

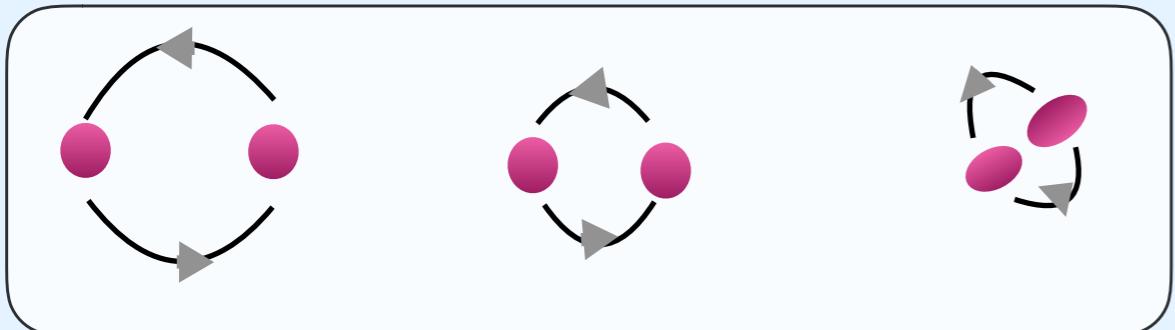
Tidal deformabilities and radii of Neutron Stars from multimessenger observations



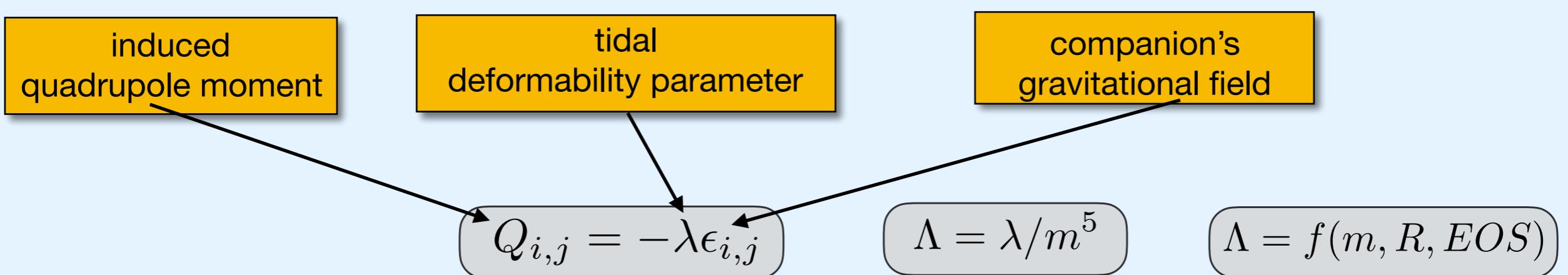
Soumi De
Syracuse University

International Workshop XLVIII on Gross Properties of Nuclei and Nuclear
Excitations
Hirschegg, Kleinwalsertal, Austria, Jan 13, 2020

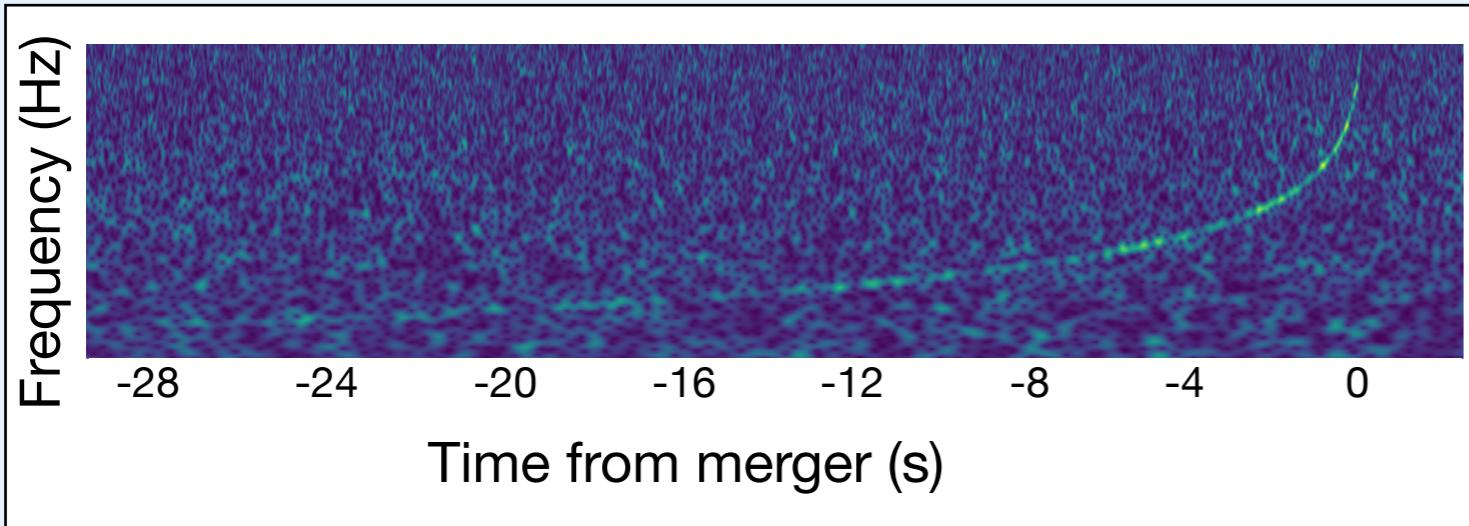
Tidal deformations in binary neutron star inspirals



Neutron stars in binaries are deformed by the gravitational field of their companion towards the end of the inspiral



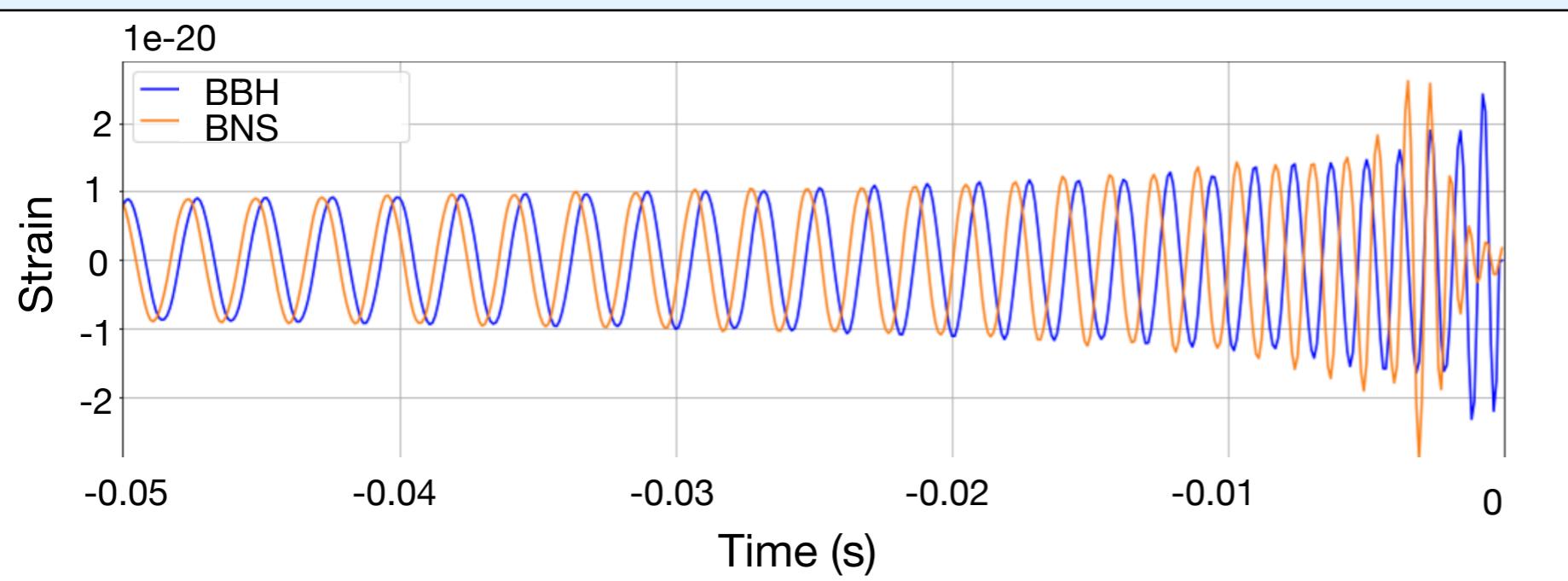
Effect of tidal deformability on gravitational waves from binary neutron stars



$$\tilde{\Lambda} = \frac{16}{13} \frac{(12q+1)\Lambda_1 + (12+q)q^4\Lambda_2}{(1+q)^5}$$

