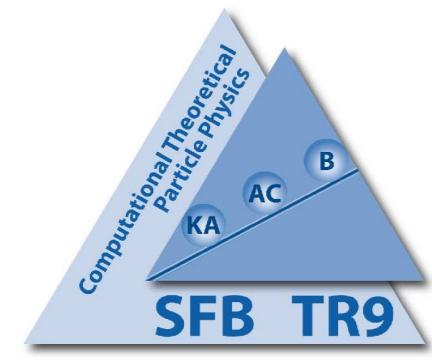
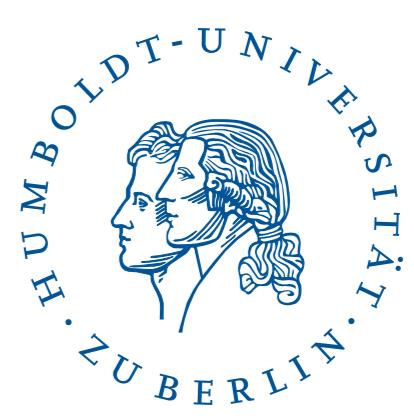


Towards the equation of state with $N_f = 2$ twisted-mass lattice QCD



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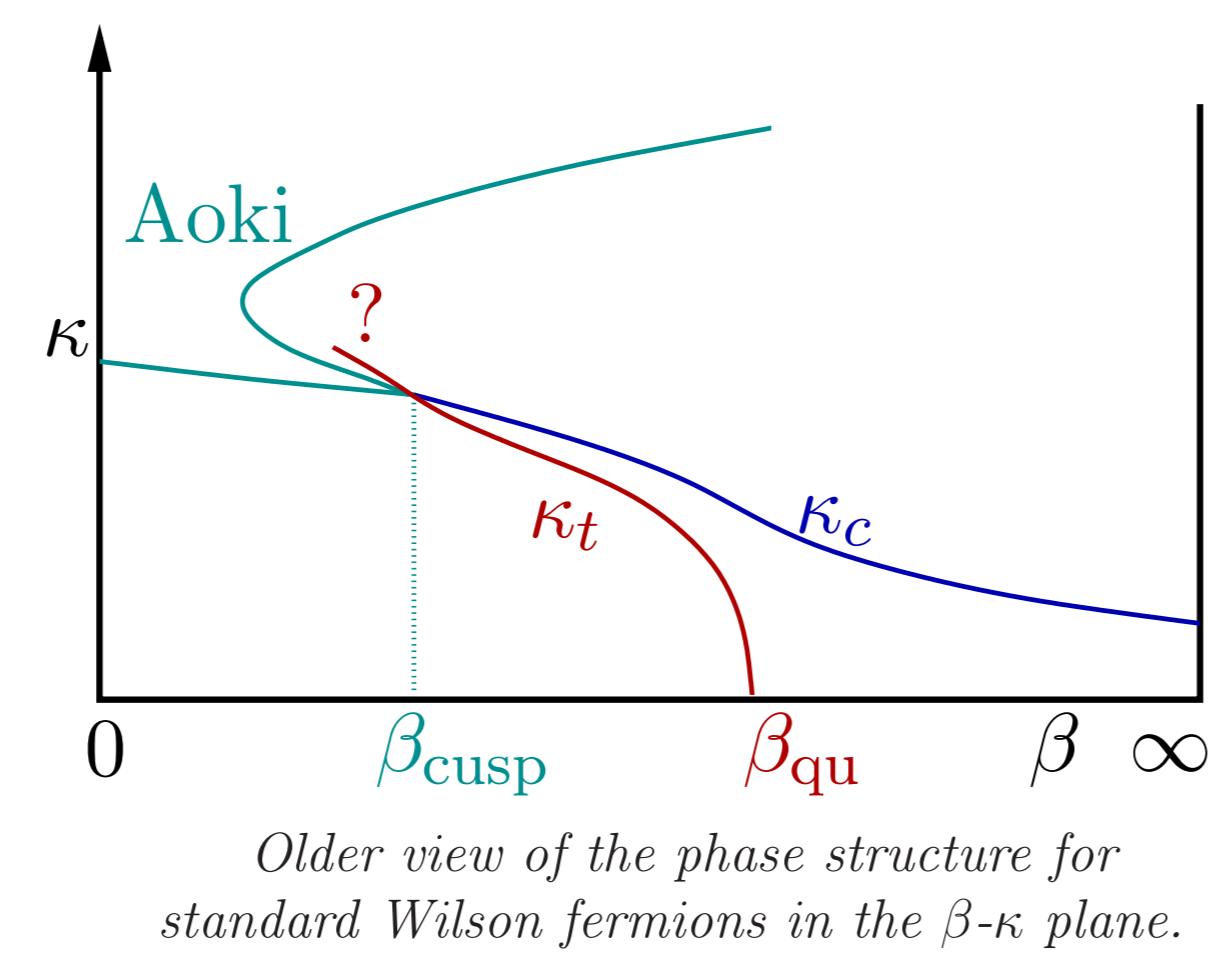
General view, the phase diagram (see [1])

