NUCLEAR CHIRAL THERMODYNAMICS and PHASES of QCD

Wolfram Weise
Technische Universität München







QCD interface with nuclear physics:

Chiral Effective Field Theory

- Nuclear Equation of State and QCD phase diagram
- Density and temperature dependence of the

Chiral (Quark) Condensate

Outlook:

New constraints from **Neutron Stars**



Part I: Prelude

QCD PHASE DIAGRAM

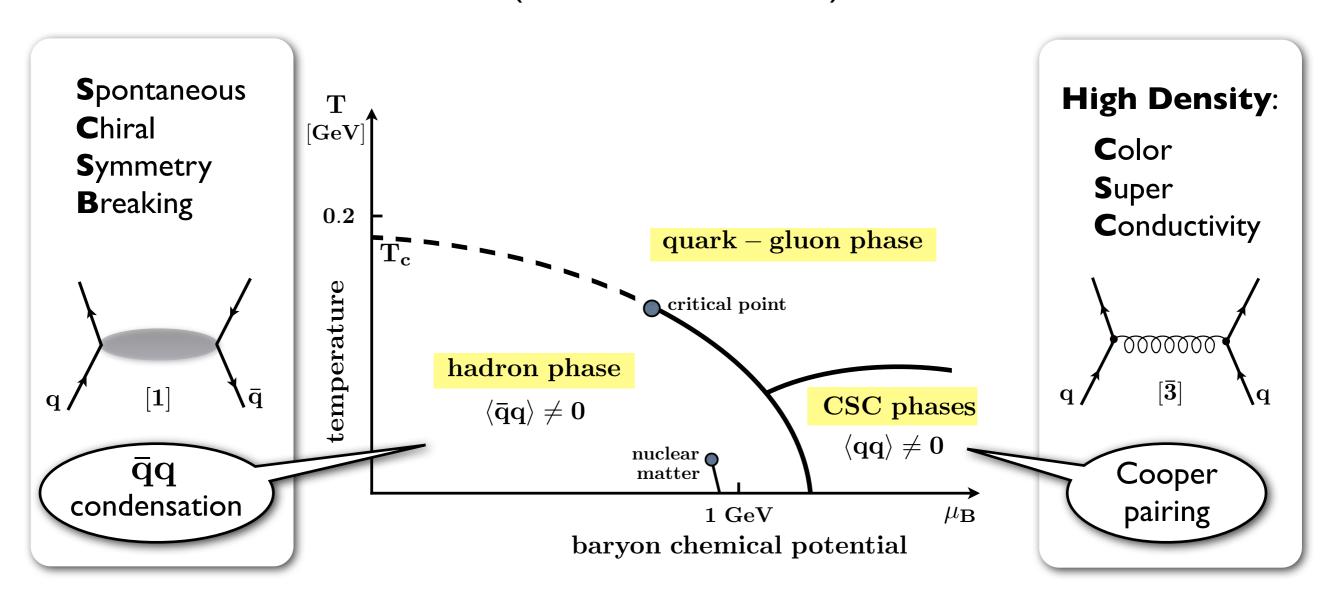
Visions & Facts





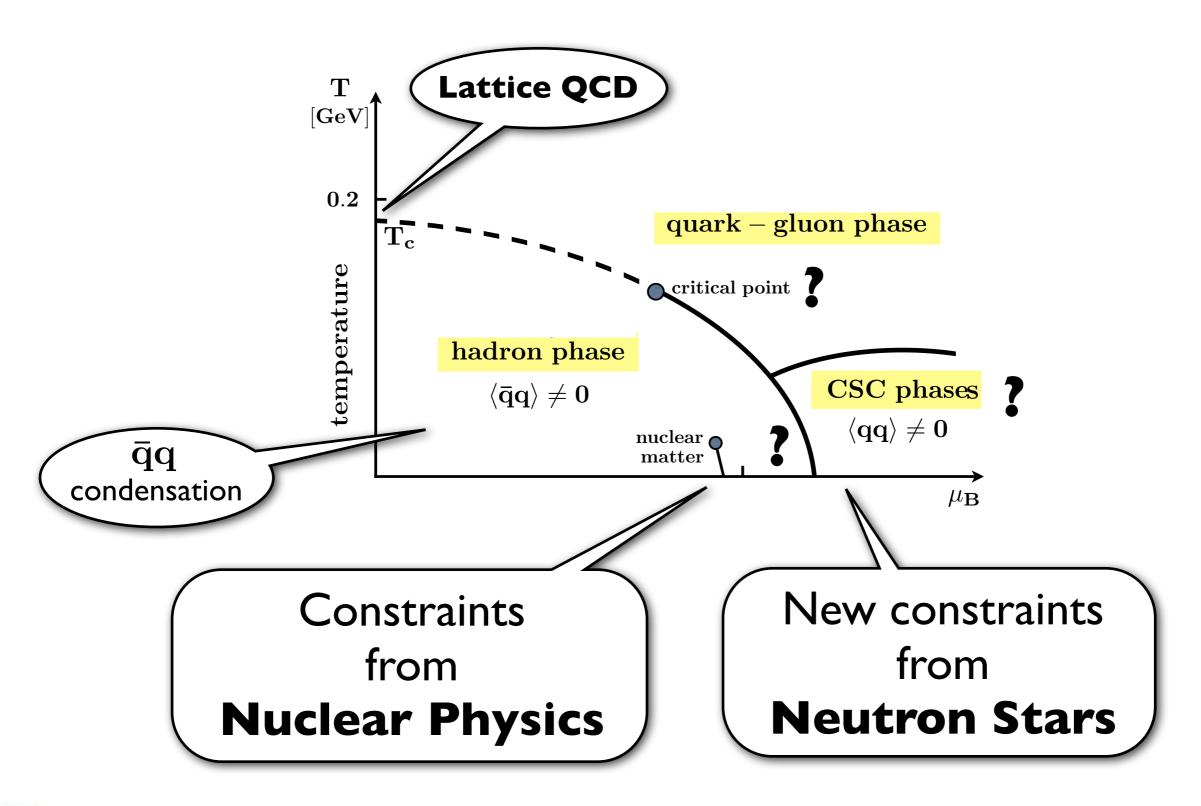
QCD PHASE DIAGRAM

(theorists' vision)





QCD PHASE DIAGRAM





MODELING the QCD PHASE DIAGRAM

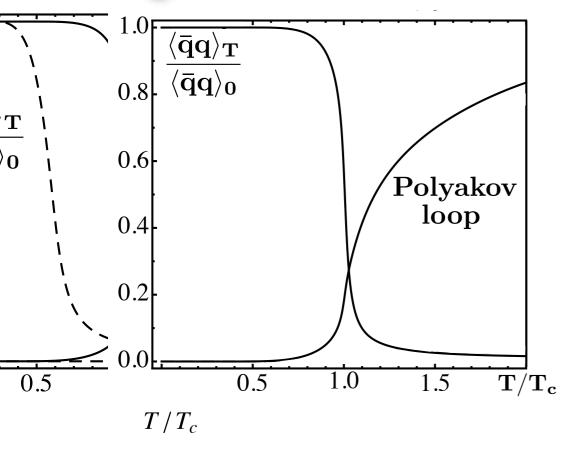


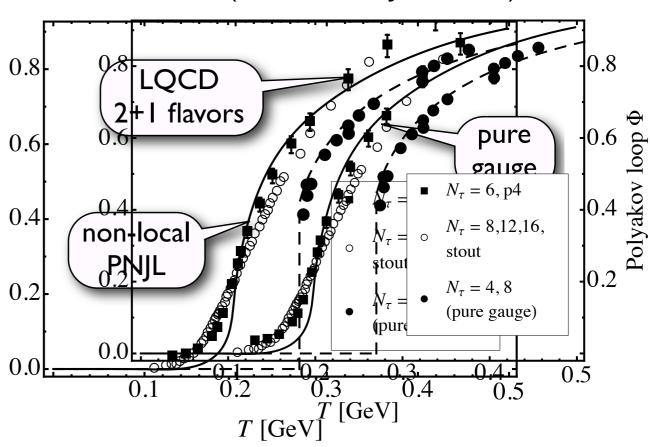
Guiding principle:

QCD symmetries and symmetry breaking patterns

 $\begin{array}{c|c} \textbf{Spontaneously broken} \\ \textbf{chiral symmetry} \\ \textbf{SU}(N_f)_R \times \textbf{SU}(N_f)_L \end{array} \begin{array}{c} \textbf{non-local} \\ \textbf{PNJL} \\ \textbf{SU}(3)_c \text{ gauge group} \\ \textbf{model} \end{array}$

chiral and deconfinement crossover transitions (3 flavor PNJL model)





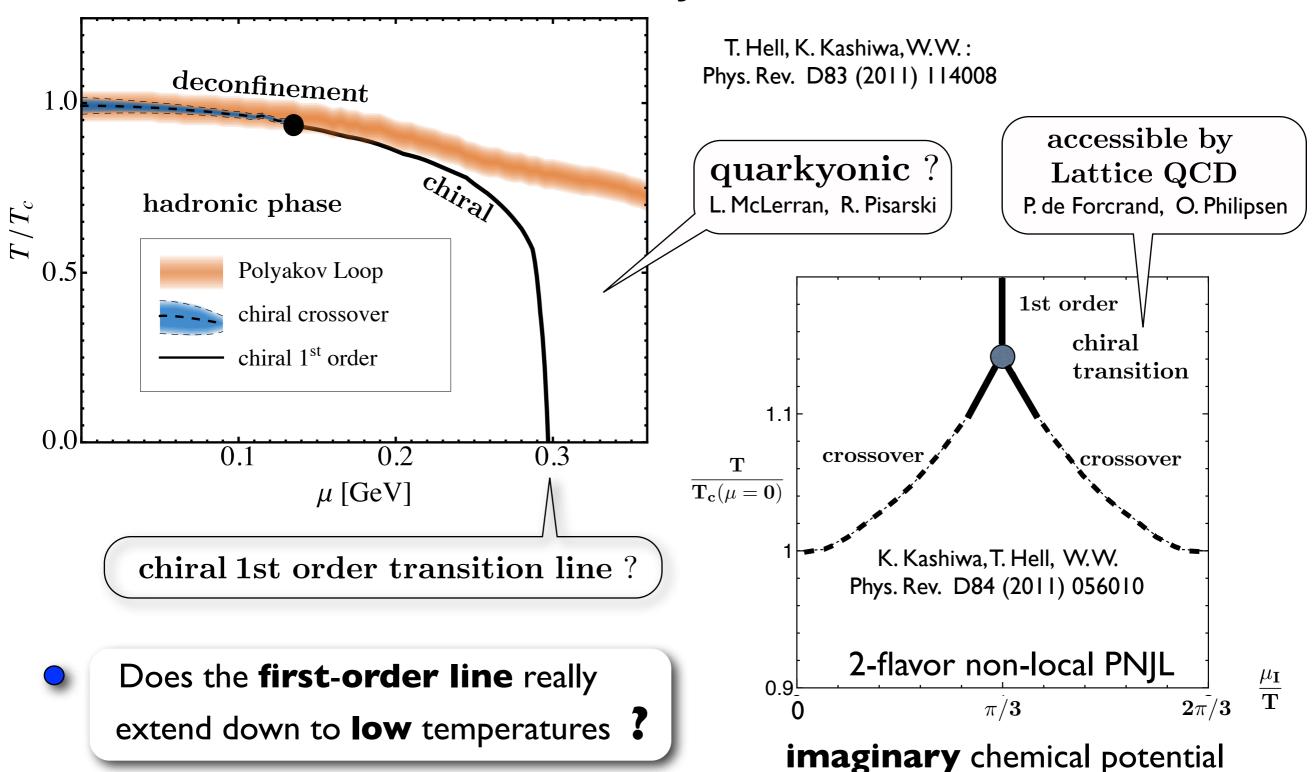
T. Hell, S. Rössner, M. Cristoforetti, W.W.: Phys. Rev. D81 (2010) 074034

lated temperature dependence of the chiral condensate \bar{q}_c in the lated temperature dependence of the chiral condensate \bar{q}_c is high left. W.W.: Phys. Rev. D83 (2011) 114008 (right) normalized to the transition temperature $T_c = 205 \, \mathrm{MeV}$ as

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PHASE DIAGRAM

Non-local 3-flavor PNJL model calculation





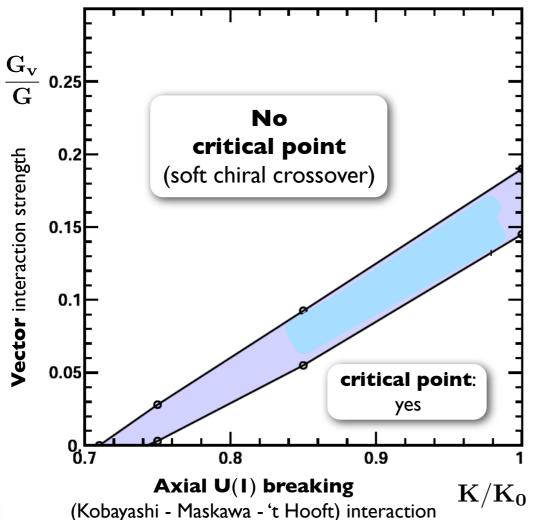


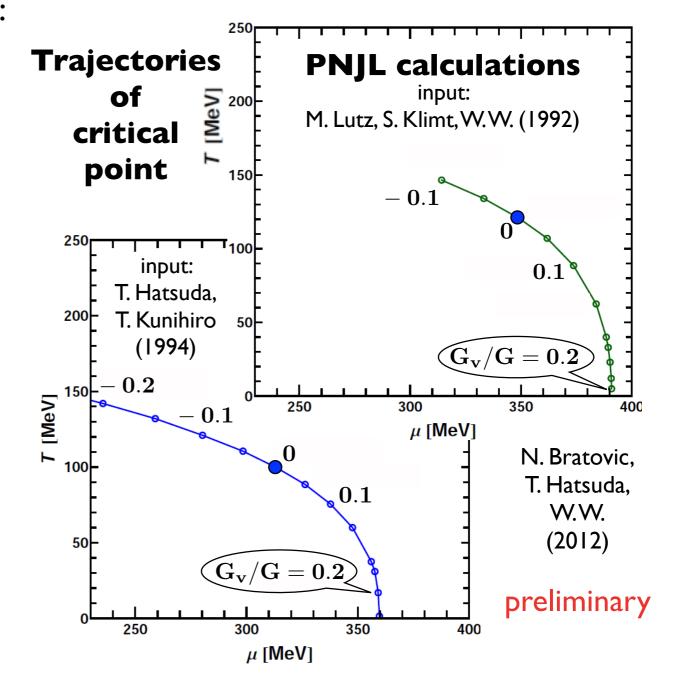
PHASE DIAGRAM (contd.)

PNJL analysis of Lattice QCD phase diagram at imaginary chemical potential suggests significant isoscalar vector term in the effective quark-quark interaction

$$\delta \mathcal{L}_V = -G_V \left(\bar{\psi} \gamma_\mu \psi \right) (\bar{\psi} \gamma^\mu \psi)$$

- Existence and location of critical point: extremely sensitive to
 - a) Strength of vector interaction
 - b) Axial U(1) breaking interaction