Enters into the phase of the emitted GW

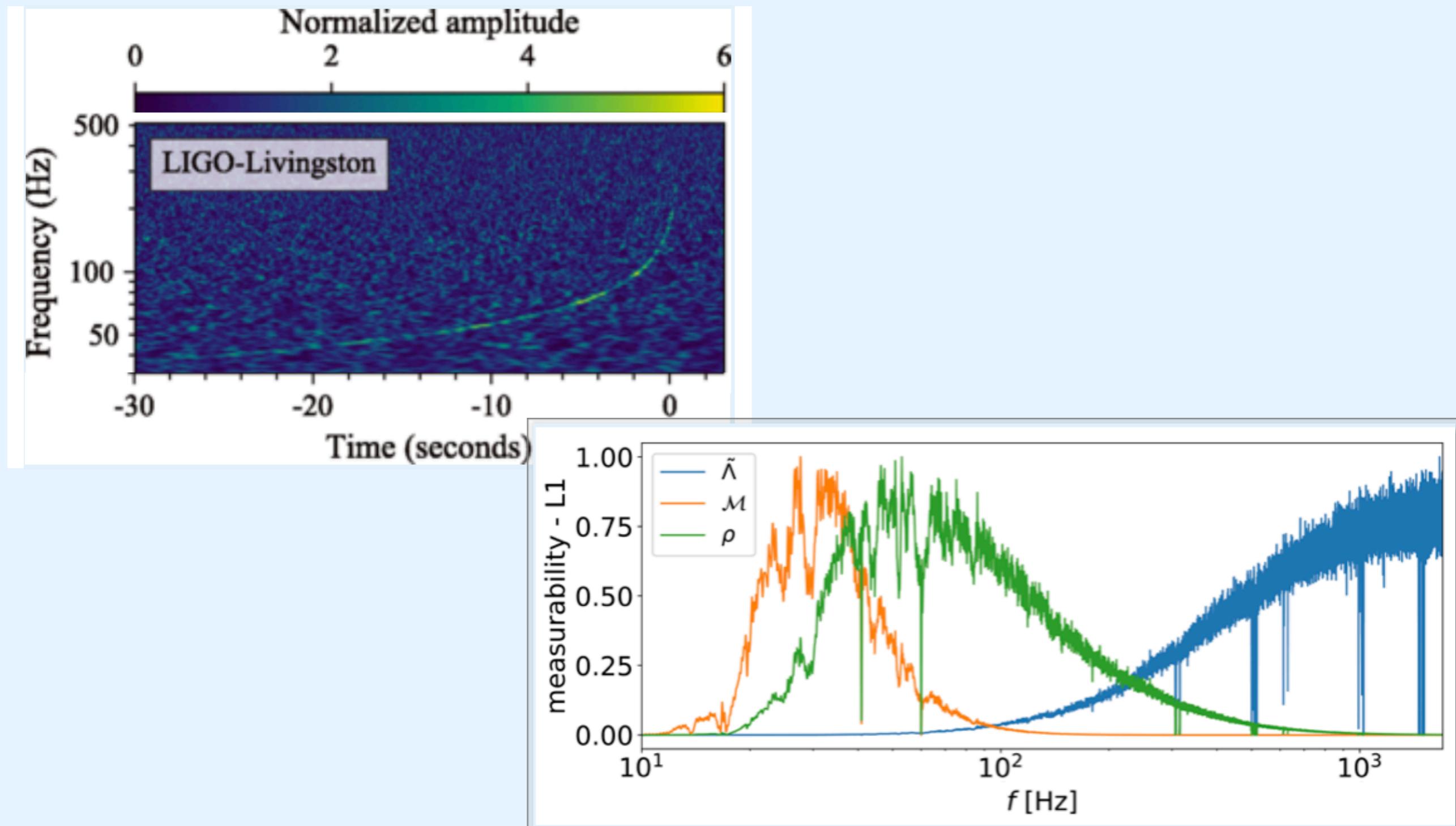
$$\Phi(t) \sim \phi_0(\mathcal{M}; t) \left[1 + \phi_1(\eta; t) \left(\frac{v}{c} \right)^2 + \dots + \phi_5(\tilde{\Lambda}; t) \left(\frac{v}{c} \right)^{10} \right]$$



$$\Lambda_{1,2} = \frac{2}{3} k_2 \left(\frac{R_{1,2} c^2}{G m_{1,2}} \right)^5$$

$$q = \frac{m_2}{m_1}$$

Measurabilities



De et al. Phys. Rev. Lett. **121**, 091102 (2018)

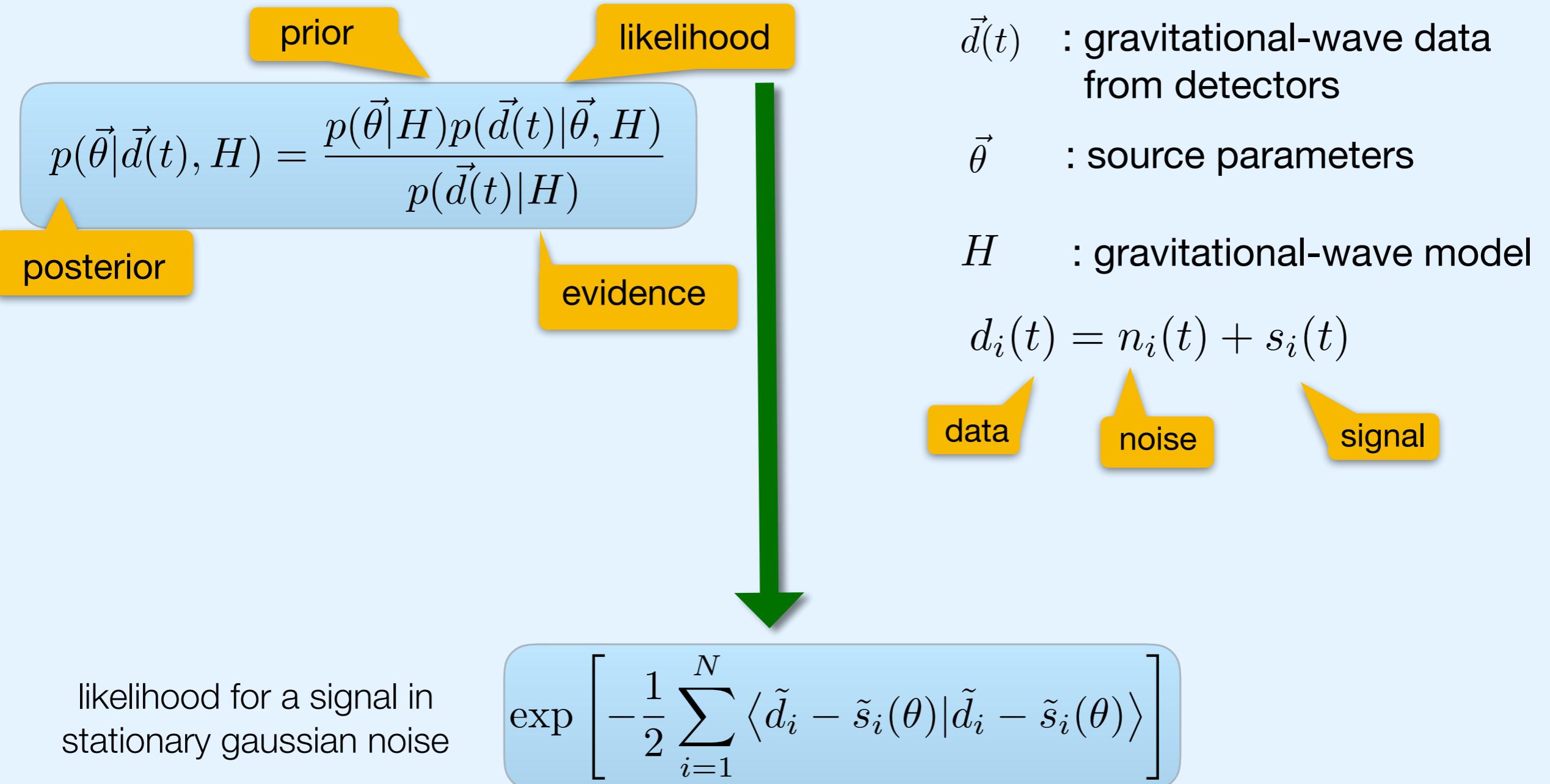
Parameters of the gravitational waveform

Parameter space $\vec{\theta}$

- Component masses : m_1, m_2
- Component spins : \vec{s}_1, \vec{s}_2
- Distance to source : d_{L}
- Source location and orientation : $\alpha, \delta, \psi, \iota$
- Coalescence time and phase : t_c, ϕ_c
- Component tidal deformabilities : Λ_1, Λ_2

Bayesian inference analyses of GW data to extract these parameters

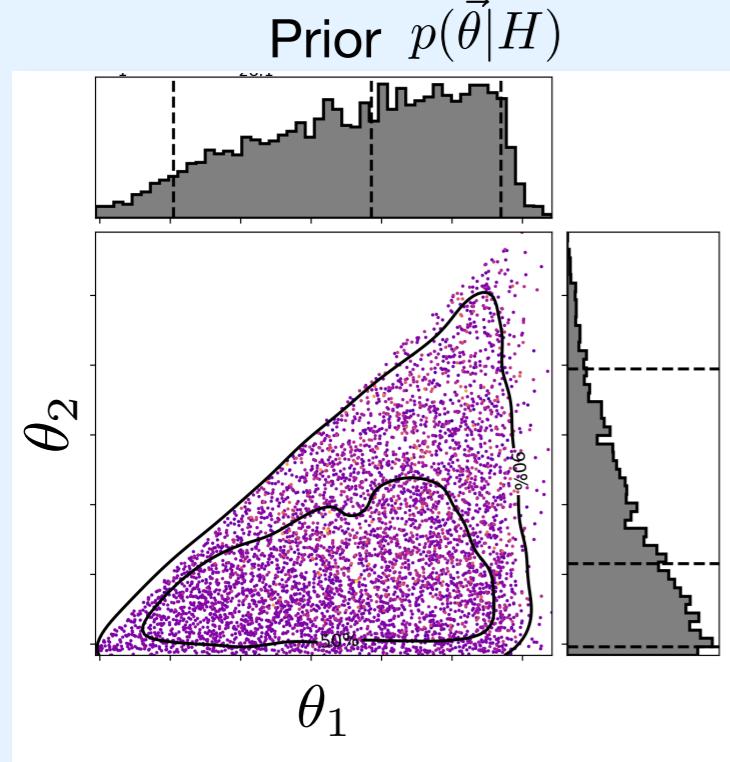
Parameter estimation of the gravitational-wave data : Bayesian framework