- Lattice QCD with $N_f = 2$ Wilson fermions modified with twisted mass term $i\mu_0\bar{\psi}\gamma_5\tau^3\psi$ and tree-level Symanzik improved gauge action is studied.

$$S_f[U, \psi, \bar{\psi}] = \sum_x \bar{\chi}(x) \left(1 - \kappa H[U] + 2i\kappa\mu_0\gamma_5\tau^3 \right) \chi(x) \quad (1)$$

$$\psi = \frac{1}{\sqrt{2}}(1 + i\gamma_5\tau^3)\chi \quad \text{and} \quad \bar{\psi} = \bar{\chi}\frac{1}{\sqrt{2}}(1 + i\gamma_5\tau^3)$$

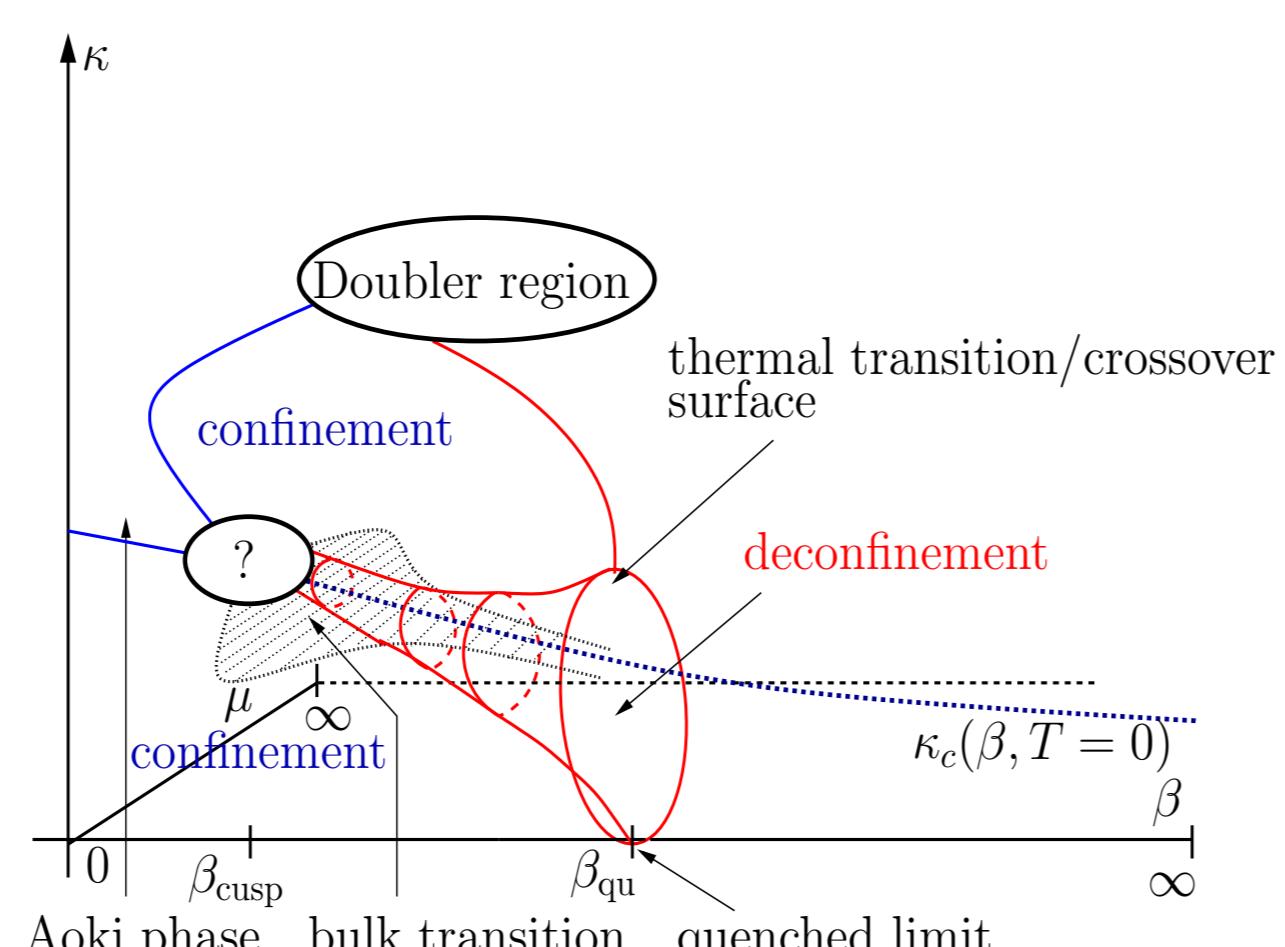
$$S_g^{\text{tlSym}}[U] = \beta \left(c_0 \sum_P \left[1 - \frac{1}{3} \text{ReTr}(U_P) \right] + c_1 \sum_R \left[1 - \frac{1}{3} \text{ReTr}(U_R) \right] \right) \quad (2)$$



