

# In-Medium SRG for Medium-Mass and Heavy Nuclei

Heiko Hergert

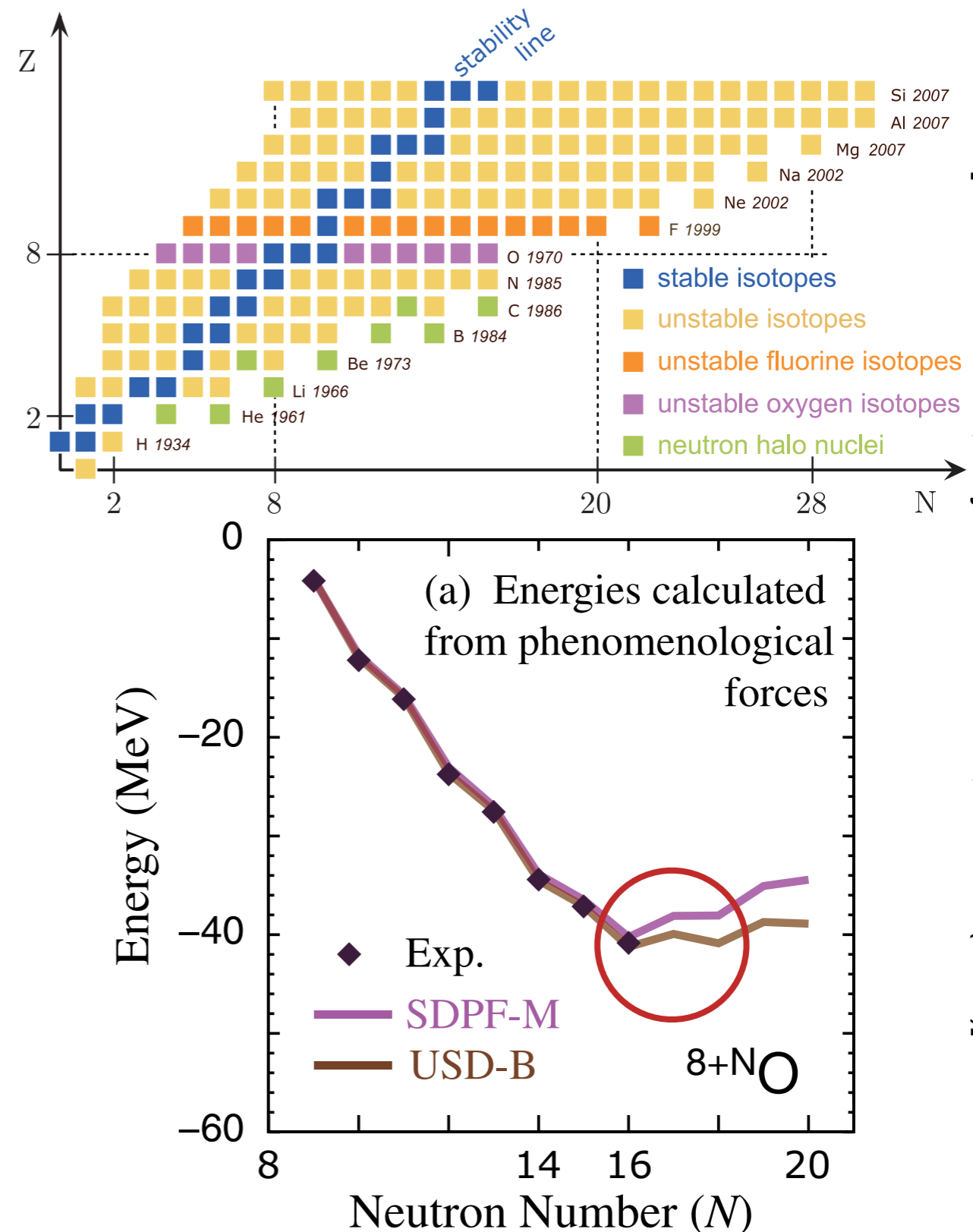
National Superconducting Cyclotron Laboratory  
Michigan State University



# Why Ab Initio Nuclear Structure?



- Nuclear Many-Body Shell Model: **oxygen drip line depends on interaction**
- empirical interactions **not interchangeable** between many-body methods
- **systematic improvements ?**
- **theoretical uncertainties ?**
- **systematic link to QCD ?**



[T. Otsuka et al., Phys. Rev. Lett. 105, 032501 (2010)]

- Ingredients
  - Nuclear Interactions from Chiral EFT
  - (Free-Space) Similarity Renormalization Group
- In-Medium SRG and Applications
  - Ground-State Results
  - IM-SRG + Shell Model for Excited States
- Next Steps
- Conclusions

Ingredients:

# Nuclear Interactions from Chiral Effective Field Theory

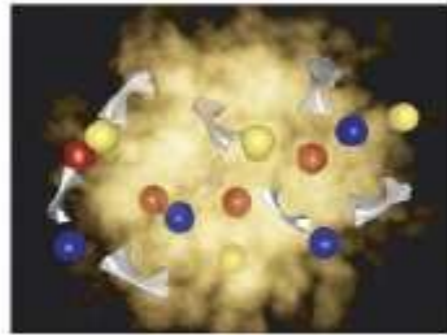
E. Epelbaum, H.-W. Hammer, and U.-G. Meissner, Rev. Mod. Phys. **81** (2009), 1773

# Scales of the Strong Interaction



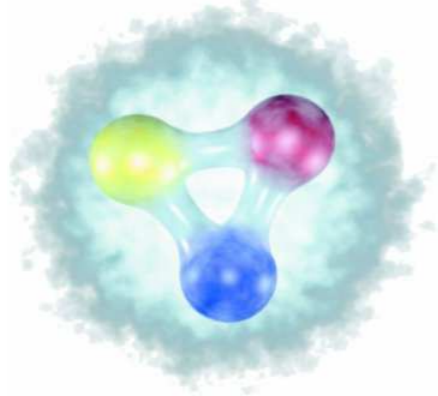
momentum transfer (resolution)

QCD



- quarks, gluons

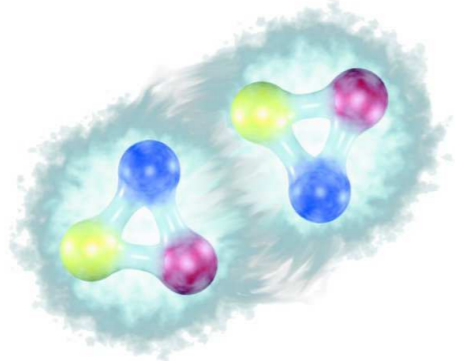
Chiral EFT



- pions, nucleons, ...

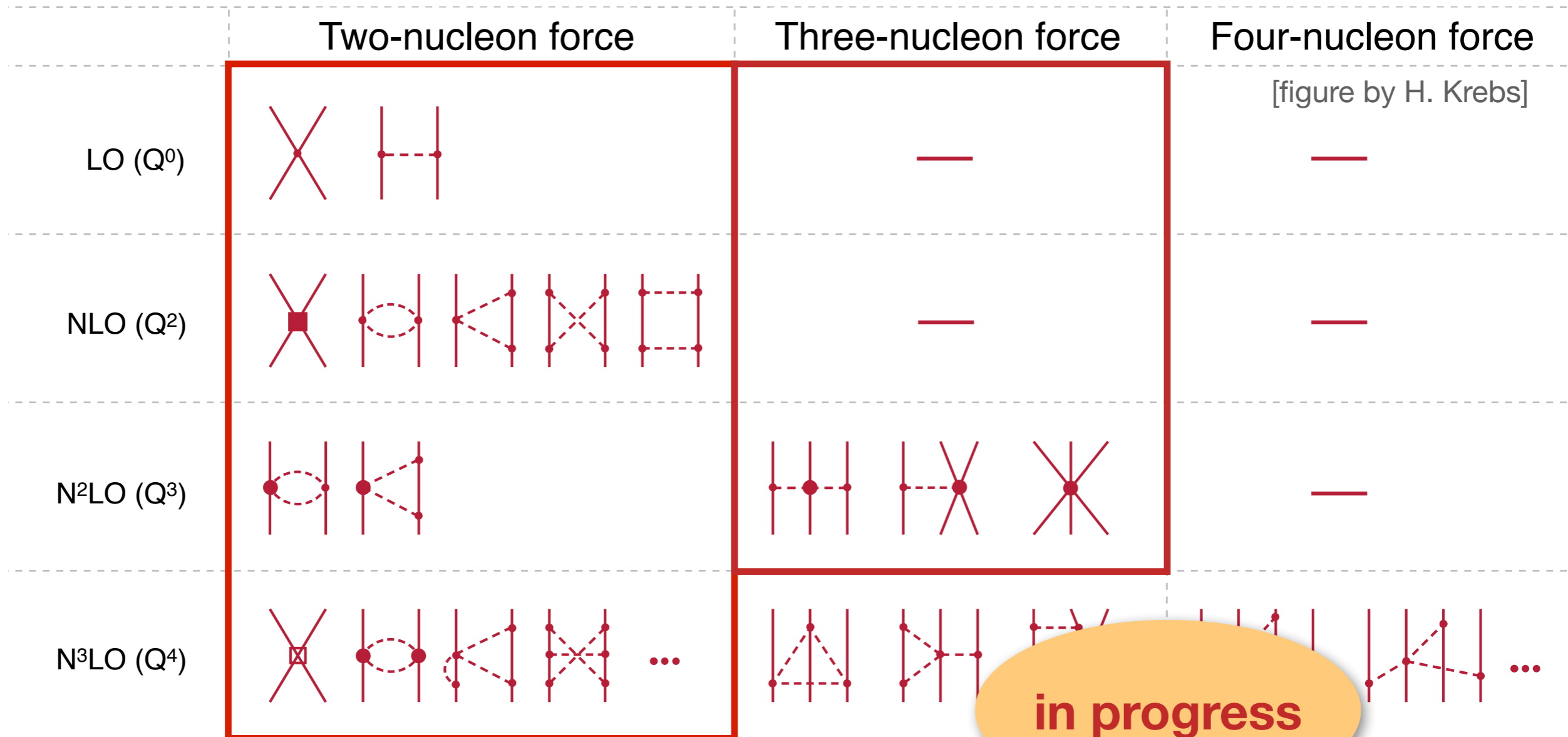
- nuclear interactions

- few-nucleon systems



chiral symmetry

# Interactions from Chiral EFT



- organization in powers  $(Q/\Lambda_\chi)^\nu$  allows **systematic improvement**
- low-energy constants **fit to NN, 3N data** (future: from Lattice QCD (?))
- **consistent** NN, 3N, ... interactions & operators (electromagnetic & weak transitions, etc.)

Ingredients:

# Similarity Renormalization Group

## Review:

S. Bogner, R. Furnstahl, and A. Schwenk, Prog. Part. Nucl. Phys. **65** (2010), 94

E. Anderson, S. Bogner, R. Furnstahl, and R. Perry, Phys. Rev. **C82** (2011), 054001

E. Jurgenson, P. Navratil, and R. Furnstahl, Phys. Rev. **C83** (2011), 034301

R. Roth, S. Reinhardt, and H. H., Phys. Rev. **C77** (2008), 064003

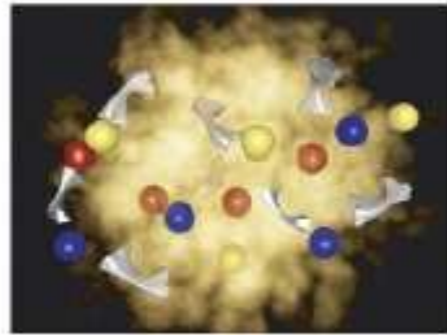
H. H. and R. Roth, Phys. Rev. **C75** (2007), 051001

# Scales of the Strong Interaction



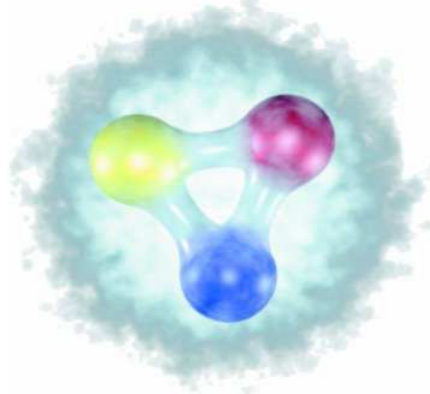
momentum transfer (resolution)

QCD



- quarks, gluons

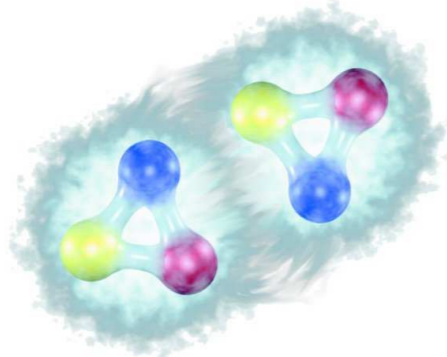
Chiral EFT



- pions, nucleons, ...

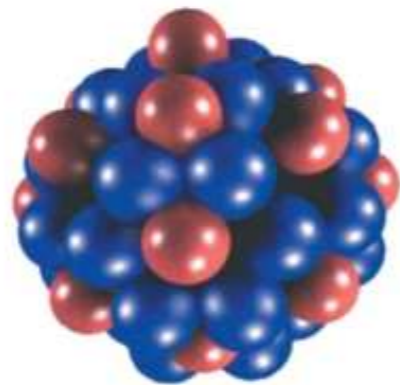
- nuclear interactions

- few-nucleon systems



- finite nuclei

- nuclear structure & reactions



chiral symmetry

(Which) Details necessary?



