

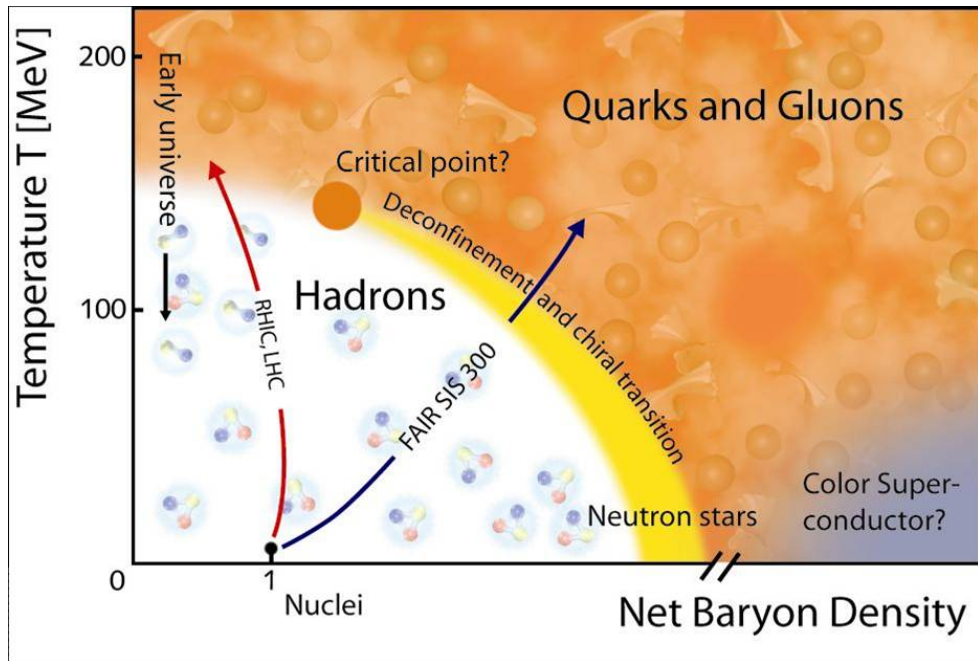
Solving the complex phase problem for QCD related systems with chemical potential

`christof.gattringer@uni-graz.at`

Ydalia Delgado Mercado, Hans Gerd Evertz, Christof Gattringer: [arXiv:1102.3096](https://arxiv.org/abs/1102.3096)

Motivation

Something we would like to understand in detail ...



Lattice QCD with chemical potential

- Grand canonical partition sum

$$Z = \int D[U] e^{-S_G[U]} \det D(\mu)^{N_f}$$

- For non-zero chemical potential μ the fermion determinant is complex

$$\det D(\mu) \in \mathbb{C}$$

and thus cannot be used as a weight factor in a MC calculation.

- For further progress in the analysis of the QCD phase diagram new approaches would be highly welcome.
- In this project we explore new strategies in simpler models.

An effective theory for QCD with chemical potential

Center symmetry and Svetitsky-Yaffe conjecture

- For pure gauge theory deconfinement can be understood as the spontaneous breaking of center symmetry. An influential idea for understanding this phase transition is the Svetitsky-Yaffe conjecture (1981):
- At T_c the critical behavior of $SU(3)$ gauge theory in $d + 1$ dimensions can be described by a d - dimensional spin system with a \mathbb{Z}_3 - invariant effective action for the local Polyakov loop $P(x) \in SU(3)$.
- Leading term of the effective action from a strong coupling expansion

$$S_c[P] = -\tau \sum_{\langle x,y \rangle} \left[\text{Tr}P(x) \text{Tr}P(y)^\dagger + \text{Tr}P(y) \text{Tr}P(x)^\dagger \right]$$

