## Heavy quarks in the quark-gluon plasma

#### Jon-Ivar Skullerud with Aoife Kelly, Dhagash Mehta, Buğra Oktay, Sinéad Ryan and others

NUI Maynooth / FASTSUM collaboration

Quarks, Gluons and Hadronic Matter under Extreme Conditions, St. Goar, 16 March 2011

## Outline

Background Quenched vs dynamical Spectral functions

#### Charmonium

Temperature dependence Reconstructed correlators Nonzero momentum Towards the physical limit

Charm diffusion

Beauty (and the beast?)

Summary and outlook

Quenched vs dynamical Spectral functions

# Background

- $J/\psi$  suppression a probe of the quark–gluon plasma?
- Heavy quarks: hard probes or thermal?
- Quenched lattice results indicate that S-waves survive well into the plasma phase
- Sequential suppression + recombination explains experimental results?
- Heavy quarks as thermometer of QGP?

Quenched vs dynamical Spectral functions

# Background

- $J/\psi$  suppression a probe of the quark–gluon plasma?
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- Quenched lattice results indicate that S-waves survive well into the plasma phase
- Sequential suppression + recombination explains experimental results?
- Heavy quarks as thermometer of QGP?
- Uncertainty about which potential to use in potential models, how to treat continuum
- How reliable are quenched lattice simulations?

Quenched vs dynamical Spectral functions

# Quenched vs dynamical

Are quenched lattice results reliable?

- $T_c^{N_f=0} \approx 1.5 T_c^{N_f=2+1}, T_c^{N_f=2} \approx T_c^{N_f=2+1}$
- No  $D \overline{D}$  threshold in quenched QCD
- Light quarks can catalyse QQ dissociation so it occurs at lower temperature
- Lower  $T_c$ , lower  $T_d$  conspire to give the same  $T_d/T_c$ ?
- Potential models indicate little change in  $T_d/T_c$

Quenched vs dynamical Spectral functions

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- Lower  $T_c$ , lower  $T_d$  conspire to give the same  $T_d/T_c$ ?
- Potential models indicate little change in  $T_d/T_c$
- Only dynamical lattice calculations can give the answer

Quenched vs dynamical Spectral functions

## Dynamical anisotropic lattices

- A large number of points in time direction required
- For  $T = 2T_c$ ,  $\mathcal{O}(10)$  points  $\Longrightarrow a_t \sim 0.025$  fm
- Far too expensive with isotropic lattices  $a_s = a_t!$

Quenched vs dynamical Spectral functions

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Quenched vs dynamical Spectral functions

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- Introduces 2 additional parameters
- Non-trivial tuning problem
   [PRD 74 014505 (2006); PRD 78, 014505 (2008)]

Quenched vs dynamical Spectral functions

# Spectral functions

•  $\rho_{\Gamma}(\omega, \overrightarrow{p})$  related to euclidean correlator  $G_{\Gamma}(\tau, \overrightarrow{p})$  according to

$$G_{\Gamma}(\tau, \overrightarrow{p}) = \int \rho_{\Gamma}(\omega, \overrightarrow{p}) \frac{\cosh[\omega(\tau - 1/2T)]}{\sinh(\omega/2T)} d\omega$$

- an ill-posed problem
- use Maximum Entropy Method to determine most likely  $\rho(\omega)$
- requires a large number of time slices to have any chance of a reliable determination
- must introduce model function  $m_0(\omega)$

Femperature dependence Reconstructed correlators Nonzero momentum Fowards the physical limit

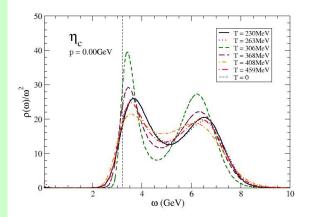
#### Simulation parameters

#### [PRD 76 194513 (2007), arXiv:1005.1209]

ξ	$a_s$ (fm)	$a_t^{-1}$ (GeV	$() m_{\pi/}$	$m_{ ho}$	Ns	$L_s$ (fm)
6.0	0.162	7.3	35 C	).54	12	1.94
		T (NA )()	<b>T</b> / <b>T</b>			_
	$N_{ au}$	T (MeV)	$T/T_c$	# (	configs	
	80	92	0.42		250	
	32	230	1.05		1000	
	28	263	1.20		1000	
	24	306	1.40		500	
	20	368	1.68		1000	
	18	408	1.86		1000	
	16	459	2.09		1000	

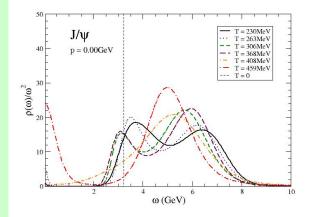
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# S-wave T dependence $(\eta_c)$



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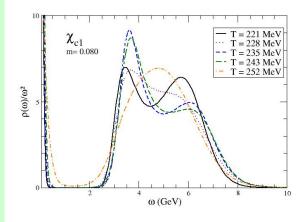
# S-wave T dependence $(J/\psi)$



 $J/\psi$  (S-wave) melts at  $T \sim 370 - 400$  MeV or  $1.7 - 1.9T_c$ ?

