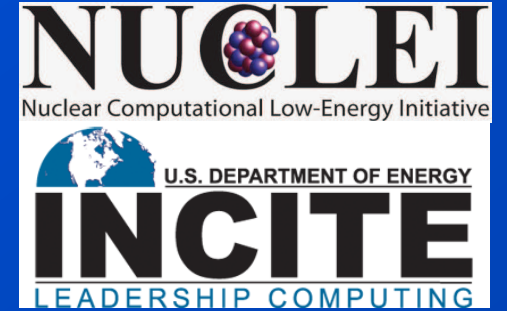


# Nuclear Properties with Chiral EFT Interactions And the Effects of Consistent Electroweak Interactions

James P. Vary, Iowa State University

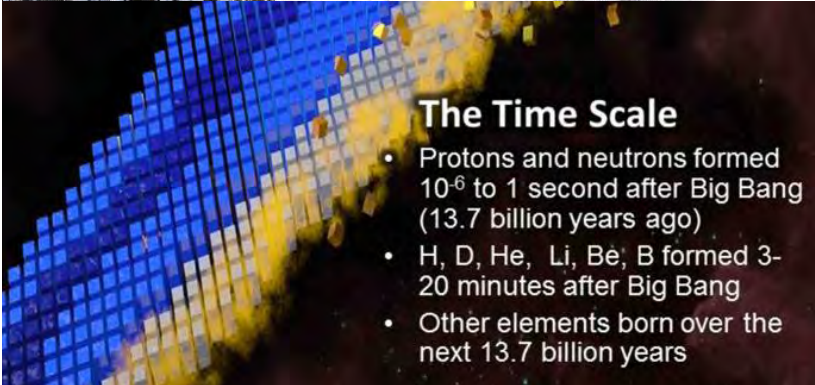
Multiparticle resonances in hadrons, nuclei and ultracold gases  
Hirschegg, January 15-19, 2018



## The Overarching Questions

- How did visible matter come into being and how does it evolve?
- How does subatomic matter organize itself and what phenomena emerge?
- Are the fundamental interactions that are basic to the structure of matter fully understood?
- How can the knowledge and technological progress provided by nuclear physics best be used to benefit society?

- NRC Decadal Study

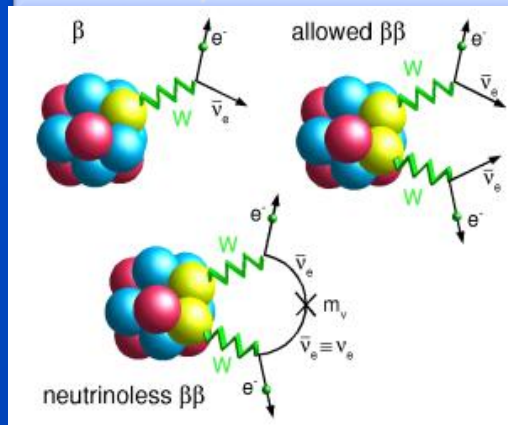


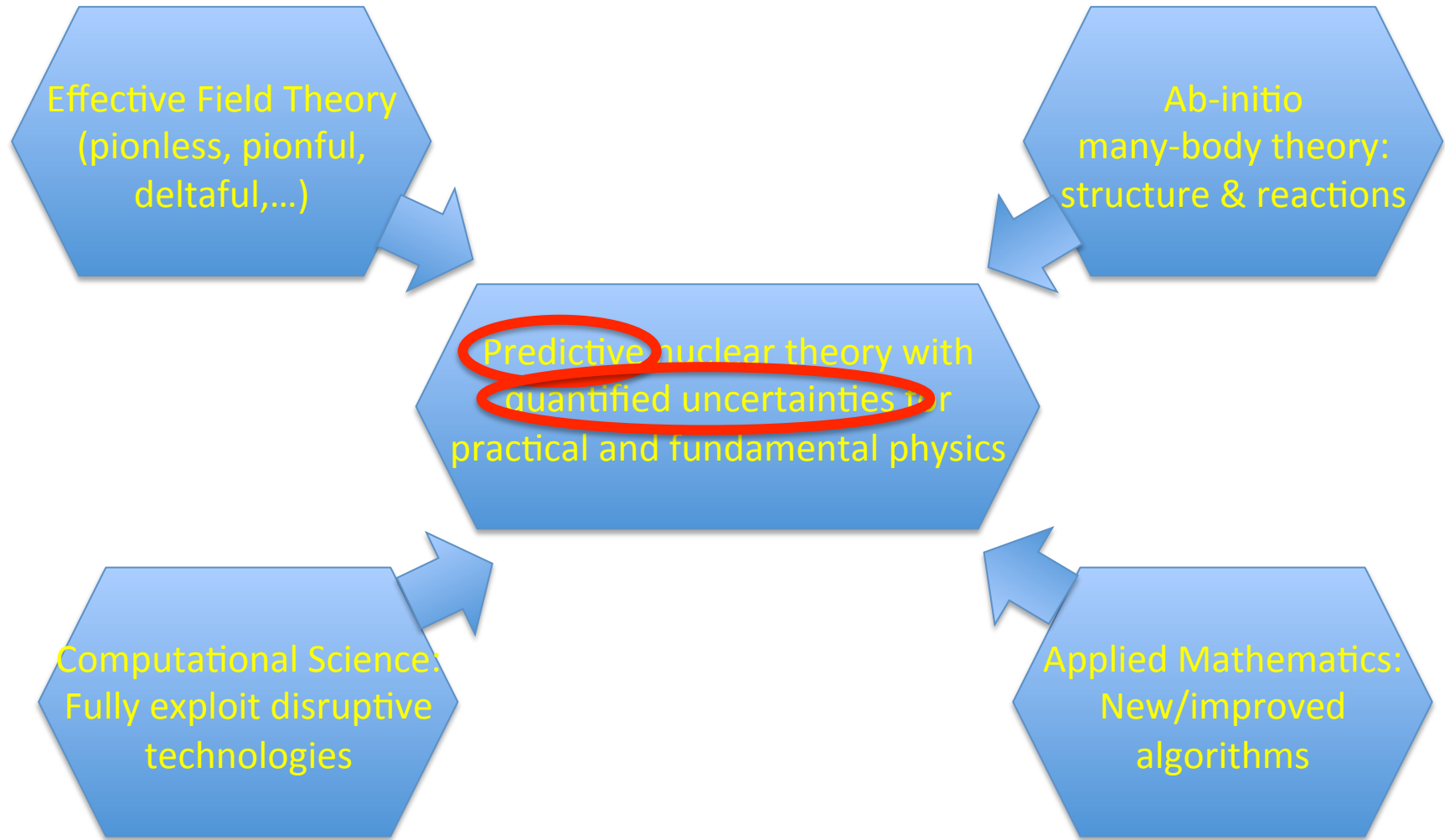
## The Time Scale

- Protons and neutrons formed  $10^{-6}$  to 1 second after Big Bang (13.7 billion years ago)
- H, D, He, Li, Be, B formed 3-20 minutes after Big Bang
- Other elements born over the next 13.7 billion years



## Topical Collaboration on Neutrinos and Fundamental Symmetries





# No-Core Configuration Interaction calculations

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Barrett, Navrátil, Vary, *Ab initio no-core shell model*, PPNP69, 131 (2013)

Given a Hamiltonian operator

$$\hat{\mathbf{H}} = \sum_{i < j} \frac{(\vec{p}_i - \vec{p}_j)^2}{2 m A} + \sum_{i < j} V_{ij} + \sum_{i < j < k} V_{ijk} + \dots$$

solve the eigenvalue problem for wavefunction of  $A$  nucleons

$$\hat{\mathbf{H}} \Psi(r_1, \dots, r_A) = \lambda \Psi(r_1, \dots, r_A)$$

- Expand eigenstates in basis states  $|\Psi\rangle = \sum a_i |\Phi_i\rangle$
  - Diagonalize Hamiltonian matrix  $H_{ij} = \langle \Phi_j | \hat{\mathbf{H}} | \Phi_i \rangle$
  - No Core Full Configuration (NCFC) – All  $A$  nucleons treated equally
  - Complete basis  $\rightarrow$  exact result
  - In practice
    - truncate basis
    - study behavior of observables as function of truncation
-