- Hopping parameter κ and twisted mass μ_0 are connected with the bare quark mass:

$$m_q = \sqrt{\frac{1}{4} \left(\frac{1}{\kappa} - \frac{1}{\kappa_c} \right)^2 + \mu_0^2} \quad (3)$$

- For maximal twist, i.e. $\kappa = \kappa_c(\beta; T=0)$, automatic $\mathcal{O}(a)$ improvement is expected. See e.g. [3, 4].
- In order to explore the phase diagram HMC simulations were performed in a wide range $\beta \approx 6/g_0^2 = 1.80, \dots, 3.90$ on linear lattice sizes $N_s = 16, 24$ and $N_\tau = 8$.
- Phase diagram for $\mu_0 = 0$ and fixed N_τ divides into three regions:
 - the **Aoki-Phase** [5] for strong coupling, i.e. at small β ,
 - a **bulk 1st order transition region** [6, 8] at intermediate β ,
 - and the **thermal transition** and scaling region at larger β .
- Acc. to Eq. (3) for $\mu_0 \neq 0$ the thermal transition line is part of a conical surface ($\kappa_t(\beta), \mu_t(\beta)$).



Present schematic view [1] of the phase diagram for tmQCD in β - κ_{μ_0} space, originally proposed by Creutz [7].

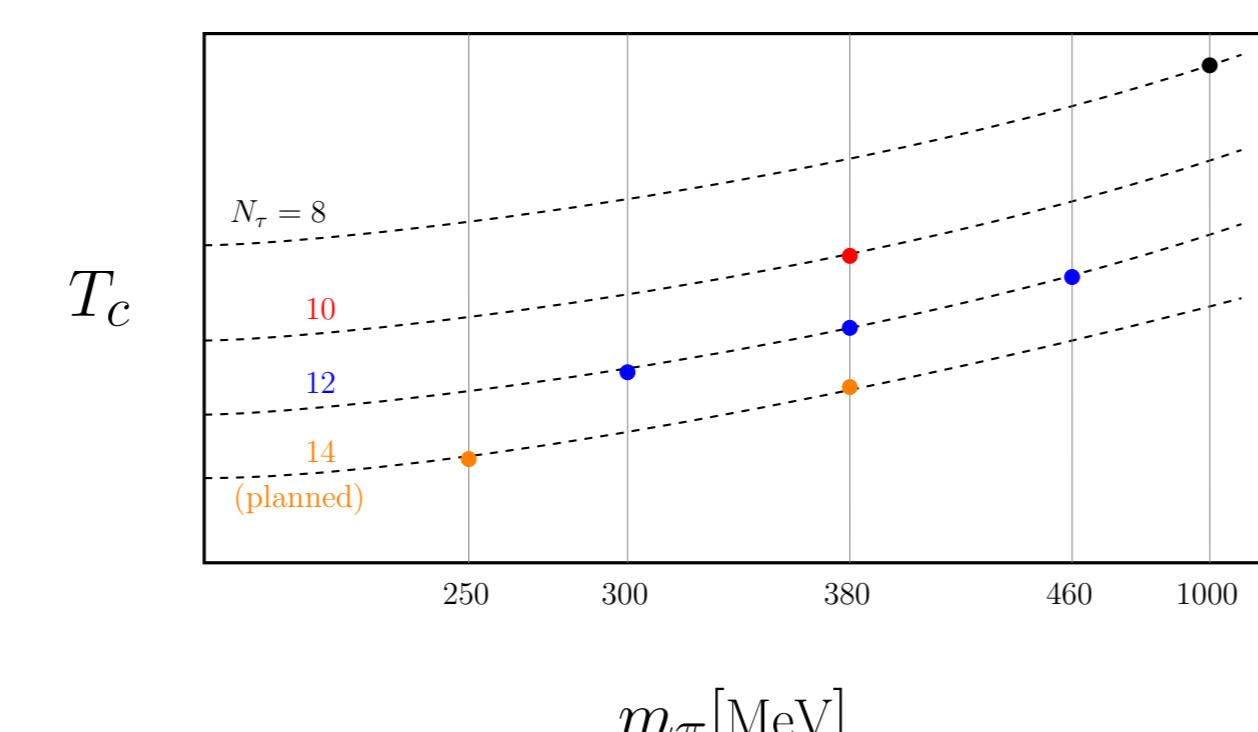
Setup for locating T_c (see [2])

- Scans in β at maximal twist.
- Keep pion mass m_π fixed while varying $T = 1/(N_\tau a(\beta))$.
- $a(\beta), r_0/a(\beta), \kappa_c(\beta)$ and $\mu_0(m_\pi, \beta)$ from ETMC data [9] (interpolations necessary).