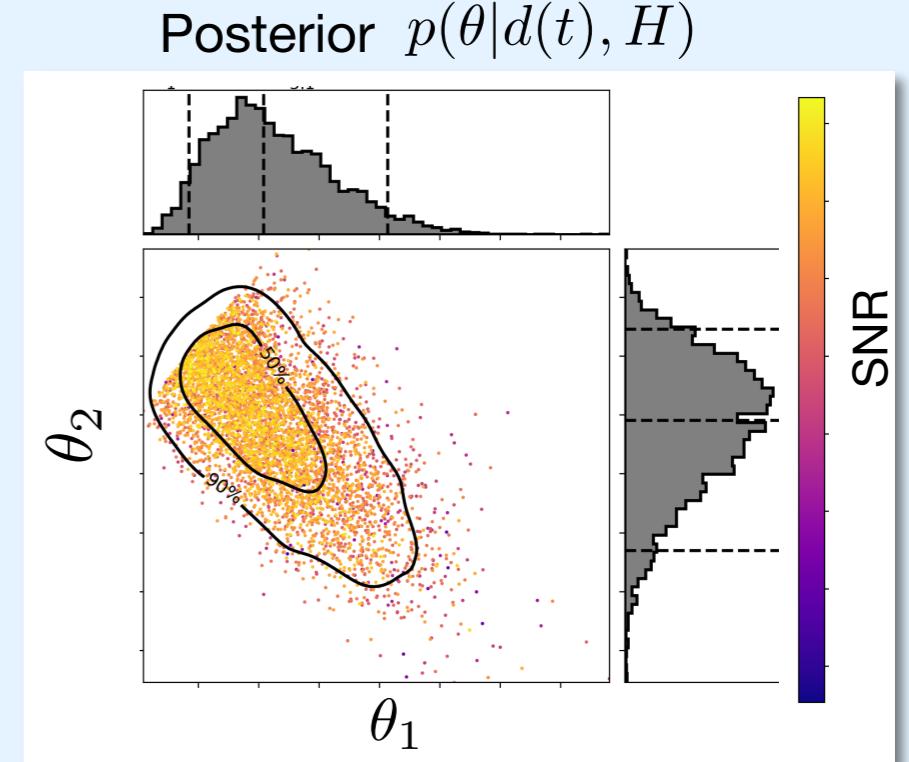
Bayesian framework

$$p(\vec{\theta}|\vec{d}(t), H) = \frac{p(\vec{\theta}|H)p(\vec{d}(t)|\vec{\theta}, H)}{p(\vec{d}(t)|H)}$$

- ✓ Data $\vec{d}(t)$ from GW-Open Science Center
- ✓ Prior distribution
- ✓ Waveform model $h(t; \alpha, \delta, m_1, m_2, \vec{s}_1, \vec{s}_2, \Lambda_1, \Lambda_2, \dots)$
- ✓ Likelihood $p(\vec{d}(t)|\vec{\theta}, H)$



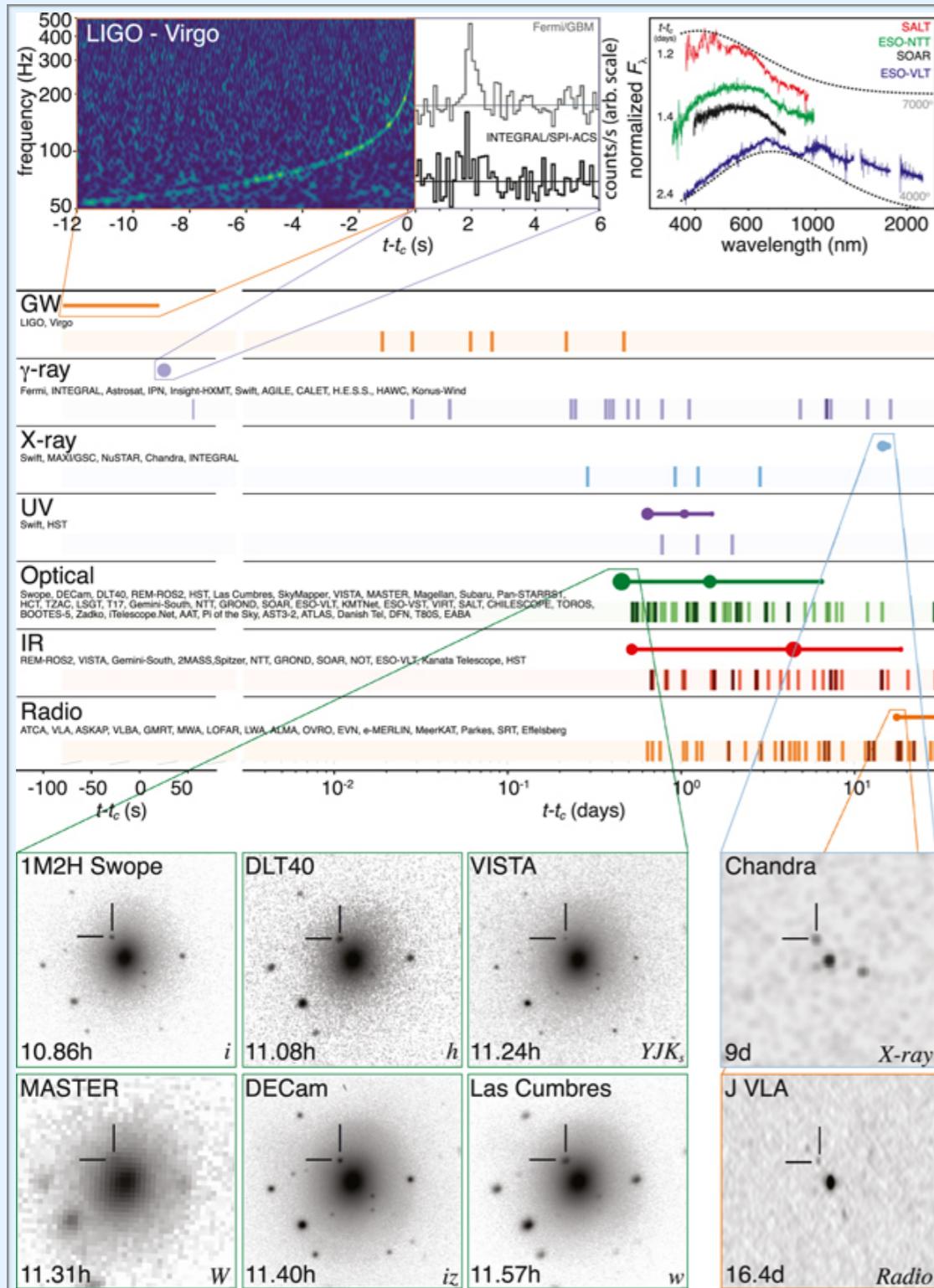
Run sampler with Markov Chain Monte Carlo techniques to find the underlying distribution



- ✓ Marginalization to find parameter estimates

$$p(\theta_n|\vec{d}(t), H) \propto \int p(\vec{d}(t)|\vec{\theta}, H)p(\vec{\theta}|H)d\theta_1\dots d\theta_{n-1}$$

Observation of GW170817



- Subsequent observation of gravitational waves and electromagnetic waves across the full spectrum
- Gave us the opportunity to use multimessenger observations to probe the physics associated with the event

Analysis of GW170817 using multimessenger observations and chiral effective field theory

arXiv:1908.10352

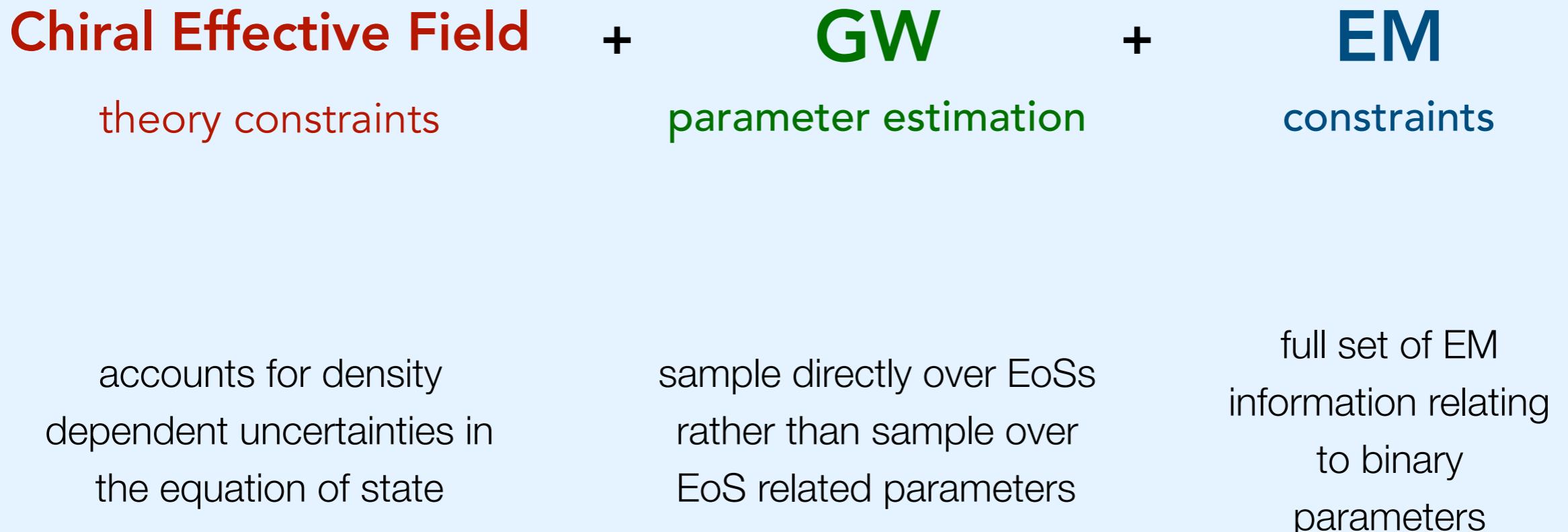
Collin Capano, Ingo Tews, Stephanie Brown, Ben Margalit, Soumi De,
Sumit Kumar, Duncan Brown, Badri Krishnan, Sanjay Reddy (2019)

What's new in the analysis?

