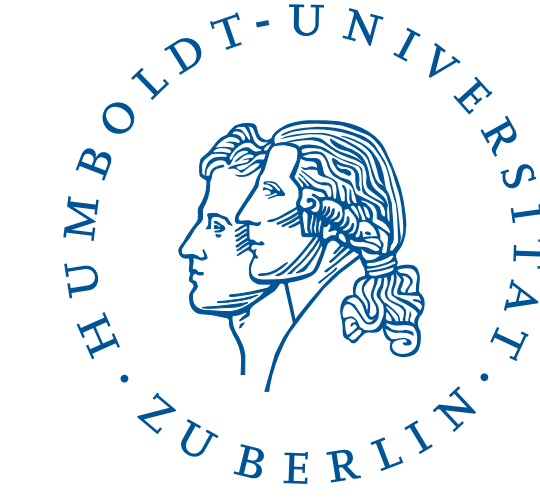
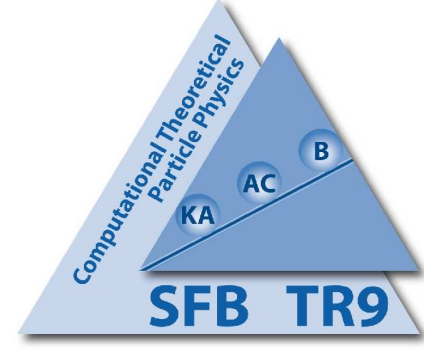


Landau gauge gluon and ghost propagators at non-zero temperature on the lattice.



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Motivation and aims

- Compute Landau gauge gluon and ghost propagators in pure $SU(3)$ gauge theory within lattice discretization.
- Study of systematic effects: lattice artifacts, Gribov copies, finite-size effects.
- Main aims:
 - provide continuum limit results within some momentum range,
 - to see, in as far gluon and ghost propagators can be used as “order parameters” for the deconfinement transition or crossover.
- Results to be compared with outcome from Dyson-Schwinger (DS) and Functional Renormalization Group (FRG) equations and to be used for tuning truncations of the systems of equations.

Setup for the $SU(3)$ lattice simulation

- Use standard Wilson plaquette action:

$$S = \beta \sum_x \sum_{\mu > \nu} \left[1 - \frac{1}{3} \Re \text{Tr} \left(U_{x\mu} U_{x+\mu;\nu} U_{x+\nu;\mu}^\dagger U_{x\nu}^\dagger \right) \right], \quad \beta = 6/g_0^2.$$

- Landau gauge fixing by maximizing the gauge functional

$$F_U[g] = \frac{1}{3} \sum_x \sum_{\mu} \Re \text{Tr} \left(g_x U_{x\mu} g_{x+\mu}^\dagger \right) \quad \text{w. r. to } g_x \in SU(3)$$

using *simulated annealing* + *overrelaxation*.

Compare random (“first”) Gribov copy (*fc*) with “best” copy (*bc*) taking global $\mathbf{Z}(3)$ -flips betw. Polyakov loop sectors into account.

- Choice of parameters:

Fix $\beta = 6.337$ corresponding to $T \equiv 1/aL_4 \simeq T_c$ for $L_4 = 12$ [1].

Fix the 3-volume $V = (L_s a)^3 = (48 a)^3 \simeq (2.7 \text{ fm})^3$.

Vary T with $L_4 = 4, 6, \dots, 18$.

Renormalize propagators at $\mu = 5 \text{ GeV}$ or 3 GeV .

Scaling, Gribov copy and finite volume tests at $T = 0.86 T_c$ and $T = 1.20 T_c$.