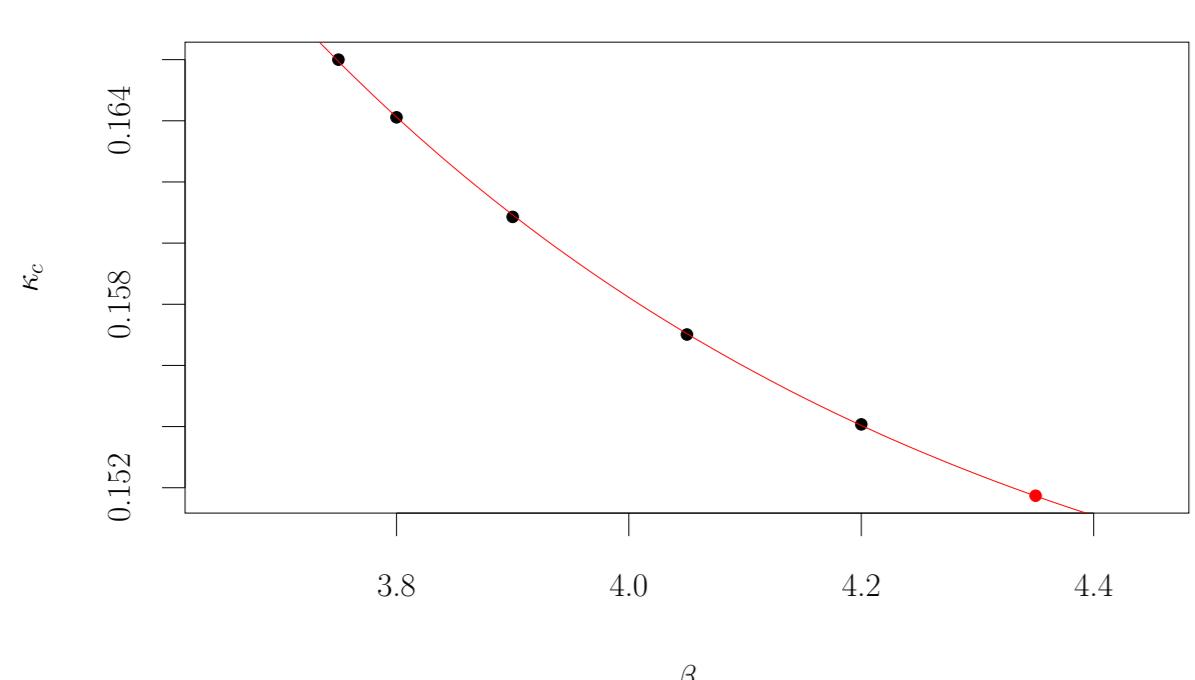
Typical Statistics:

- $N_\tau = 10$: O(4k) measurements per β -value
- $N_\tau = 12$: O(3k-10k) measurements per β -value
- Have scanned T -range $\approx [0.9 \dots 1.05] T_c$ For EoS: → wider range needed: $T \in [0.8 \dots 1.7] T_c$ (partly available at $m_\pi = 400$ MeV)
- Need also further $T = 0$ points (preliminary data available at $\beta = 4.35$, only one mass yet)