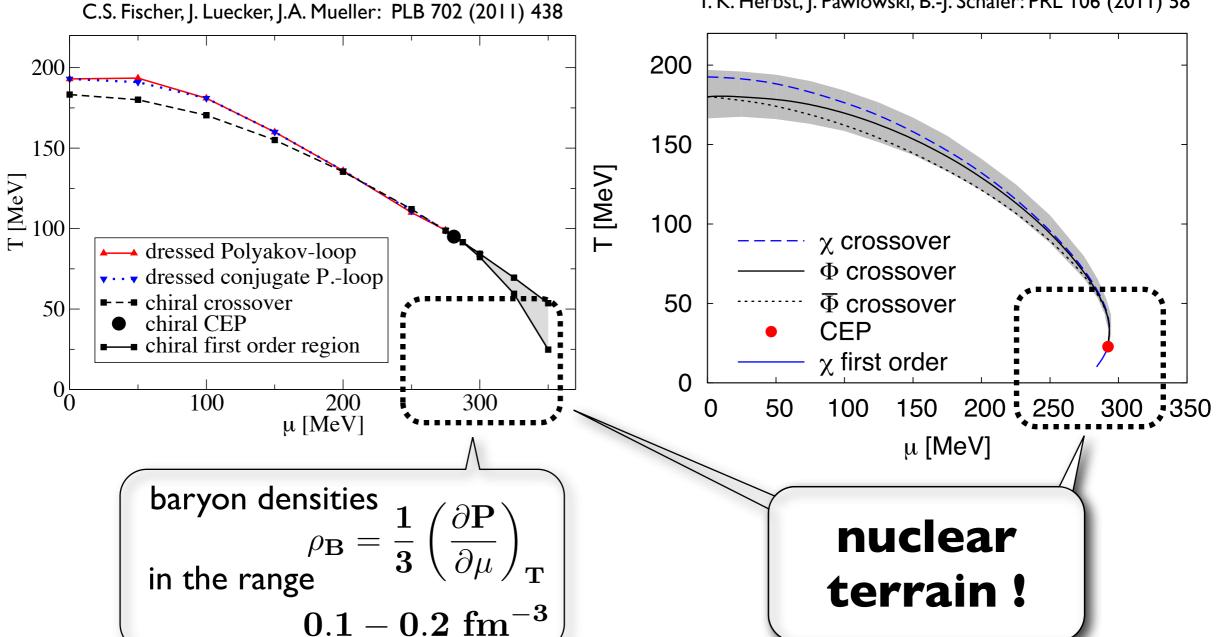


PHASE DIAGRAM (contd.)

Dyson - Schwinger QCD

Polyakov - Quark-Meson model

T. K. Herbst, J. Pawlowski, B.-J. Schäfer: PRL 106 (2011) 58



Quarks are not the relevant active quasiparticles at low temperatures and at baryon chemical potentials $\mu_{\rm B} \lesssim 1~{
m GeV}$





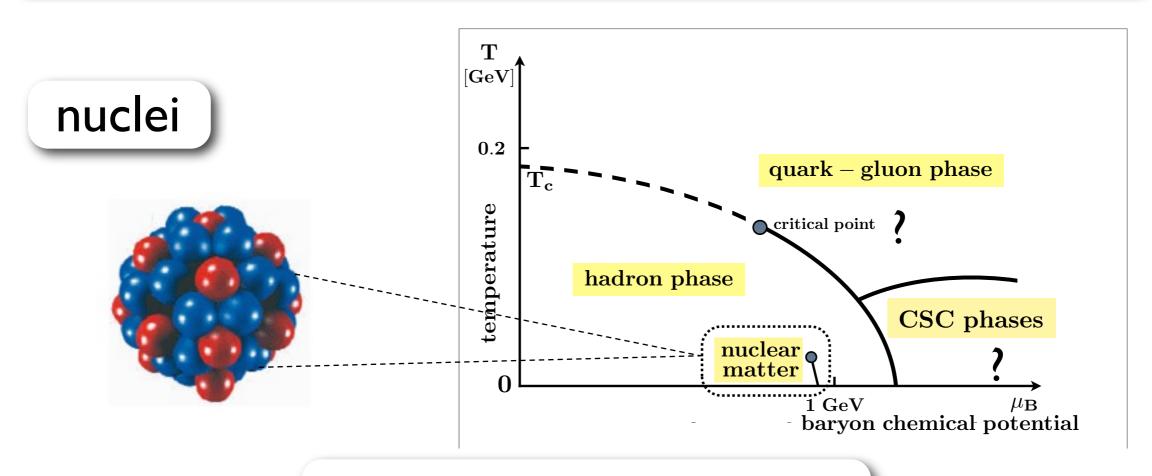
Part II:

NUCLEAR CHIRAL THERMODYNAMICS





NUCLEAR MATTER and QCD PHASES



Scales in nuclear matter:

momentum scale:

Fermi momentum

- NN distance:
- energy per nucleon:
- compression modulus:

$$m k_F \simeq 1.4~fm^{-1} \sim 2m_\pi$$

$$d_{
m NN} \simeq 1.8~{
m fm} \simeq 1.3~{
m m}_\pi^{-1}$$

$$\mathbf{E}/\mathbf{A} \simeq -\mathbf{16} \ \mathbf{MeV}$$

$$\mathbf{K} = (\mathbf{260} \pm \mathbf{30}) \; \mathbf{MeV} {\sim} \; \mathbf{2m_{\pi}}$$





PIONS and NUCLEI in the context of LOW-ENERGY QCD

- CONFINEMENT of quarks and gluons in hadrons
- Spontaneously broken CHIRAL SYMMETRY

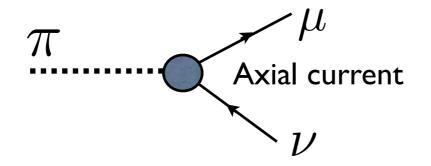
LOW-ENERGY / LOW-TEMPERATURE QCD:

Effective Field Theory of weakly interacting

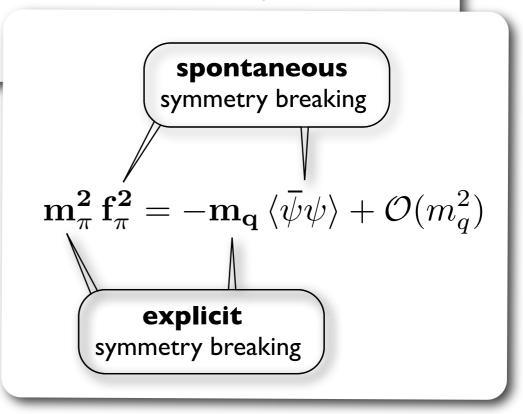
Nambu-Goldstone Bosons (PIONS)

representing QCD at (energy and momentum) scales

$$\mathbf{Q} << \mathbf{4}\pi\,\mathbf{f}_\pi \sim \,\mathbf{1}\,\mathbf{GeV}$$



$$\mathbf{f}_{\pi} = \mathbf{92.4\,MeV}$$





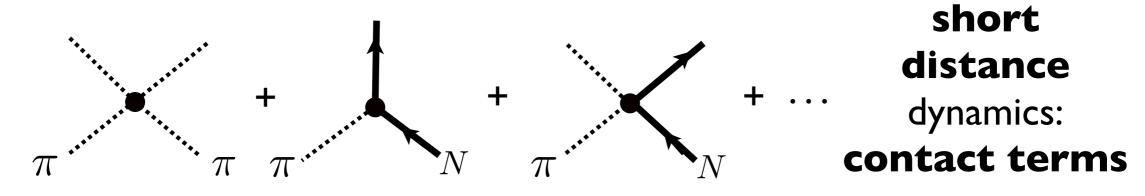


CHIRAL EFFECTIVE FIELD THEORY

- Systematic framework at interface of QCD and Nuclear Physics
- Interacting systems of PIONS (light / fast) and NUCLEONS (heavy / slow):

$$\mathcal{L}_{eff} = \mathcal{L}_{\pi}(U, \partial U) + \mathcal{L}_{N}(\Psi_{N}, U, ...)$$
$$U(x) = exp[i\tau_{a}\pi_{a}(x)/f_{\pi}]$$

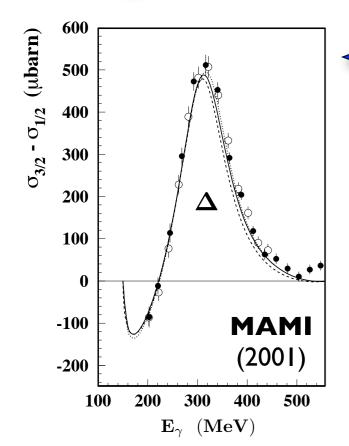
Construction of Effective Lagrangian: Symmetries