Chiral Effective Field + GW + EM

theory constraints parameter estimation constraints

What's new in the analysis?



Equations of state from Chiral Effective Field theory

- Chiral Effective Field Theory at low density, provides a description of matter in terms of nucleons and pions
- It is expected to provide a good description of the equation of state of matter up to $1 - 2 n_{\text{sat}}$
- The equations of state constrained by the Chiral EFT are characterized by density-dependent theoretical uncertainty estimates
- The EoSs also ensure $c_{\text{sound}} < c_{\text{light}}$ and $M > 1.9M_{\odot}$

Gravitational wave parameter estimation method

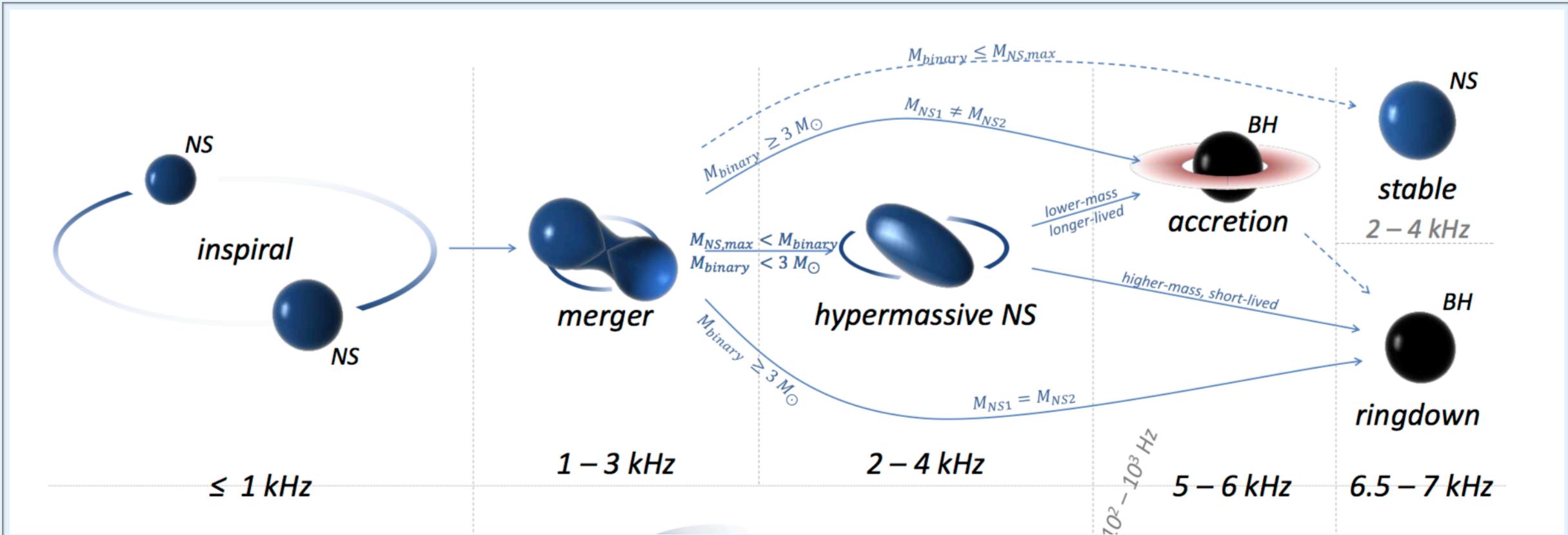
Variable parameter space $\vec{\theta}$ in this analysis

- Component masses : m_1, m_2
 - Component spins : \vec{s}_1, \vec{s}_2
 - Distance to source : d_L
 - Source location and orientation : $\alpha, \delta, \psi, \iota$
 - Coalescence time and phase : t_c, ϕ_c
 - ~~Component tidal deformabilities : Λ_1, Λ_2~~
 - Equation of state constrained by Chiral EFT, indexed by :
 i + m_1, m_2 Each EoS is described by a
(m, R, Λ) dataset
- \downarrow
- R_1, R_2 and Λ_1, Λ_2
- by-passes the need of a universal relation

Gravitational wave parameter estimation method

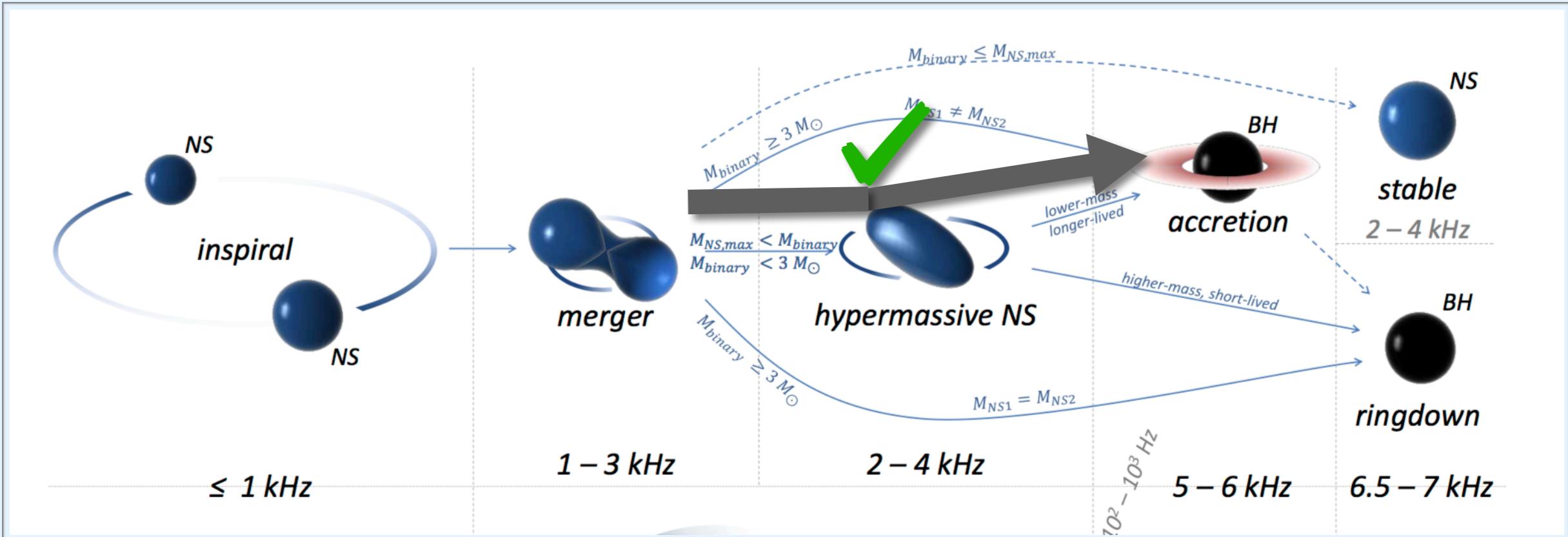
- Analyzed GW data between 20 - 2048 Hz
- Used TaylorF2 waveform model (Better waveform models did not produce an significant difference in results)
- Assumed low spins $\chi < 0.05$
- Performed Bayesian inference analysis using two families of EoSs generated with a uniform prior on $R_{1.4M_\odot}$ using Chiral EFT

