Clusters in Nuclear Matter and in Heavy Nuclei Stefan Typel



TECHNISCHE UNIVERSITÄT DARMSTADT



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Nuclear Equation of State and Neutron Stars

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Outline



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Clusters in Nuclear Matter

- Theoretical Methods and Applications
- Descriptions in Different Ranges of Density and Temperature
- Unified Description with Generalized Relativistic Density Functional
- Compact Star Matter

Clusters in Heavy Nuclei

- Application of gRDF to Heavy Nuclei
- α Particle Correlations at Surface of Sn Nuclei
- Consequences

Conclusions



Introduction

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Correlations and Clusters



- essential in strongly interacting systems
- different types of correlations
 - two-, three-, ..., many-body correlations
 - short-range/long-range correlations
 - configuration-space/momentum-space correlations (e.g. ²H vs. NN pairing)

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different systems

- nuclear matter (different types of pairing, light clusters in dilute matter)
- light nuclei (clusters in structure and nuclear reactions)
 - \Rightarrow development of structure models with clusters as explicit degrees of freedom (D. M. Brink's α -particle model, resonating group method, ...)

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 - \Rightarrow development of structure models with clusters as explicit degrees of freedom
 - (D. M. Brink's α -particle model, resonating group method, . . .)
- problems
 - effective theoretical description of cluster formation and dissolution in nuclear matter
 - effects of clusters in heavy nuclei



Clusters in Nuclear Matter

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Theoretical Methods and Applications



ab-initio approaches

- realistic nuclear interaction (e.g. fitted to nucleon-nucleon scattering data)
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- effective in-medium interaction
- quasiparticles
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applications

- equation of state/phase diagram
 - \Rightarrow description of astrophysical objects
 - (neutron stars, their mergers, core-collapse supernovae)
- simulation of heavy-ion collisions

Descriptions in Different Ranges of Density and Temperature



exact low-density limits

- finite temperatures: virial equation of state
- zero temperature: Lee-Yang type expansions

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important distinction

- nuclear/baryonic matter (only nucleons/baryons, no Coulomb interaction)
- compact star matter (with leptons and Coulomb interaction, charge-neutral system)
- \Rightarrow different properties and phase diagrams

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application in astrophysical simulations

- wide ranges in density, temperature, isospin asymmetry
- combination of various effective methods
- energy density functionals with clusters
- \Rightarrow unified equation of state

Description at Very Low Densities



finite temperature, exact limit

\Rightarrow virial equation of state (VEOS)

- (E. Beth and G. Uhlenbeck, Physica 3(1936) 729, Physica 4 (1937) 915;
- C. J. Horowitz and A. Schwenk, NPA 776 (2006) 55, $\ldots)$
- **expansion** of pressure in powers of fugacities $z_i = \exp(\mu_i/T)$

$$\rho = TV\left(\sum_{i} \frac{g_{i}}{\lambda_{i}^{3}} z_{i} + \sum_{ij} \frac{b_{ij}}{\lambda_{i}^{3/2} \lambda_{j}^{3/2}} z_{i} z_{j} + \dots\right) \quad \text{with thermal wavelength} \quad \lambda_{i} = \left[2\pi/(m_{i}T)\right]^{1/2}$$

and virial coefficients $g_i, b_{ij}, \ldots \Rightarrow$ limitation $n_i \lambda_i^{-3} \ll 1$

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• only two-body correlations relevant at lowest densities, encoded in $b_{ij} = \frac{1 + \delta_{ij}}{2} \frac{\lambda_i^{3/2} \lambda_j^{3/2}}{\lambda_{ij}^3} \int dE \exp\left(-\frac{E}{T}\right) D_{ij}(E) \pm \delta_{ij} \frac{g_i}{2E^{5/2}} \quad \lambda_{ij} = \left\{2\pi/[(m_i + m_j)T]\right\}^{1/2}$ with 'density of states' $D_{ij}(E) = \sum_k g_k^{(ij)} \delta(E - E_k^{(ik)}) + \sum_l \frac{g_l^{(ij)}}{\pi} \frac{d\delta_l^{(ij)}}{dE}$ $\Rightarrow \text{ contributions from bound states and continuum,}$ denends only on bound state energies $E^{(ik)}$ and phase shifts $\delta^{(ij)}$ (experiment)

depends only on bound-state energies $E_k^{(lk)}$ and phase shifts $\delta_l^{(lj)}$ (experiment!) (not independent! Levinson theorem)

Description at Low Densities I



simplification of VEOS

- \Rightarrow nuclear statistical equilibrium (NSE)
 - consider nucleons and all nuclei (ground and excited states)
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extension of VEOS