## Basis expansion $\Psi(r_1, \dots, r_A) = \sum a_i \Phi_i(r_1, \dots, r_A)$

---

- Many-Body basis states  $\Phi_i(r_1, \dots, r_A)$  Slater Determinants
- Single-Particle basis states  $\phi_\alpha(r_k)$  with  $\alpha = (n, l, s, j, m_j)$
- Radial wavefunctions: Harmonic Oscillator (HO), natural orbitals, Woods-Saxon, Coulomb-Sturmian, Complex Scaled HO, Berggren, . . .
- $M$ -scheme: Many-Body basis states eigenstates of  $\hat{J}_z$

$$\hat{J}_z |\Phi_i\rangle = M |\Phi_i\rangle = \sum_{k=1}^A m_{ik} |\Phi_i\rangle$$

- $N_{\max}$  truncation: Many-Body basis states satisfy

$$\sum_{\alpha \text{ occ.}}^A (2n + l)_\alpha \leq N_0 + N_{\max}$$

$N_{\max}$  runs from zero to computational limit.  
( $N_{\max}, \hbar\Omega$ ) fix HO basis

- Alternatives:

- Full Configuration Interaction (single-particle basis truncation)
  - Importance Truncation Roth, PRC79, 064324 (2009)
  - No-Core Monte-Carlo Shell Model Abe *et al*, PRC86, 054301 (2012)
  - SU(3) Truncation Dytrych *et al*, PRL111, 252501 (2013)
-

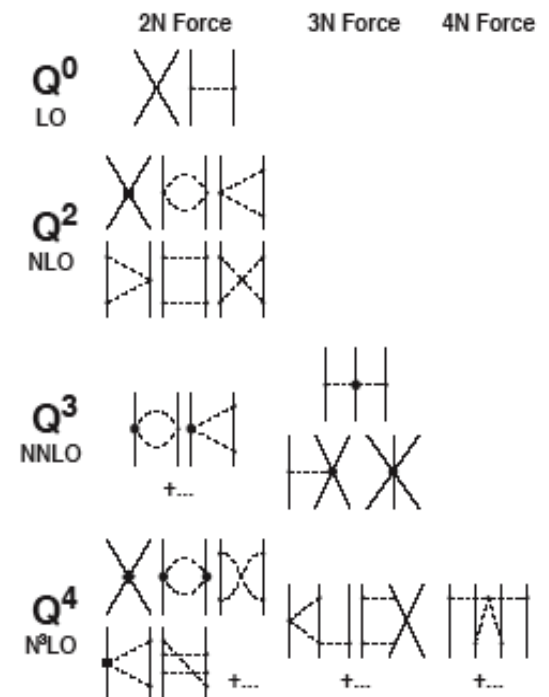
# Nuclear interaction

Nuclear potential not well-known,  
though in principle calculable from QCD

$$\hat{H} = \hat{T}_{\text{rel}} + \sum_{i < j} V_{ij} + \sum_{i < j < k} V_{ijk} + \dots$$

In practice, alphabet of realistic potentials

- Argonne potentials: AV8', AV18
  - plus Urbana 3NF (UIX)
  - plus Illinois 3NF (IL7)
- Bonn potentials
- Chiral NN interactions
  - plus chiral 3NF, ideally to the same order
- JISP16
- Daejeon16
- ...

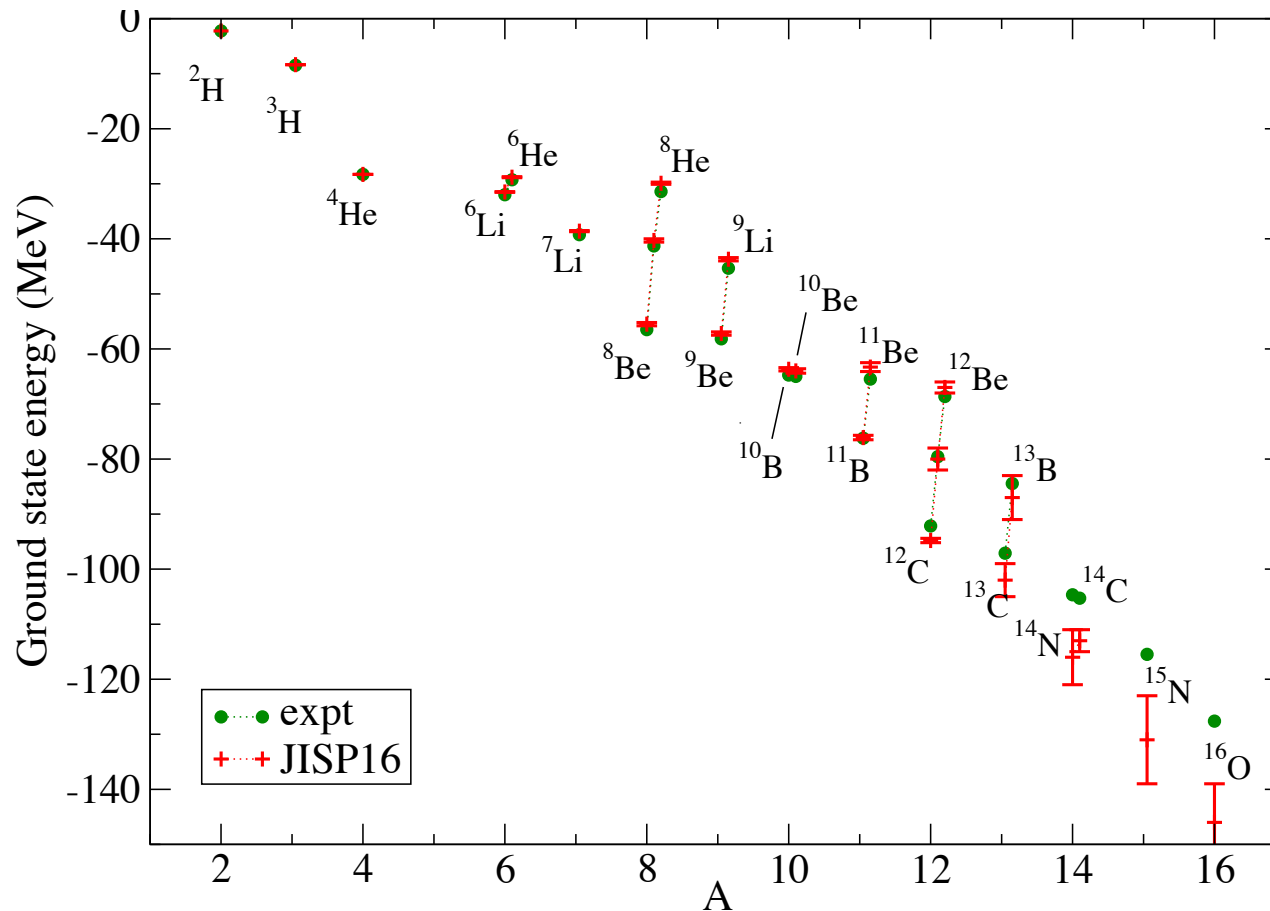


Major development during the past 5-10 years:  
High-precision ab initio calculations now used to  
“discover” the correct strong NN+NNN interaction

# Ground state energy of *p*-shell nuclei with JISP16

Compare theory and experiment for 24 nuclei

Maris, Vary, IJMPE22, 1330016 (2013)



- $^{10}\text{B}$  – most likely JISP16 produces correct  $3^+$  ground state, but extrapolation of  $1^+$  states not reliable due to mixing of two  $1^+$  states
- $^{11}\text{Be}$  – expt. observed parity inversion within error estimates of extrapolation
- $^{12}\text{B}$  and  $^{12}\text{N}$  – unclear whether gs is  $1^+$  or  $2^+$  (expt. at  $E_x = 1$  MeV) with JISP16