Explicit $\Delta(1230)$ DEGREES of FREEDOM

Large spin-isospin polarizabilty of the Nucleon



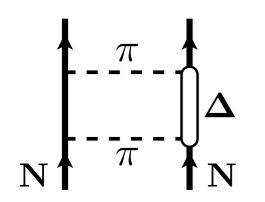
example: polarized Compton scattering

$$eta_\Delta = rac{g_A^2}{f_\pi^2(M_\Delta-M_N)} \sim 5\, ext{fm}^3 \ M_\Delta - M_N \simeq 2\,\, m_\pi << 4\pi\, f_\pi$$
 (small scale)

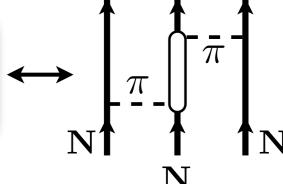


N. Kaiser, S. Gerstendörfer, W.W., NPA637 (1998) 395

N. Kaiser, S. Fritsch, W.W., NPA750 (2005) 259



$$egin{aligned} \mathbf{\Delta} & \mathbf{V_c(r)} = -rac{9\,\mathbf{g_A^2}}{32\pi^2\,\mathbf{f_\pi^2}}\,eta_{\mathbf{\Delta}}\,rac{\mathbf{e^{-2m_\pi r}}}{\mathbf{r^6}}\,\mathbf{P(m_\pi r)} \end{aligned} egin{aligned} lacksquare$$



 $1\,\mathrm{fm}$

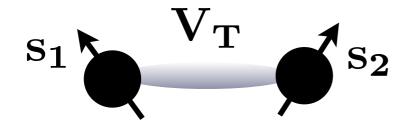
strong 3-body interaction

J. Fujita, H. Miyazawa (1957) Pieper, Pandharipande, Wiringa, Carlson (2001)

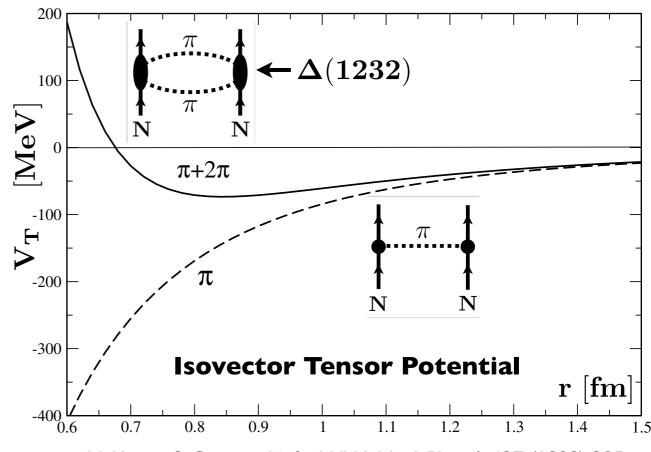


Important pieces of the CHIRAL NUCLEON-NUCLEON INTERACTION

ISOVECTORTENSOR FORCE

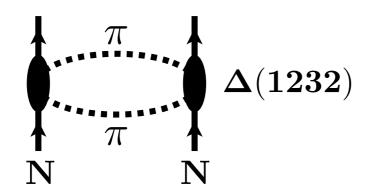


o note: **no** ρ meson



N. Kaiser, S. Gerstendörfer, W.W.: Nucl. Phys. A 637 (1998) 395

CENTRAL ATTRACTION from TWO-PION EXCHANGE



lacksquare note: **no** σ boson

Van der WAALS - like force:

$$\mathbf{V_c}(\mathbf{r}) \propto -rac{\exp[-2\mathbf{m_\pi r}]}{\mathbf{r^6}}\mathbf{P}(\mathbf{m_\pi r})$$

... at intermediate and long distance





CHIRAL DYNAMICS and the **NUCLEAR MANY-BODY PROBLEM**

N. Kaiser, S. Fritsch, W.W. (2002 - 2005)

Small scales:

$$\mathbf{k_F} \sim \mathbf{2}\,\mathbf{m_\pi} \sim \mathbf{M_\Delta} - \mathbf{M_N} << 4\pi\,\mathbf{f_\pi}$$

- PIONS (and DELTA isobars) as explicit degrees of freedom
- IN-MEDIUM CHIRAL PERTURBATION THEORY

pion exchange processes in presence of filled Fermi sea

2nd order TENSOR force + nucleon's SPIN-ISOSPIN polarizability

short-distance dynamics: N contact interactions (incl. resummations)





IN-MEDIUM CHIRAL PERTURBATION THEORY

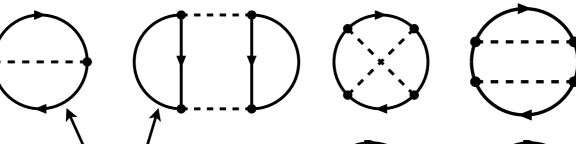
Loop expansion of (In-Medium) Chiral Perturbation Theory \updownarrow Systematic expansion of ENERGY DENSITY $\mathcal{E}(\mathbf{k_F})$ in powers of Fermi momentum [modulo functions $\mathbf{f_n}(\mathbf{k_F}/\mathbf{m_\pi})$] (works for $\mathbf{k_F} << 4\pi\,\mathbf{f_\pi} \sim 1\,\mathrm{GeV}$)

Finite nuclei ←→ energy density functional
J.W. Holt, N. Kaiser, W.W.: Eur. Phys. J. A 47 (2011) 128

many quantitatively successful applications throughout the nuclear chart

e.g. P. Finelli et al.: Nucl. Phys. A 770 (2007) I

Nuclear thermodynamics: compute free energy density



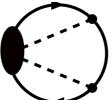
(3-loop order)

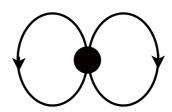
N. Kaiser, S. Fritsch, W.W. (2002-2004)

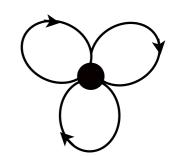


nucleon propagators incl. Pauli blocking











NUCLEAR MATTER

- In-medium ChPT 3-loop $(\pi, \mathbf{N}, \boldsymbol{\Delta})$
- Input parameters: two contact terms
- basically: analytic calculation
- Output:

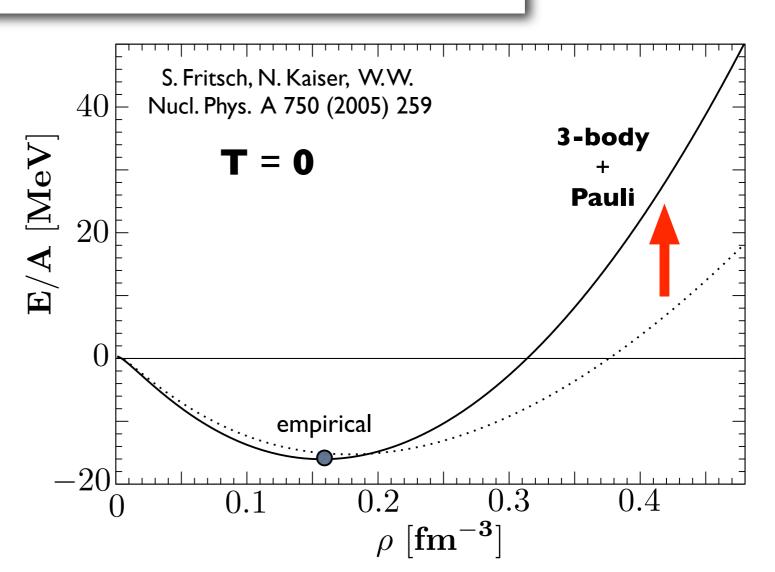


$${
m E_0/A} = -16\,{
m MeV}$$
 , $ho_0 = 0.16\,{
m fm^{-3}}$, ${
m K} = 290\,{
m MeV}$



- $\blacktriangleright \quad \text{Asymmetry energy} \quad \mathbf{A}(\mathbf{k_F^0}) = \mathbf{34\,MeV}$
- Quasiparticle interaction and Landau parameters

J.W. Holt, N. Kaiser, W.W. Nucl. Phys. A 870 (2011) 1, arXiv: 1111.1924 [nucl-th] (NPA (2012), in print)