## Basic Concept

**continuous unitary transformation** of the Hamiltonian to band-diagonal form w.r.t. a given “uncorrelated” many-body basis

- evolved Hamiltonian

$$H(s) = U(s)HU^\dagger(s) \equiv T + V(s)$$

- flow equation:

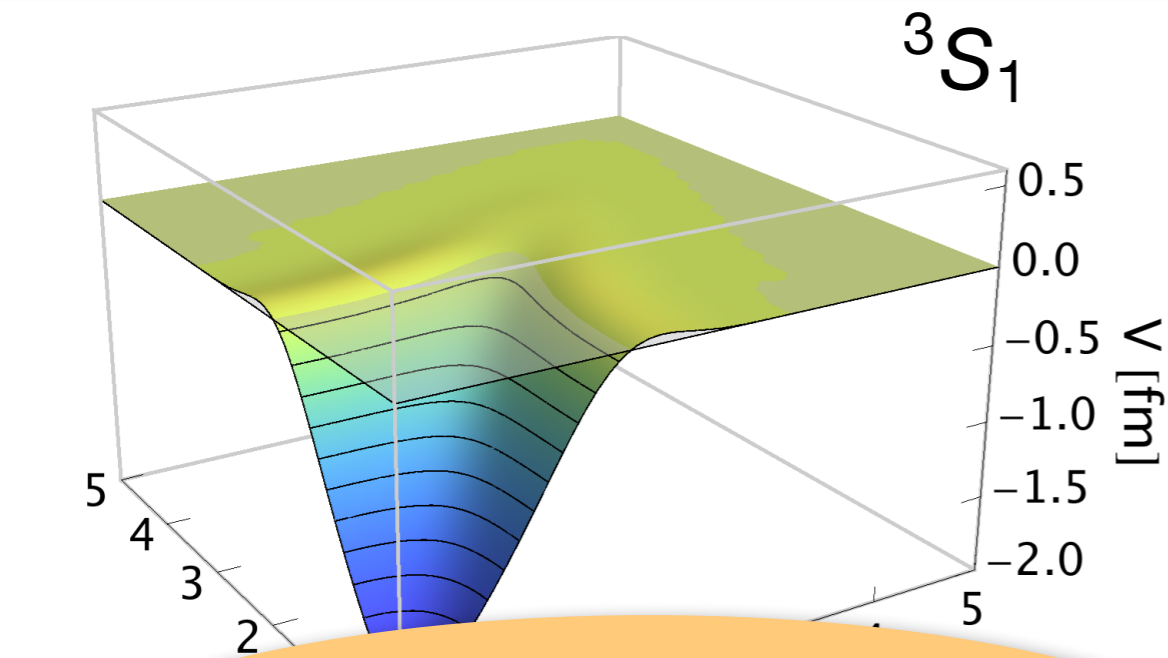
$$\frac{d}{ds}H(s) = [\eta(s), H(s)], \quad \eta(s) = \frac{dU(s)}{ds}U^\dagger(s) = -\eta^\dagger(s)$$

- choose  $\eta(s)$  to achieve desired behavior, e.g. decoupling of momentum or energy scales
- **consistently evolve observables** of interest

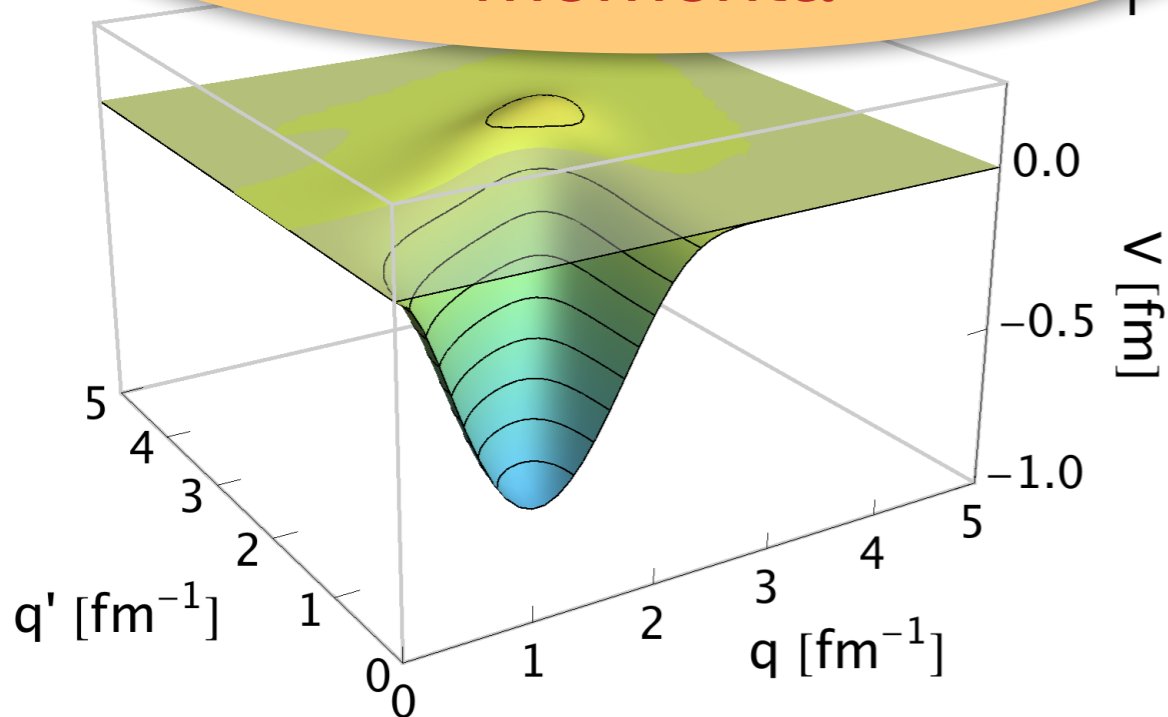
# SRG in Two-Body Space



momentum space matrix elements



lowering resolution scale  $\lambda$   
 $\Leftrightarrow$  decoupling of low and high momenta

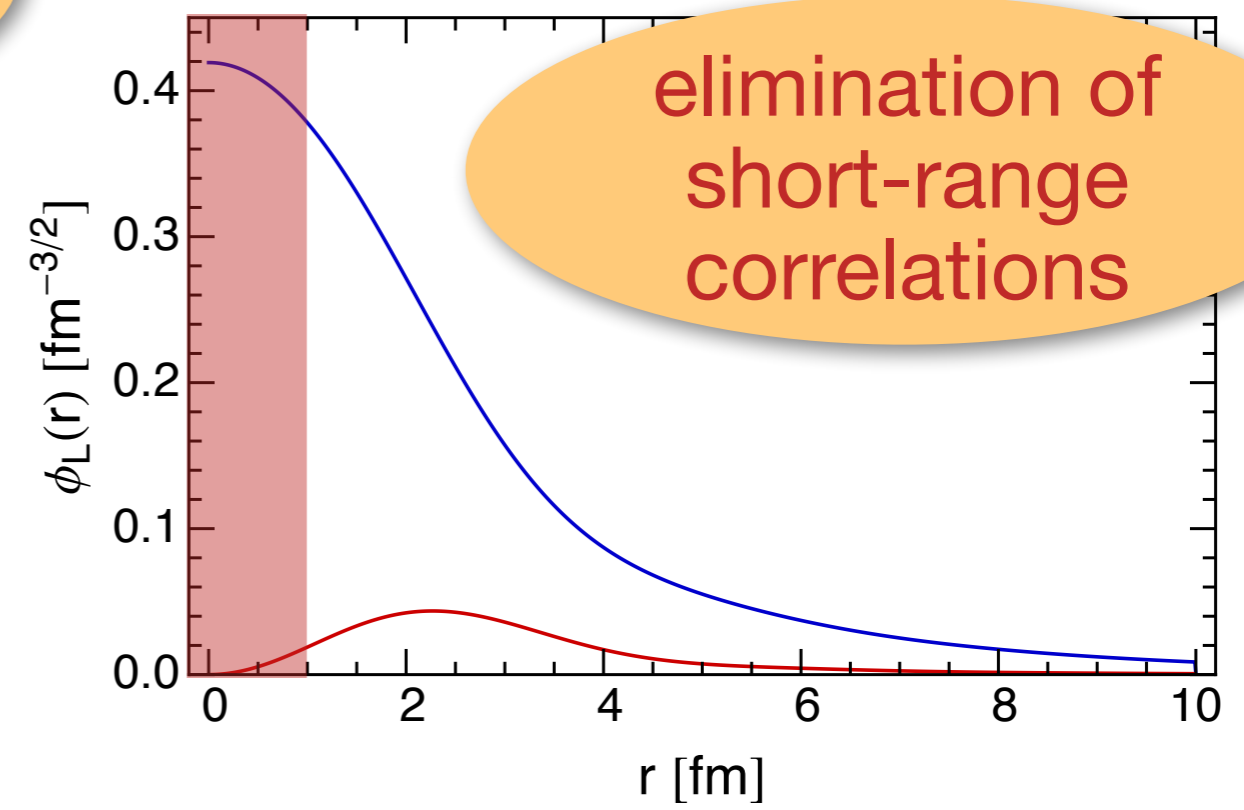


$\lambda = 1.8 \text{ fm}^{-1}$

$$\eta(\lambda) = 2\mu [T_{\text{rel}}, H(\lambda)]$$

$$\lambda = s^{-1/4}$$

deuteron wave function



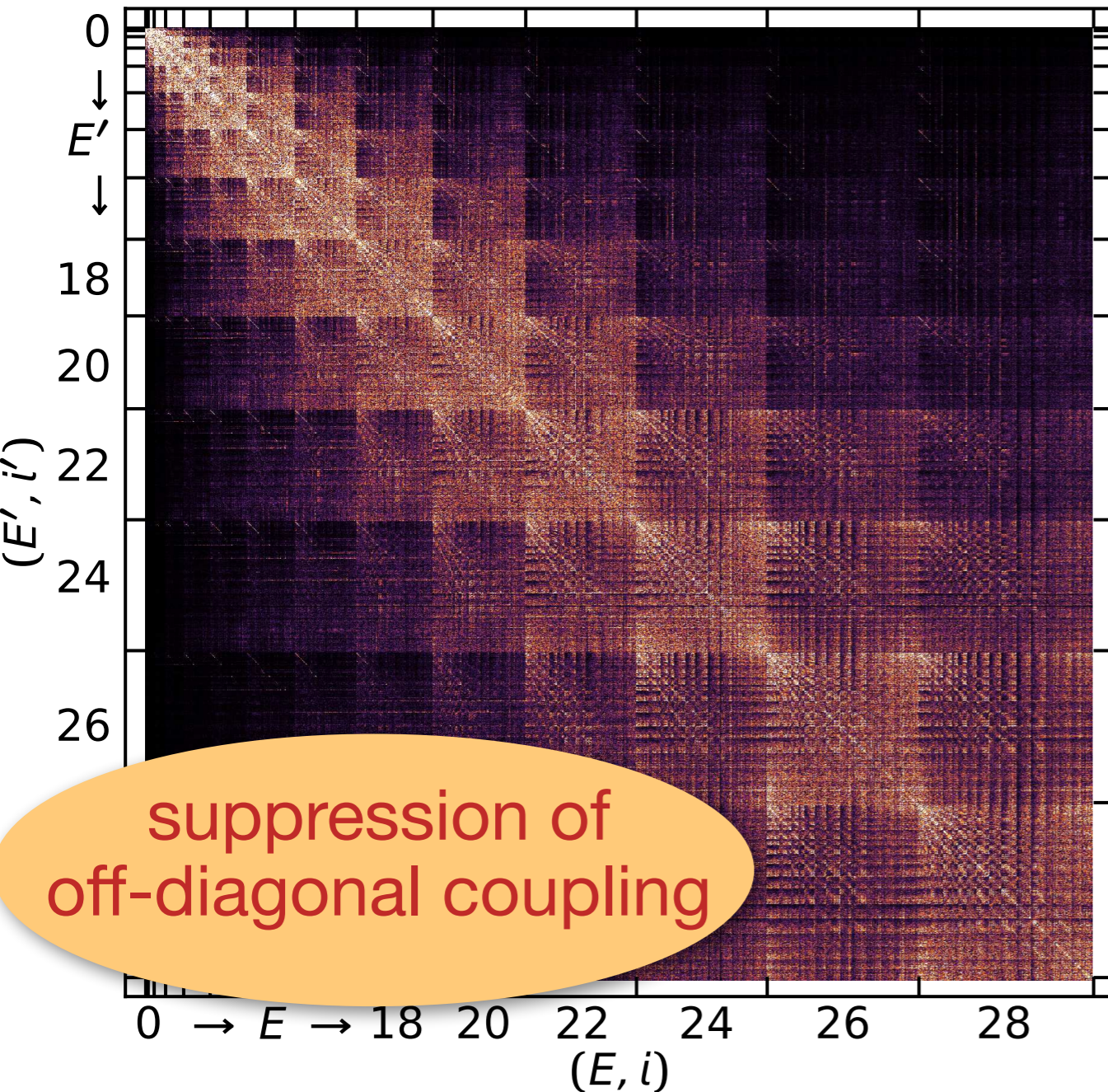
elimination of short-range correlations

# SRG in Three-Body Space



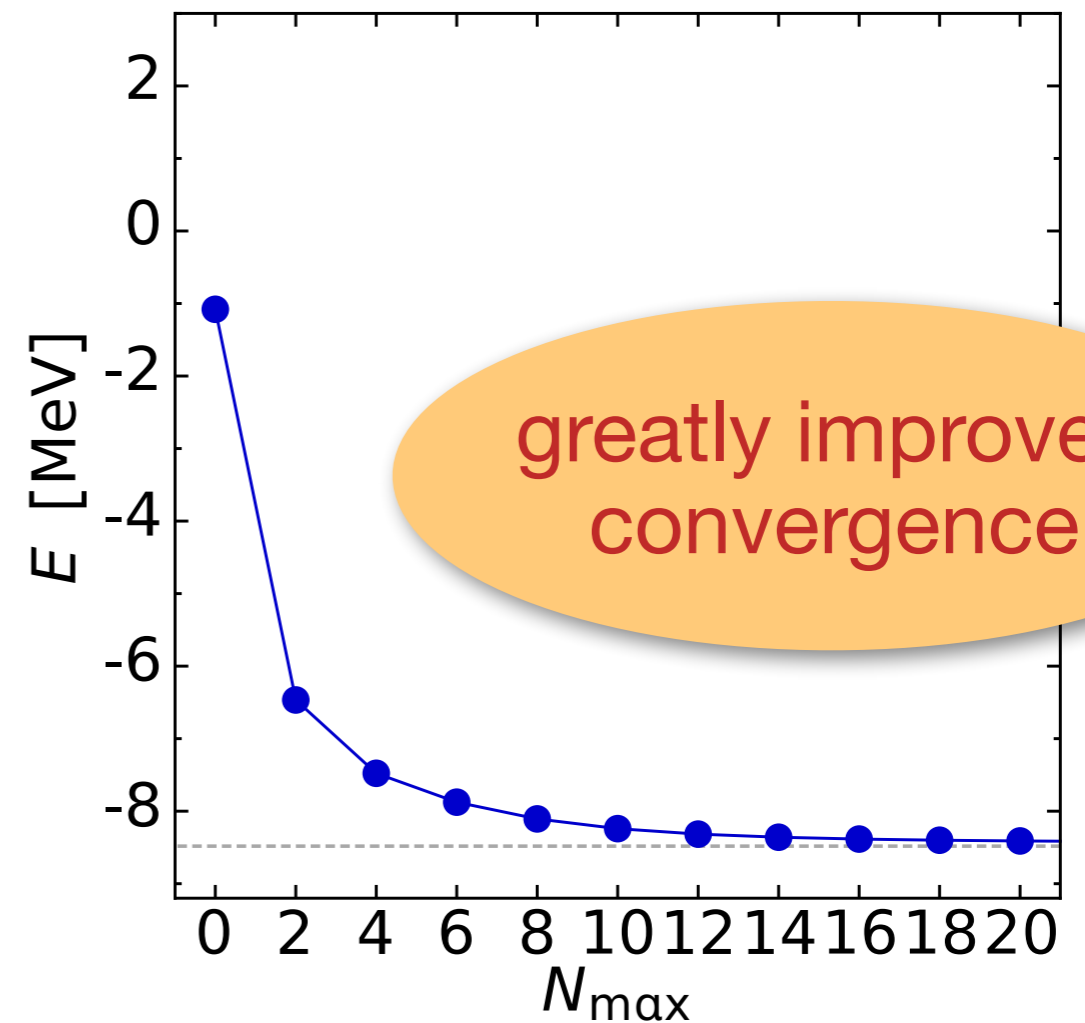
## 3B Jacobi-HO Matrix Elements

$$J^\pi = \frac{1}{2}^+, T = \frac{1}{2}, \hbar\Omega = 28 \text{ MeV}$$



$$\lambda = 1.33 \text{ fm}^{-1}$$

## ${}^3\text{H}$ ground-state (NCSM)



[figures by R. Roth, A. Calci, J. Langhammer]