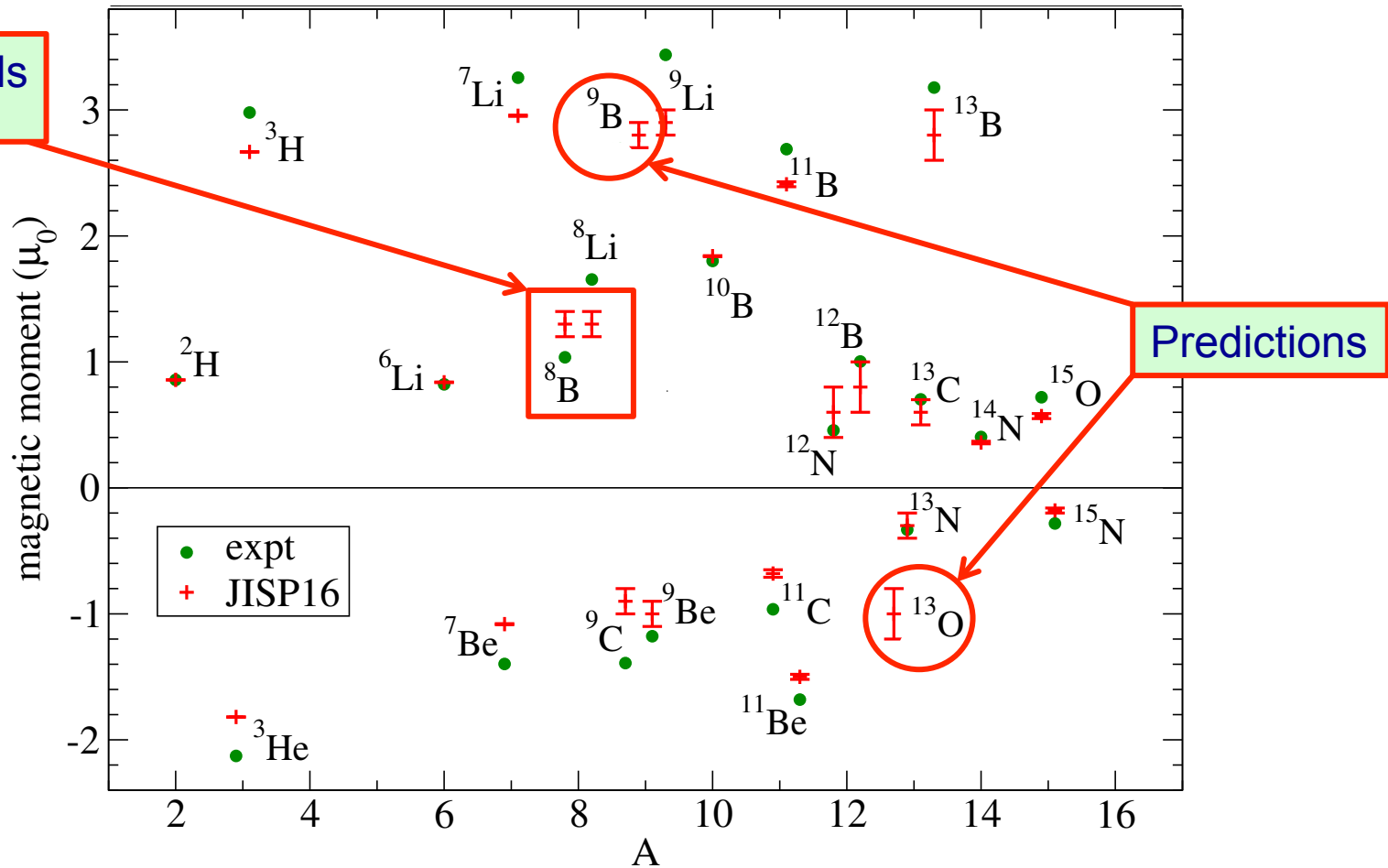
Many additional nuclear states: A.M. Shirokov, V.A. Kulikov, P. Maris and J.P. Vary, in Nucleon-Nucleon and Three-Nucleon Interactions, edited by L.D. Blokhintsev and I.I. Strakovsky, (Nova Science, 2014), Chapter 8, p. 231.

# Ground state magnetic moments with JISP16

Compare theory and experiment for 23 magnetic moments Maris, Vary, IJMPE22, 1330016 (2013)

$$\mu = \frac{1}{J+1} \left( \langle \mathbf{J} \cdot \mathbf{L}_p \rangle + 5.586 \langle \mathbf{J} \cdot \mathbf{S}_p \rangle - 3.826 \langle \mathbf{J} \cdot \mathbf{S}_n \rangle \right) \mu_0$$

Only 8B needs reduction



● Good agreement with data, given that we do not have any meson-exchange currents

GFMC with AV18 + IL7 interaction  
 15 magnetic moments compared between theory and experiment  
 Two-body currents tend to enhance the magnetic moments (exc.  $^8\text{B}$ )