Lattice observables for $T > 0$:

- Lattice gauge potential

$$A_\mu(x + \hat{\mu}/2) = \frac{1}{2ia g_0} (U_{x\mu} - U_{x\mu}^\dagger) |_{\text{tracless}}.$$

- Transverse gluon propagator in momentum space

$$D_T = \frac{1}{(d-2)8} \left\langle \sum_{i=1}^3 A_i^a(q) A_i^a(-q) - \frac{q_4^2}{q^2} A_4^a(q) A_4^a(-q) \right\rangle.$$

- Longitudinal gluon propagator

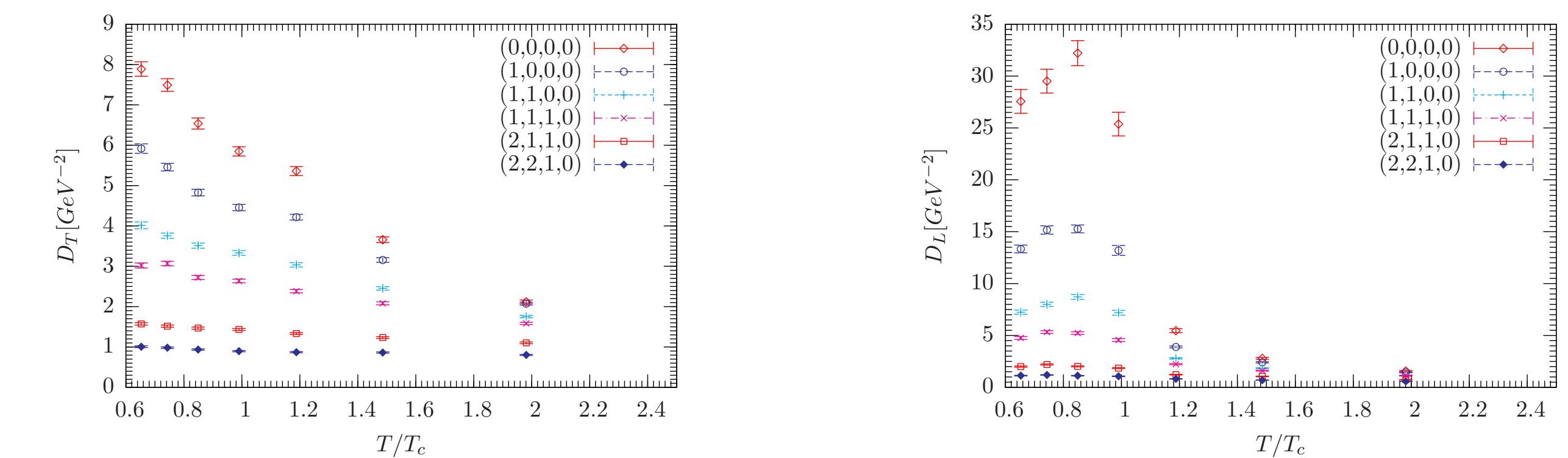
$$D_L = \frac{1}{8} \left(1 + \frac{q_4^2}{q^2} \right) \langle A_4^a(q) A_4^a(-q) \rangle.$$

- Ghost propagator from Faddeev-Popov operator $M = -\partial D$

$$G^{ab}(q) = a^2 \sum_{x,y} \langle e^{-2\pi i k \cdot (x-y)/L} [M^{-1}]_{xy}^{ab} \rangle = \delta^{ab} G(q),$$

- Momenta: $q_\mu(k_\mu) = \frac{2}{a} \sin\left(\frac{\pi k_\mu}{L_\mu}\right)$, $k_\mu \in (-L_\mu/2, L_\mu/2]$. Here put $q_4 = 0$.

Result: gluon propagators vs. T at lowest q



⇒ Transverse propagator depends smoothly on T .