Constraints from electromagnetic observations of GW170817



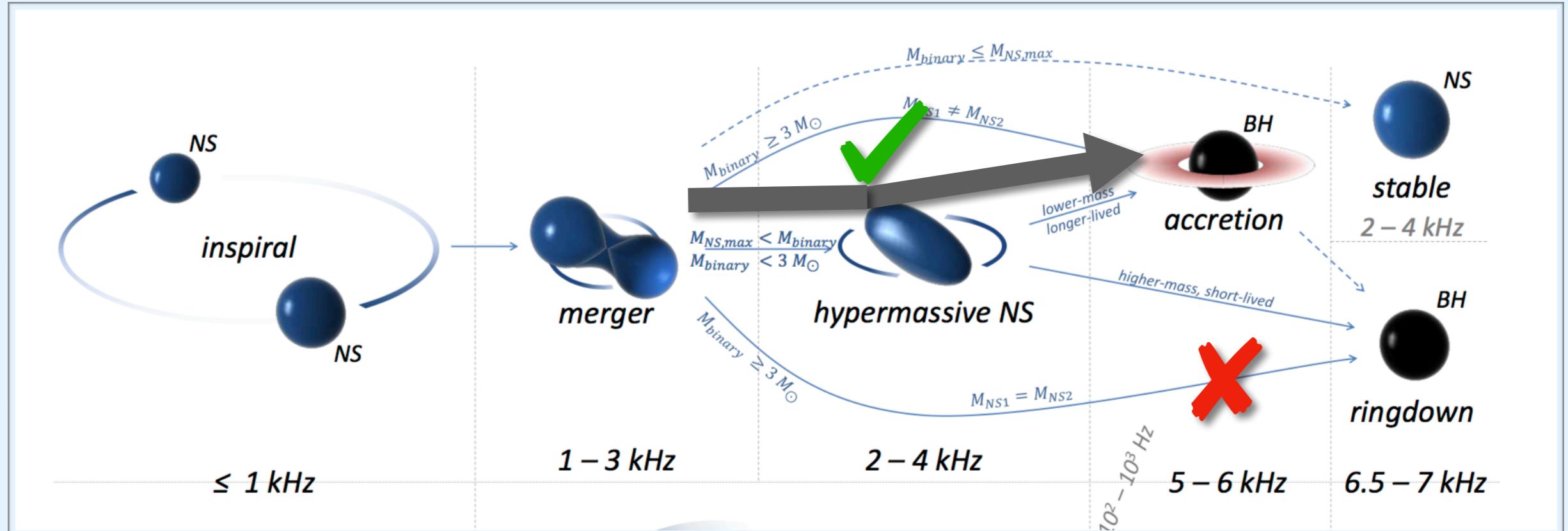
Bartos et al. (2013) CQG 30, 123001

Constraints from electromagnetic observations of GW170817



Bartos et al. (2013) CQG 30, 123001

Constraints from electromagnetic observations of GW170817



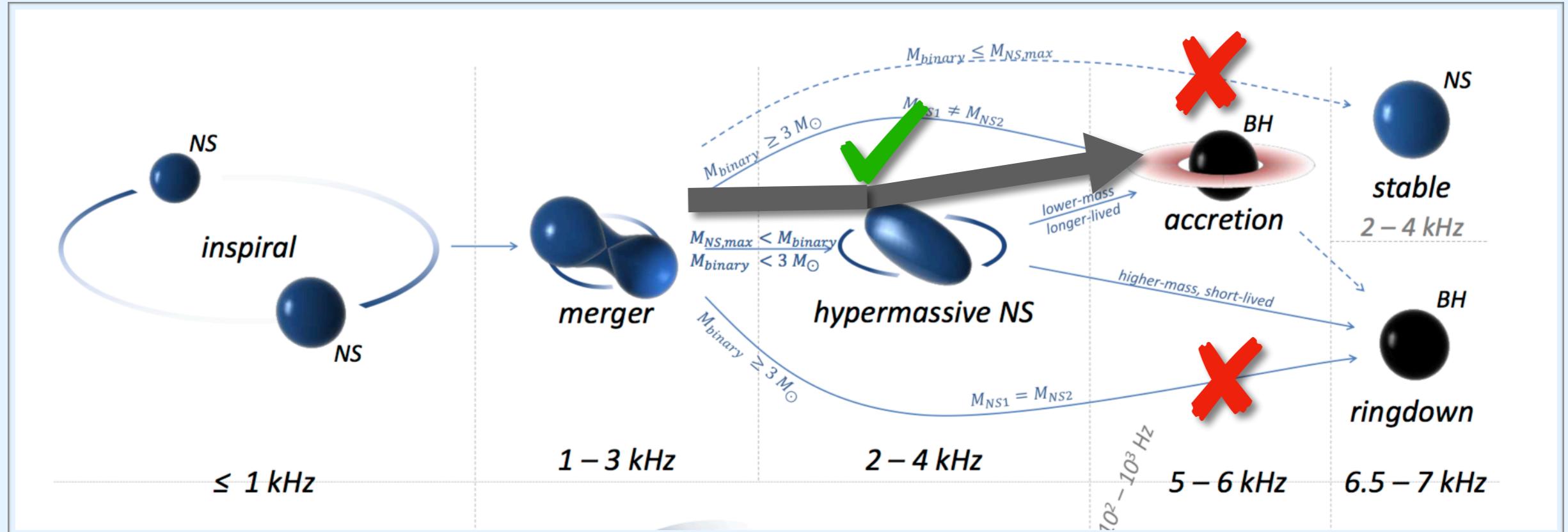
Bartos et al. (2013) CQG 30, 123001

- Total mass in agreement with **No prompt collapse to BH**

$$M_{\text{total}} < M_{\text{thres}} \approx M_{\text{max}} \left(2.380 - 3.606 \frac{GM_{\text{max}}}{c^2 R_{1.6M_{\odot}}} \right) \pm 0.05 M_{\odot}$$

Bauswein et al. PRL 111, 131101 (2013)

Constraints from electromagnetic observations of GW170817



Bartos et al. (2013) CQG 30, 123001

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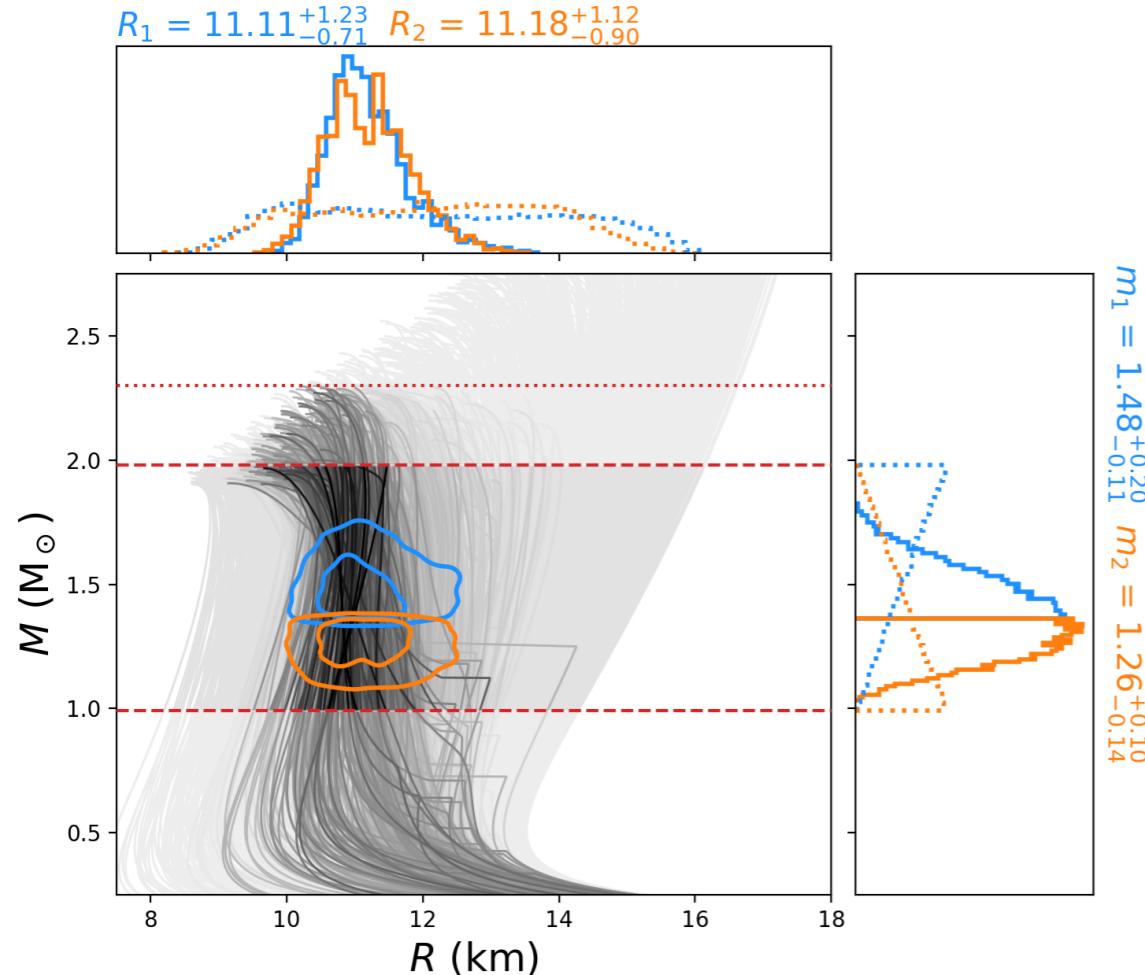
Bauswein et al. PRL 111, 131101 (2013)