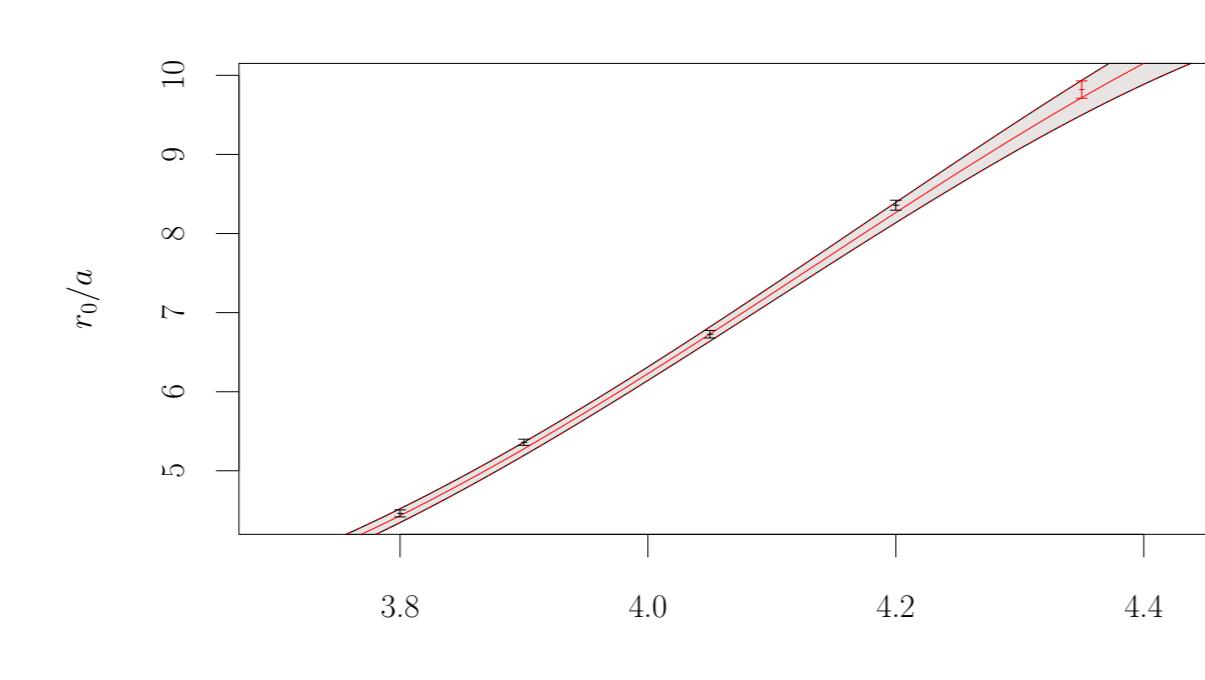
Simulated Pion Masses:



Interpolation of κ_c :



Interpolation of r_0/a :



Locating the thermal transition or crossover (cf. talk by L. Zeidlewicz)

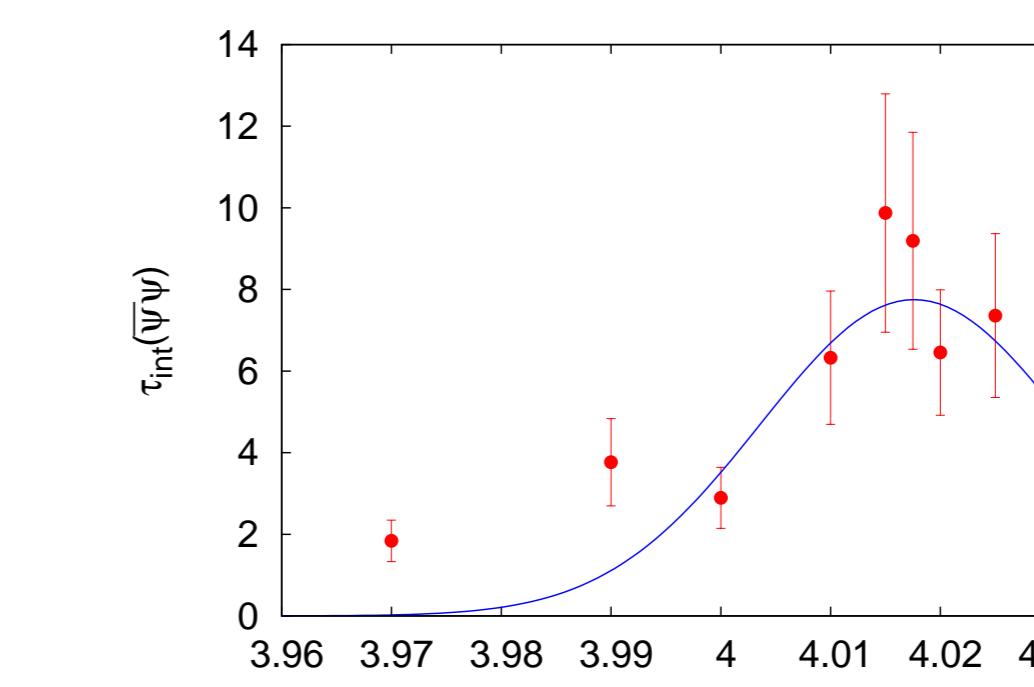
Observables

Polyakov-Loop, Plaquette, $\langle \bar{\psi}\psi \rangle$, Pionnorm, their variances and integrated autocorrelation times τ_{int} . Mostly weak crossover signals as expected in intermediate mass range.

