QCD phase transitions and quark spectral functions Christian S. Fischer, Jan Lücker and Jens Müller JLU Gießen

JUSTUS-LIEBIG-UNIVERSITÄT GIESSEN

Abstract

We investigate the QCD phase diagram by solving the Dyson-Schwinger equations of QCD at finite temperature and chemical potential. From the quark propagator we extract order parameters for the chiral and the deconfinement transition. Moreover we investigate the propagation of quarks in the quark-gluon plasma by analyzing its spectral functions.

Order parameters 1: chiral symmetry breaking

We extract the quark condensate from the quark propagator:

$$egin{aligned} &\langlear\psi\psi
angle &= Z_2 N_c T \sum_n \int rac{d^3 p}{(2\pi)^3} ext{Tr}_D \left[S(ec p^2, \omega_n)
ight] \end{aligned}$$

The truncation scheme



For the quark-gluon vertex we use a phenomenological *ansatz*. For the gluon propagator we

large for broken chiral symmetry
 small for (approximately) restored chiral symmetry

Order parameters 2: deconfinement

For confinement/deconfinement we use dual condensates [7, 8, 9].

$$\Sigma_n = \int rac{darphi}{2\pi} e^{-iarphi n} \langle ar{\psi}\psi
angle_arphi,$$

where $\langle \bar{\psi}\psi \rangle_{\varphi}$ is the condensate with generalised, U(1)-valued quark boundary conditions $\psi(\vec{x}, 1/T) = e^{i\varphi}\psi(\vec{x}, 0)$. The quark loop is evaluated at $\varphi = \pi$ to break Roberge-Weiss symmetry.

• Σ_{+1} , the dressed Polyakov loop equals the ordinary Polyakov loop for $m o\infty$ • Σ_{-1} , the conjugated dressed Polyakov loop, differs from Σ_{+1} for $\mu>0$

► use lattice results for the quenched propagator (yellow dot in gluon DSE) and

ullet add the quark loop to account for unquenching effects $(N_f=2)$

Behaviour of the order parameters



Dressed Polyakov loop can be calculated from DSEs [4] and FRGs [5].
Physical quark masses → both transitions are crossovers at µ = 0
Both transitions in the same temperature regime, pseudo-critical temperatures may differ

Ansatz for quark spectral function

 $egin{aligned}
ho_{\pm}(\omega,|p|) &= 2\piig[Z_1\delta(\omega\mp E_1)+Z_2\delta(\omega\pm E_2)ig]\ &+ ext{ continuum part from Landau damping} \end{aligned}$

Results 1: chiral transition



• The pseudocritical temperature at $\mu = 0$ is $T_c \approx 183 MeV$ • At quark chemical potential $\mu \approx 280 MeV$ the crossover turns into a first order phase transition, indicating a critical endpoint

Dispersion relations at $\mu = 0$ (quenched)



- In the confined phase, the condensate is nearly independent of the boundary condition
- In the deconfined phase a "valley" developes
- \blacktriangleright For larger values of μ the valley becomes steeper, developing a first order phase transition in φ

Results 3: the phase diagram



Literature

- [1] J. A. Mueller, C. S. Fischer, D. Nickel, Eur. Phys. J. **C70** (2010) 1037-1049. [arXiv:1009.3762 [hep-ph]].
- [2] C. S. Fischer, A. Maas and J. A. Muller, Eur. Phys. J. C **68** (2010) 165 [arXiv:1003.1960

- E₁: quark, E₂: plasmino
 quantitative agreement with lattice results [10]
- ${\scriptstyle \blacktriangleright}\, p \rightarrow 0$ agree with analytical results from HTL
- ${\scriptstyle
 ho}\, p > T$ plasmino dies out
- dispersion relations stay timelike due to continuum part



► The deconfinement temperatures for quarks and antiquarks differ for moderate $\mu \rightarrow$ similar to PQM studies, Ref. [6]

At densities close to the critical endpoint deconfinement and chiral transition coincide

 So far, no sign of a quarkyonic phase
 At large µ baryonic contributions play a role → better truncation scheme needed [hep-ph]].

[3] C. S. Fischer, J. A. Mueller, Phys. Rev. D80 (2009) 074029. [arXiv:0908.0007 [hep-ph]].
[4] C. S. Fischer, Phys. Rev. Lett. 103 (2009) 052003. [arXiv:0904.2700 [hep-ph]].
[5] J. Braun, L. M. Haas, F. Marhauser, J. M. Pawlowski, Phys. Rev. Lett. 106 (2011) 022002. [arXiv:0908.0008 [hep-ph]].

[6] T. K. Herbst, J. M. Pawlowski and B. J. Schaefer, Phys. Lett. B 696 (2011) 58 [arXiv:1008.0081 [hep-ph]].

[7] C. Gattringer, Phys. Rev. Lett. **97** (2006) 032003. [hep-lat/0605018].
[8] F. Synatschke, A. Wipf, C. Wozar, Phys. Rev. **D75** (2007) 114003. [hep-lat/0703018].

[9] E. Bilgici, F. Bruckmann, C. Gattringer, C. Hagen, Phys. Rev. D77 (2008) 094007. [arXiv:0801.4051 [hep-lat]].

[10] F. Karsch, M. Kitazawa, Phys. Lett. **B658** (2007) 45-49. [arXiv:0708.0299 [hep-lat]].