\Rightarrow generalized (cluster) Beth-Uhlenbeck approach

(G. Röpke, L. Münchow, and H. Schulz, NPA 379 (1982) 536,
M. Schmidt, G. Röpke, and H. Schulz, Ann. Phys. 202 (1990) 57,
G. Röpke, N.-U. Bastian et al., NPA 897 (2013) 70,
N.-U. Bastian et al., arXiv:1804.10178)

- quantum statistical description with thermodynamic Green's functions
- part of interaction included in self-energies of quasiparticles
- modified second virial coefficient
 - \Rightarrow dependence on particle-pair momentum,

correction factor in continuum contribution

 \Rightarrow suppression of cluster formation with increasing density

Description at Low Densities – Low-Temperature Limit I



pure neutron matter

exact limit at $T = 0 \Rightarrow$ Lee-Yang type expansion (T. D. Lee and C. N. Yang, Phys. Rev. 105 (1957) 1119, H.-W. Hammer and R. J. Furnstahl, NPA 678 (2000) 277)

$$\frac{E}{N} = \frac{3}{5} \frac{k_n^2}{2m_n} \left[1 + \frac{10}{9\pi} \zeta + \frac{4}{21\pi^2} (11 - 2\ln 2) \zeta^2 + \dots \right]$$

 $\zeta = a_{nn}k_n$ with s-wave scattering length a_{nn} and Fermi momentum k_n

 \Rightarrow small radius of convergence ($a_{nn} \approx -18.8$ fm)

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nuclear matter

condensation of (bosonic) clusters expected, does not stop at α condensation

 \Rightarrow increase of cluster size

(no Coulomb interaction \rightarrow no size limit)

- ⇒ coexistence of low-/high-density phases ('liquid-gas phase transition')
- \Rightarrow effect on symmetry energy



Description at Low Densities – Low-Temperature Limit II



compact star matter

- charge neutral system (nucleons + leptons) in β equilibrium
- ► phase transition to solid crystal (Coulomb correlations essential), driven by plasma parameter $\Gamma = Z_{\text{ion}}^{5/3} e^2 / (a_e T) \approx 175$ with $a_e = [3n_e/(4\pi)]^{1/3}$



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 - geometric picture (finite size of particles)
 - \Rightarrow excluded-volume mechanism
 - applications to compact star matter
 (M. Hempel and J. Schaffner-Bielich, NPA 837 (2010) 210;
 Banik et al., ApJ. Suppl. 214 (2014) 22;
 T. Fischer et al., EPJ A 50 (2014) 46; M. Hempel, PRC 91 (2015) 055897)
 - generalized formulation, different interpretation (S. Typel, EPJ A 52 (2016) 16)



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 - generalized formulation, different interpretation (S. Typel, EPJ A 52 (2016) 16)
 - medium modification of cluster properties
 - ⇒ mass shifts
 - ► action of Pauli principle ⇒ blocking of states
 - density, temperature, momentum dependence





Description at High Densities



- baryon density n above n_{sat}
 - \Rightarrow no clusters as degrees of freedom
 - \Rightarrow only single baryons (nucleons, hyperons, ...)
- microscopic models (e.g. Brueckner HF)
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- no explicit correlations between baryons
- \Rightarrow ideal mixture of Fermion gases
- ⇒ step function in single-particle momentum distributions at zero temperature

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k^F_{min} k^F_{maj} majority minority -1/k⁴

experiments: nucleon knockout from nuclei in inelastic electron scattering

(O. Hen et al. (CLAS Collaboration), Science 346 (2014) 614)

 \Rightarrow no sharp cut-off, high-momentum tail

Unified Description at All Densities and Temperatures



- energy density functionals: various types
 - nonrelativistic (e.g. Skyrme, Gogny) or relativistic/covariant
 - often derived from mean-field models in different approximations (Hartree, Hartree-Fock, Hartree-Fock-Bogoliubov)
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here: generalized relativistic density functional (gRDF)

- nucleons, clusters (= many-body correlations) and mesons as degrees of freedom in grand-canonical ensemble
- minimal coupling of nucleons (free or bound) to mesons
- ► quasiparticles with effective mass m^{*}_i = m_i − S_i and effective chemical potential μ^{*}_i = μ_i − V_i
- effective interaction by meson exchange with density dependent couplings \Rightarrow vector (V_i) and scalar (S_i) potentials
- medium dependent masses of clusters

Generalized Relativistic Density Functional



generalisation of relativistic mean field model

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extended set of particle species