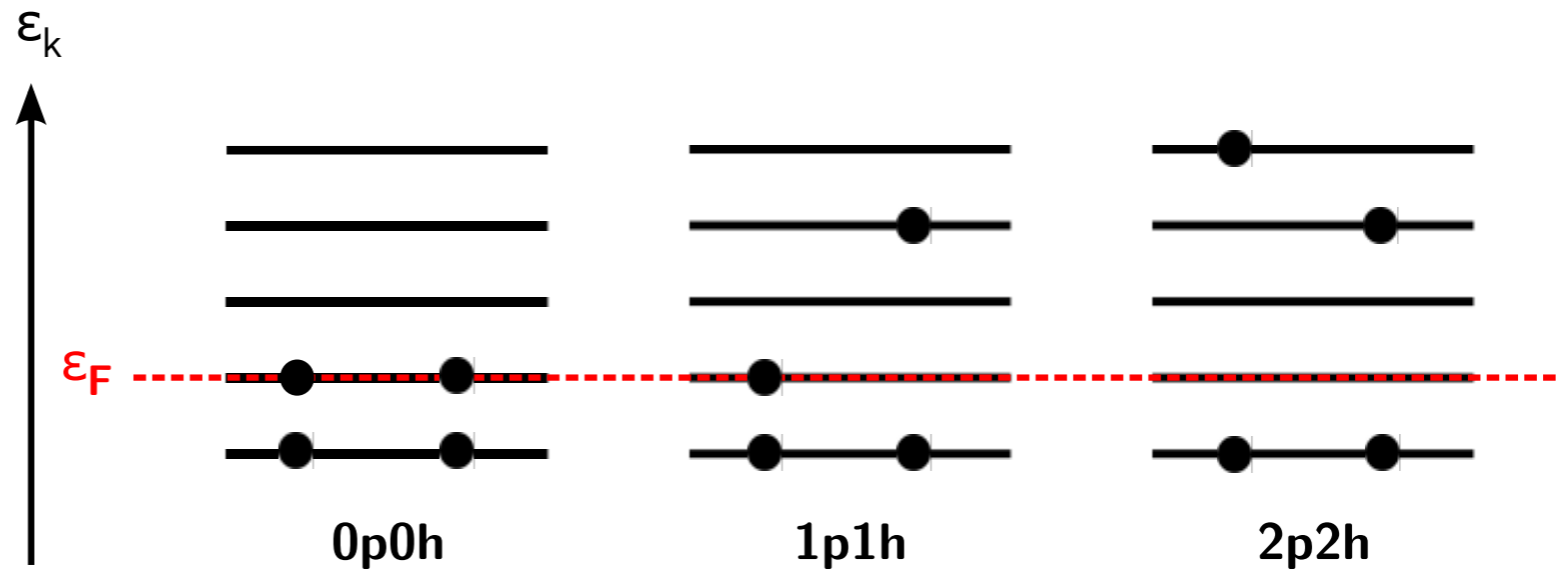
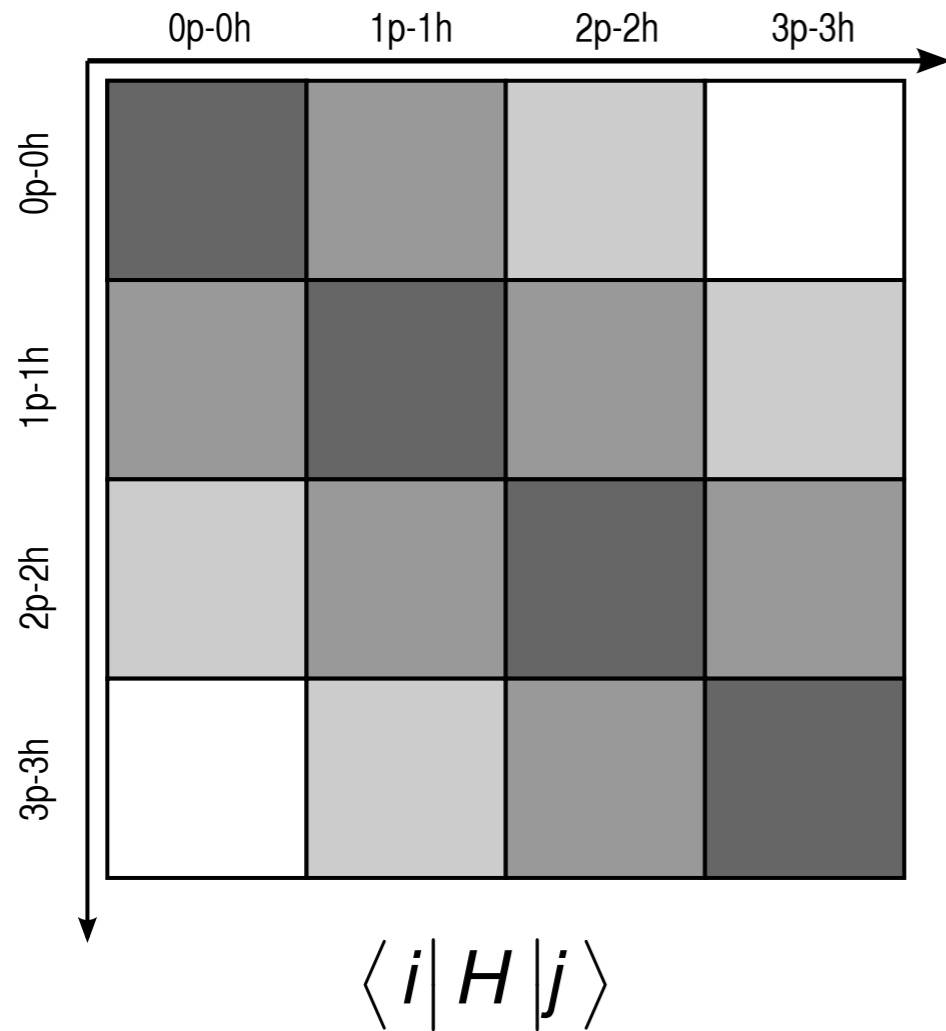
# In-Medium SRG

S. K. Bogner, H. H., T. Morris, A. Schwenk, and K. Tsukiyama, to appear in Phys. Rept.

H. H., S. K. Bogner, S. Binder, A. Calci, J. Langhammer, R. Roth, and A. Schwenk, Phys. Rev. C **87**, 034307 (2013)

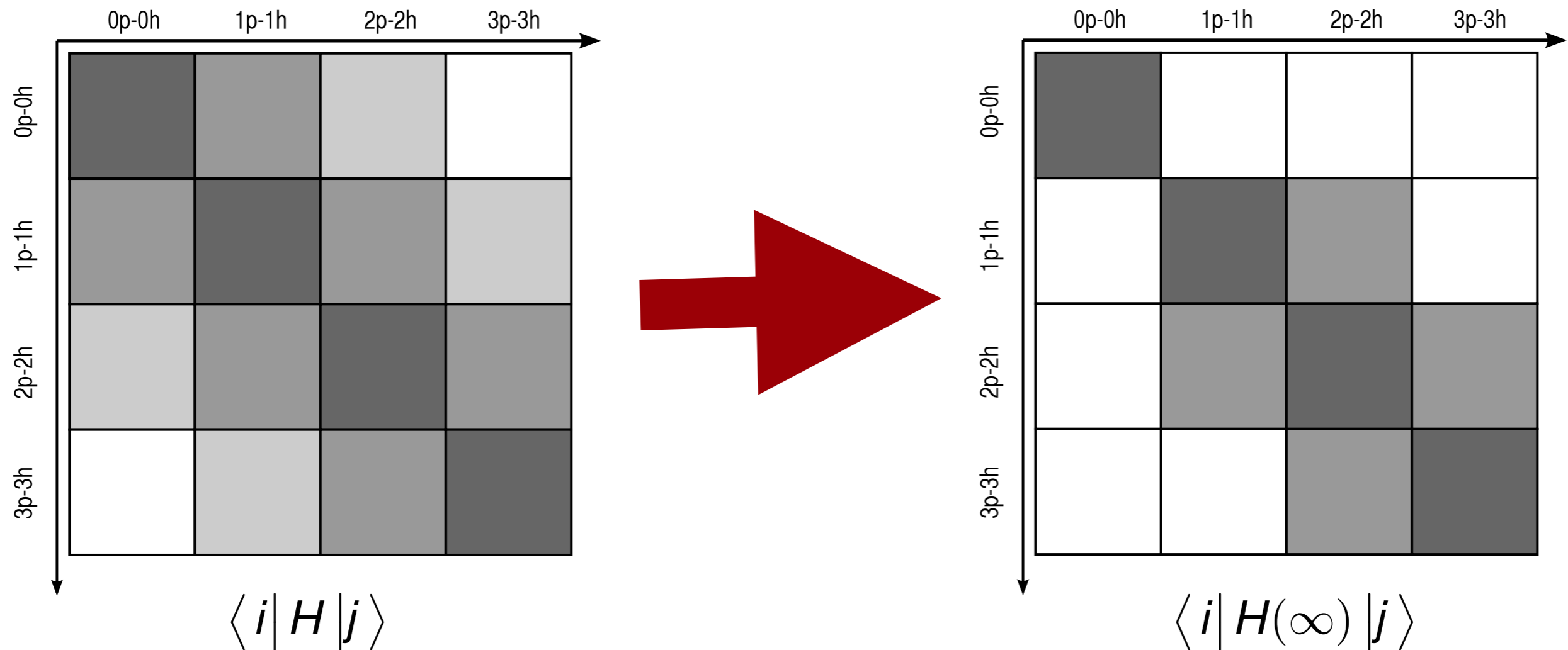
K. Tsukiyama, S. K. Bogner, and A. Schwenk, Phys. Rev. Lett. **106**, 222502 (2011)

# Decoupling in A-Body Space



excitations **relative** to reference state:  
 → **normal-ordering**

# Decoupling in A-Body Space



**aim:** decouple reference state  $|\Phi\rangle$   
(0p-0h) from excitations

# Normal Ordering



- **second quantization:**  $A_{I_1 \dots I_N}^{k_1 \dots k_N} = a_{k_1}^\dagger \dots a_{k_N}^\dagger a_{I_N} \dots a_{I_1}$

- particle- and hole density matrices:

$$\lambda_I^k = \langle \Phi | A_I^k | \Phi \rangle \longrightarrow n_k \delta_I^k, \quad n_k \in \{0, 1\}$$

$$\xi_I^k = \lambda_I^k - \delta_I^k \longrightarrow -\bar{n}_k \delta_I^k \equiv -(1 - n_k) \delta_I^k$$

- define **normal-ordered operators** recursively:

$$A_{I_1 \dots I_N}^{k_1 \dots k_N} = : A_{I_1 \dots I_N}^{k_1 \dots k_N} : + \lambda_{I_1}^{k_1} : A_{I_2 \dots I_N}^{k_2 \dots k_N} : + \text{singles} \\ + \left( \lambda_{I_1}^{k_1} \lambda_{I_2}^{k_2} - \lambda_{I_2}^{k_1} \lambda_{I_1}^{k_2} \right) : A_{I_3 \dots I_N}^{k_3 \dots k_N} : + \text{doubles} + \dots$$