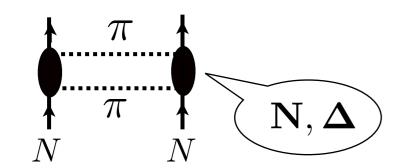


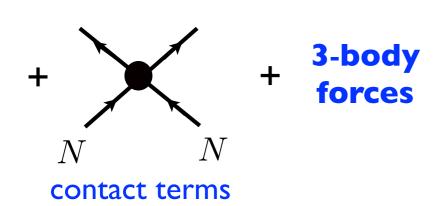
NUCLEAR THERMODYNAMICS

NUCLEAR CHIRAL (PION) DYNAMICS

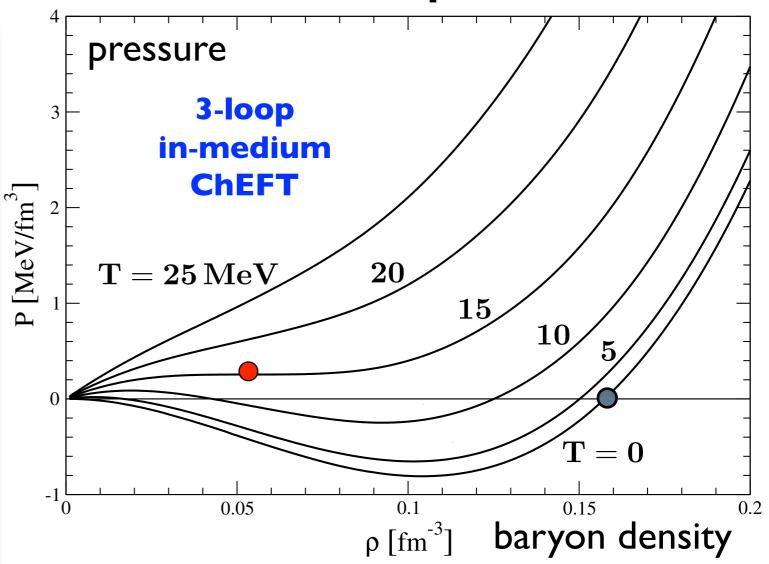
BINDING & SATURATION:

Van der Waals + Pauli





nuclear matter: equation of state



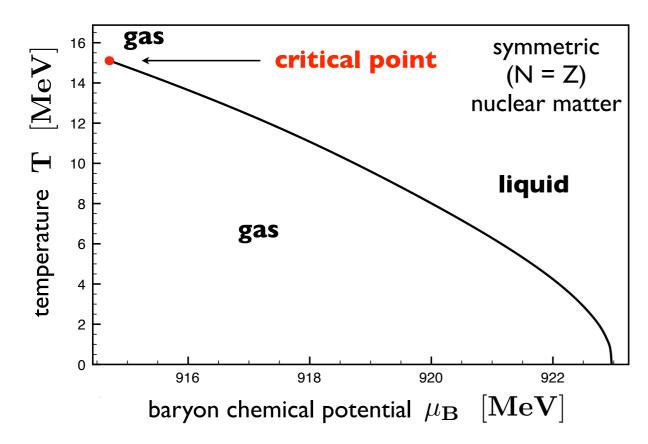
Liquid - Gas Transition at Critical Temperature $T_c = 15 \text{ MeV}$

(empirical: $T_c = 16 - 18 \text{ MeV}$)





PHASE DIAGRAM of NUCLEAR MATTER

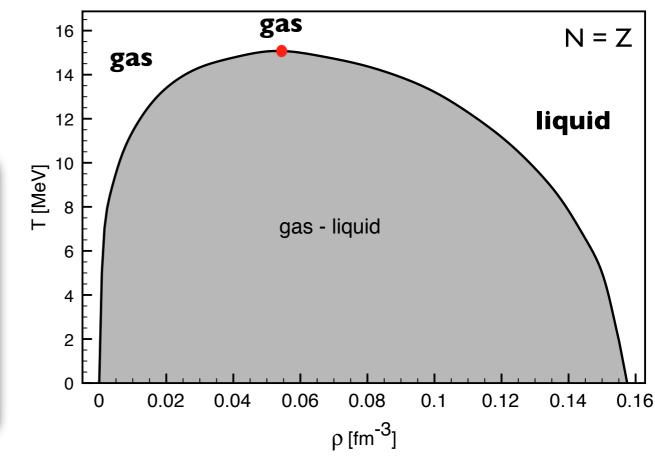


- Pion-nucleon dynamics incl. delta isobars
- Short-distanceNN contact terms
- Three-body forces

In-medium
 chiral effective field theory
 (3-loop calculation of free energy density)

S. Fritsch, N. Kaiser, W.W.: NPA 750 (2005) 259

S. Fiorilla, N. Kaiser, W.W. arXiv:1111.3688 [nucl-th], Nucl. Phys. A (2012), in print

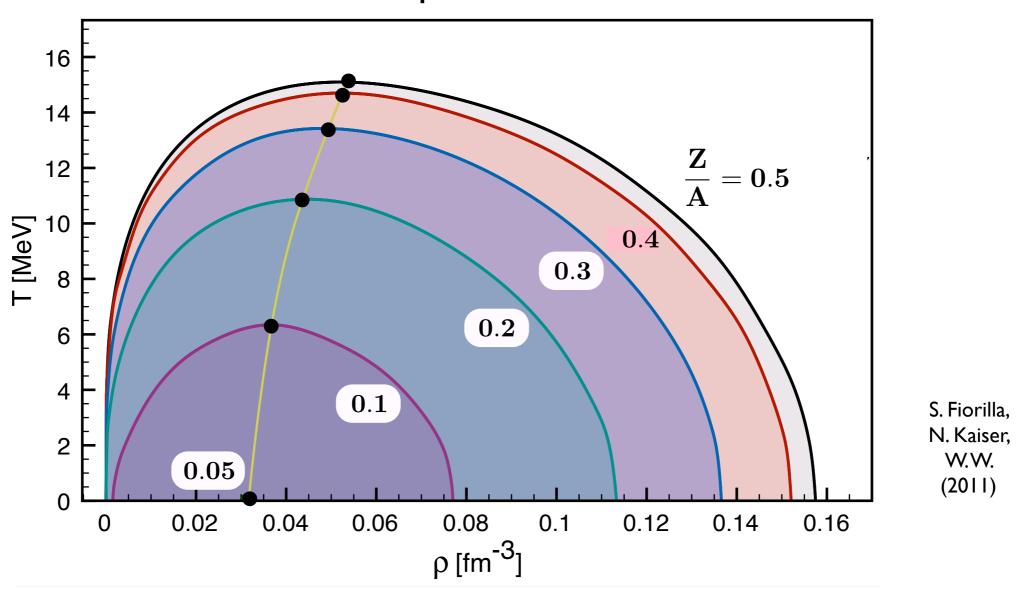






PHASE DIAGRAM of NUCLEAR MATTER

Trajectory of CRITICAL POINT for asymmetric matter as function of proton fraction Z/A

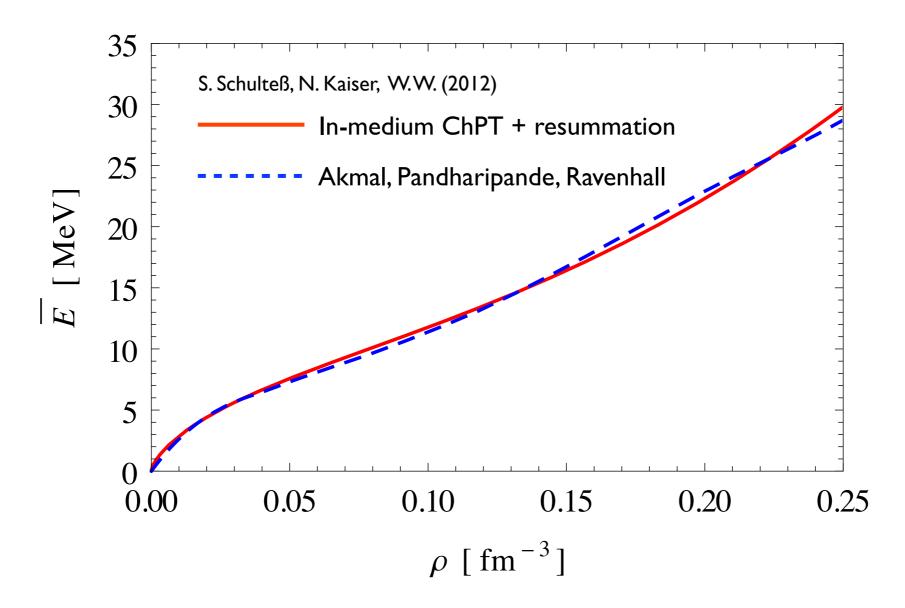


... determined almost entirely by isospin dependent (one- and two-) pion exchange dynamics