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#### **P**-waves



P-waves melt at T < 250 MeV or  $1.2T_c$ ?

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#### Reconstructed correlators

Reconstructed correlator is defined as

$$G_r(\tau; T, T_r) = \int_0^\infty \rho(\omega; T_r) K(\tau, \omega, T) d\omega$$

where K is the kernel

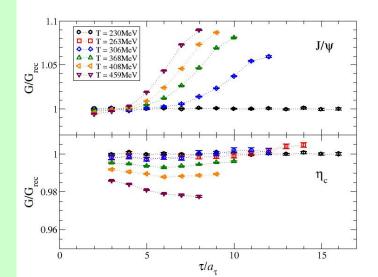
$$K(\tau, \omega, T) = rac{\cosh[\omega(\tau - 1/2T)]}{\sinh(\omega/2T)}$$

If  $\rho(\omega; T) = \rho(\omega; T_r)$  then  $G_r(\tau; T, T_r) = G(\tau; T)$ 

We use  $N_{\tau} = 32$  as our reference temperature

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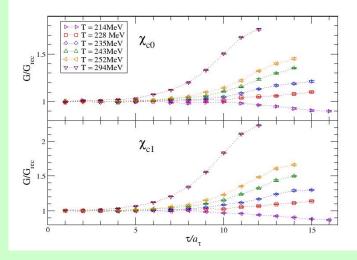
#### S-waves



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#### P-waves



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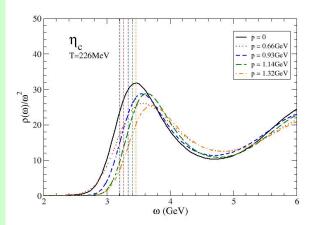
### Nonzero momentum

[With MB Oktay, arXiv:1005.1209]

- Charmonium is produced at nonzero momentum
- Transverse momentum (and rapidity) distributions important to distinguish between models
- Momentum dependent binding?
- Gives an additional window to transport properties

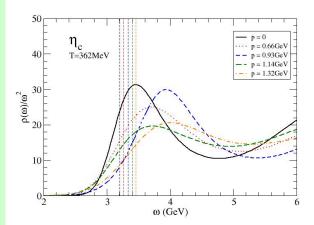
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## Nonzero momentum results $\eta_c, 12^3 \times 32$



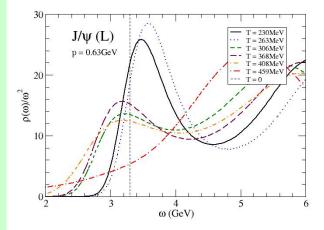
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## Nonzero momentum results $\eta_c, 12^3 \times 20$



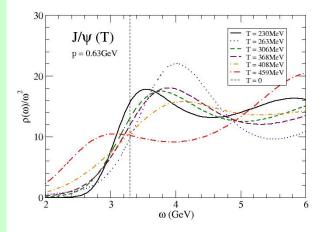
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#### Transverse vs longitudinal



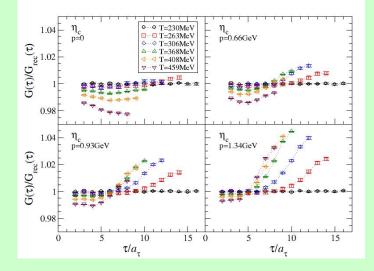
Temperature dependence Reconstructed correlators **Nonzero momentum** Towards the physical limit

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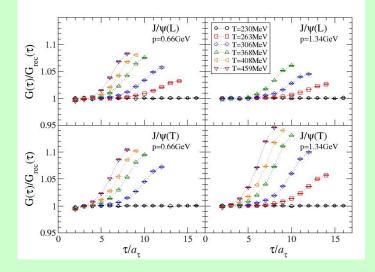
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#### Reconstructed correlators



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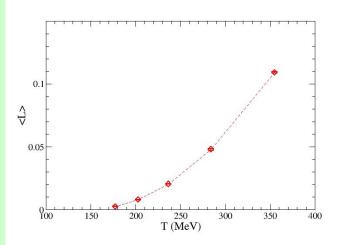
## Towards the physical limit