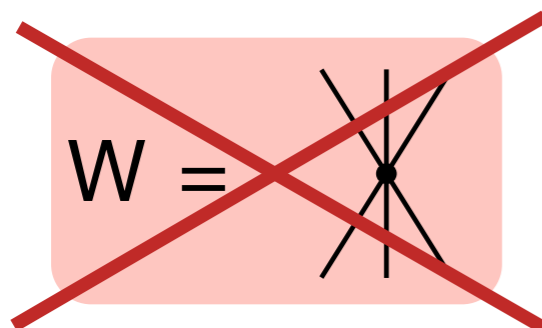
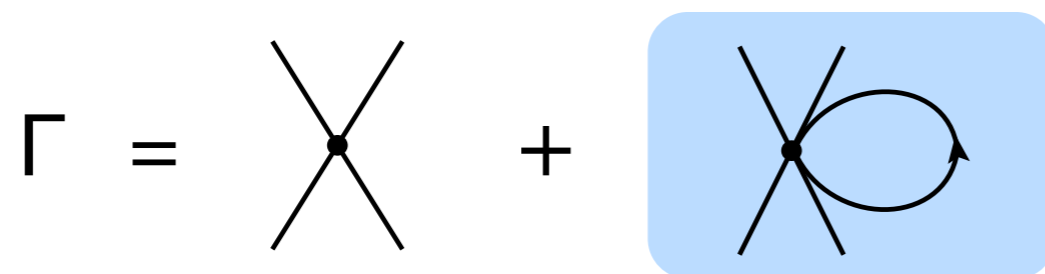
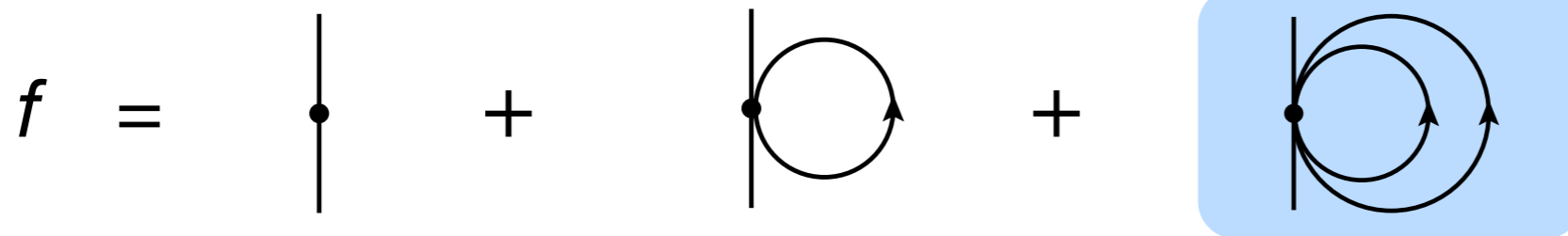
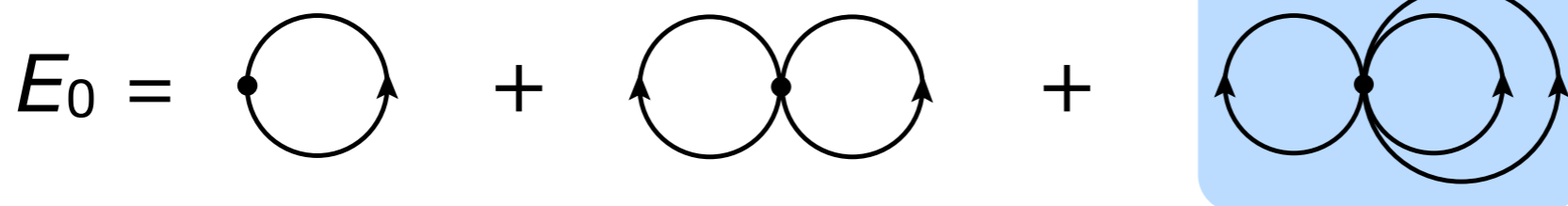
- **algebra is simplified** significantly because

$$\langle \Phi | : A_{I_1 \dots I_N}^{k_1 \dots k_N} : | \Phi \rangle = 0$$

- **Wick's theorem** gives simplified expansions (**fewer terms!**) for products of normal-ordered operators

## Normal-Ordered Hamiltonian

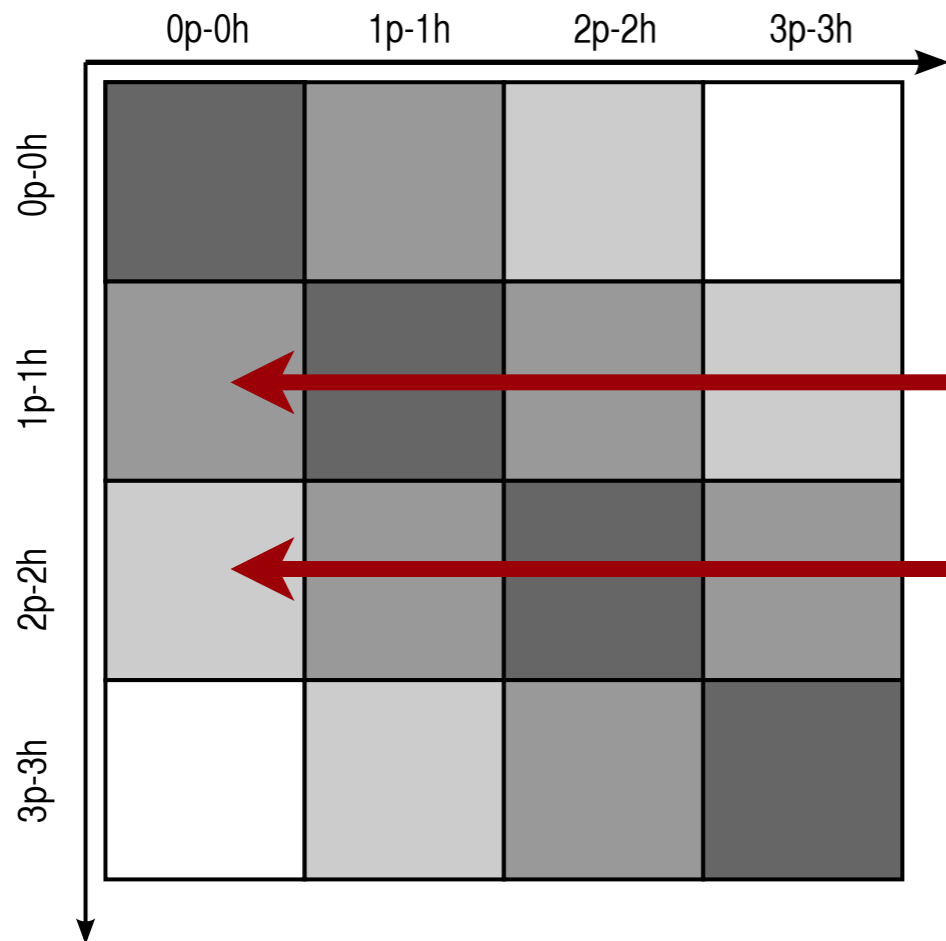
$$H = E_0 + \sum_{kl} f_l^k : A_l^k : + \frac{1}{4} \sum_{klmn} \Gamma_{mn}^{kl} : A_{mn}^{kl} : + \frac{1}{36} \sum_{ijklmn} W_{lmn}^{ijk} : A_{lmn}^{ijk} :$$



two-body formalism with  
in-medium contributions from  
three-body interactions



# Choice of Generator



$$\langle \begin{smallmatrix} p \\ h \end{smallmatrix} | H | \Psi \rangle = \sum_{kl} f_l^k \langle \Psi | : A_p^h :: A_l^k : | \Psi \rangle = -n_h \bar{n}_p f_h^p$$

$$\langle \begin{smallmatrix} pp' \\ hh' \end{smallmatrix} | H | \Psi \rangle = \sum_{klmn} \Gamma_{mn}^{kl} \langle \Psi | : A_{pp'}^{hh'} :: A_{mn}^{kl} : | \Psi \rangle \sim \Gamma_{hh'}^{pp'}$$

- define off-diagonal Hamiltonian (suppressed by IM-SRG flow):

$$H^{od} \equiv f^{od} + \Gamma^{od}, \quad f^{od} \equiv \sum_{ph} f_h^p : A_h^p : + \text{H.c.}, \quad \Gamma^{od} \equiv \sum_{pp'hh'} \Gamma_{hh'}^{pp'} : A_{hh'}^{pp'} : + \text{H.c.}$$

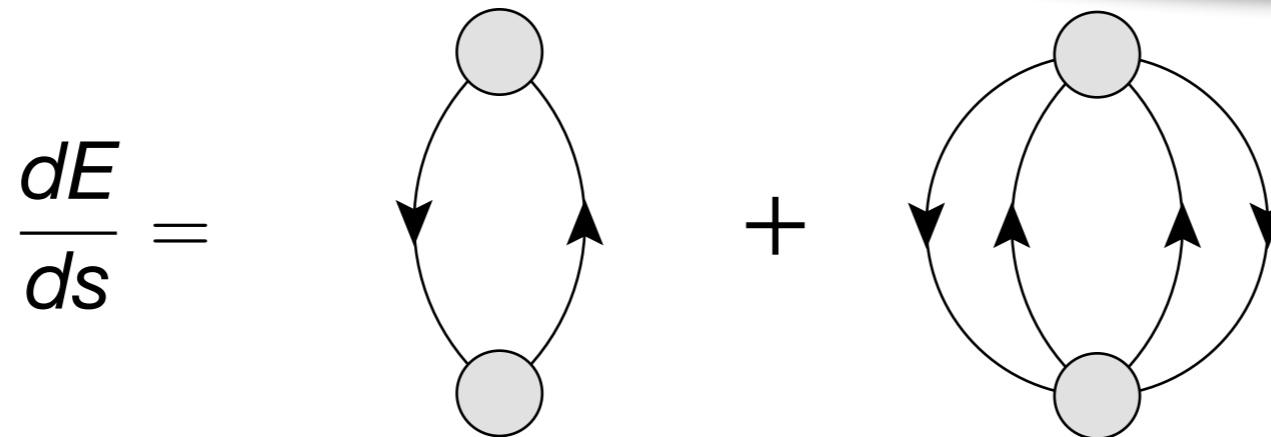
- construct generator, e.g.,  $\eta^l = [H^d, H^{od}]$  (Wegner-type)

# IM-SRG(2) Flow Equations

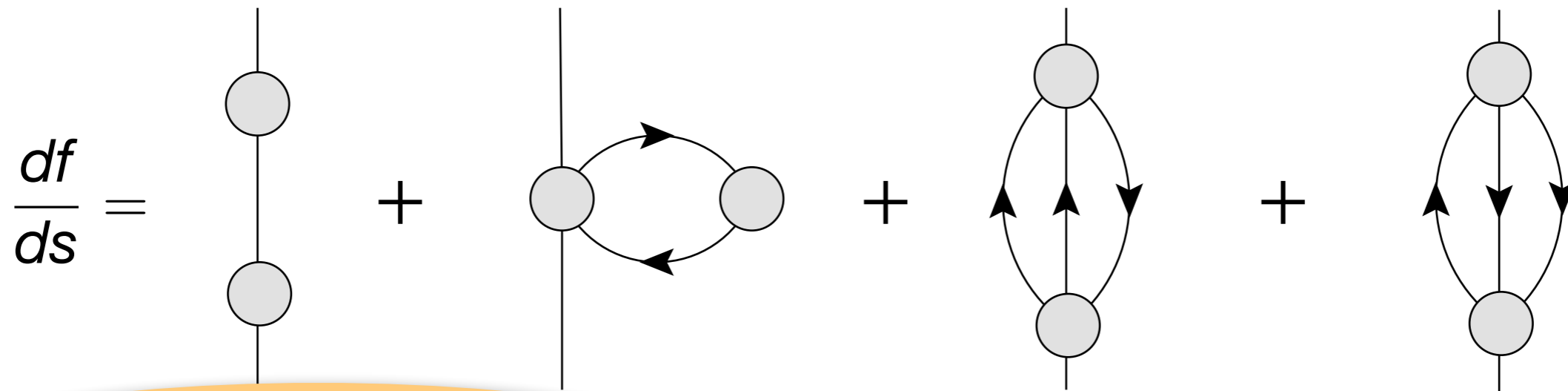


## 0-body Flow

~ 2nd order MBPT for  $H(s)$



## 1-body Flow



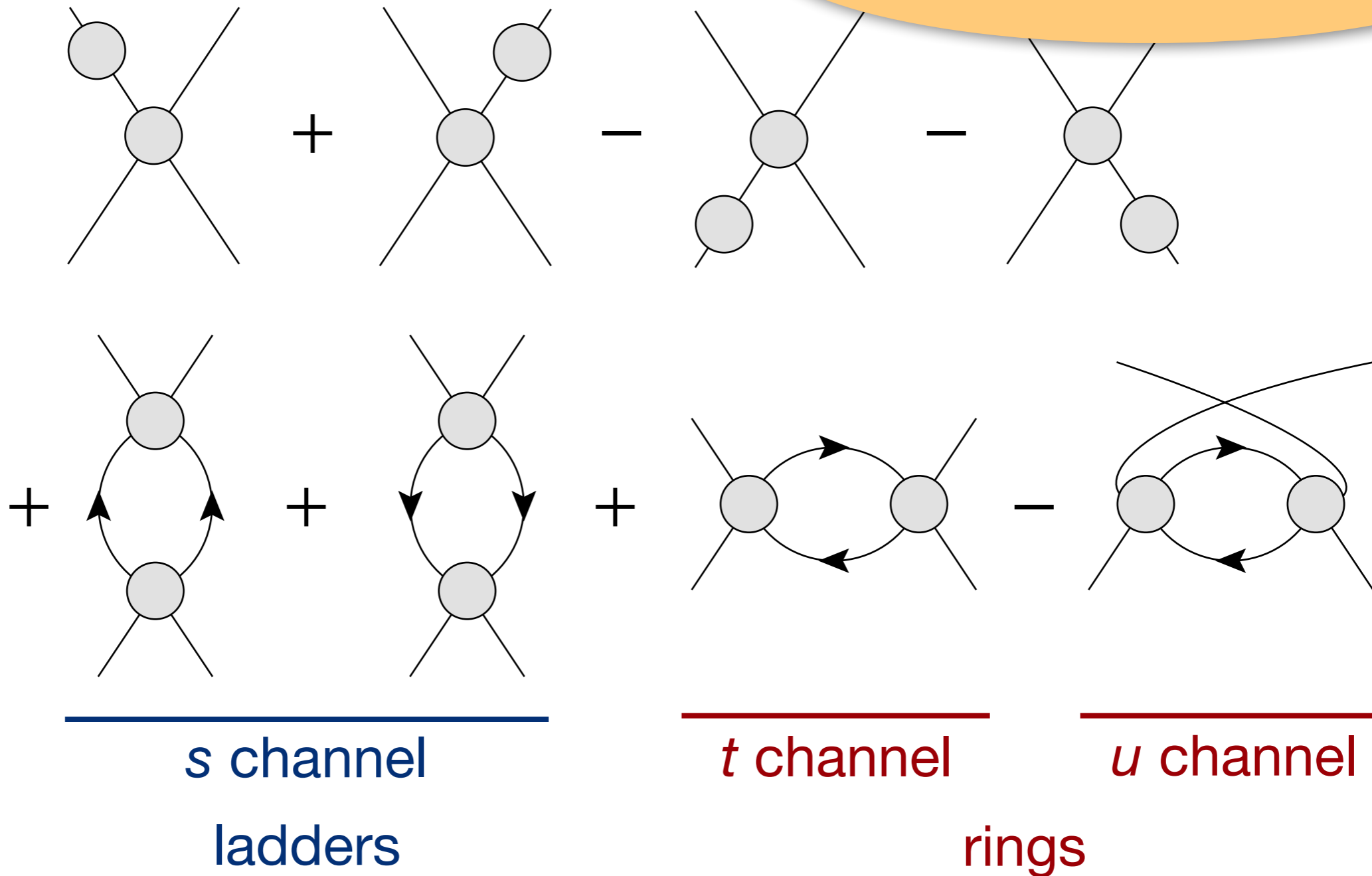
IM-SRG(2): truncate ops.  
at two-body level