NEUTRON MATTER

In-medium chiral effective field theory (3-loop) with resummation of short distance contact terms (large nn scattering length, $a_{\rm s}=19~{
m fm}$) N. Kaiser, Nucl. Phys. A 860 (2011) 370



perfect agreement with sophisticated many-body calculations

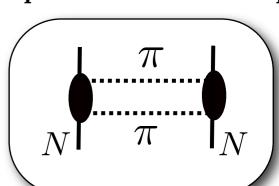


CHIRAL CONDENSATE at finite BARYON DENSITY

- Order parameter of spontaneously broken chiral symmetry in QCD
- Hellmann Feynman theorem: $\langle \Psi | \bar{\mathbf{q}} \mathbf{q} | \Psi \rangle = \langle \Psi | \frac{\partial \mathcal{H}_{\mathbf{QCD}}}{\partial \mathbf{m_q}} | \Psi \rangle = \frac{\partial \mathcal{E}(\mathbf{m_q}; \rho)}{\partial \mathbf{m_q}}$

$$\mathbf{m_q} \frac{\partial \mathbf{M_N}}{\partial \mathbf{m_q}}$$

in-medium chiral effective field theory



$$rac{\langle \mathbf{ar{q}q}
angle_{
ho}}{\langle \mathbf{ar{q}q}
angle_{
m 0}} = \mathbf{1} - rac{
ho}{\mathbf{f_{\pi}^2}} \left[rac{\hat{\sigma_{\mathbf{N}}}}{\mathbf{m_{\pi}^2}} \left(\mathbf{1} - rac{\mathbf{3\,p_F^2}}{\mathbf{10\,M_N^2}} + \ldots
ight) + rac{\partial}{\partial \mathbf{m_{\pi}^2}} \left(rac{\mathbf{E_{int}(p_F)}}{\mathbf{A}}
ight)
ight]$$

(free) Fermi gas of nucleons

nuclear interactions (dependence on pion mass)





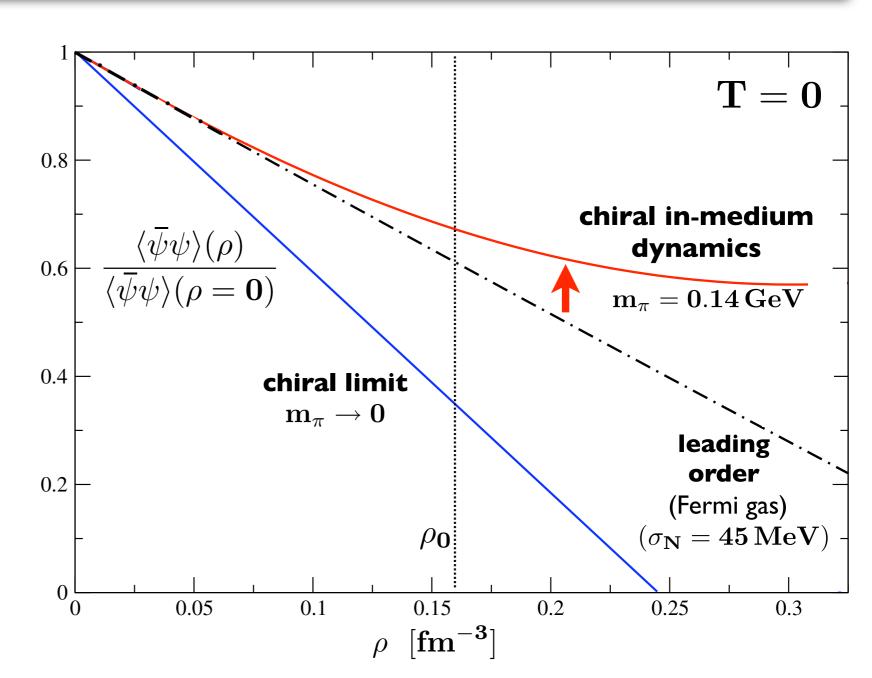
CHIRAL CONDENSATE: DENSITY DEPENDENCE

In-medium
Chiral
Effective
Field Theory

(NLO 3-loop)

constrained by realistic nuclear equation of state

N. Kaiser, Ph. de Homont, W.W. Phys. Rev. C 77 (2008) 025204



- Substantial change of symmetry breaking scenario between chiral limit $m_{\bf q}=0$ and physical quark mass $m_{\bf q}\sim 5\,MeV$
- Nuclear Physics would be very different in the chiral limit!





CHIRAL CONDENSATE:

DENSITY and TEMPERATURE DEPENDENCE

• Free energy density

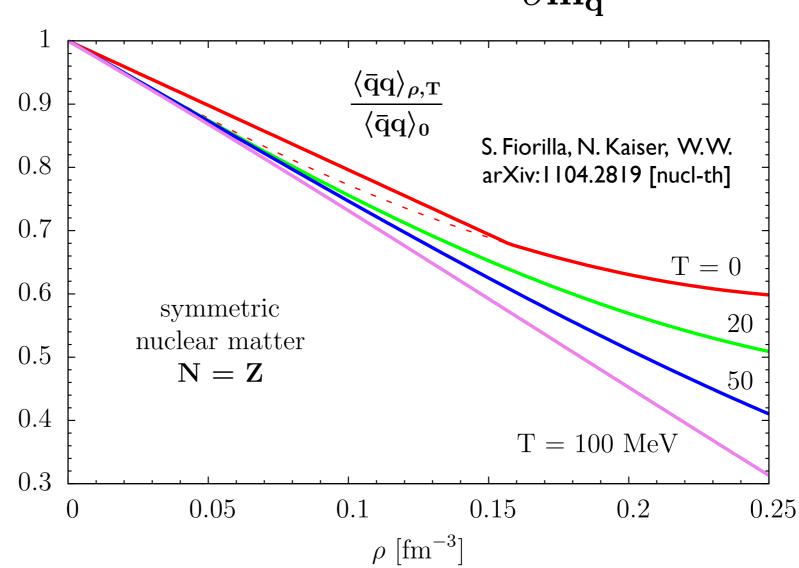
$$\mathcal{F}(\mathbf{m}_{\mathbf{q}}; \rho, \mathbf{T})$$

In-medium
Chiral
Effective
Field Theory

(NLO 3-loop)

constrained by realistic nuclear equation of state

$$\langle \mathbf{\Psi} | \mathbf{ar{q}} \mathbf{q} | \mathbf{\Psi}
angle_{
ho, \mathbf{T}} = rac{\partial \mathcal{F}(\mathbf{m_q};
ho, \mathbf{T})}{\partial \mathbf{m_q}}$$