Center breaking terms from the fermion determinant

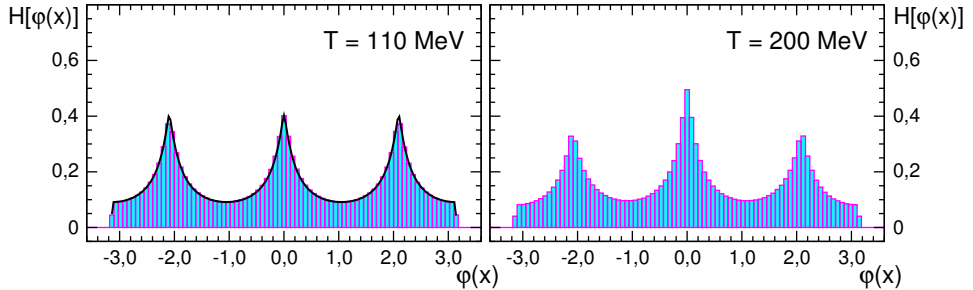
- The fermion determinant breaks the center symmetry explicitly.
- The leading center symmetry breaking can be calculated using hopping expansion, including the chemical potential μ .

$$S[P] = -\tau \sum_{\langle x,y \rangle} \left[\text{Tr}P(x) \text{Tr}P(y)^\dagger + \text{Tr}P(y) \text{Tr}P(x)^\dagger \right] \\ - \kappa \sum_x \left[e^\mu \text{Tr}P(x) + e^{-\mu} \text{Tr}P(x)^\dagger \right]$$

- For non-vanishing chemical potential μ the action is complex and the model has a complex phase problem.

Reduction to the center degrees of freedom

- The distribution of $\text{Tr}P(x)$ is dominated by the center.



- Reduction to the center elements: $\text{Tr}P(x) \rightarrow L(x) \in \{1, e^{\pm i2\pi/3}\}$

$$S = -\tau \sum_{\langle x,y \rangle} \left[L(x)L(y)^* + L(y)L(x)^* \right] - \kappa \sum_x \left[e^\mu L(x) + e^{-\mu} L(x)^* \right]$$

Contains leading center symmetric and center symmetry breaking terms.

Flux representation for the effective theory

Flux representation - I

- Identity for the nearest neighbor term:

$$e^{\tau[L(x)L(x+\hat{\nu})^*+c.c.]} = C \sum_{b_{x,\nu}=-1}^{+1} B^{|b_{x,\nu}|} (L(x)L(x+\hat{\nu}))^{b_{x,\nu}}$$

- Identity for the magnetic term:

$$e^{\kappa[e^{\mu} L(x) + e^{-\mu} L(x)^*]} = \sum_{s_x=-1}^{+1} M_{s_x} L(x)^{s_x}$$

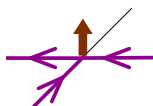
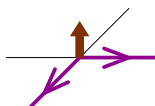
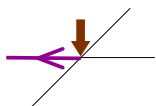
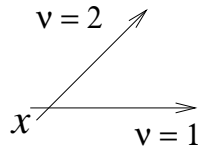
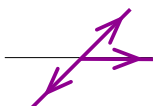
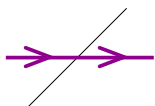
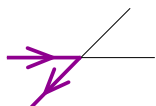
- C, B and M_s are real and positive functions of τ, κ and μ .
- Remaining sums over the center elements at each site:

$$\sum_{L \in \{1, e^{\pm i2\pi/3}\}} L^n = 3T(n) \quad \text{with} \quad T(n) = \delta_{0, n \bmod 3}$$

Flux representation - II

- Representation with dimers $b_{x,\nu} \in \{-1, 0, 1\}$ and monomers $s_x \in \{-1, 0, 1\}$

$$Z = \sum_{\{b,s\}} W(b, s) \prod_x T \left(\sum_{\nu} [b_{x,\nu} - b_{x-\hat{\nu},\nu}] + s_x \right) \quad \text{with} \quad W(b, s) > 0$$

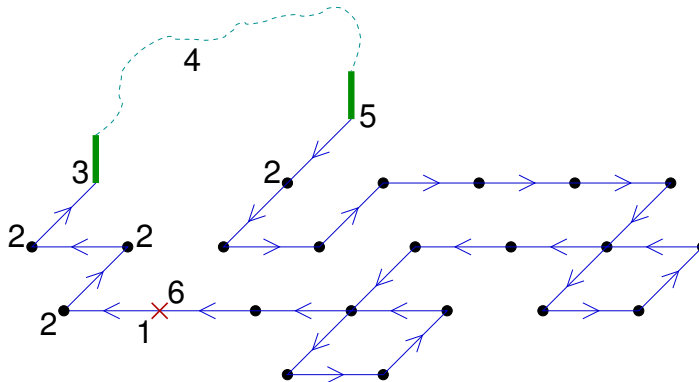


Monomers: $\downarrow \uparrow$

Dimers: $\rightleftarrows \quad \nearrow \searrow$

Update with a Prokof'ev – Svistunov worm algorithm

- We start the worm at some point (Position 1).
- The worm may decide to change dimers (Pos. 2) or insert monomers (3).
- Insertion of a monomer is followed by a random hop (4) and another monomer insertion (5).
- The worm closes when it reaches its starting point (6).