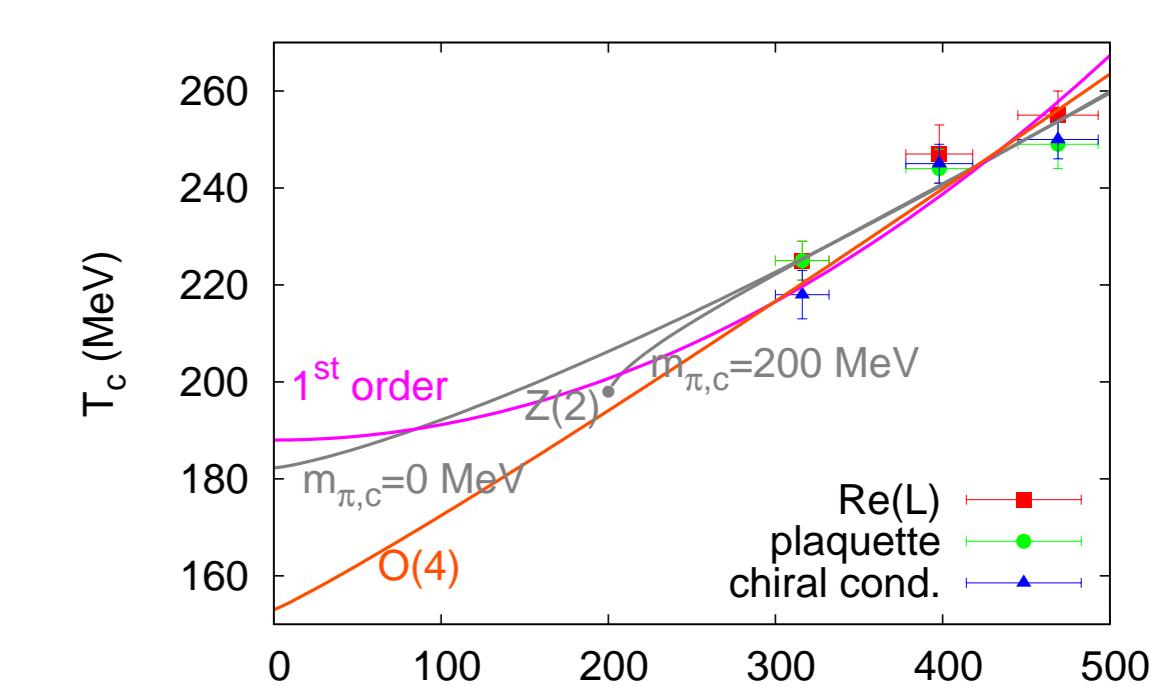
Result

$m_\pi [\text{MeV}]$	≈ 320	≈ 400	≈ 470	≈ 400
N_τ	12	12	12	10
$\beta_c(\langle \bar{\psi}\psi \rangle)$	3.9400(96)	4.0150(11)	4.0298(11)	3.8804(04)
$T_c [\text{MeV}]$	218(5)	245(4)	250(4)	239(7)

Example: τ_{int} for $\langle \bar{\psi}\psi \rangle$:



T_c versus pion mass m_π



EoS: trace anomaly

- Use the integral method (see e.g. [10]) to get $p(T)$ and $\epsilon(T)$ from trace anomaly

$$\begin{aligned} \frac{\epsilon - 3p}{T^4} &= -\frac{T}{V} \left\langle \frac{d \ln Z}{d \ln a} \right\rangle_{\text{sub}} \\ &= \left(a \frac{d\beta}{da} \right) \left(c_0 \langle \text{ReTr} U_P \rangle_{\text{sub}} + c_1 \langle \text{ReTr} U_R \rangle_{\text{sub}} \right. \\ &\quad \left. + \underbrace{\frac{\partial \kappa_c}{\partial \beta} \langle H \rangle_{\text{sub}}}_{\approx 0 \text{ neglected so far}} + \left(2a\mu_0 \frac{\partial \kappa_c}{\partial \beta} + \frac{\partial(a\mu_0)}{\partial \beta} \right) \langle \bar{\psi}\psi \rangle_{\text{sub}} \right) \end{aligned}$$

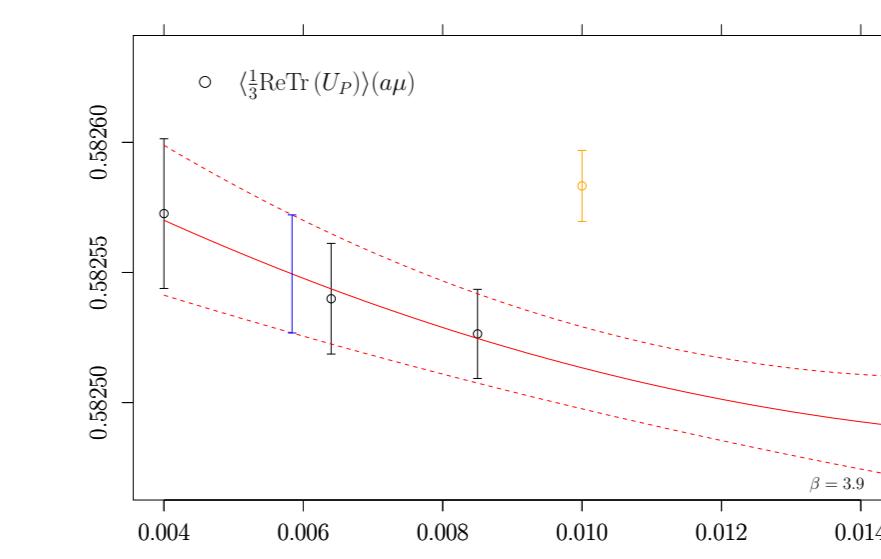
where:

$$\langle \dots \rangle_{\text{sub}} \equiv \langle \dots \rangle_{T>0} - \langle \dots \rangle_{T=0}$$