# IM-SRG(2) Flow Equations



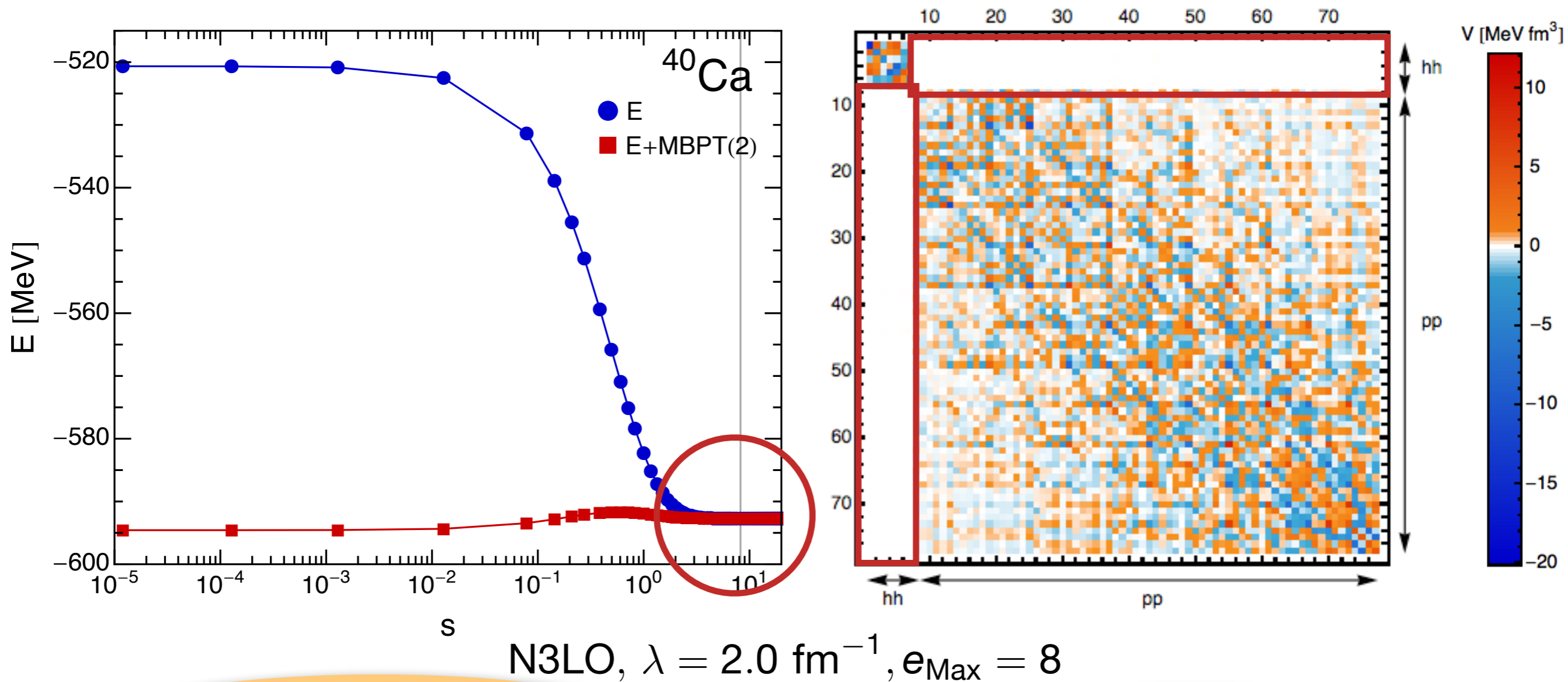
## 2-body Flow

$$\frac{d\Gamma}{ds} =$$



**$O(N^6)$  scaling**  
 (before particle/hole distinction)

# Decoupling



non-perturbative  
resummation of MBPT series  
(correlations)

off-diagonal couplings  
are rapidly driven to zero

Applications:

# Ground-State Results

H. H., in preparation

H. H., S. Bogner, T. Morris, S. Binder, A. Calci, J. Langhammer, R. Roth, Phys. Rev. C **90**, 041302 (2014)

H. H., S. Binder, A. Calci, J. Langhammer, and R. Roth, Phys. Rev. Lett **110**, 242501 (2013)

H. H., S. K. Bogner, S. Binder, A. Calci, J. Langhammer, R. Roth, and A. Schwenk, Phys. Rev. C **87**, 034307 (2013)

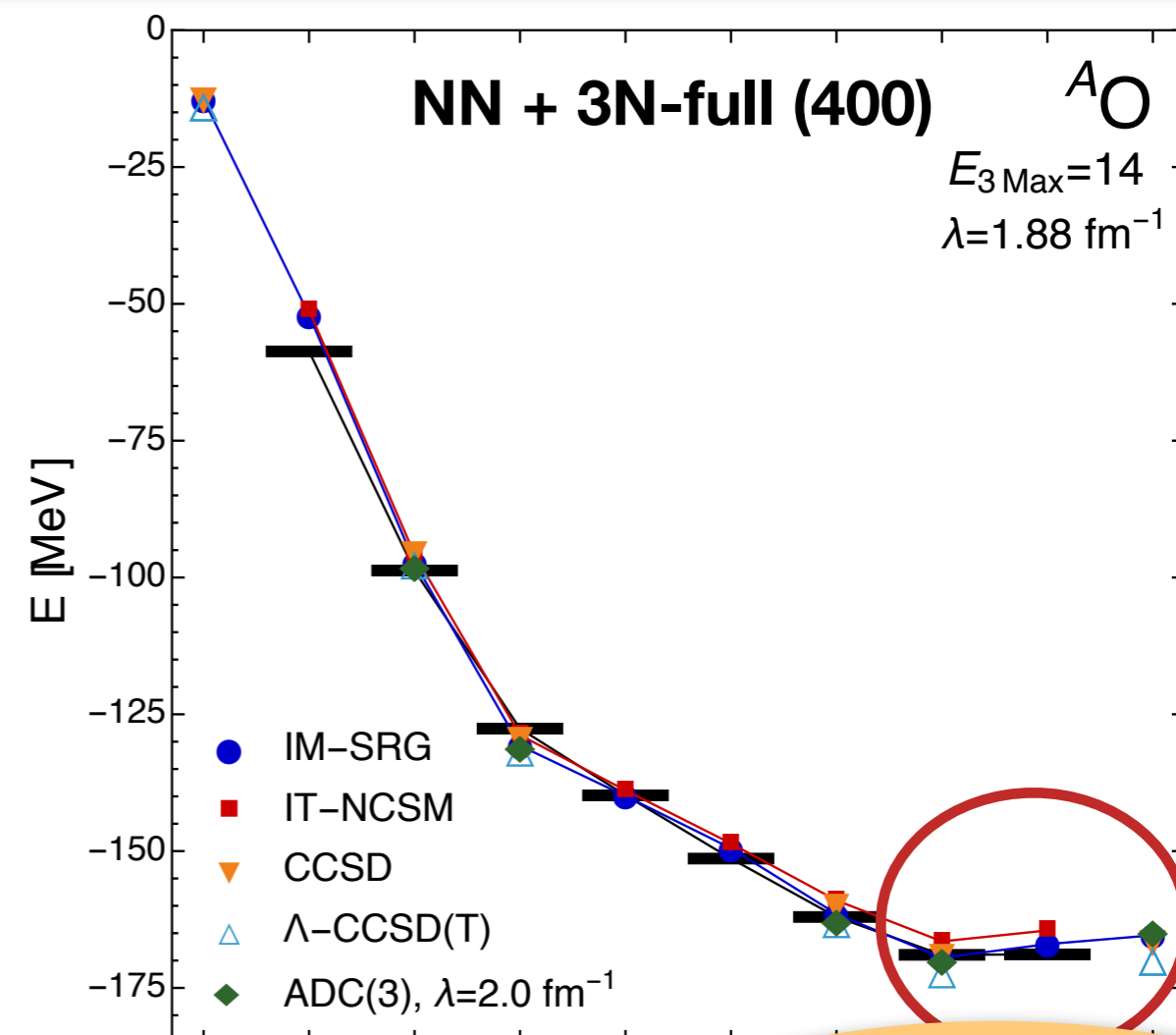
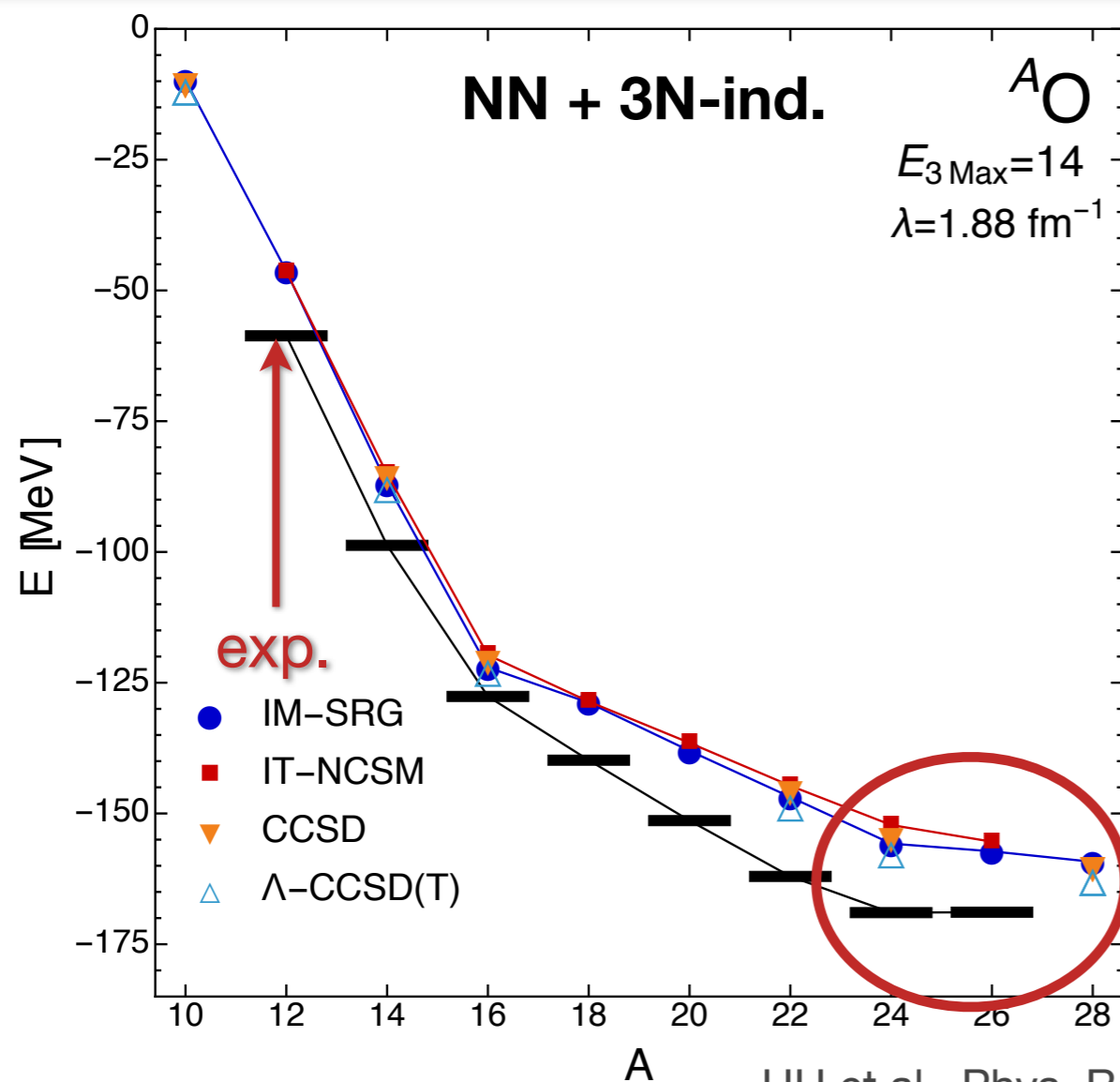
## Initial Hamiltonian

- NN: chiral interaction at  $N^3LO$  (Entem & Machleidt)
- 3N: chiral interaction at  $N^2LO$  ( $c_{D,CE}$  fit to  ${}^3H$  energy & beta decay)

## SRG-Evolved Hamiltonians

- **NN + 3N-induced:** start with initial NN Hamiltonian, keep two- and three-body terms
- **NN + 3N-full:** start with initial NN + 3N Hamiltonian, keep two- and three-body terms

# Results: Oxygen Chain



HH et al., Phys. Rev. Lett. **110**, 242501 (2013)