- Max NS mass in agreement with **No long-lived NS**

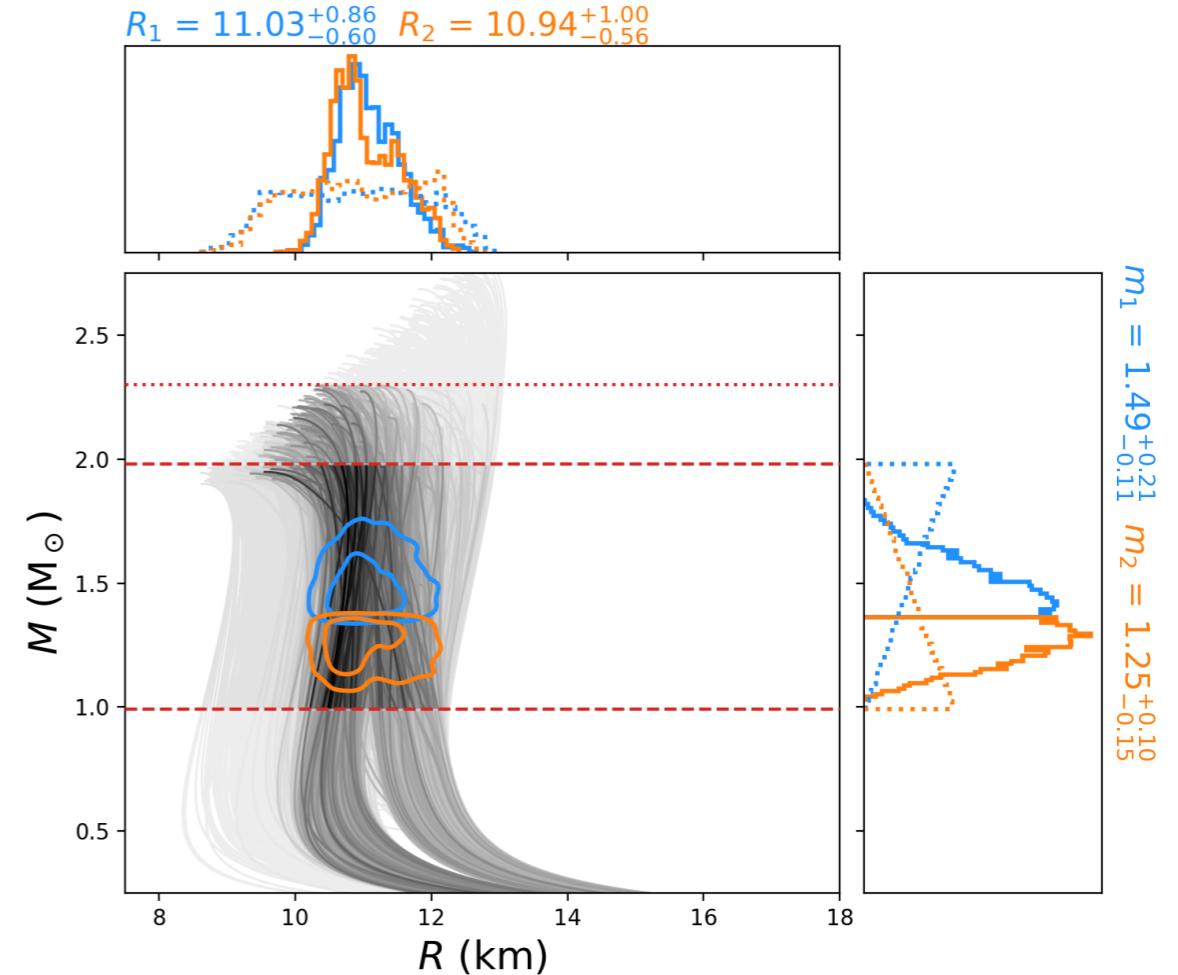
$$M_{\text{max}} < 2.3M_{\odot}$$

Analysis of GW170817

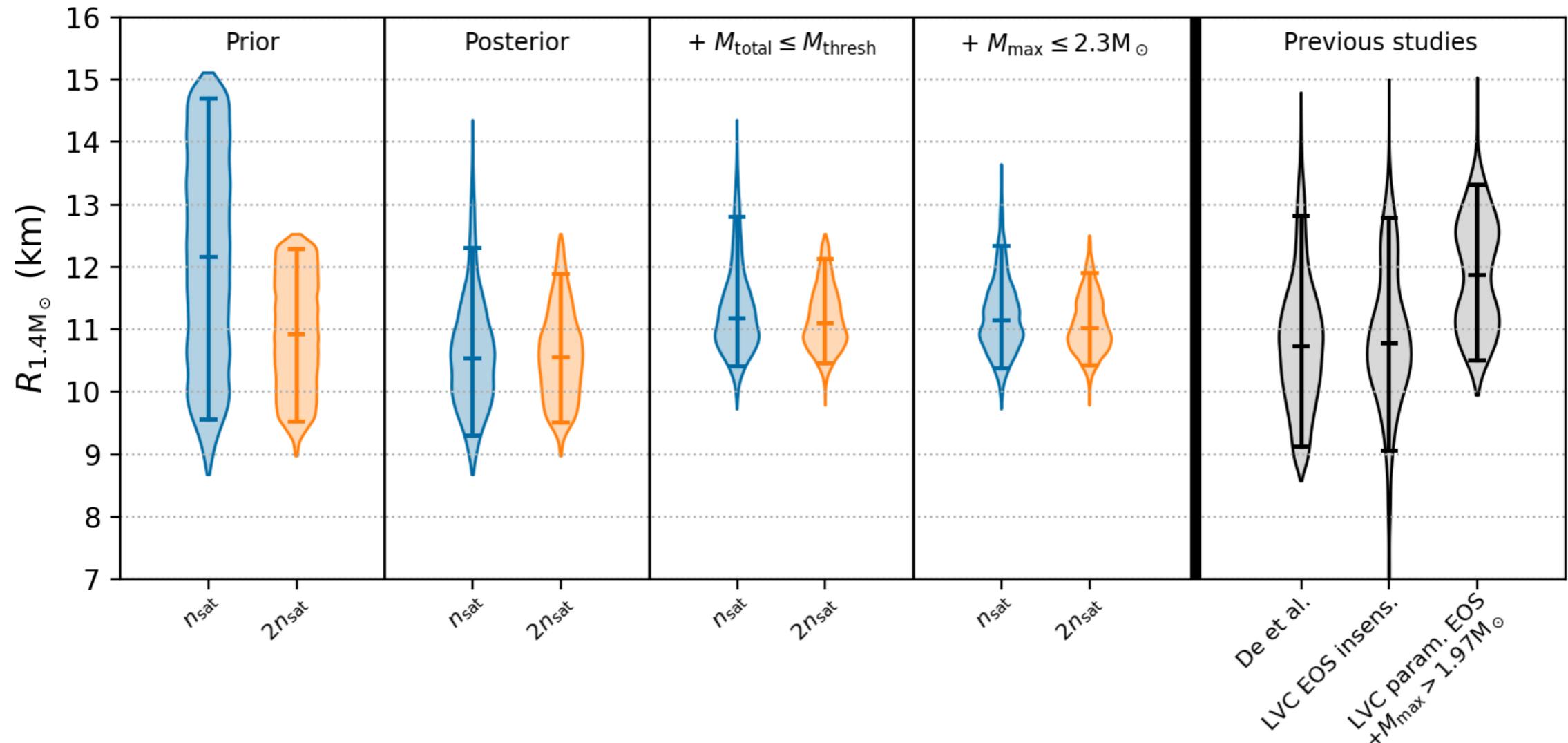
Constrained till n_{sat}



Constrained till $2n_{\text{sat}}$



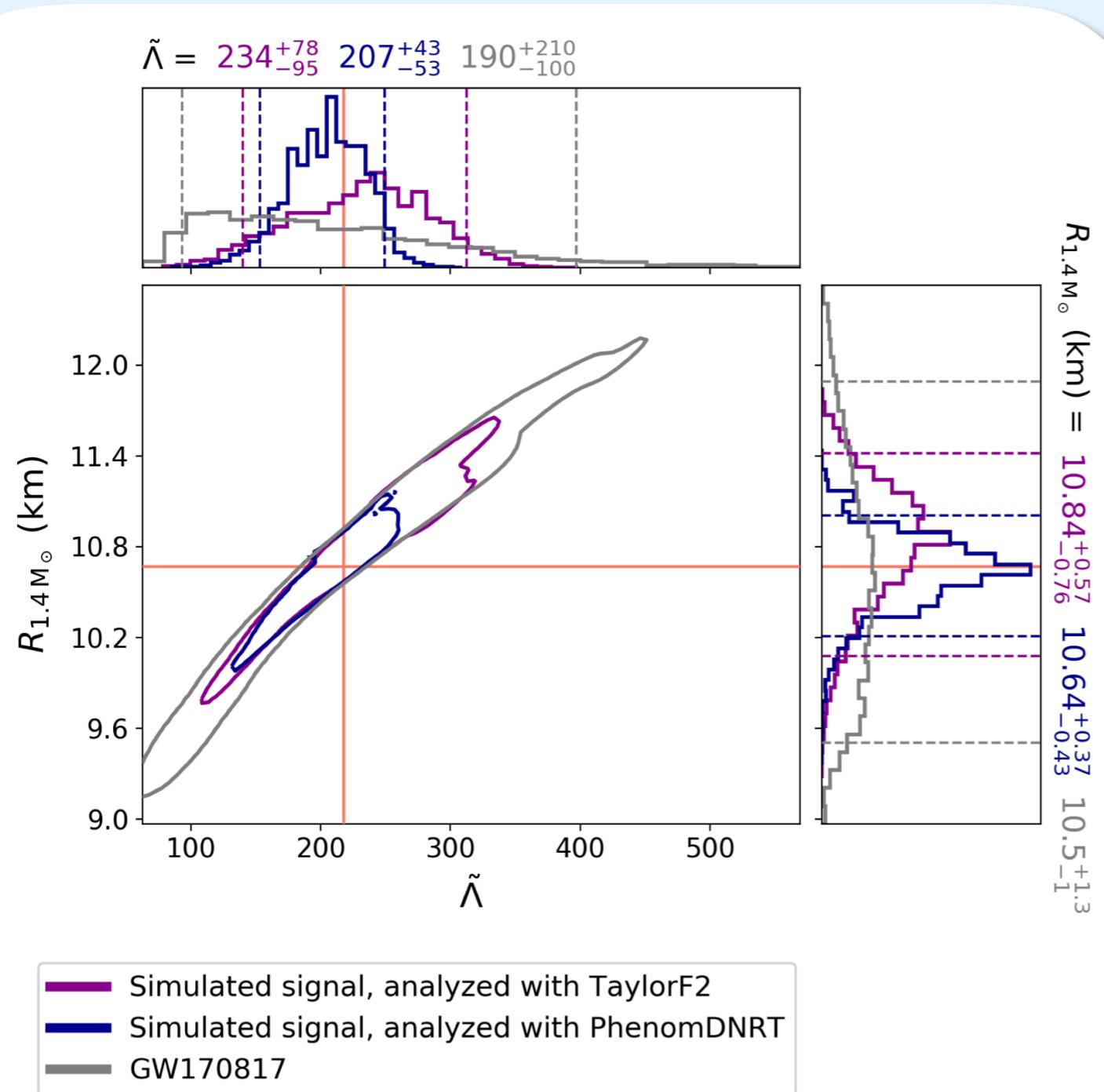
Radius measurements from analyses of GW170817