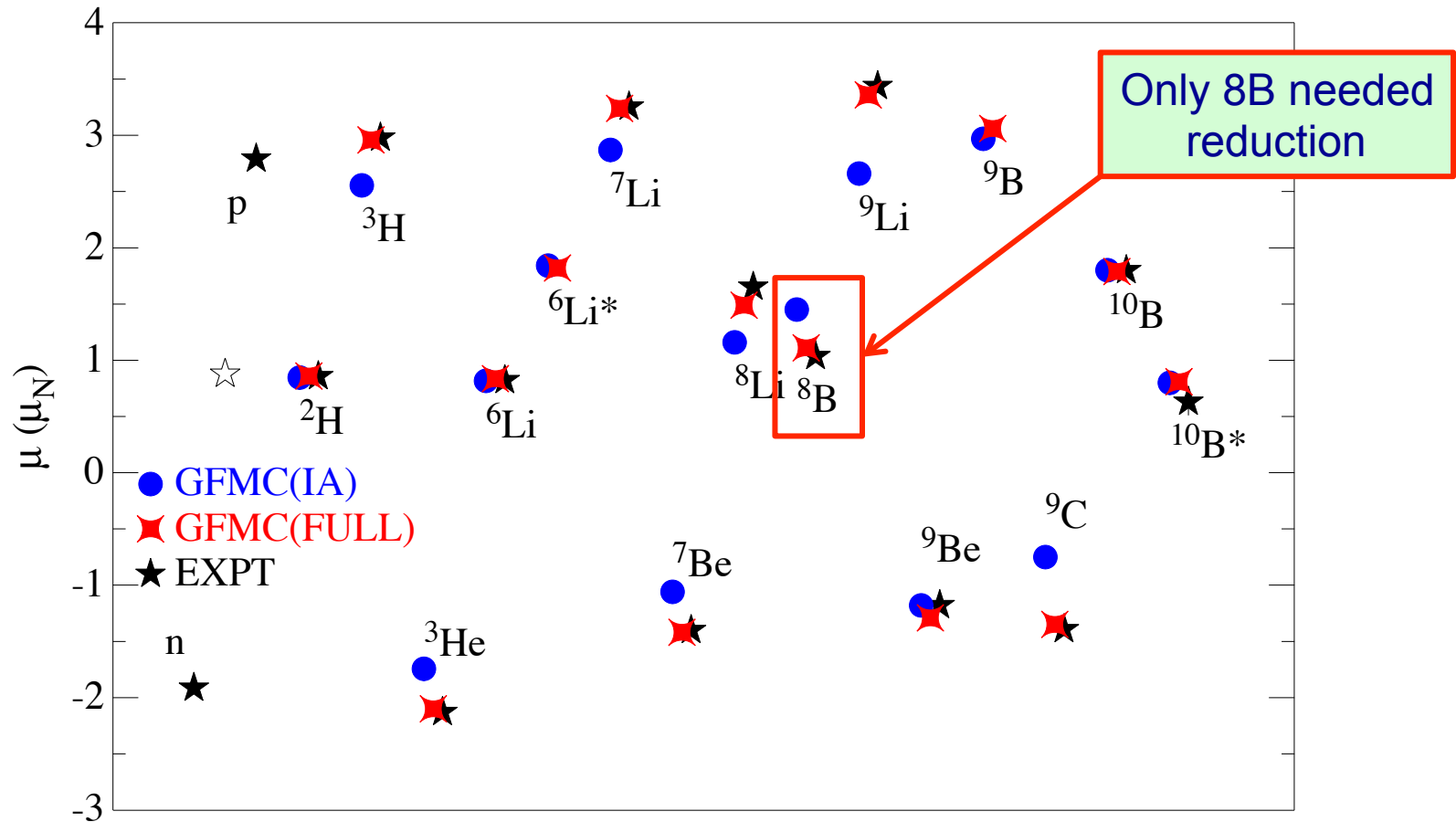
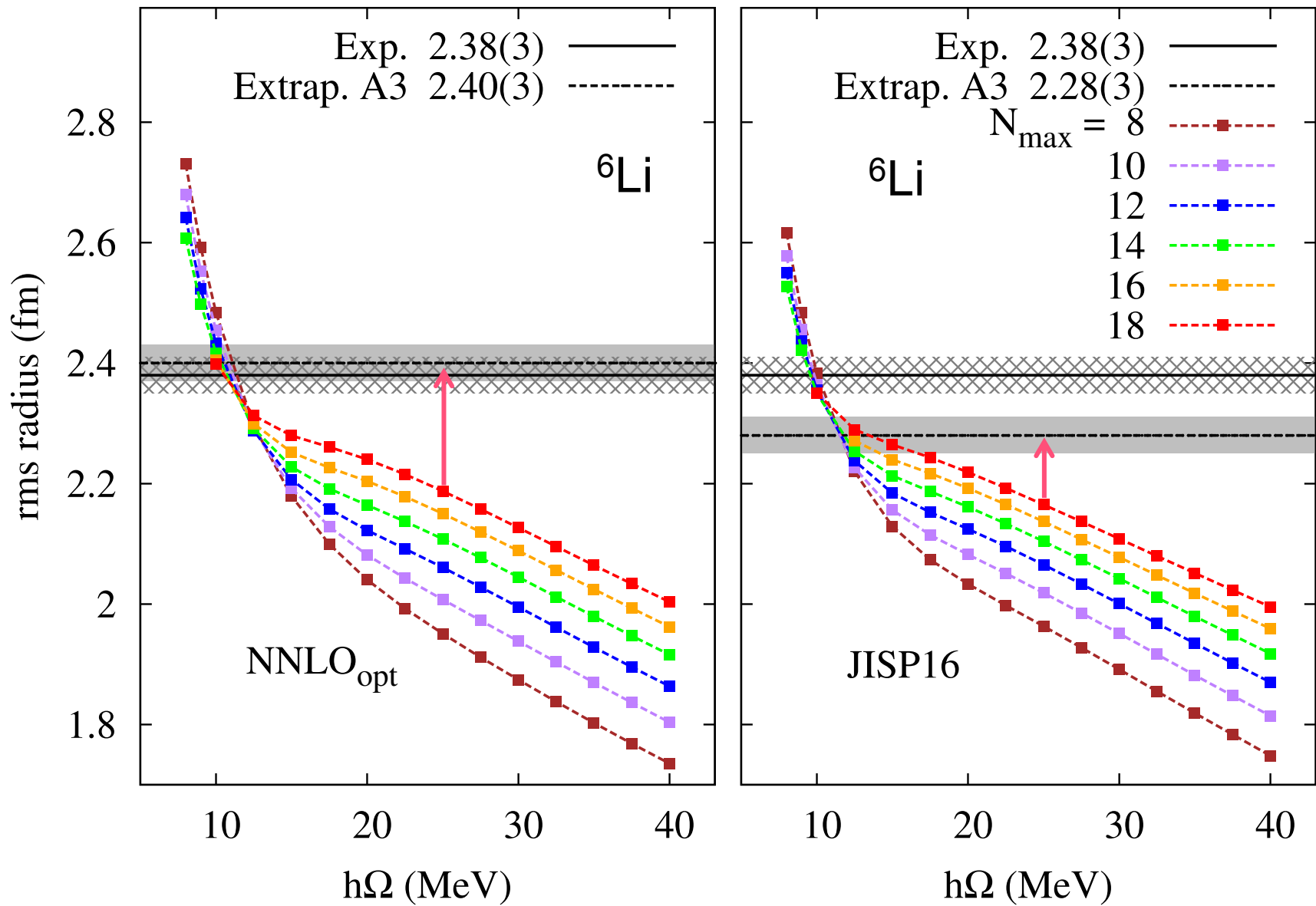


Figure from Rocco Schiavilla

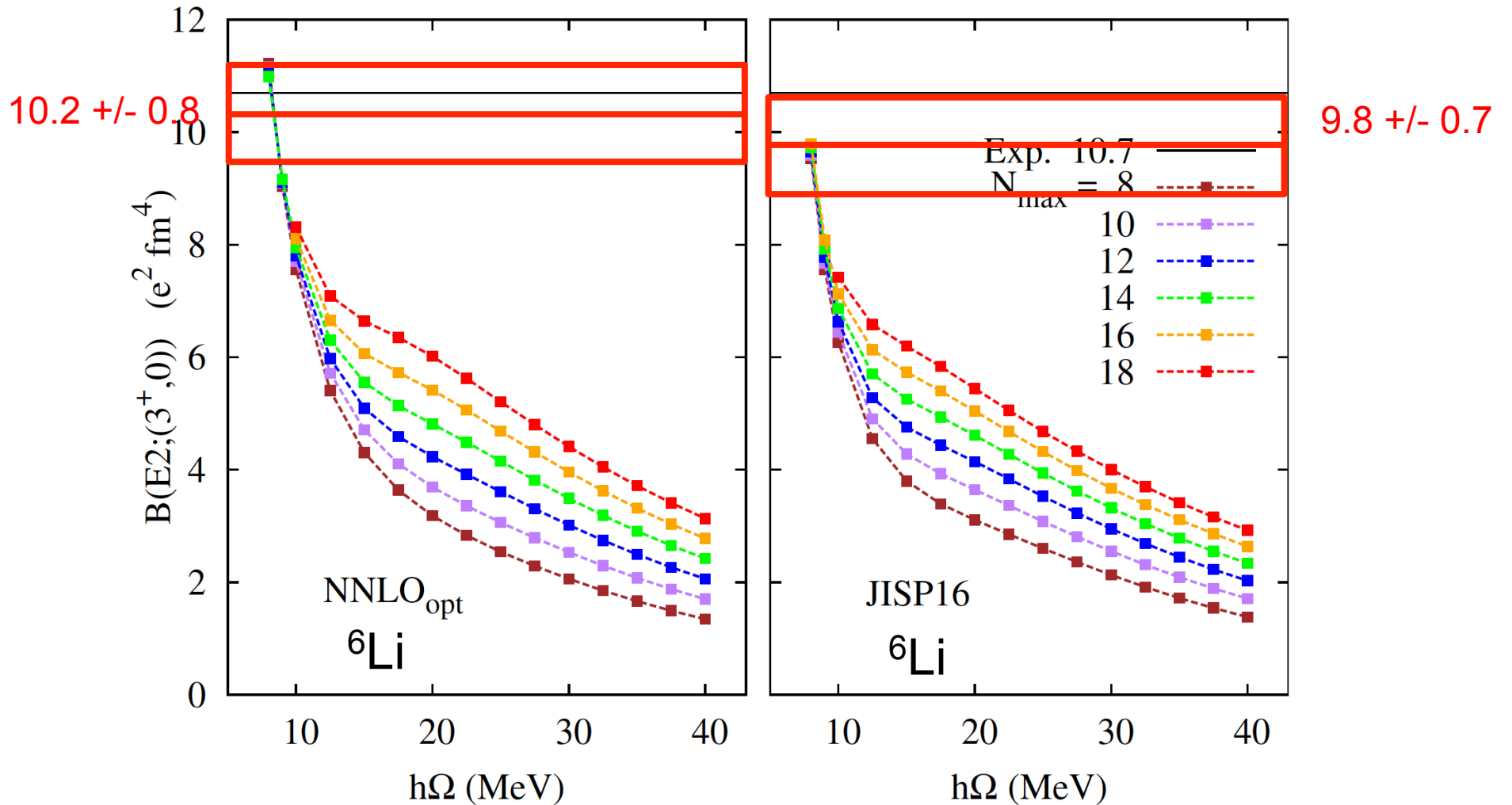




Ik Jae Shin, Youngman Kim, Pieter Maris, James P. Vary, Christian Forssen, Jimmy Rotureau and Nicolas Michel, J. Phys. G., Nuc. Part. Phys. 44, 075103 (2017); arXiv:1605.02819