Anisotropic clover-improved Wilson fermions, 2+1 flavours [HadSpec Collab, PRD **79** 034502 (2009)]

ξ	<i>as</i> (	fm) $a_t^{-1}$ (	(GeV)	$m_{\pi}/m_{ ho}$	Ns	L <sub>s</sub> (f	m)
3.5	0.	122	5.68	0.45	24	2.	.93
-		<b>T</b> (NA ) ()			<i>c</i> :		-
_	$N_{ au}$	T (MeV)	$T/T_c$	# con	tigs	used	_
	160	35	0.2		—		
	32	177	1.0		242	38	
	28	203	1.1		306	100	
	24	237	1.3		259	57	
	20	284	1.6		625	539	
	16	355	2.0		289	102	

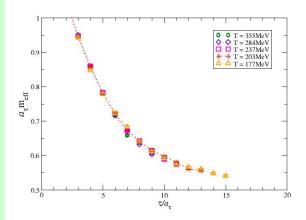
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### Polyakov loop



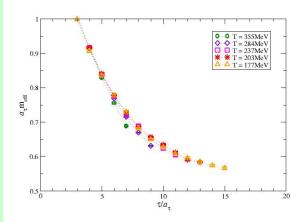
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### Pseudoscalar effective mass



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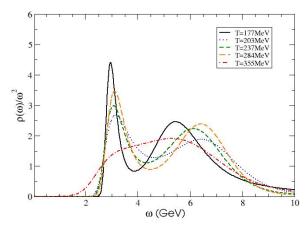
### Vector effective mass



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## Pseudoscalar spectral function

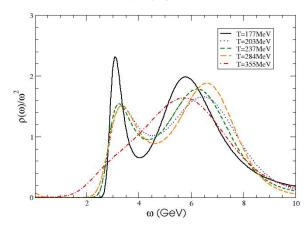
 $\eta_{c} (a_{\tau}m_{c} = 0.087)$ 



Temperature dependence Reconstructed correlators Nonzero momentum Towards the physical limit

## Vector spectral function

 $J/\psi (a_{\tau}m_{c} = 0.087)$ 



# Charm diffusion

How fast do charm quarks thermalise? The heavy quark diffusion constant *D* is given by

$$D = \frac{1}{\chi^{00}} \lim_{\omega \to 0} \frac{\rho_V(\omega)}{\omega} \,,$$

 $ho_V$  is the spectral function of the conserved-current operator  $V_i$ 

$$\chi^{00} = \frac{1}{T} \int \langle V_0(\overrightarrow{x},t) V_0(\overrightarrow{0},t) d^3 x$$

# Charm diffusion

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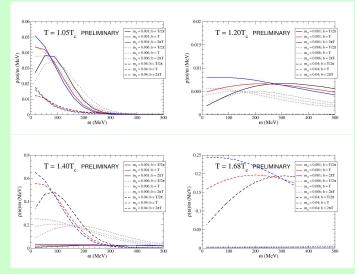
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$$\chi^{00} = \frac{1}{T} \int \langle V_0(\overrightarrow{x}, t) V_0(\overrightarrow{0}, t) d^3 x$$

Preliminary results using default model  $m(\omega) = m_0 \omega (b + \omega)$ 

#### Results



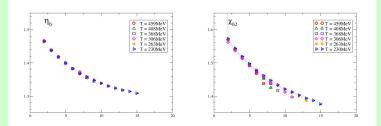
# Beauty (and the beast?)

[See also talk by Gert Aarts]

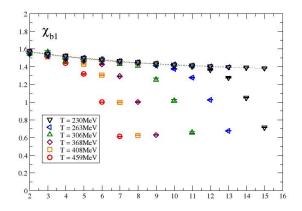
- Many b quarks will be produced at ALICE
- $T_d^{\Upsilon} \sim 5 T_c$  hard to do on the lattice
- $\chi_b$  melts at  $T_d^{\chi_b} \lesssim 1.2 T_c$ ?
- Use NRQCD and relativistic action, compare two approaches

### Results from relativistic beauty

- Used the same action as for charm (and light quarks)
- Used both point and derivative operators for P-waves



#### Operator dependence



Derivative operators better behaved — smaller constant mode?

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# Summary

- Charmonium S-waves survive to  $T \sim 1.6 2T_c$
- P-waves melt at  $T < 1.3 T_c$
- Significant momentum dependence in reconstructed correlators
- Transverse vector correlators are more sensitive to temperature and momentum
- Charm diffusion feasible from lattice simulations
- Relativistic beauty results compatible with NRQCD
- Simulations on finer lattices with realistic quark content underway