- nucleons, electrons, muons, photons, hyperons (optional), ...
- ▶ light nuclei (²H, ³H, ³He, ⁴He) and heavy nuclei (A > 4)
 - bindung energies from mass tables
 ⇒ shell effects included, full distribution, not only average heavy nucleus
- two-nucleon scattering states
 - \Rightarrow consistency with virial EoS at low densities

excited states of nuclei

temperature dependent degeneracy factors with density of states

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medium dependence of particle properties

quasiparticles with mass shifts (coupling to mesons, effective Pauli principle)

(S. Typel et al., Phys. Rev. C 81 (2010) 015803; M. D. Voskresenskaya et al., Nucl. Phys. A 887 (2012) 42;
M. Hempel et al., Phys. Rev. C 91 (2015) 045805; S. Typel, arXiv:1504.01571; H. Pais et al., arXiv:1612.07022;
H. Pais et al. Nuovo Cim. C 39 (2016) 393; S. Typel, J. Phys. G 45 (2018) 114001)

Mass Shifts I



concept applies to composite particles: clusters

- light and heavy nuclei
- nucleon-nucleon correlations in continuum
 - \Rightarrow medium dependent resonances

effective change of masses/binding energies

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- light and heavy nuclei
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- effective change of masses/binding energies
- two major contributions $\Delta m_i = \Delta m_i^{\text{strong}} + \Delta m_i^{\text{Coul}}$
 - strong shift $\Delta m_i^{\text{strong}} = \Delta m_i^{\text{meson}} + \Delta m_i^{\text{Pauli}}$
 - effects of strong interaction (coupling to mesons)
 - Pauli exclusion principle: blocking of states in the medium
 - \Rightarrow reduction of binding energies
 - \Rightarrow cluster dissolution at high densities: Mott effect
 - \Rightarrow replaces traditional excluded-volume mechanism

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 - electromagnetic shift Δm_i^{Coul} (in compact star matter)
 - electron screening of Coulomb field
 - \Rightarrow increase of binding energies

Mass Shifts II



- light nuclei and NN scattering states
 - parametrization from G. Röpke simplified and modified for high densities and temperatures
 - scattering states: mass shifts as for deuteron



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 - ► dependence of $\Delta m_i^{\text{Pauli}}$ on temperature *T* and effective density $n_i^{\text{eff}} = \frac{2}{A_i} \left[Z_i Y_q + N_i (1 - Y_q) \right] n_b$ \Rightarrow asymmetry of medium
 - ► Δm^{Coul} in Wigner-Seitz approximation
 - full coupling of nucleons in clusters to meson fields



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 - ► △*m*^{Coul} in Wigner-Seitz approximation
 - full coupling of nucleons in clusters to meson fields
- heavy nuclei
 - heuristic parametrization



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extension with DZ31 masses

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neutronisation with increasing baryon density

Compact Star Matter Zero Temperature



- formation of neutron star crust, sequence of ions, phase transitions
- more neutron rich at higher densities, approaching neutron drip density
- 'pasta phases' before transition to uniform matter



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Compact Star Matter Light Clusters



- formation and dissolution with increasing density
- temperature dependence



Compact Star Matter Heavy Clusters



- full distribution of nuclei with shell effects (ground and excited states)
- single-nucleus approximation (SNA) not sufficient



Light Clusters in Heavy-Ion Collisions



emission of light nuclei

 determination of density and temperature of source

S. Kowalski et al. PRC 75 (2007) 014601 J. Natowitz et al. PRL 104 (2010) 202501 R. Wada et al. PRC 85 (2012) 064618

 thermodynamic conditions as in neutrinosphere of core-collapse supernovae

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- thermodynamic conditions as in neutrinosphere of core-collapse supernovae
- ► particle yields \Rightarrow chemical equilibrium constants $K_c[i] = n_i / (n_p^{Z_i} n_n^{N_i})$ L. Qin et al., PRL 108 (2012) 172701
- mixture of ideal gases not sufficient
- new data from INDRA collaboration H. Pais et al., arXiv:1911.10849



M. Hempel, K. Hagel, J. Natowitz, G. Röpke, S. Typel, PRC C 91 (2015) 045805



very reasonable nuclear matter parameters n_{sat} = 0.149 fm⁻³, a_V = 16.02 MeV, K = 242.7 MeV, J = S₀ = 31.67 MeV, L = 55.04 MeV (S. Typel et al., PRC 81 (2010) 015803)