=> Apply extrapolation method:

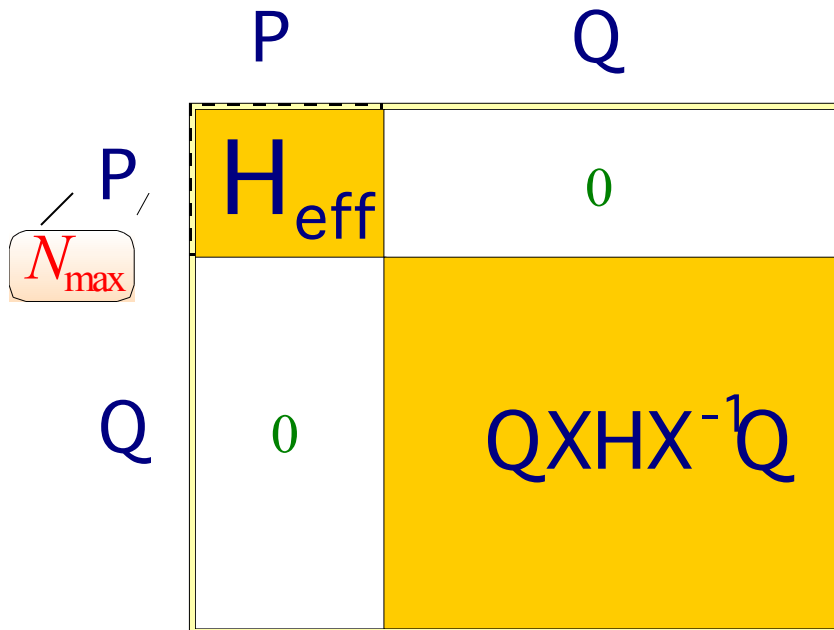
D. Odell, T. Papenbrock and L. Platter, Phys. Rev. C93, 044331(2016)



Ik Jae Shin, Youngman Kim, Pieter Maris, James P. Vary, Christian Forssen, Jimmy Rotureau and Nicolas Michel, J. Phys. G., Nuc. Part. Phys. 44, 075103 (2017); arXiv:1605.02819

# Effective Hamiltonian in the NCSM

## Okubo-Lee-Suzuki (OLS) renormalization scheme



$$H : E_1, E_2, E_3, \dots, E_{d_P}, \dots, E_\infty$$

$$H_{\text{eff}} : E_1, E_2, E_3, \dots, E_{d_P}$$

$$QXHx^{-1}P = 0$$

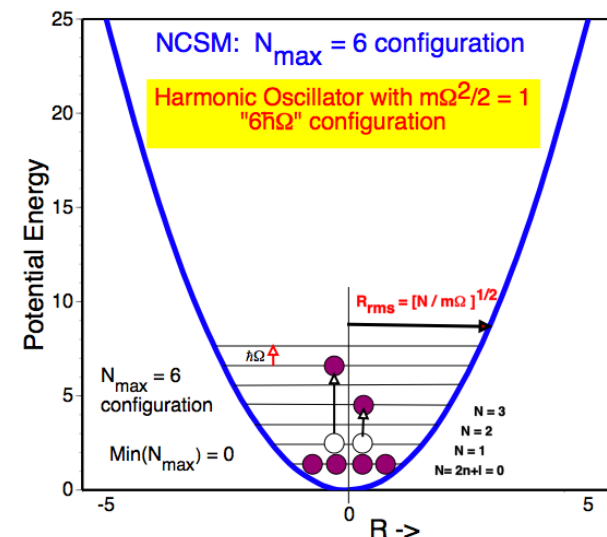
$$H_{\text{eff}} = PXHX^{-1}P$$

model space dimension

unitary  $X = \exp[-\arctan h(\omega^+ - \omega)]$

- $n$ -body cluster approximation,  $2 \leq n \leq A$
- $H_{\text{eff}}^{(n)}$   $n$ -body operator
- Two ways of convergence:
  - For  $P \rightarrow 1$   $H_{\text{eff}}^{(n)} \rightarrow H$
  - For  $n \rightarrow A$  and fixed  $P$ :  $H_{\text{eff}}^{(n)} \rightarrow H_{\text{eff}}$

Adapted from Petr Navratil



## Outline of the OLS process

$$UHU^\dagger = U[T + V]U^\dagger = H_d$$

$$H_{\text{eff}} = U_{OLS}HU_{OLS}^\dagger = PH_{\text{eff}}P = P[T + V_{\text{eff}}]P$$

$$U^P = PUP$$

$$\tilde{U}^P = P\tilde{U}^P P = \frac{U^P}{\sqrt{U^{P\dagger}U^P}}$$

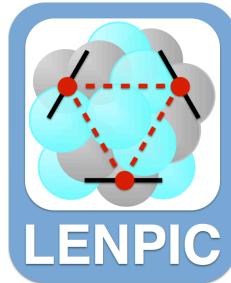
$$H_{\text{eff}} = \tilde{U}^{P\dagger}H_d\tilde{U}^P = \tilde{U}^{P\dagger}UHU^\dagger\tilde{U}^P = P[T + V_{\text{eff}}]P$$

$$O_{\text{eff}} = \tilde{U}^{P\dagger}UOU^\dagger\tilde{U}^P = P[O_{\text{eff}}]P$$

$$U_{OLS} = \tilde{U}^{P\dagger}U$$

# Calculation of three-body forces at N<sup>3</sup>LO

Low  
Energy  
Nuclear  
Physics  
International  
Collaboration



J. Golak, R. Skibinski,  
K. Tolponicki, H. Witala



E. Epelbaum, H. Krebs



A. Nogga



R. Furnstahl



S. Binder, A. Calci, K. Hebeler,  
J. Langhammer, R. Roth



P. Maris, J. Vary



H. Kamada

## Goal

Calculate matrix elements of 3NF in a partial-wave decomposed form which is suitable for different few- and many-body frameworks

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## Challenge

Due to the large number of matrix elements, the calculation is extremely expensive.