Capano, Tews, Brown, Margalit, **SD** et al (2019), arXiv:1908.10352

Posterior samples available at <https://github.com/sugwg/gw170817-eft-eos>

Prospects for improving measurements



- Using simulated signals at SNR ~ 100 we find $\sim 2.9\times$ improvement in measurement uncertainty
- Better waveform models further improves the measurements by 2x
- Gravitational waves alone will be able to constrain upper and lower bounds of tidal deformability and radii for loud SNR signals
- Low SNR signals would need combination of information from GW + EM + nuclear theory

Astrophysical source distinguishability using gravitational waves

- Perform Bayesian Inference analyses with BNS and BBH template waveforms models
- Compute Evidence for each model marginalizing over the full parameter space

$$\mathcal{Z} = \int p(\vec{d}(t) | \vec{\theta}, H) p(\vec{\theta} | H) d\theta_1 \dots d\theta_n$$

- Compute Bayes factor comparing BNS to BBH model $\mathcal{B} = \mathcal{Z}_{\text{BNS}} / \mathcal{Z}_{\text{BBH}}$

$\mathcal{B} \sim 1.0$ for GW170817. The event cannot distinguish between BNS and BBH mergers

$\mathcal{B} \sim 10^5$ for simulated GW170817 in aLIGO design sensitivity, SNR~100. Gravitational waves can reliably distinguish between BNS and BBH mergers for high-SNR observations.

Summary

- We analyzed GW170817 combining GW and EM observations and Chiral Effective Field theory
- We have provided the most stringent constraints on radii of neutron stars to date
- With Advanced LIGO design sensitivity
 - there are prospects of improving the current constraints by more than two times using gravitational waves alone
 - gravitational waves can be used to distinguish between BNS and BBH events

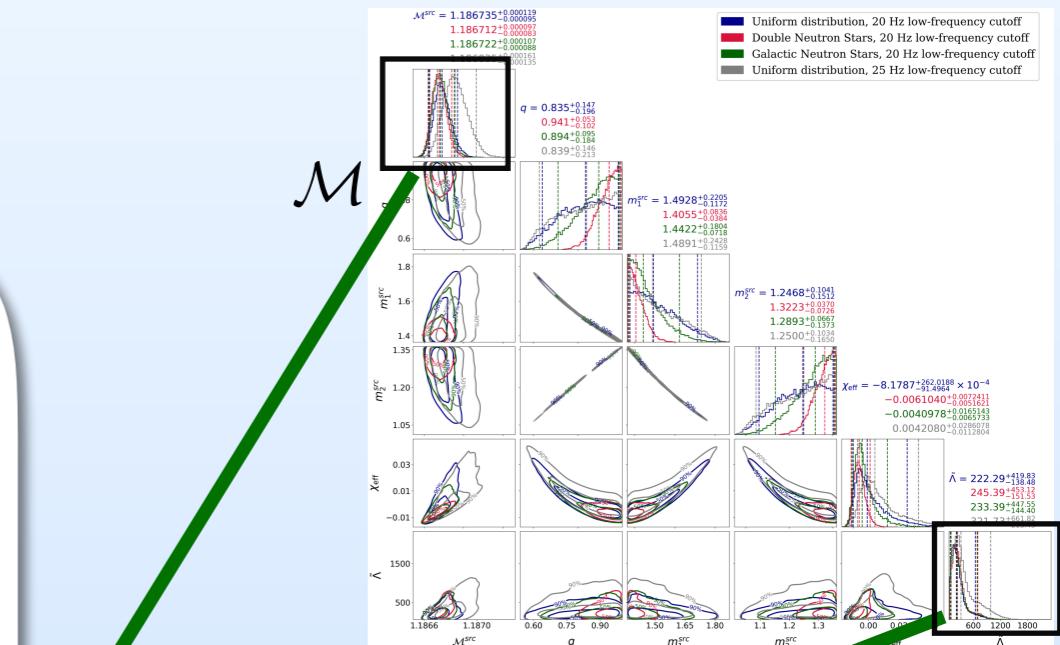
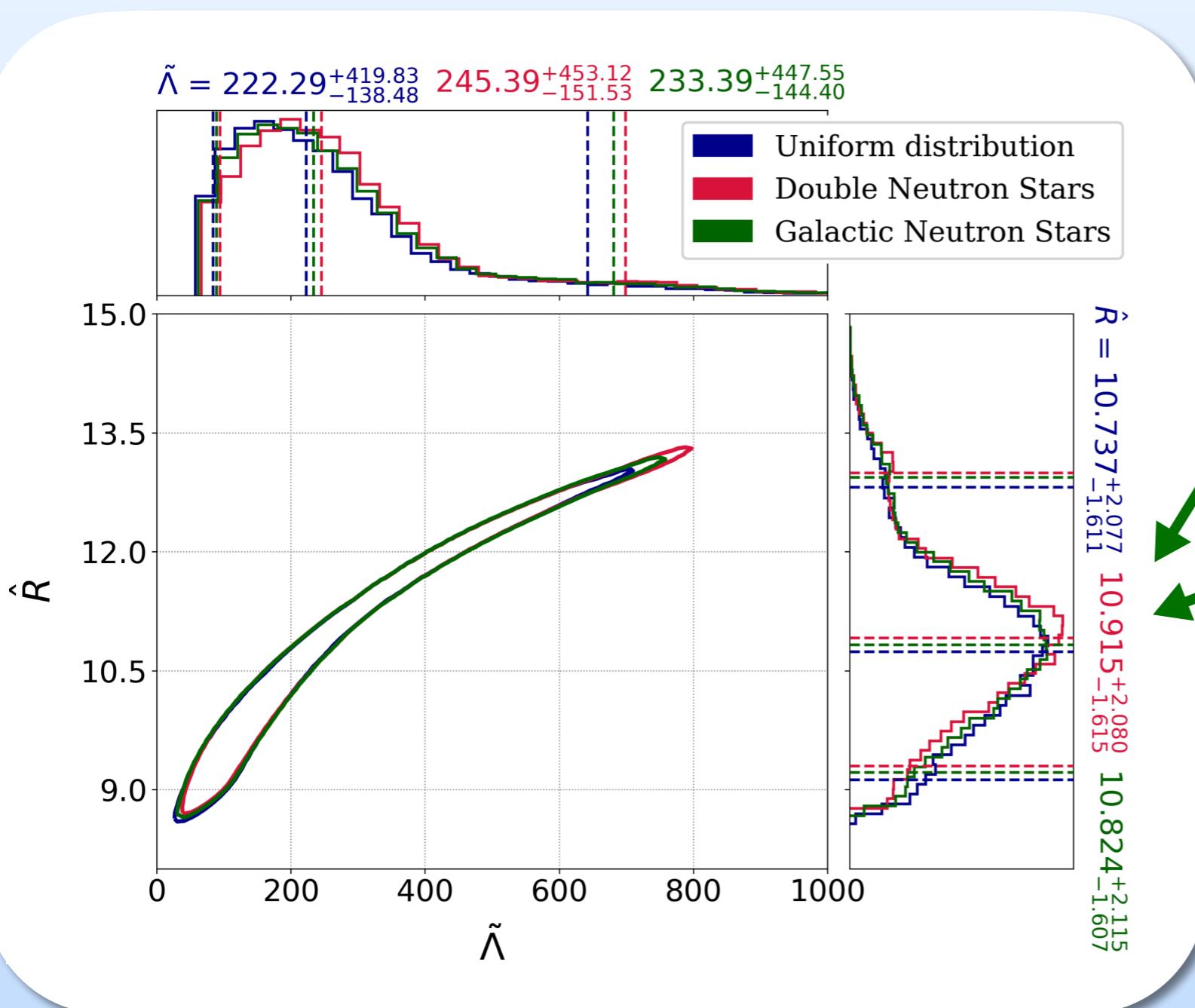
Thank you for your attention!