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 100
 Excluded Tax 4²Mas

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- symmetry energy consistent with unitary gas constraint (E. E. Kolomeitsev et al., Astrophys. J. 848 (2017) 105)
- neutron star mass-radius relation consistent with maximum mass and recent mass-radius constraints

$$M_{\rm max} = 2.42 \ {\rm M}_{\odot}, R_{1.4} = 13.2 \ {\rm km}$$



(MeV)



Clusters in Heavy Nuclei

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Application of Generalized Relativistic Density Functional to Heavy Nuclei



- extension to zero temperature
- simplified nuclear structure calculation (S. Typel, PBC 89 (2014) 064321)
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 - nucleons and α particles as degrees of freedom
 - Thomas-Fermi approximation for nucleons
 - explicit α-particle wave function in WKB approximation
 - finite size of α particles

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- application to chain of Sn isotopes
 - α particle distribution at surface of nuclei



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 - reduced probability of *α* occurrence with increasing neutron excess

(consistent with trend of α particle reduced widths in (d, ^6Li) pickup reactions

on Sn nuclei, A. A. Cowley, Phys. Rev. C 93 (2016) 054329)



Study of α Particle Correlations at Surface of Sn Nuclei I



quasifree (p,pα) knockout reactions on Sn nuclei

- experimental signatures:
 - dependence of cross sections on neutron excess
 - Iocalisation of α particles at surface
 - \Rightarrow broad momentum distribution

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experiments at RCNP, Osaka (E461)

- targets: stable ^{112–124}Sn nuclei
- beam: 392 MeV protons, 100 pnA
- proton detection: Grand Raiden
- α detection: LAS
- first experiment (June 2015): failure of some detectors
- second experiment (February 2018): successful



(a)

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Study of α Particle Correlations at Surface of Sn Nuclei II



• quasifree (p,p α) knockout reactions on Sn nuclei

- experiment
 - Spectrometer setting: $\theta_{lab}(p) = 45.3 \text{ deg}, \theta_{lab}(\alpha) = 60 \text{ deg}$
 - momentum coverage: $Q_{\alpha} \leq$ 80 MeV/c
 - analysis: Junki Tanaka and Yang Zaihong

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- theory
 - distorted-wave eikonal model in impulse approximation
 - \Rightarrow factorization of cross section
 - α particle distribution from GRDF
 - proton optical potential from global Dirac phenomenology (S. Hama et al., Phys. Rev. C 41 (1990) 2737)
 - elastic proton-α cross section
 (K. Yoshida et al., Phys. Rev. C 94 (2016) 044604)
 - scaled α particle optical potential (M. Nolte et al., Phys. Rev. C 36 (1987) 1312)
 - correction for experimental acceptances

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Consequences of α Particle Correlations at Surface of Heavy Nuclei



Sn nuclei

- reduction of neutron skin thickness
- no effect for np-symmetric nuclei and very neutron-rich nuclei
- strong effect

mass number	112	116	120	124
rel. change	-44%	-31%	-23%	-15%



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Sn nuclei

- reduction of neutron skin thickness
- no effect for np-symmetric nuclei and very neutron-rich nuclei
- strong effect

mass number	112	116	120	124
rel. change	-44%	-31%	-23%	-15%

- ²⁰⁸Pb nucleus
 - expected reduction of neutron skin thickness: 0.018 fm



Consequences of α Particle Correlations at Surface of Heavy Nuclei



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- ²⁰⁸Pb nucleus
 - expected reduction of neutron skin thickness: 0.018 fm
 - ⇒ affects correlation of neutron skin thickness with slope parameter *L* of symmetry energy (no α cluster effects in conventional energy density functional calculations)
 - \Rightarrow systematic uncertainty

X. Viñas et al., Eur. Phys. J. A 50 (2014) 27





Conclusions

January 16, 2020 | Darmstädter Haus, Hirschegg, Kleinwalsertal, Austria | S. Typel | 62

Summary and Outlook



- cluster formation essential feature in nuclei and nuclear matter
- effective theoretical description with generalized relativistic density functional
- applications
 - equation of state of nuclear and compact star matter
 - clusters at surface of heavy nuclei

Summary and Outlook



- cluster formation essential feature in nuclei and nuclear matter
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 - equation of state of nuclear and compact star matter
 - clusters at surface of heavy nuclei
- new parametrisation of nucleon-meson couplings
 - scalar density dependence to avoid problems at zero baryon density
 - inclusion of tensor couplings (project with Diana Alvear, visitor in Erasmus+ program)
- clusters as effective means to describe correlations above saturation density (project with Stefano Burrello, Alexander-von-Humboldt fellowship)
- improvement of mass shifts for light and heavy nuclei, effects of momentum dependence?