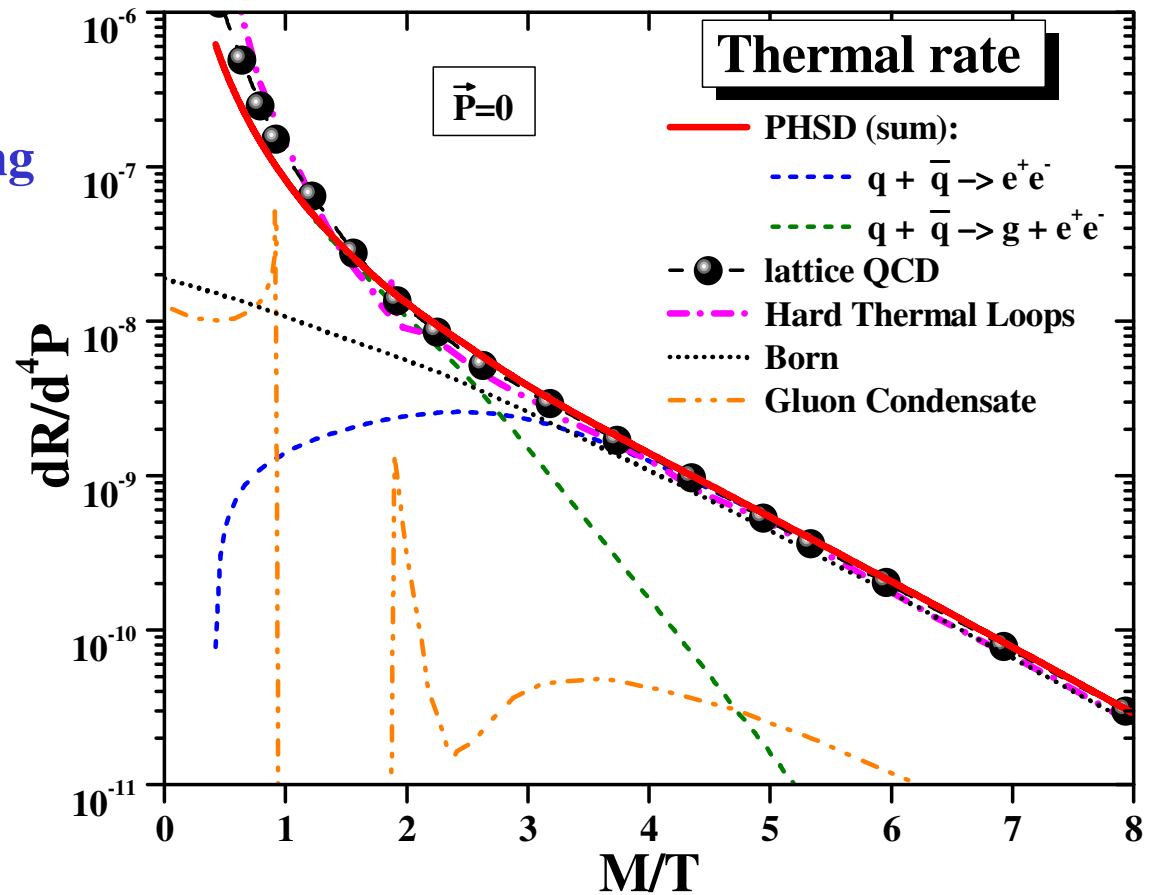
- relaxation time approach in DQPM: $\frac{\sigma_0(T)}{T} \approx \frac{2}{9} \frac{e^2 n_q(T)}{M_q(T) \Gamma_q(T) T}$

Thermal dilepton rates in lattice QCD

Dileptons from dynamical off-shell quark and gluon interactions, LO and NLO in the coupling

- Qualitative agreement of dynamical quasiparticles, lattice QCD and HTL

back to back lepton pairs





Summary

- **PHSD** provides a consistent description of **off-shell parton dynamics** in line with the **lattice QCD equation of state** (from the BMW collaboration)
- **PHSD versus experimental observables:**
 - enhancement of meson m_T slopes (at top SPS and RHIC)
 - strange antibaryon enhancement (at SPS)
 - partonic emission of high mass dileptons at SPS and RHIC
 - enhancement of collective flow v_2 with increasing energy
 - reasonable description of v_2, v_3, v_4 at RHIC vs p_T
 - quark number scaling of v_2 (at RHIC)
 - ...
- **PHSD results for infinite systems in equilibrium:**
 - minimum of η/s close to T_c and maximum of ζ/s close to T_c
 - pronounced** maximum of ζ/η close to T_c
 - minimum of σ/T close to T_c



PHSD group



Wolfgang Cassing (Giessen Univ.)
Volodya Konchakovski (Giessen Univ.)
Olena Linnyk (Giessen Univ.)
Thorsten Steinert (Giessen Univ.)



Elena Bratkovskaya (FIAS & ITP Frankfurt Univ.)
Vitalii Ozvenchuk (HGS-HIRE, FIAS & ITP Frankfurt Univ.)
Rudy Marty (FIAS, Frankfurt Univ.)
Hamza Berrehrah (FIAS, Frankfurt Univ.)
Daniel Cabrera (ITP&FIAS, Frankfurt Univ.)



External Collaborations:

SUBATECH, Nantes Univ. :

Jörg Aichelin
Christoph Hartnack
Pol-Bernard Gossiaux



Texas A&M Univ.:

Che-Ming Ko



JINR, Dubna:

Vadim Voronyuk
Viatcheslav Toneev



Kiev Univ.:

Mark Gorenstein