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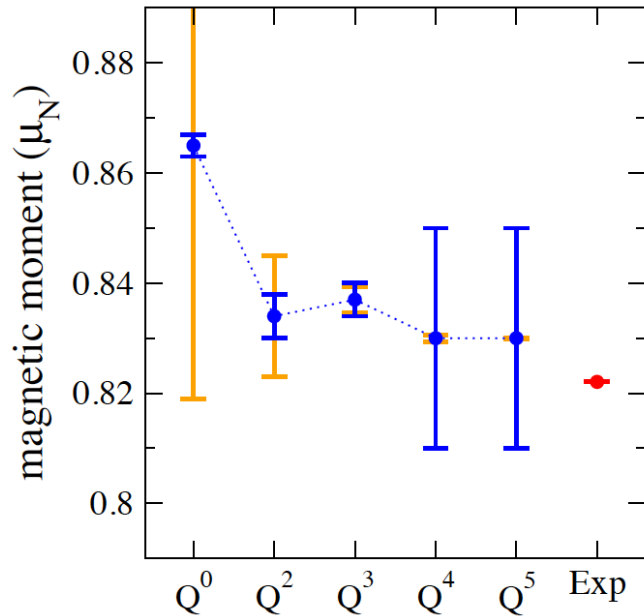
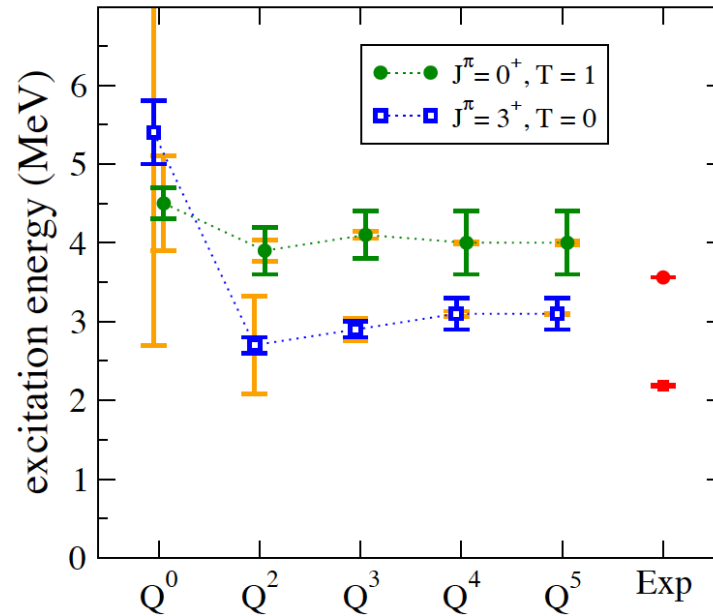
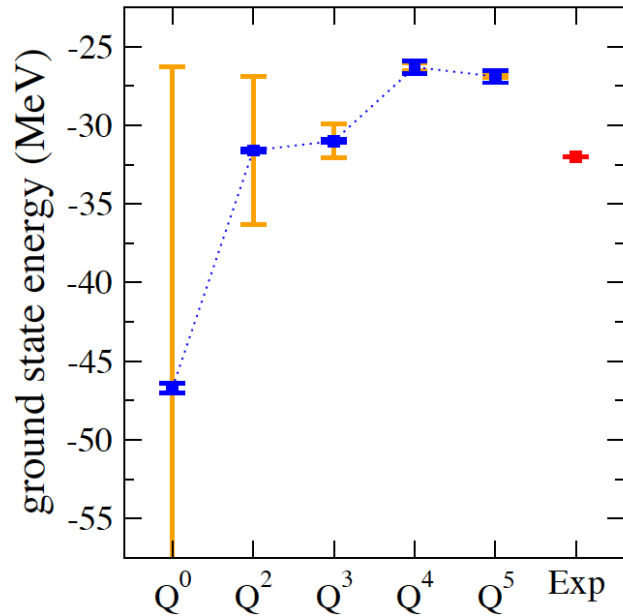
## Strategy

Develop an efficient code which allows to treat arbitrary local 3N interactions.  
(Krebs and Hebeler)

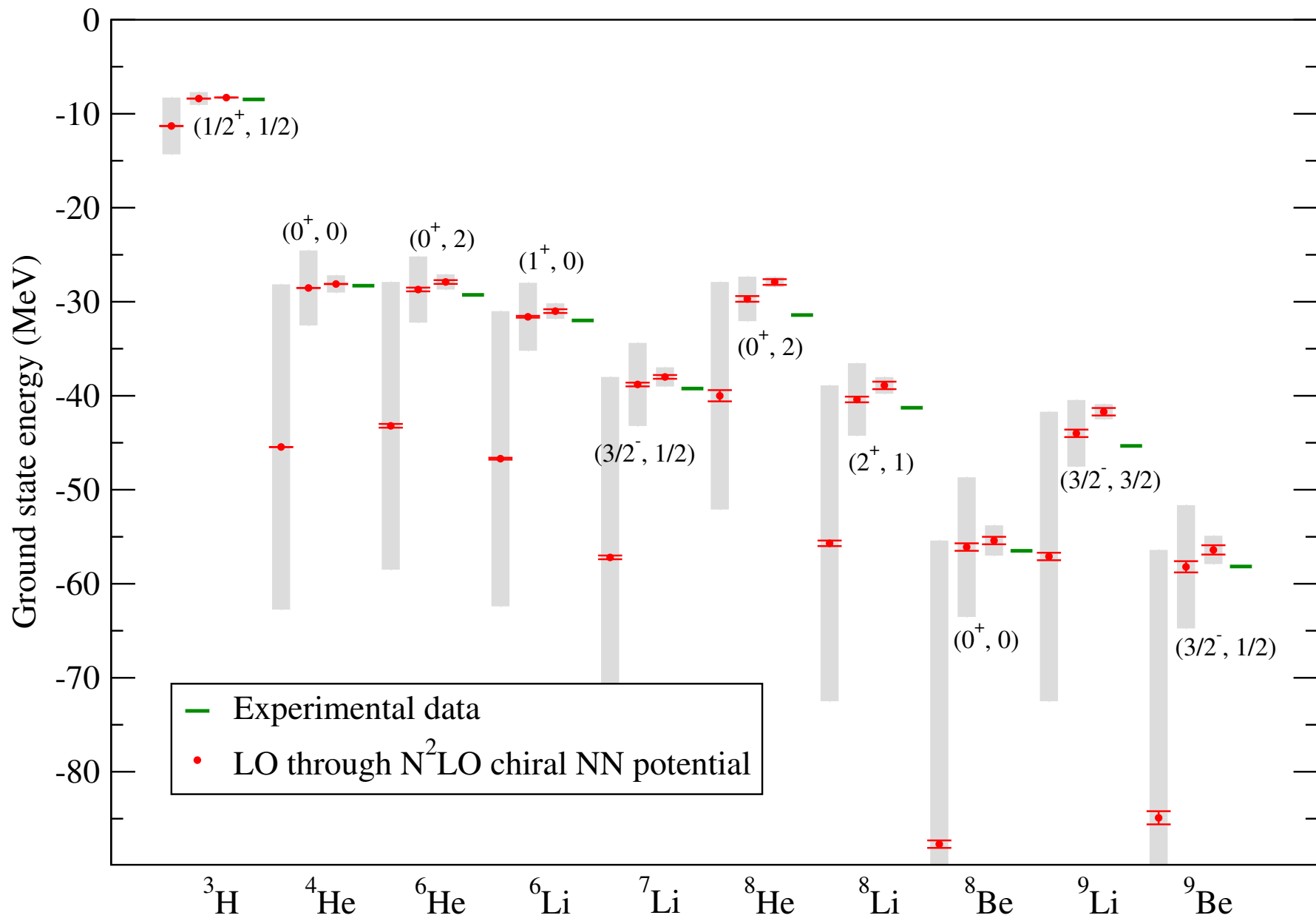
Initial LENPIC Collaboration results: Chiral NN results for  ${}^6\text{Li}$  by Chiral order