A. Cipollone et al., Phys. Rev. Lett. **111**, 242501 (2013) (C. Barbieri et al.)

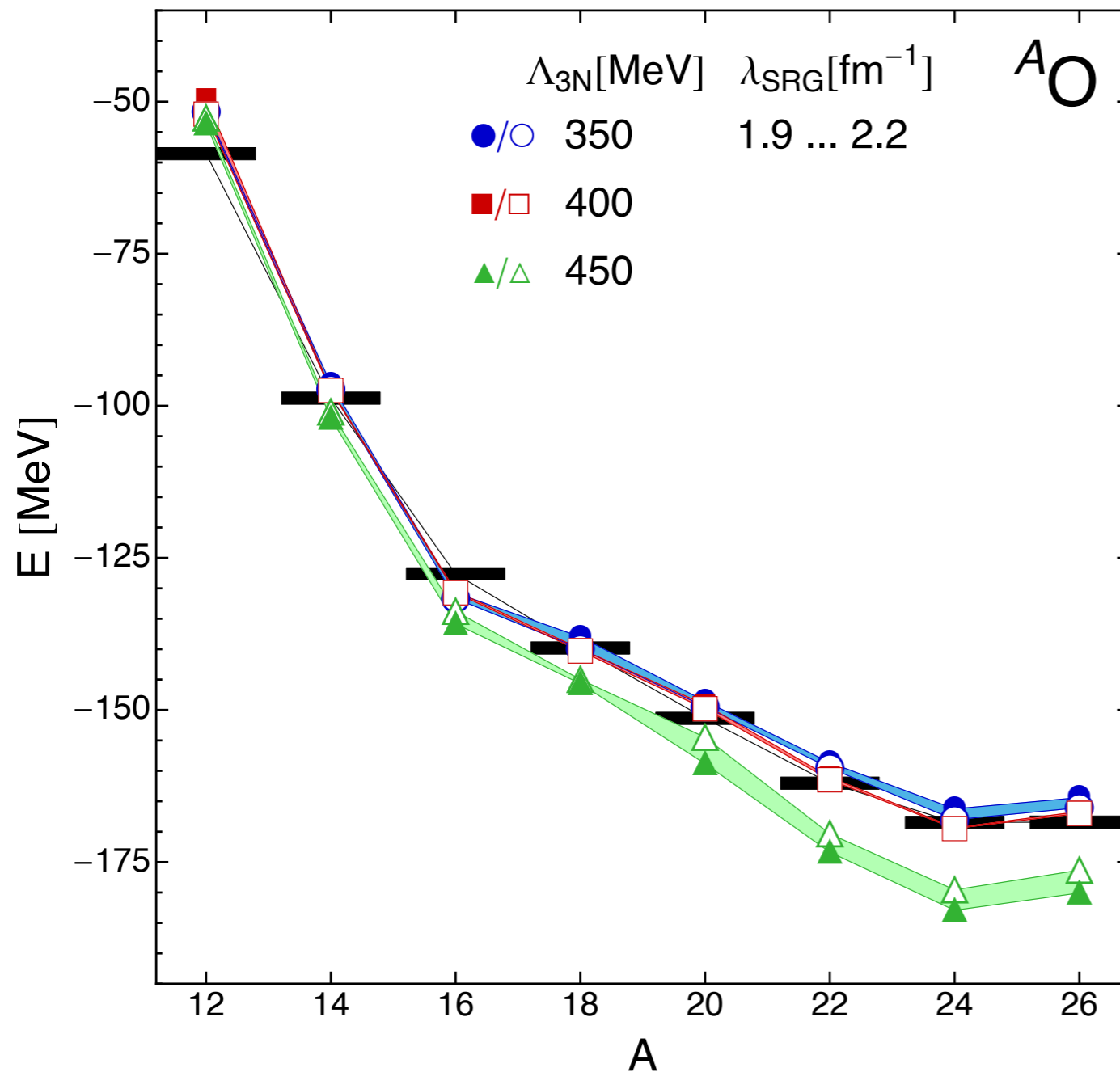
**3N interaction completely fixed by  $A \leq 4$  data**

- **Multi-Reference IM-SRG** with number-projected Hartree-Fock-Bogoliubov as reference state (**pairing correlations**)
- **consistent results from different many-body methods**

# Variation of Scales



## NN + 3N-full



- variation of **initial 3N cutoff only**
- diagnostics for chiral interactions
- **dripline at  $A=24$  is robust under variations**
- **(leading) continuum effects too small to bind  $^{26}\text{O}$**

Phys. Rev. Lett. **110**, 242501 (2013)

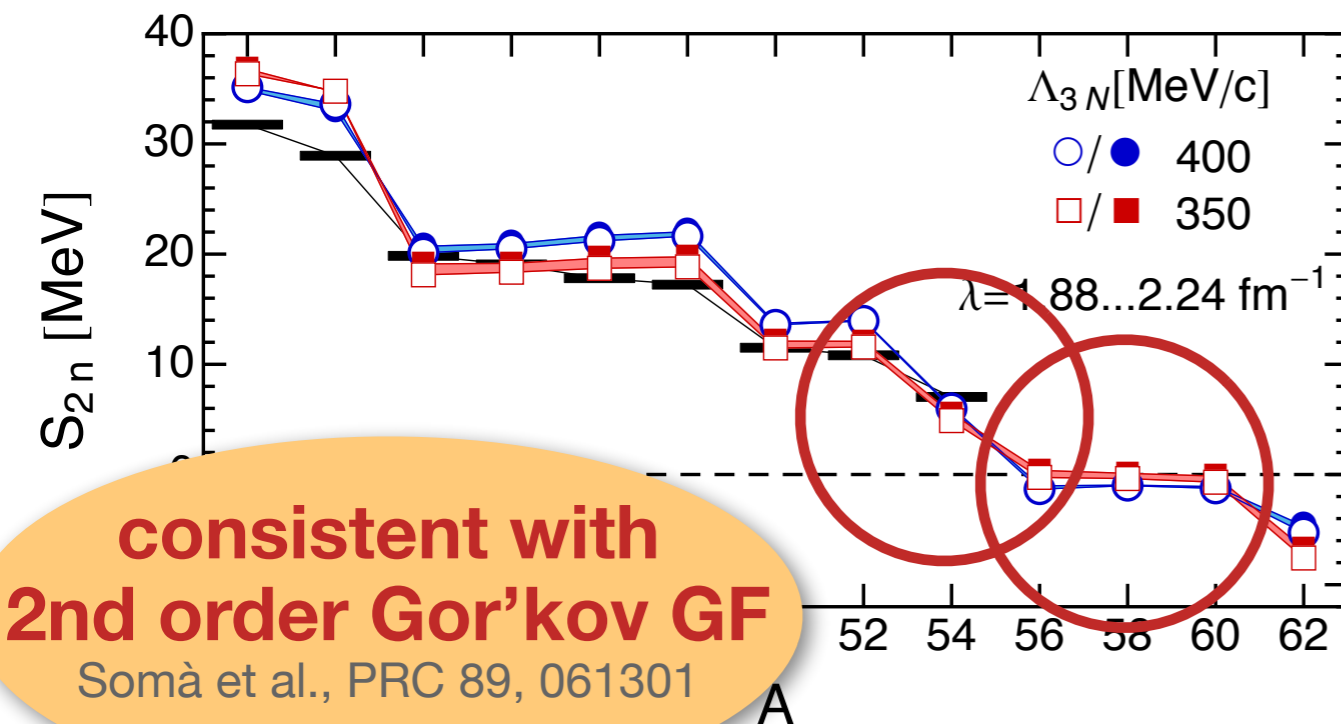
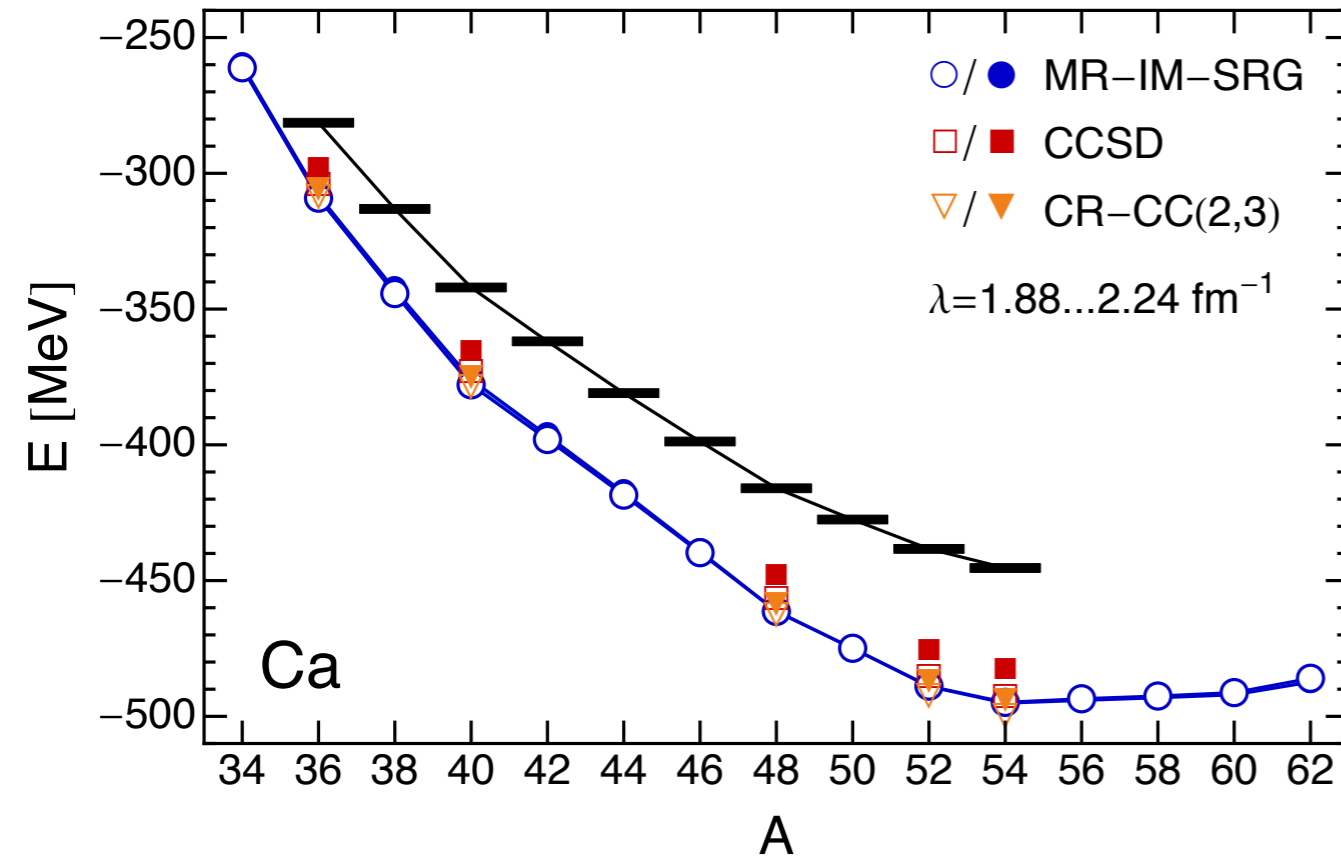


# Two-Neutron Separation Energies



PRC 90, 041302(R) (2014)

NN + 3N-full (400)



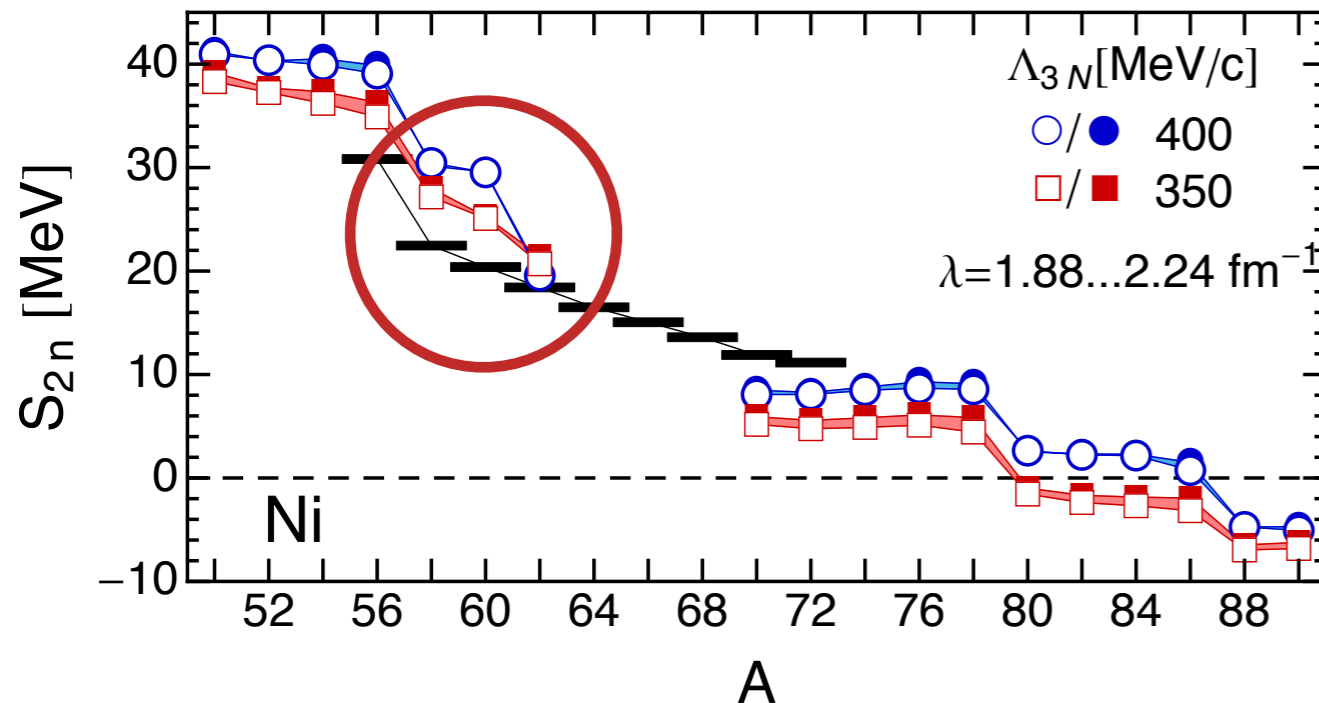
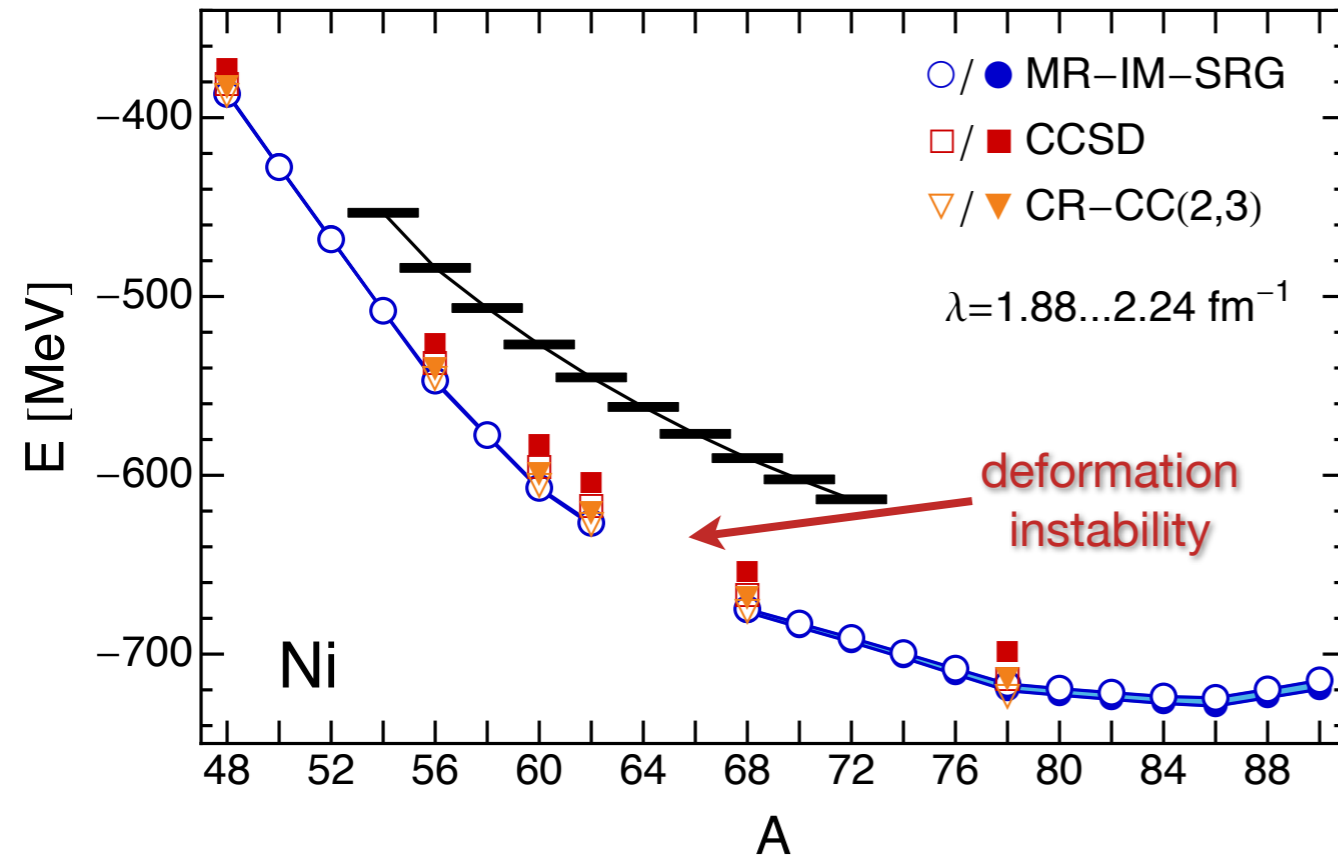
- differential observables ( $S_{2n}$ , spectra,...) filter out interaction components that cause overbinding
- predict flat trends for g.s. energies/ $S_{2n}$  beyond  $^{54}\text{Ca}$  - await experimental data
- $^{52}\text{Ca}$ ,  $^{54}\text{Ca}$  robustly magic due to 3N interaction
- no continuum coupling yet, other  $S_{2n}$  uncertainties  $< 1\text{MeV}$