No indication of first order chiral phase transition for

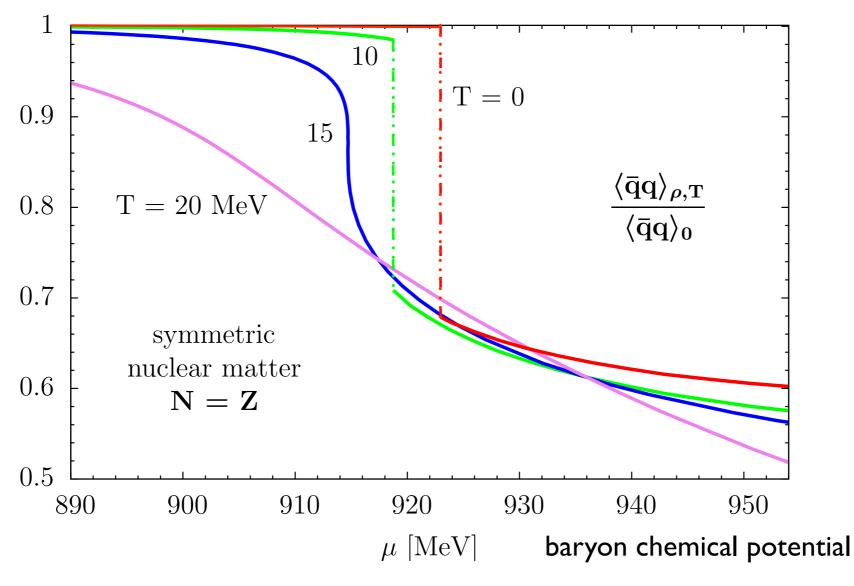
$$\rho \lesssim 2 \rho_0$$
, $T \lesssim 100 \,\mathrm{MeV}$





CHIRAL CONDENSATE:

Dependence on TEMPERATURE and BARYON CHEMICAL POTENTIAL



S. Fiorilla, N. Kaiser, W.W. (2011)

- Liquid-gas phase transition leaves its signature also in chiral condensate
- but: **no** tendency toward **chiral first order transition** in the range

$$\mu_{\mathbf{B}} \lesssim 1 \; \mathbf{GeV}$$



Outlook:

New Constraints from NEUTRON STARS





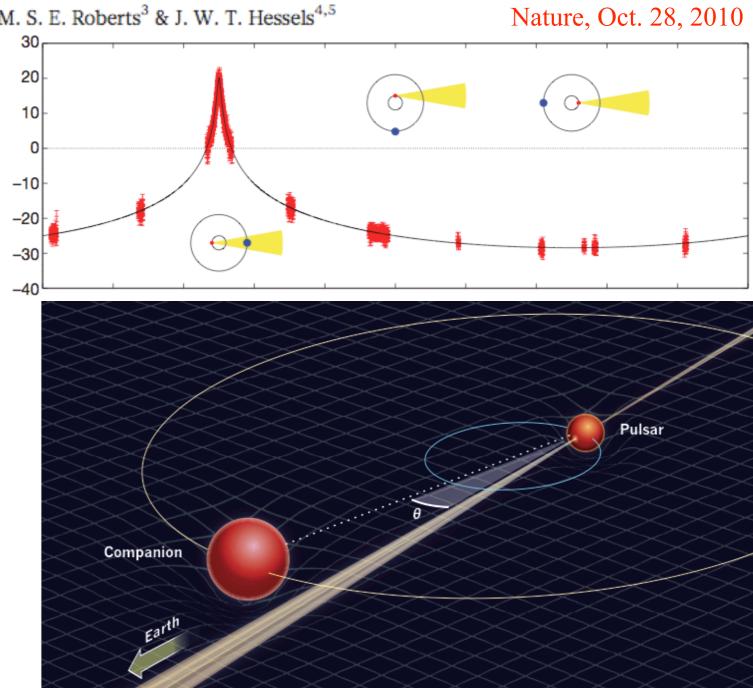
A two-solar-mass neutron star measured using Shapiro delay

P. B. Demorest¹, T. Pennucci², S. M. Ransom¹, M. S. E. Roberts³ & J. W. T. Hessels^{4,5}

direct measurement of neutron star mass from increase in travel time near companion

J1614-2230 most edge-on binary pulsar known (89.17°) + massive white dwarf companion (0.5 M_{sun})

heaviest neutron star with 1.97±0.04 M_{sun}

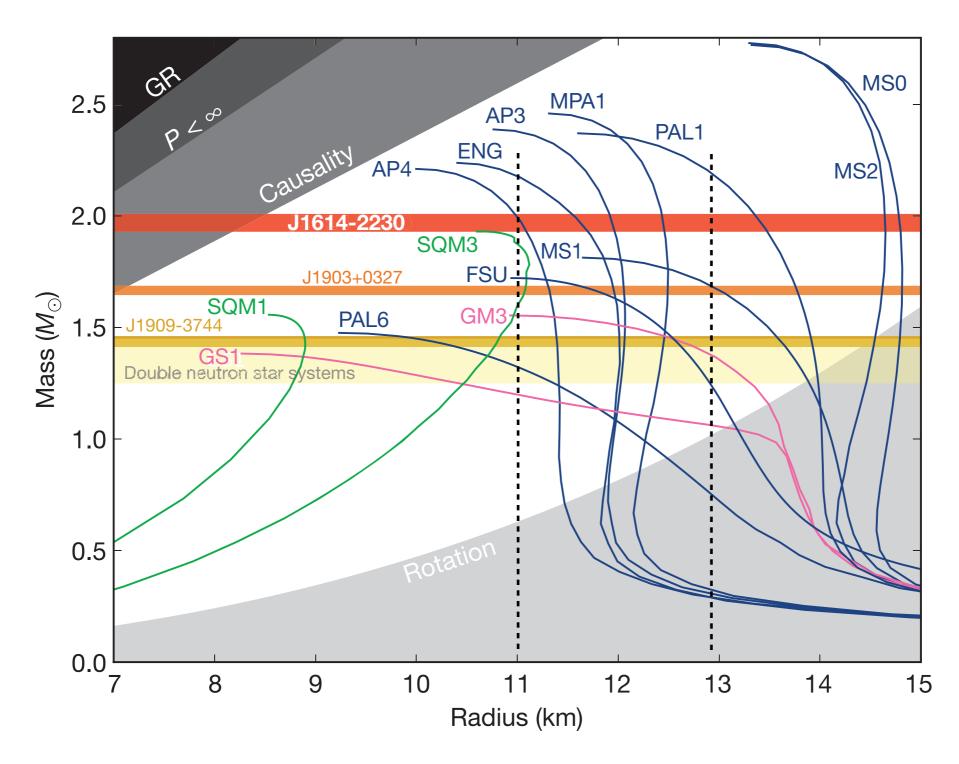




TWO-SOLAR-MASS NEUTRON STAR

... observed using Shapiro delay

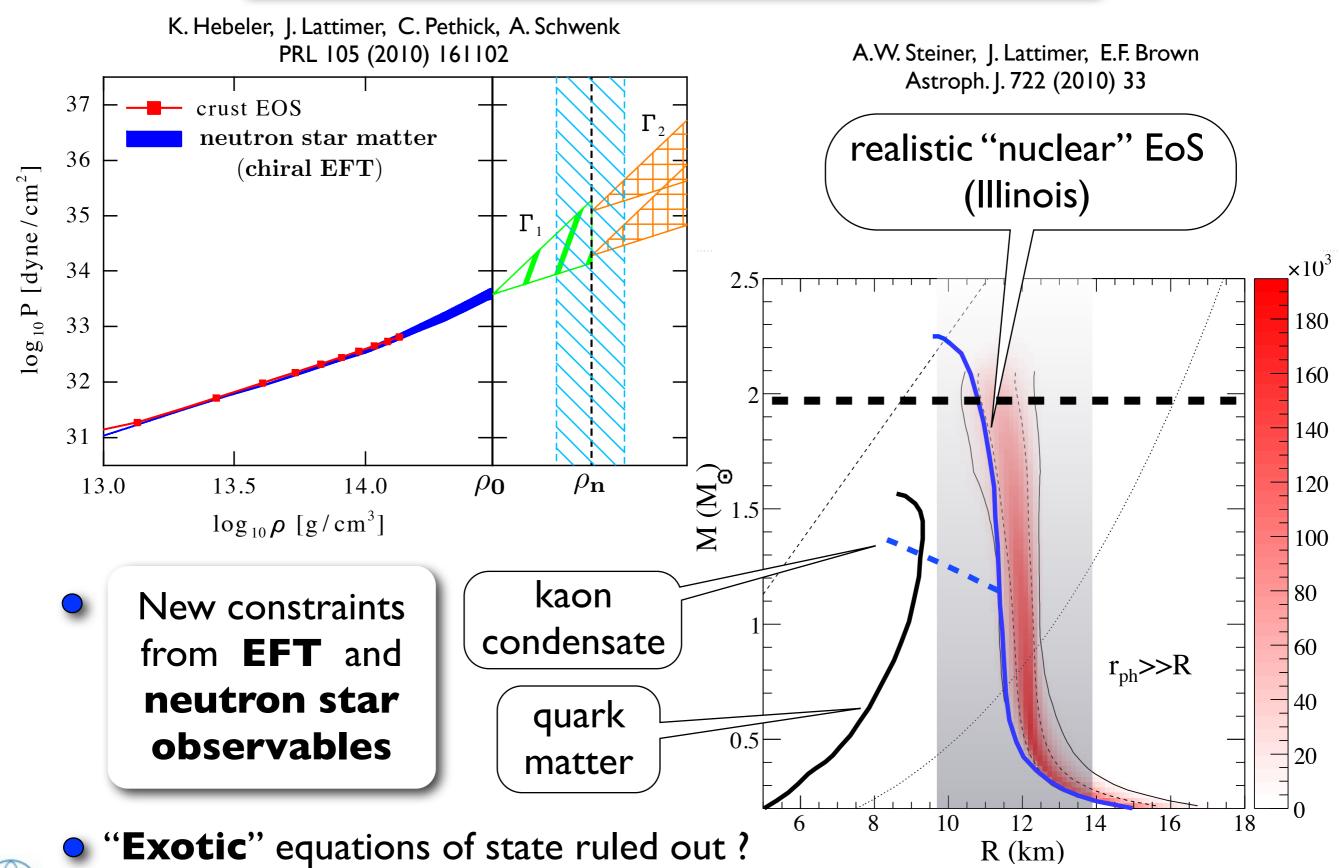
P.B. Demorest et al., Nature 467 (2010) 1081