Orange: Chiral order uncertainties; Blue/Green: Many-body method uncertainties

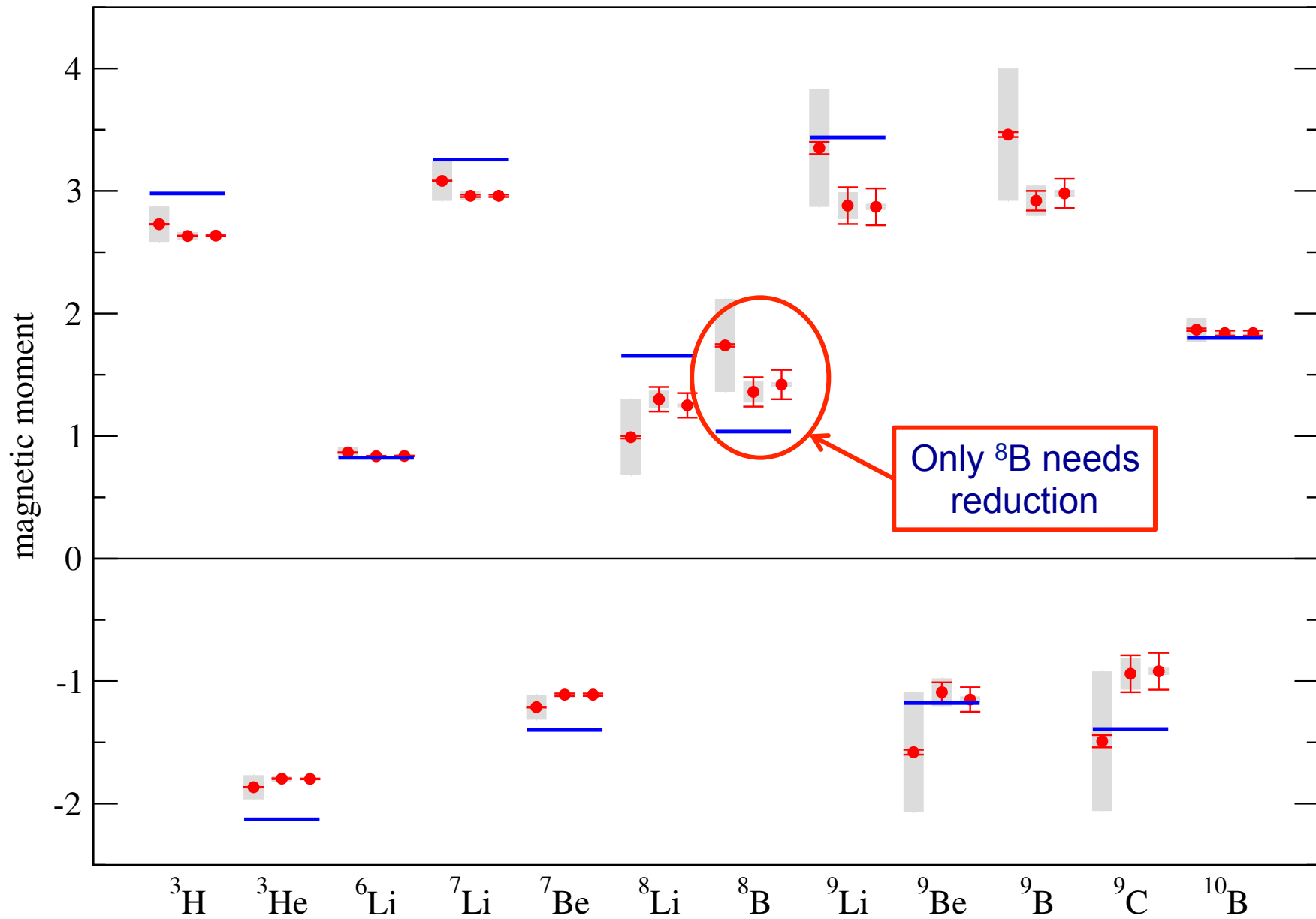
S. Binder, et al, Phys. Rev. C **93**, 044002 (2016); arXiv:1505.07218



- Ground state energy  
similar behavior as  ${}^3\text{H}$  and  ${}^4\text{He}$
- Open question:
  - Are chiral uncertainty estimates applicable to excitation energies?
- Need compatible chiral 3NFs
- Need chiral expansion conserved current operator



Preliminary LENPIC results with Chiral NN only and  $R = 1.0$  fm, IA for operator  
S. Binder, et al., LENPIC Collaboration, in preparation





Consider two nucleons as a model problem with  $V = \text{LENPIC}$  chiral NN solved in the harmonic oscillator basis with  $\hbar\Omega = 5, 10$  and  $20$  MeV. Also, consider the role of an added harmonic oscillator quasipotential

Hamiltonian #1  $H = T + V$

Hamiltonian #2  $H = T + U_{\text{osc}}(\hbar\Omega_{\text{basis}}) + V$

Other observables:

Root mean square radius	R
Magnetic dipole operator	M1
Electric dipole operator	E1
Electric quadrupole moment	Q
Electric quadrupole transition	E2
Gamow-Teller	GT
Neutrinoless double-beta decay	$M(0\nu 2\beta)$

Dimension of the “full space” is 400 for all results depicted here

Deuteron gs energy:  
truncation vs OLS

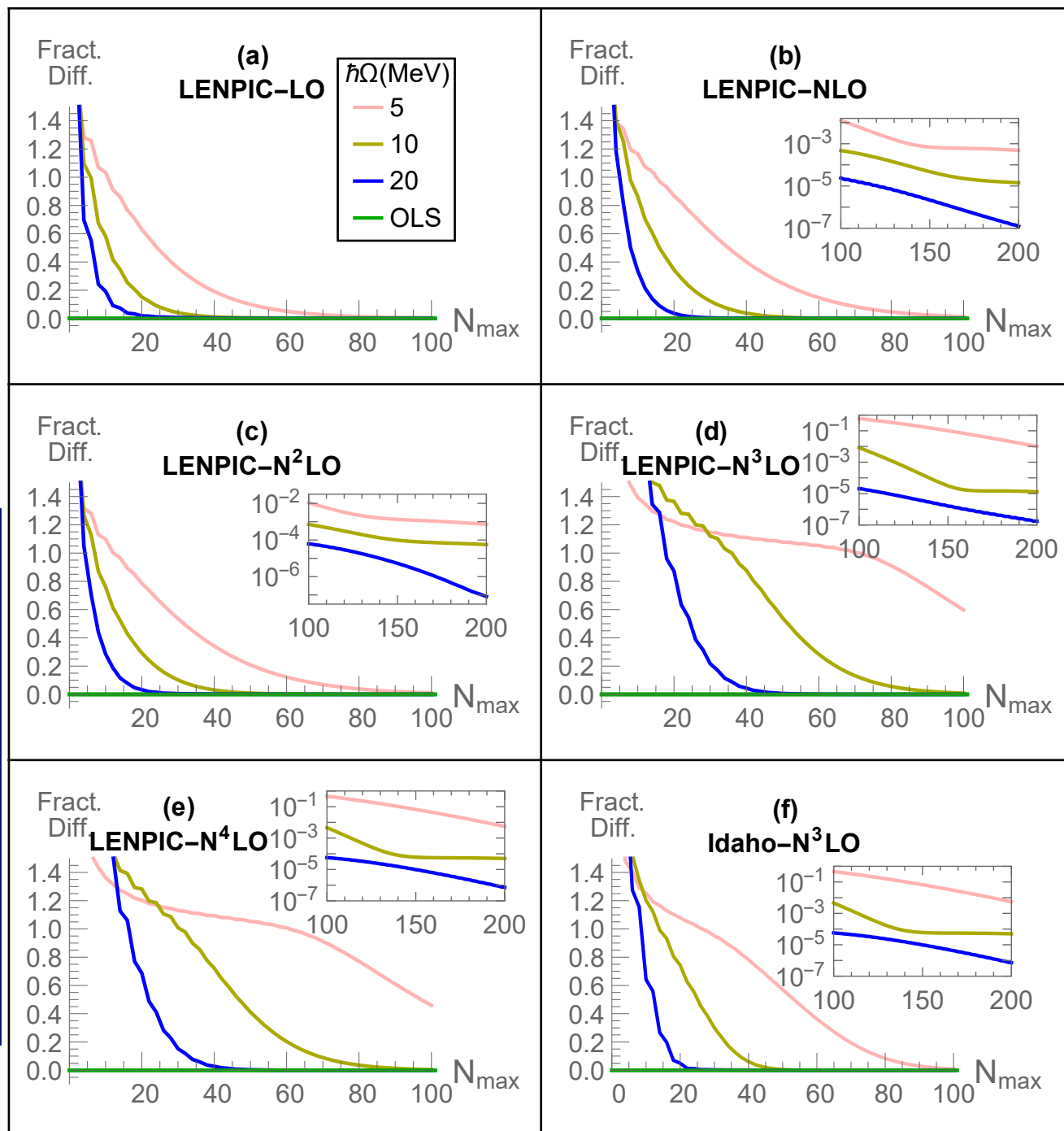
$$\text{Fract. Diff.} = \frac{E_{\text{model}} - E_{\text{exact}}}{|E_{\text{exact}}|}$$

Insets: Semilog plots  
of high  $N_{\text{max}}$  region

OLS gives exact results  
for all cases (green lines  
at Fract. Diff. = 0)