# Two-Neutron Separation Energies



PRC 90, 041302(R) (2014)

## NN + 3N-full (400)

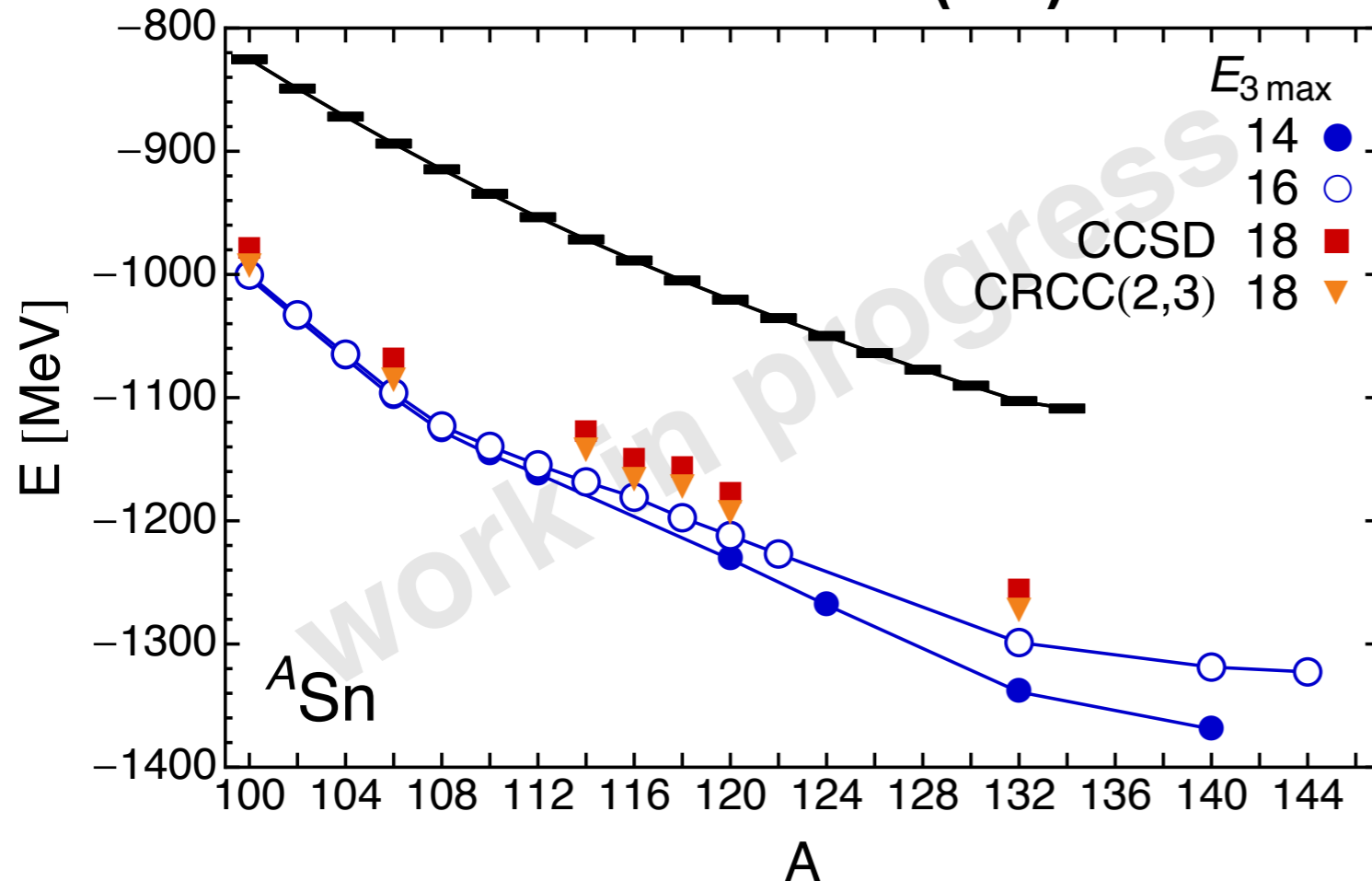


- flat trends for g. s. energies and  $S_{2n}$  (similar to Ca)
- deformation instability in  $^{64,66}\text{Ni}$  calculations - issue with “shell” structure
- further evidence from 3N cutoff variation
- no continuum coupling yet, other  $S_{2n}$  uncertainties  $< 1\text{MeV}$

# The *Ab Initio* Mass Frontier: Tin



NN + 3N-full (400)



$E_{3\max}$	memory (float) [GB]
14	5
16	~20
18	100+

- systematics of overbinding similar to Ca/Ni
- not converged with respect to 3N matrix element truncation:

$$e_1 + e_2 + e_3 \leq E_{3\max}$$

( $e_{1,2,3}$  : SHO energy quantum numbers)

- need technical improvements to go further

Applications:

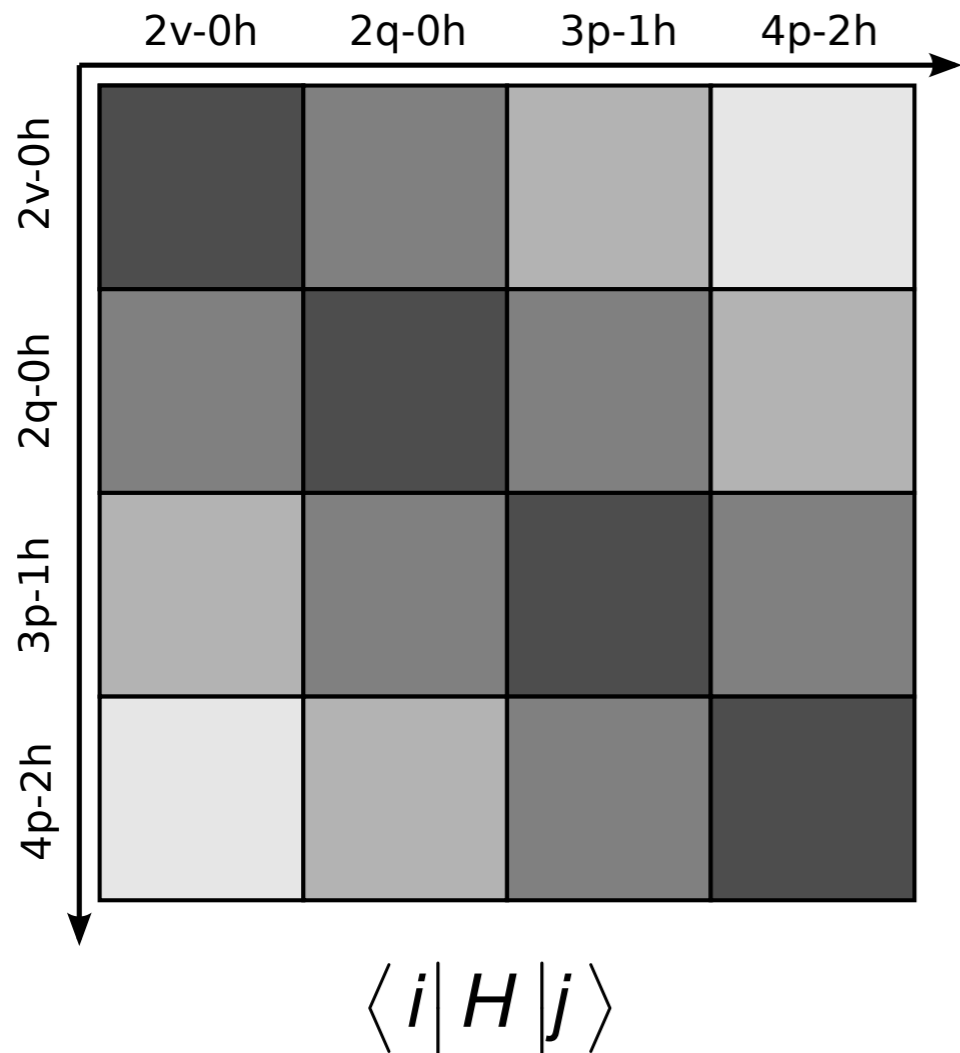
# IM-SRG + Shell Model for Excited States

S. K. Bogner, H. H., J. D. Holt, A. Schwenk, in preparation

S. K. Bogner, H. H., J. D. Holt, A. Schwenk, S. Binder, A. Calci, J. Langhammer, R. Roth,  
Phys. Rev. Lett. 113, 142501 (2014)

K. Tsukiyama, S. K. Bogner, and A. Schwenk, Phys. Rev. C **85**, 061304(R) (2012)

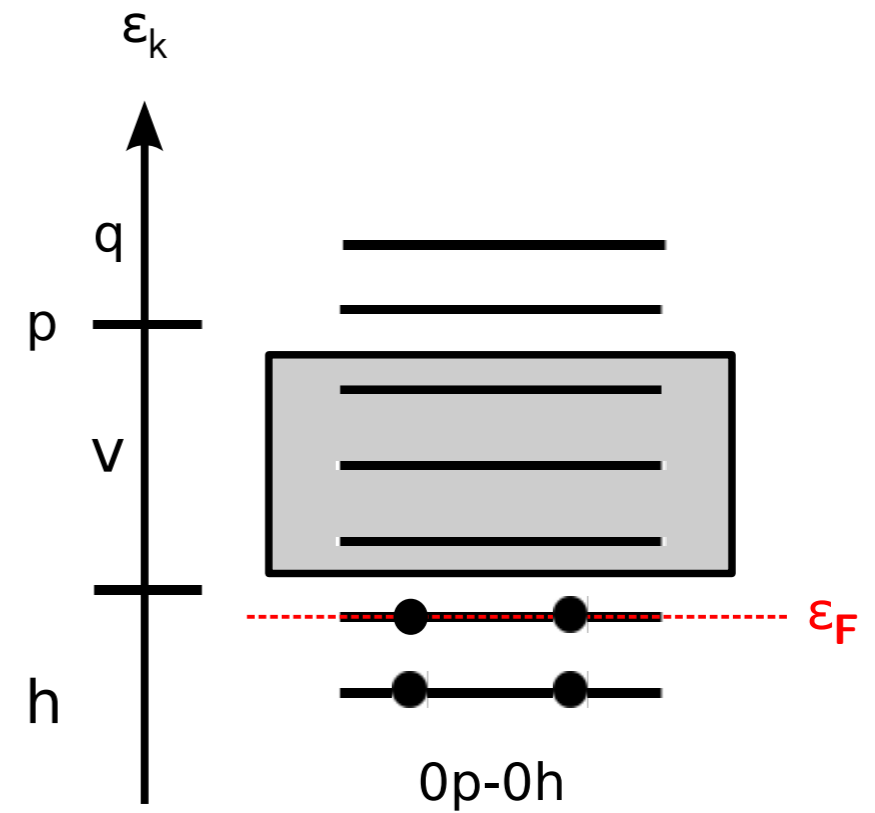
# Valence Space Decoupling



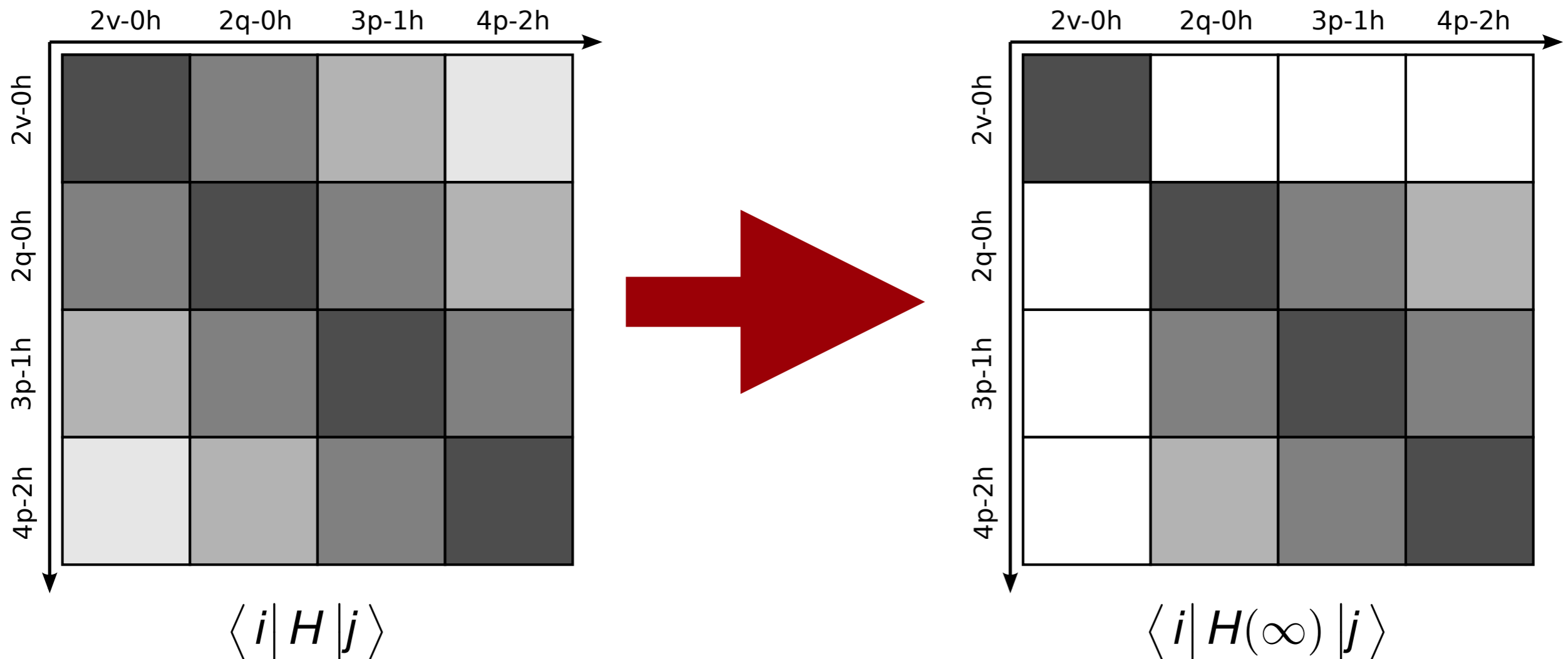
non-valence  
particle states

valence  
particle states

hole states  
(core)