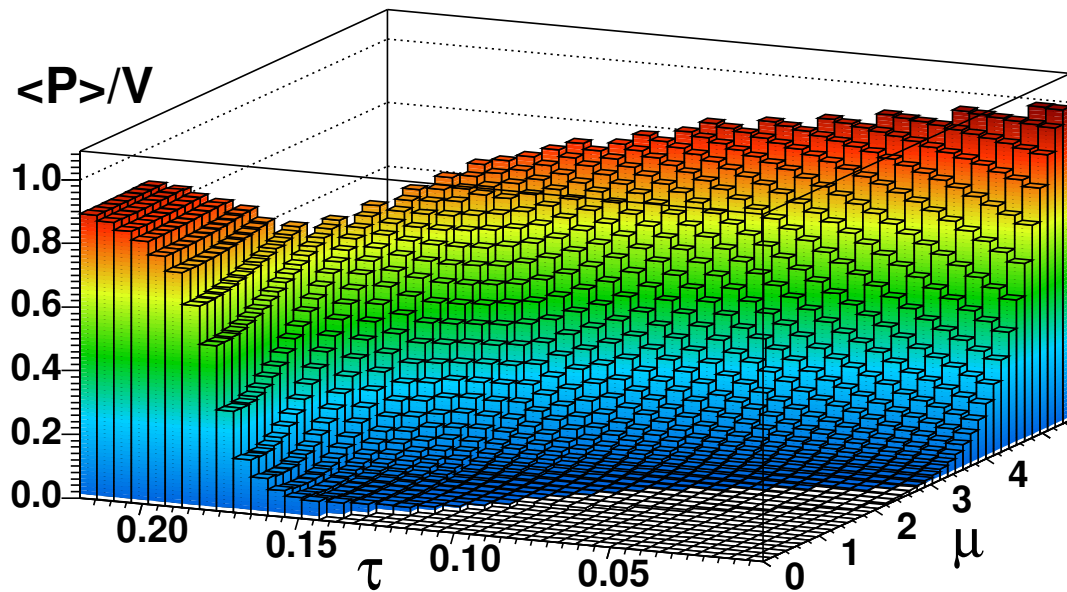


Numerical analysis

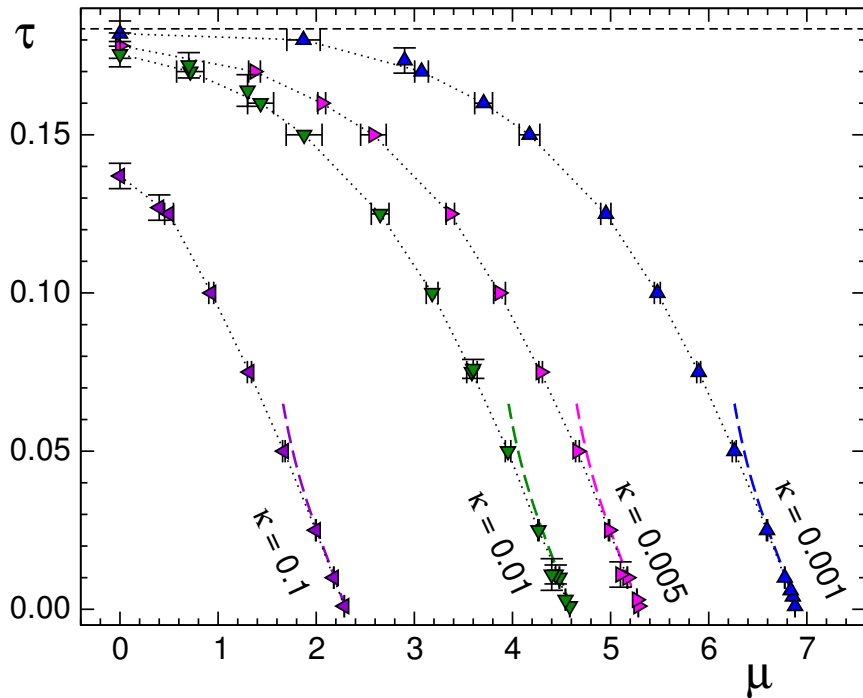
- The worm algorithm is ergodic and efficient also for large κ and μ .
- As observables we consider the internal energy E , the heat capacity C , the expectation value of the Polyakov loop $\langle P \rangle$ and the corresponding susceptibility χ_P .
- All our observables can be expressed in terms of the dimer- and monomer occupation numbers and their fluctuations.
- We work on lattice with sizes 16^3 to 72^3 with statistics of 10^5 to 10^6 worms.
- Comparison of the results for small τ with low temperature expansion.