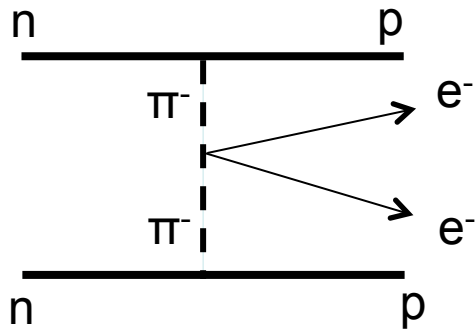
Convergence patterns  
sensitive to chiral order

Even unbound cases  
(Fract. Diff. > 1.0)  
are accurately treated  
with OLS



## Consider a 2-body contribution within EFT to $0\nu\beta\beta$ -decay at $N^2\text{LO}$

G. Prézeau, M. Ramsey-Musolf and P. Vogel, Phys. Rev. D 68, 034016 (2003)



$$M^0 = \langle \Psi_{A,Z+2} | \sum_{ii} \frac{R}{r_{ij}} [F_1(x_{ij}) \vec{\sigma}_i \vec{\sigma}_j + F_2(x_{ij}) T_{ij}] \tau_i^+ \tau_j^+ | \Psi_{A,Z} \rangle$$

$$F_1(x) = (x - 2)e^{-x}, \quad F_2(x) = (x + 1)e^{-x}, \quad x = m_\pi |\vec{r}|$$

$$T_{ij} = 3\vec{\sigma}_i \cdot \hat{r}_{ij} \vec{\sigma}_j \cdot \hat{r}_{ij} - \vec{\sigma}_i \vec{\sigma}_j$$

Regulator applied to  $0\nu\beta\beta$ -decay operator for consistency with LENPIC interaction

$$f\left(\frac{r}{R}\right) = \left(1 - \exp\left(-\frac{r^2}{R^2}\right)\right)^6$$

$R = 1.0$  fm for these results

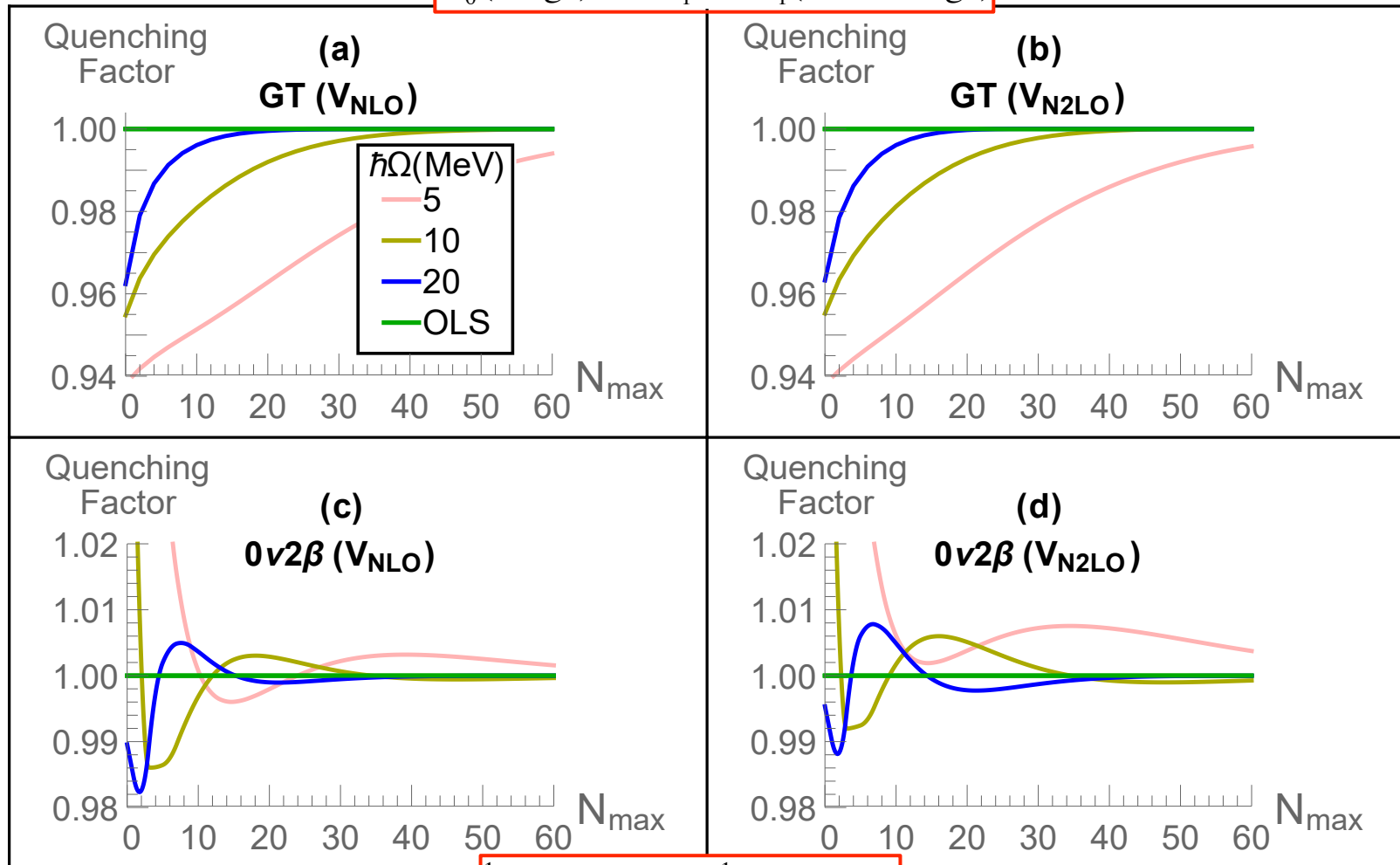
Additional operators under development – stay tuned

Two nucleons in a Harmonic Oscillator trap with trap  $\hbar\Omega =$  basis  $\hbar\Omega$   
 LENPIC Chiral NN interaction at N<sup>2</sup>LO with R = 1.0 fm  
 Comparison of GT and 0ν2β-decay matrix elements from truncation with Exact/OLS

Recast Fract. Diff. (FD) results as a Quenching Factor (QF)

$$QF = \text{Exact/Model} = 1 - (\text{FD} \times |\text{Exact}|)/\text{Model}$$

$${}^1S_0(nn \text{ gs}) \rightarrow {}^3S_1 - {}^3D_1(\text{deuteron gs})$$



$${}^1S_0(nn \text{ gs}) \rightarrow {}^1S_0(pp \text{ gs})$$

Outlook:

Implement in finite nuclei:

Perform benchmark  $A=6$  calculations with UNC group (underway)

Evaluate/save density matrices (static and transition)  
and use them to evaluate consistent OLS'd or SRG'd observables

Expand treatment to wider range of EW operators within Chiral EFT  
at NLO & N2LO

Extend to 3-body interactions with OLS or SRG on operators at the 3-  
body level

Extend to medium weight nuclei with “Double OLS” approach

## Partial list of projects underway – keep on your radar screens

Iteratively improved natural orbitals (with Notre Dame Univ)

Ab initio nuclear reactions

multiple-scattering with realistic 1-body density matrices (with Ohio Univ)

non-perturbative time-dependent Coulomb excitation (with IMP-Lanzhou)

Benchmarking neutrinoless double-beta decay (with UNC)

Consistent electroweak operators (LENPIC)

moments

transitions

double-beta decay

} This talk

Valence effective interactions (with S. Korea, France, Russia)

Artificial Neural Network developments and applications to NCSM (with LBNL)

Collaborators at Iowa State University  
Members of NUCLEI and Topical Collaboration Teams

Robert Basili (grad student)  
Weijie Du (grad student)  
Matthew Lockner (grad student)  
Pieter Maris  
Soham Pal (grad student)  
Shiplu Sarker (grad student)

New faculty position at Iowa State in Nuclear Theory  
Supported, in part, by the Fundamental Interactions  
Topical Collaboration

**Watch for the Advertisement**