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Extra Slides

Measurement of the neutron star radius (1)

De, Finstad, Lattimer, Brown, Berger,
Biwer (2018)

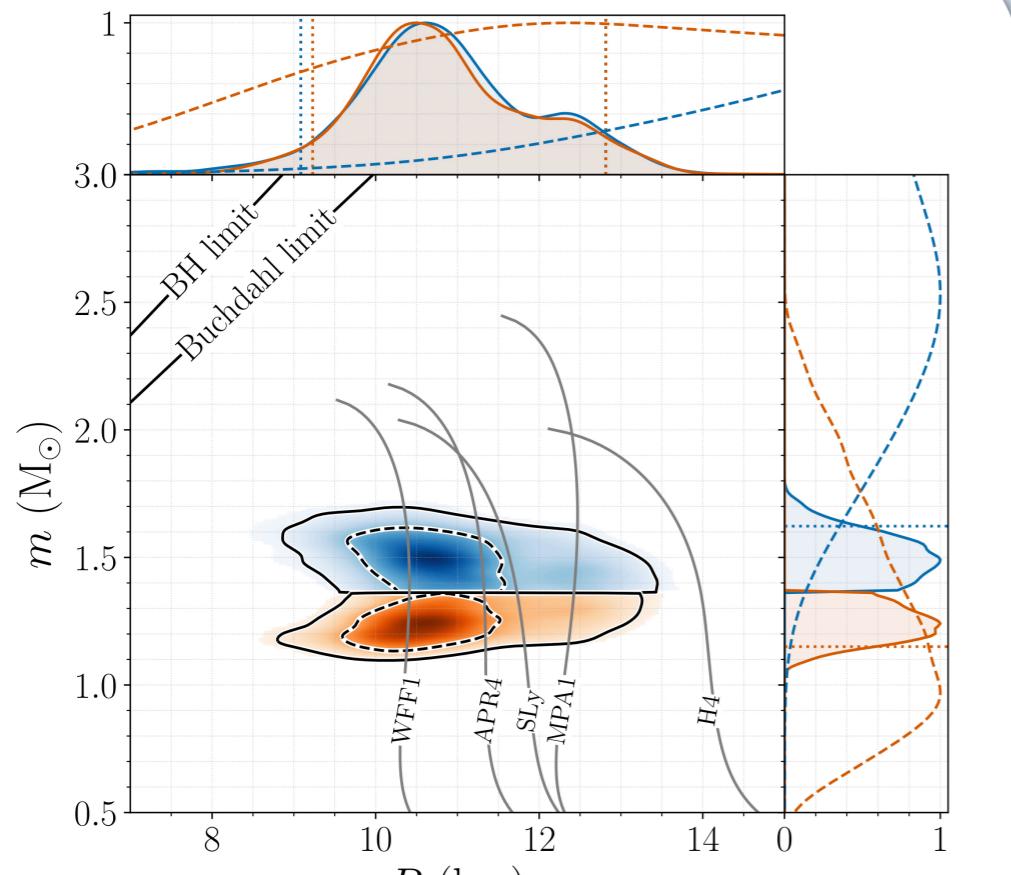


$$\hat{R} = (11.2 \pm 0.2) \frac{\mathcal{M}}{M_{\odot}} \left(\frac{\tilde{\Lambda}}{800} \right)^{1/6}$$

Measurement of the neutron star radius (2)

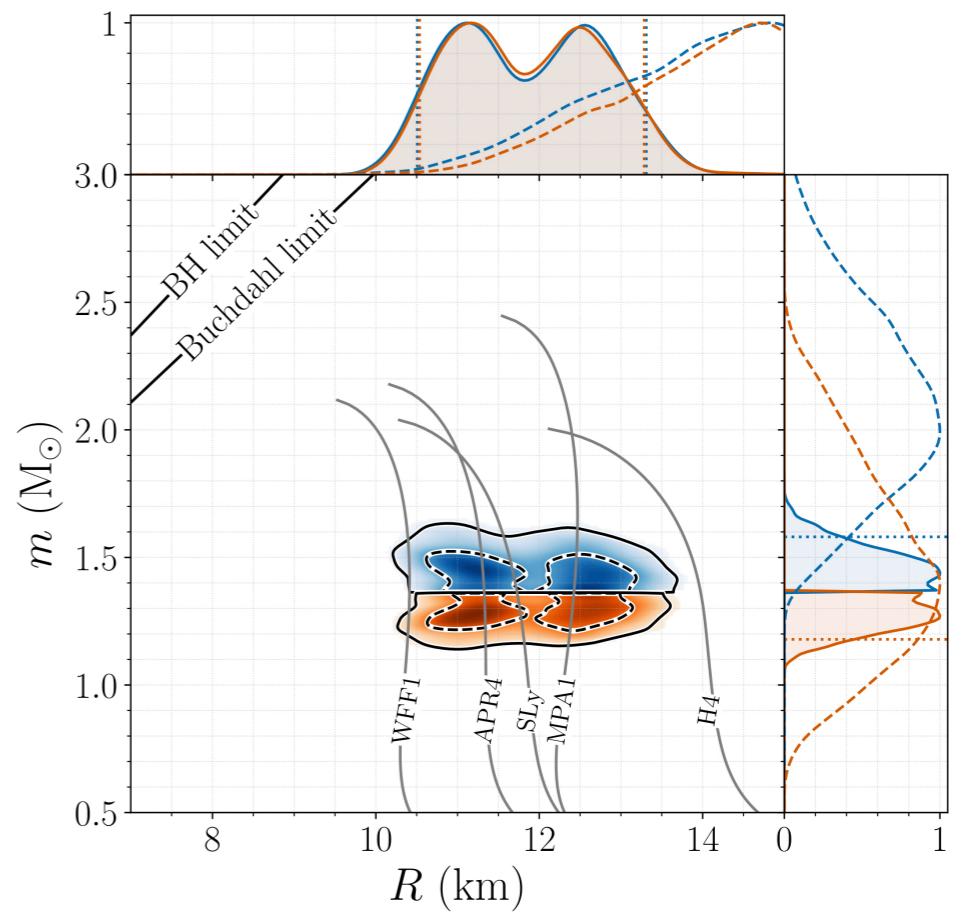
LIGO-Virgo Collaboration (2018)

EOS-insensitive relations :



$$(R_1, R_2) = (10.8_{-1.7}^{+2.0}, 10.7_{-1.5}^{+2.1})$$

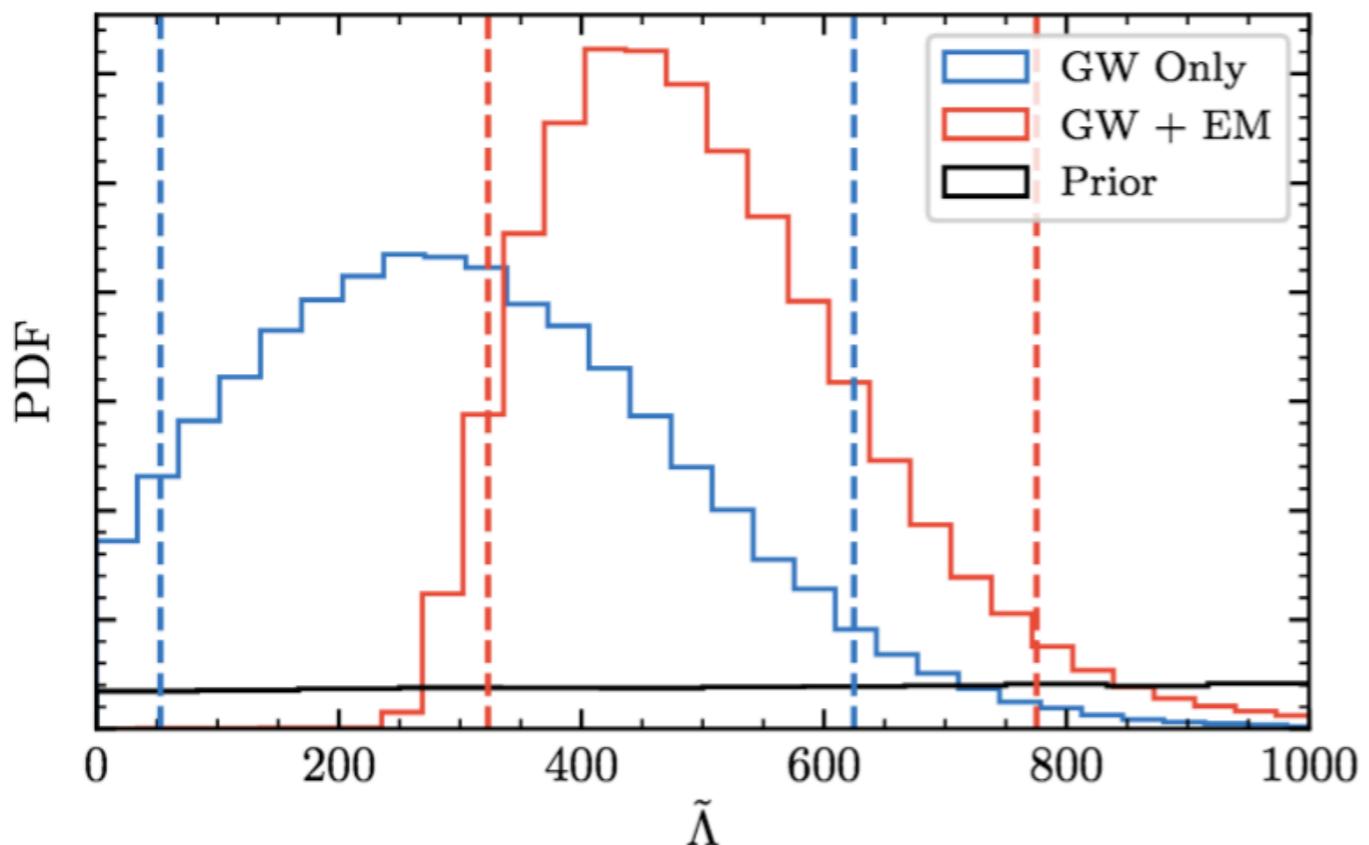
Parameterized EOS :



$$(R_1, R_2) = (11.9_{-1.4}^{+1.4}, 11.9_{-1.4}^{+1.4})$$

Measurement of the neutron star radius (3)

Radice, Dai (2018)



- EM data incorporation showed $\tilde{\Lambda} < 300$ excluded
 - Used $\hat{R}(\mathcal{M}, \tilde{\Lambda})$ expression from De et al to get
- GW : $R_{1.4} = 11.3^{+1.5}_{-2.8} \pm 0.2$ km
- GW + EM : $R_{1.4} = 12.2^{+1.0}_{-0.8} \pm 0.2$ km