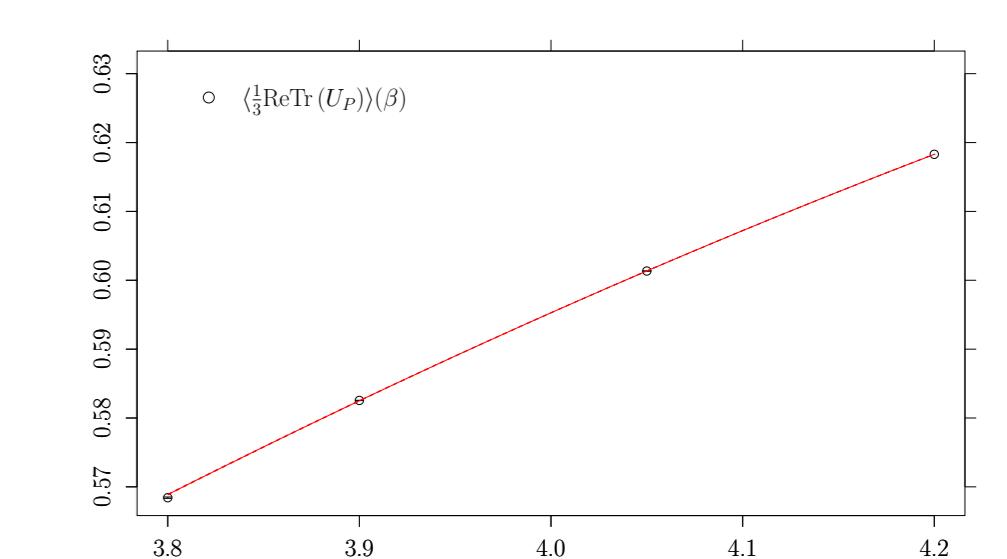
- Have to use interpolation for $(\frac{r_0}{a}) (\beta)$ to get:

$$\left(a \frac{d\beta}{da} \right) = - \left(\frac{r_0}{a} \right) \left(\frac{d(\frac{r_0}{a})}{d\beta} \right)^{-1}$$

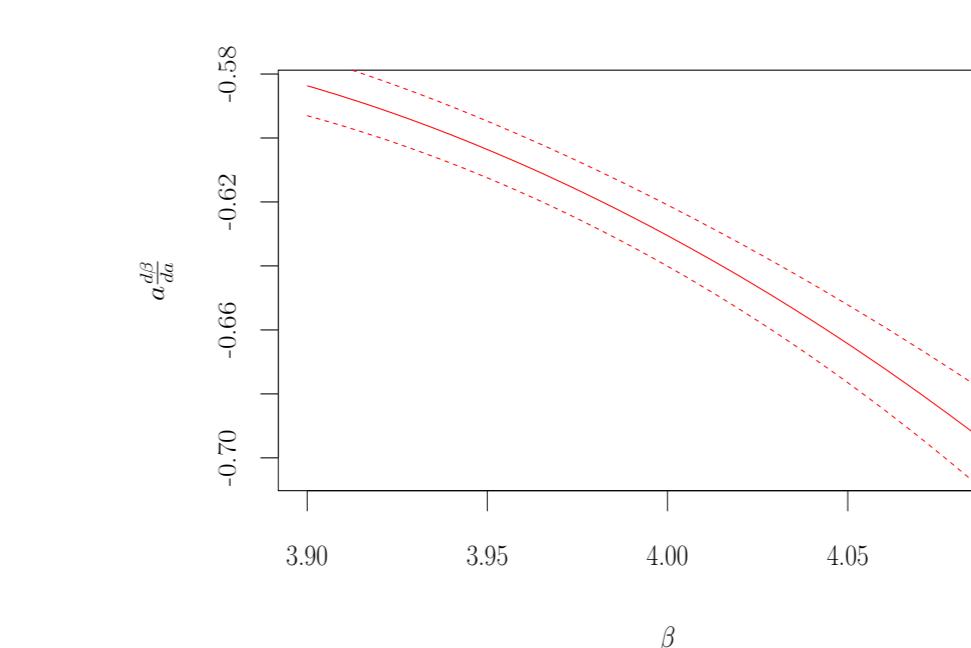
Plaquette interpolation in $a\mu_0$:



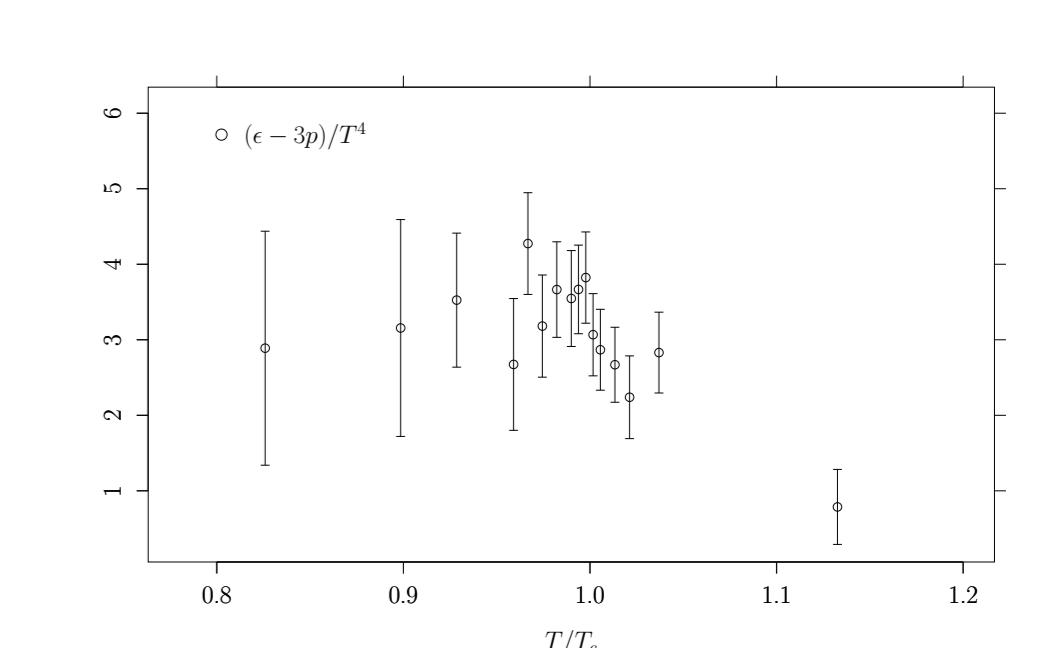
Plaquette interpolation in β :



Interpolation of beta-function:



Trace anomaly (preliminary):



Conclusions

- The finite temperature transition has been located down to $m_\pi = O(300\text{MeV})$. Our results at $N_\tau = 10, 12$ support a weak crossover behaviour.
- The chiral limit is consistent with $O(4)$ universality, but does not exclude other scenarios.
- Do not see a splitting of chiral and deconfinement transitions.
- First (and preliminary) results for the trace anomaly presented.

Acknowledgement:

We appreciate support by the ETM Collaboration. Large-scale computations were carried out at INFN and HLRN Berlin, Hannover. We thank Elena Garcia Ramos for providing new data for $T = 0$ prior to publication. F.B. and M.M.P. acknowledge support by DFG via GK 1504 and SFB/TR 9.

References

- [1] E.-M. Ilgenfritz, K. Jansen, M.P. Lombardo, M. Müller-Preussker, M. Petschlies, O. Philipsen, L. Zeidlewicz [tmfT Collaboration], Phys. Rev. D80 (2009) 094502, arXiv:0905.3112 [hep-lat].
- [2] F. Burger, E.-M. Ilgenfritz, M. Kirchner, M.-P. Lombardo, M. Müller-Preussker, O. Philipsen, C. Urbach, L. Zeidlewicz [tmfT Collaboration], arXiv:1102.4530 [hep-lat].
- [3] F. Farchioni *et al.*, PoS LAT2005 (2006) 072, arXiv: hep-lat/0509131.
- [4] C. Urbach [ETM Collaboration], PoS LAT2007 (2007) 022, arXiv:0710.1517 [hep-lat].
- [5] S. Aoki, Phys. Rev. D30 (1984) 2653.
- [6] G. Münster, JHEP 0409, 035 (2004), arXiv: hep-lat/0407006.
- [7] M. Creutz, Phys. Rev. D76 (2007) 054501, arXiv:0706.1207 [hep-lat].
- [8] S.R. Sharpe, J.M.S. Wu, Phys. Rev. D71 (2005) 074501, arXiv: hep-lat/0411021.
- [9] ETM Collaboration (R. Baron *et al.*), arXiv:0911.5061 [hep-lat].
- [10] C. DeTar, U. M. Heller, Eur. Phys. J. A 41 (2009) 405, arXiv:0905.2949 [hep-lat].