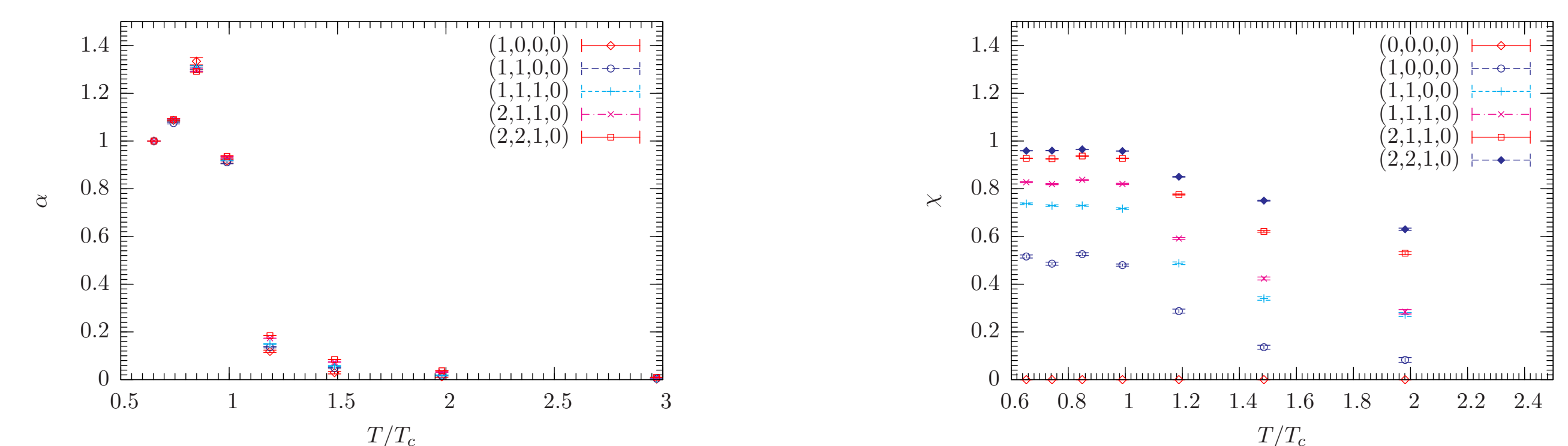
⇒ Long. propagator at low momenta has a characteristic behaviour around T_c .

⇒ Ghost propagator almost independent of T (not shown here).

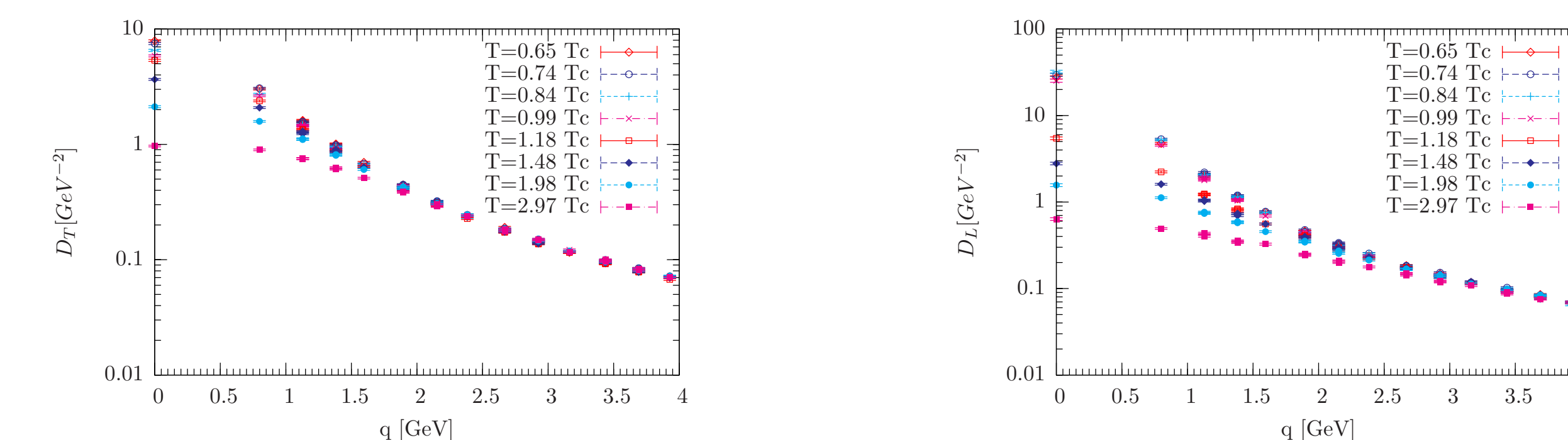
“Order parameters” from longitud. propagator