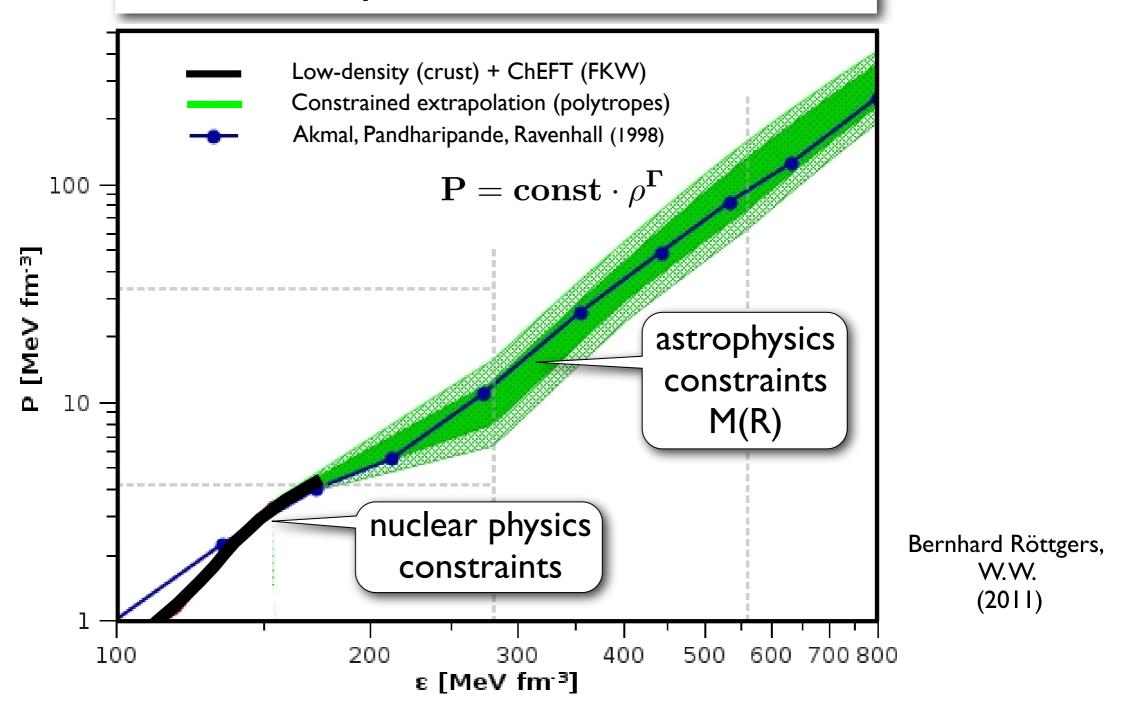
News from NEUTRON STARS





NEUTRON STAR MATTER

Equation of State



Including new neutron star constraints plus
 Chiral Effective Field Theory at lower density





SUMMARY

- Exploration of QCD phase diagram:
 progress concerning basic symmetry breaking patterns
 - Lattice QCD (restricted to small quark chemical potentials)
 - Models (PNJL, PQM) (but: nuclear physics constraints missing)
 - Dyson-Schwinger QCD (-- same problem --)
- Nuclear thermodynamics based on

In-medium Chiral Effective Field Theory

Fermi liquid \leftrightarrow interacting Fermi gas (1st order transition)

No indication of first order chiral transition in the range

$$ho \lesssim 2 \,
ho_0 \,, \,\,\, \mathbf{T} \lesssim 100 \, \mathrm{MeV}$$

- Major challenge: design QCD phase diagram that is consistent with established hadronic and nuclear physics
- New dense & cold matter constraints from neutron stars:
 - Mass radius relation; observation of two-solar-mass n-star
 - "Non-exotic" equation of state works best!





The End

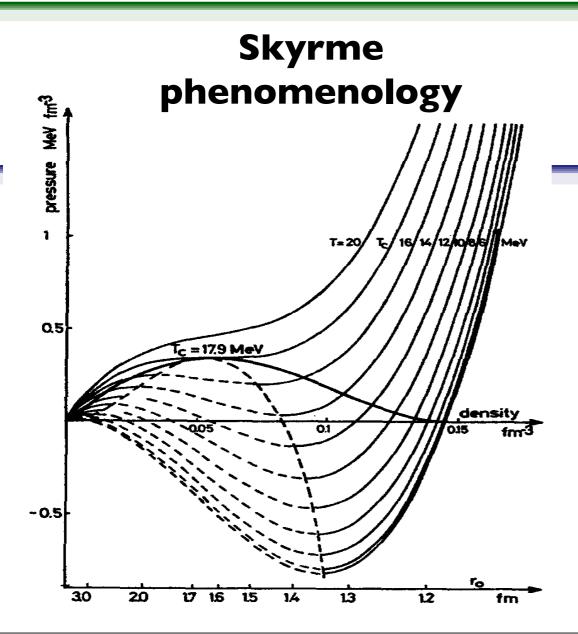
thanks to:

Nino Bratovic Jeremy Holt Salvatore Fiorilla Norbert Kaiser

Thomas Hell Kouji Kashiwa

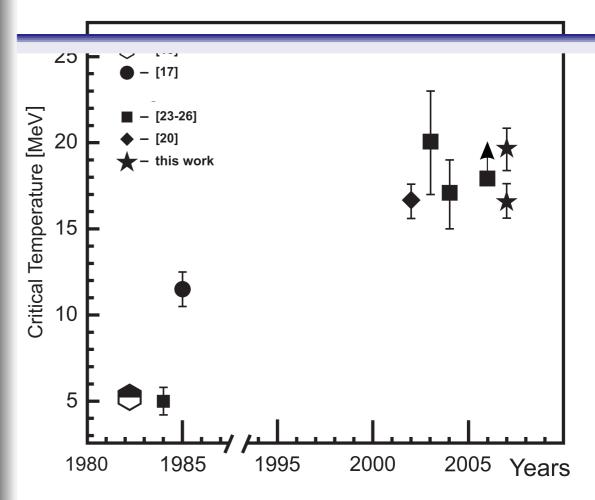


NUCLEAR THERMODYNAMICS



G. Sauer, H. Chandra, U. Mosel Nucl. Phys. A 264 (1976) 221

Multifragmentation and fission analysis



V.A. Karnaukhov et al.: Phys. Atom. Nucl. 71 (2008) 2067