Results

Order parameter

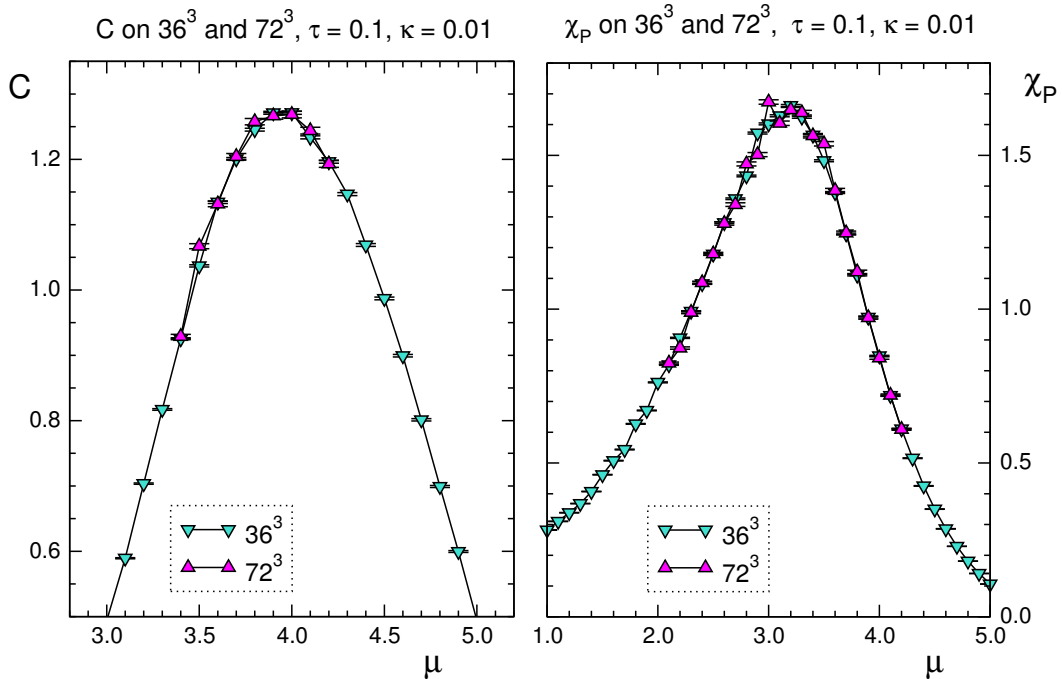


Exploring the phase diagram

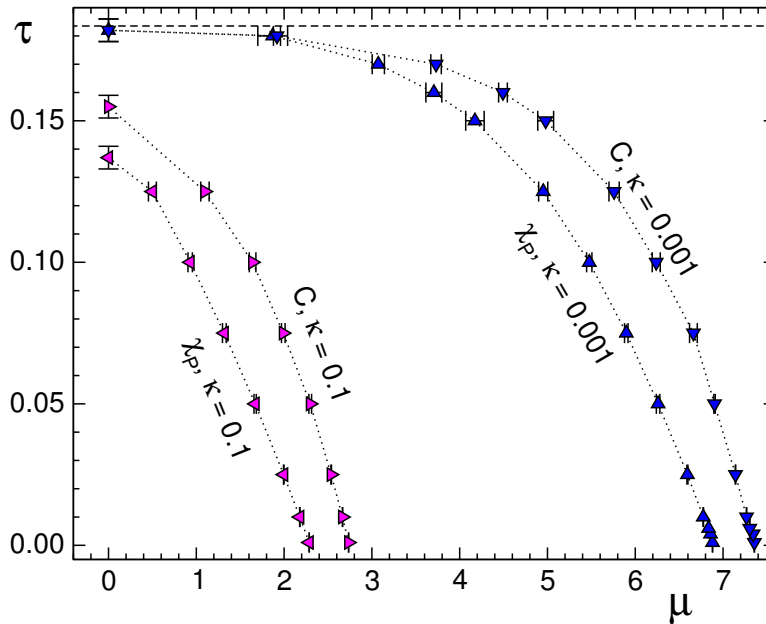


Volume scaling

For all values of κ we studied, no volume scaling of the second derivatives of the free energy C and χ_P was observed \Rightarrow Crossover.



Strips of crossover behavior



If QCD is reduced to center degrees of freedom \Rightarrow Crossover transitions.

Developments and perspectives

More realistic models:

- Action for effective theory with SU(3)-valued Polykov loops:

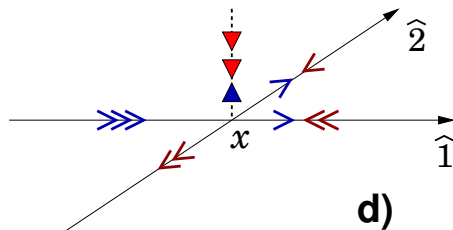
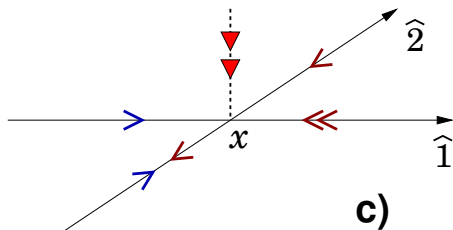
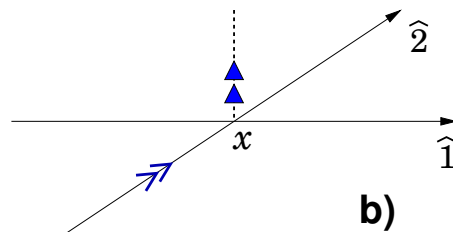
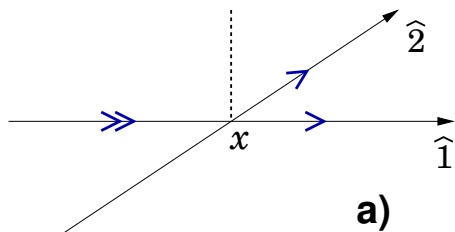
$$\begin{aligned} S[P] = & -\tau \sum_{\langle x,y \rangle} \left[\text{Tr}P(x) \text{Tr}P(y)^\dagger + \text{Tr}P(y) \text{Tr}P(x)^\dagger \right] \\ & -\kappa \sum_x \left[e^\mu \text{Tr}P(x) + e^{-\mu} \text{Tr}P(x)^\dagger \right] \end{aligned}$$

- Flux representation:

$$Z = \sum_{\{l, \bar{l}, s, \bar{s}\}} \mathcal{W}[l, \bar{l}, s, \bar{s}] \prod_x T \left(\sum_\nu [(l_{x,\nu} - \bar{l}_{x,\nu}) - (l_{x-\hat{\nu},\nu} - \bar{l}_{x-\hat{\nu},\nu})] + (s_x - \bar{s}_x) \right)$$

$$\mathcal{W}[l, \bar{l}, s, \bar{s}] \text{ real and positive ; } l_{x,\nu}, \bar{l}_{x,\nu}, s_x, \bar{s}_x \in \mathbb{N}_0$$

Examples of admissible flux and monomer vertices



Generalized PS worm algorithms can be applied for a Monte Carlo calculation.

Summary

- We study QCD reduced to the leading center symmetric and center symmetry breaking terms.
- The complex phase problem is solved by using a flux representation with dimers and monomers, which is suitable for a PS worm algorithm.
- We map out the phase diagram using E , $\langle P \rangle$, C and χ_P .
- If QCD is reduced to the center degrees of freedom only crossover transitions remain.
- Work in progress: Generalizations to theories closer to QCD.
- The new models can and should be used for testing various approaches to finite μ (expansions, complex Langevin, reweighting ...).