$$\alpha = \frac{D_L(0, T) - D_L(q, T)}{D_L(0, \frac{2}{3}T_c) - D_L(q, \frac{2}{3}T_c)}$$

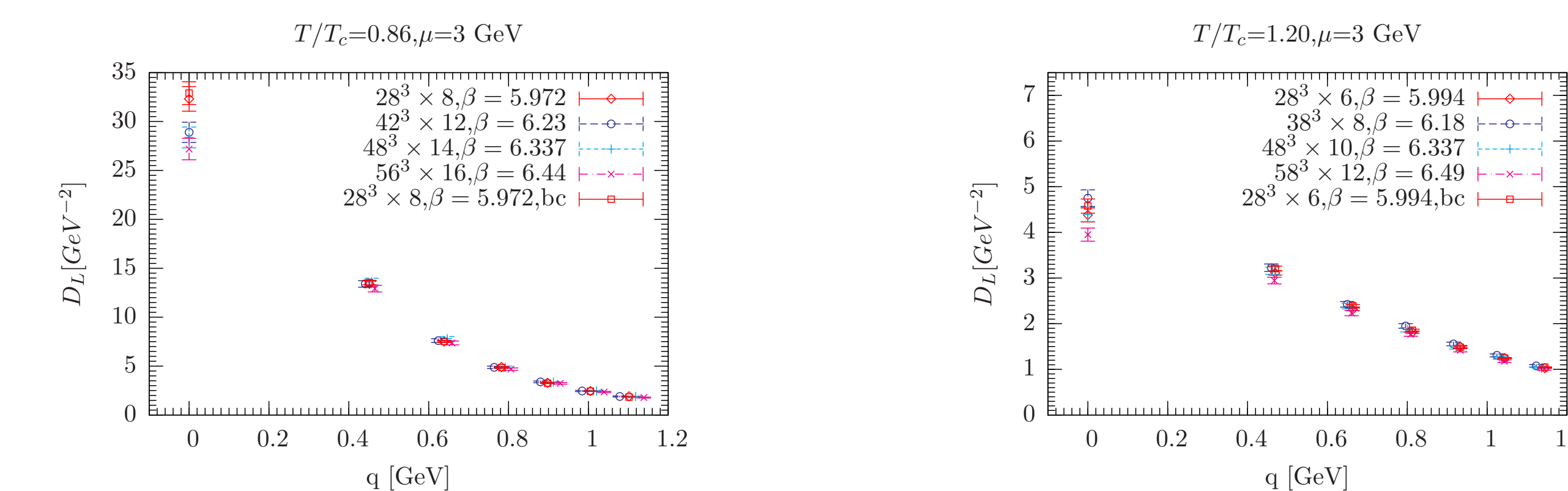
$$\chi = [D_L(0, T) - D_L(q, T)] / D_L(0, T)$$



Result: gluon propagators for various T



Scaling and Gribov copy test (here for D_L)



⇒ Scaling violation, Gribov copy effects visible at small q .

⇒ Effects for transversal gluon propagator even stronger.

Conclusions

- Longitudinal Landau gauge gluon propagator provides “order parameter(s)” for finite-temperature transition.
- Lattice artifacts, Gribov copy effects reasonably small at $q > 0.8 \text{ GeV}$.
- Finite-volume effects have been studied up to $(64a)^3 \simeq (3.6 \text{ fm})^3$. Negligible effect above $\simeq 0.5 \text{ GeV}$.

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