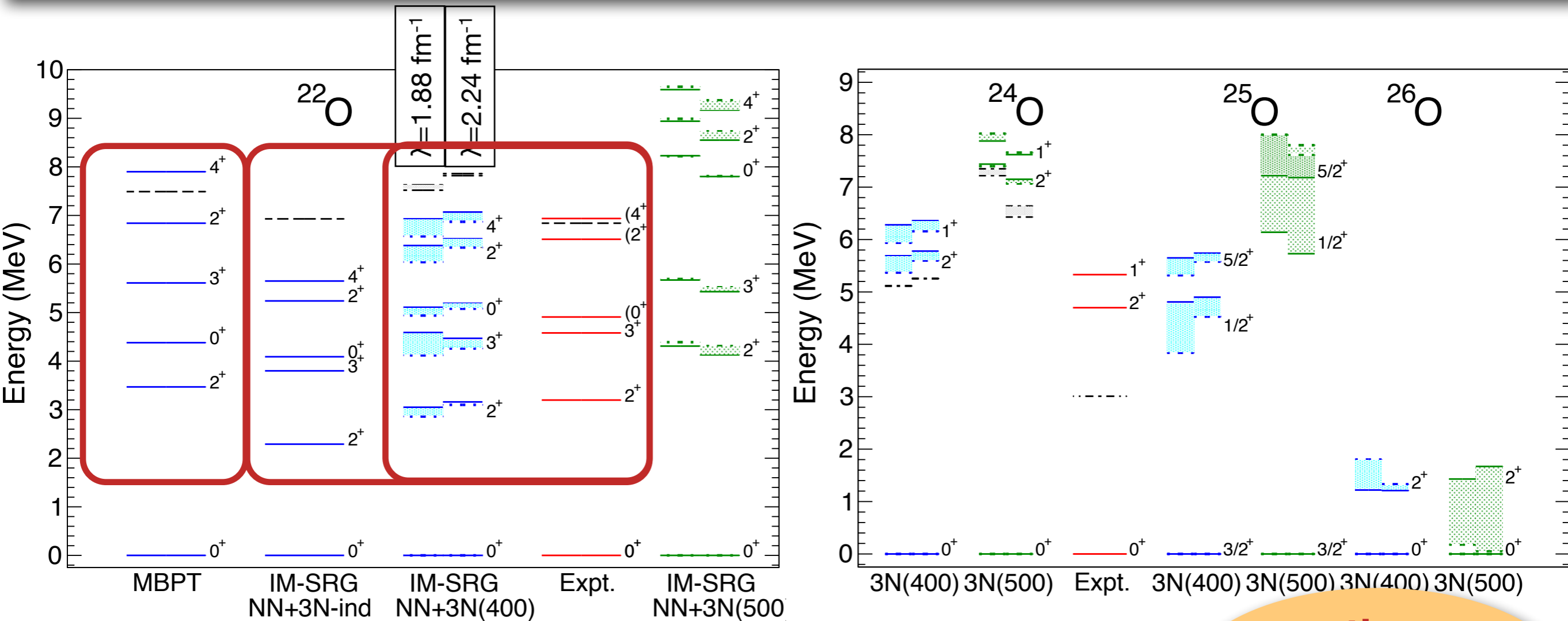
# Valence Space Decoupling



- construct generator from off-diagonal Hamiltonian

$$\{H^{od}\} = \{f_{h'}^h, f_{p'}^p, f_h^p, f_v^q, \Gamma_{hh'}^{pp'}, \Gamma_{hv}^{pp'}, \Gamma_{vv'}^{pq}\} \& \text{H.c.}$$

# From Oxygen...



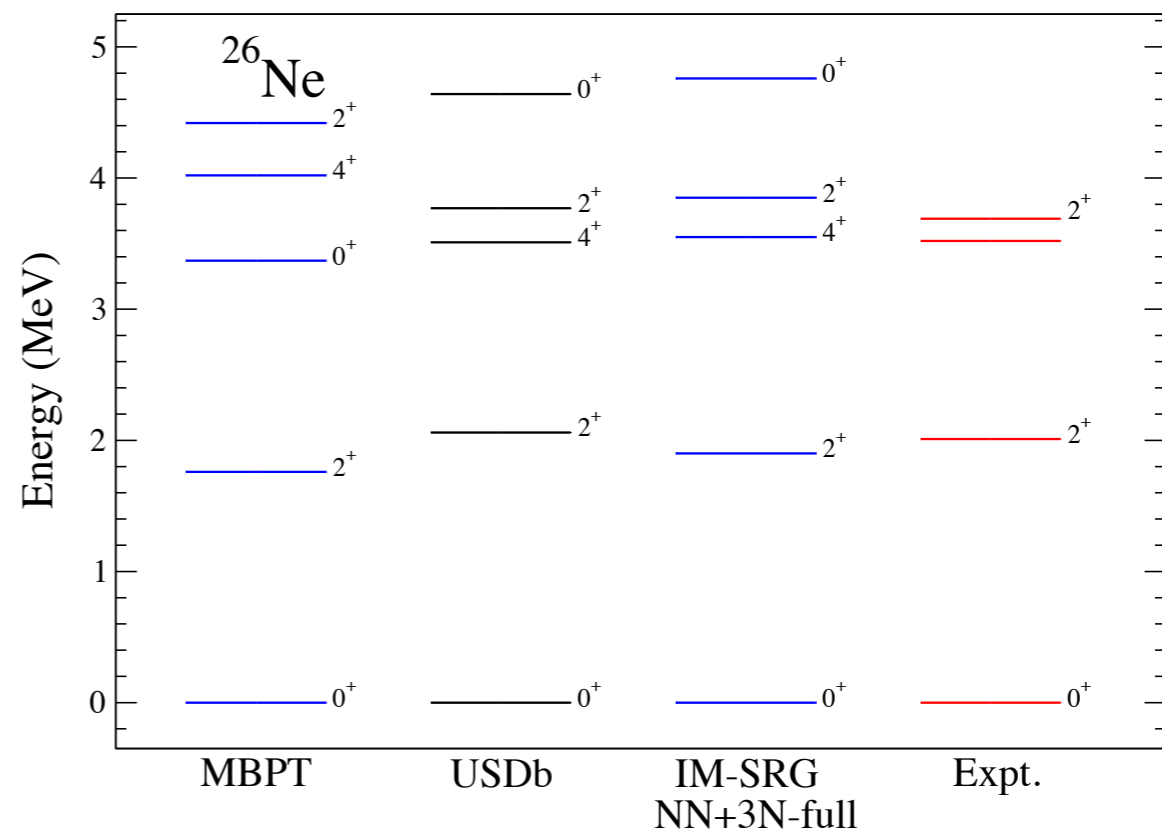
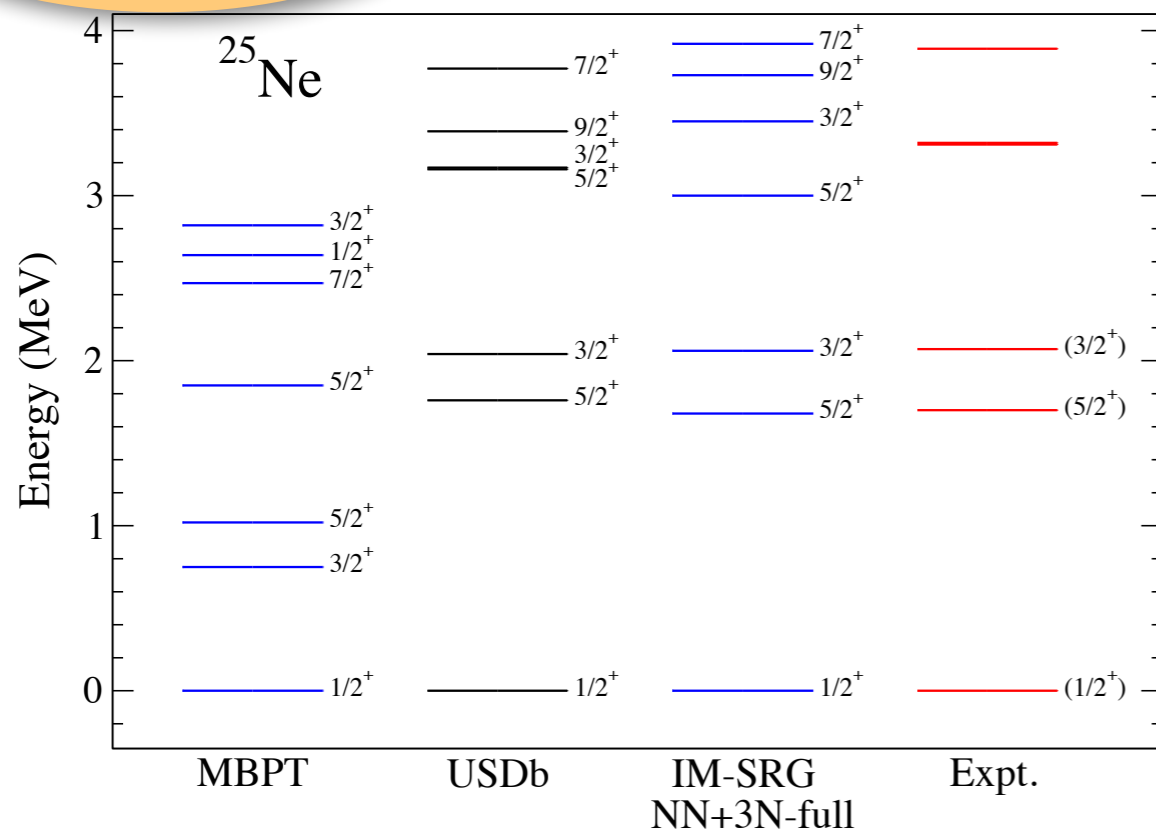
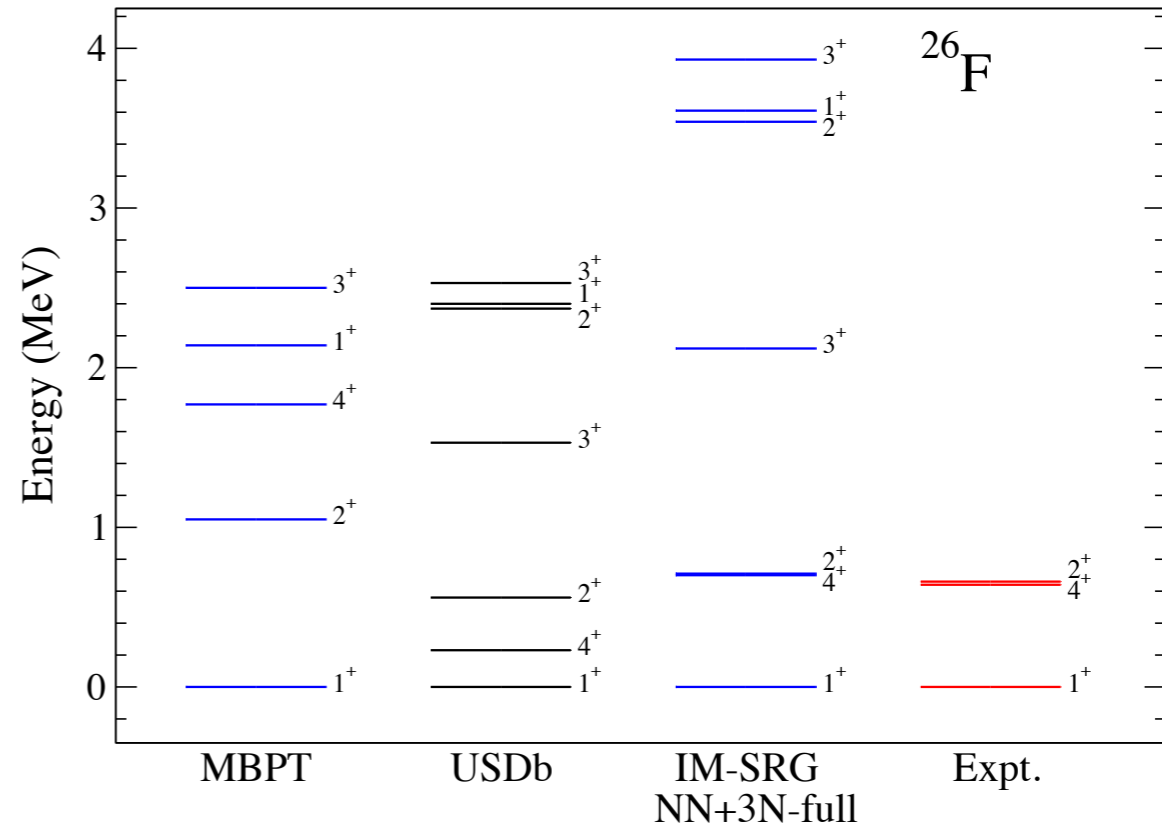
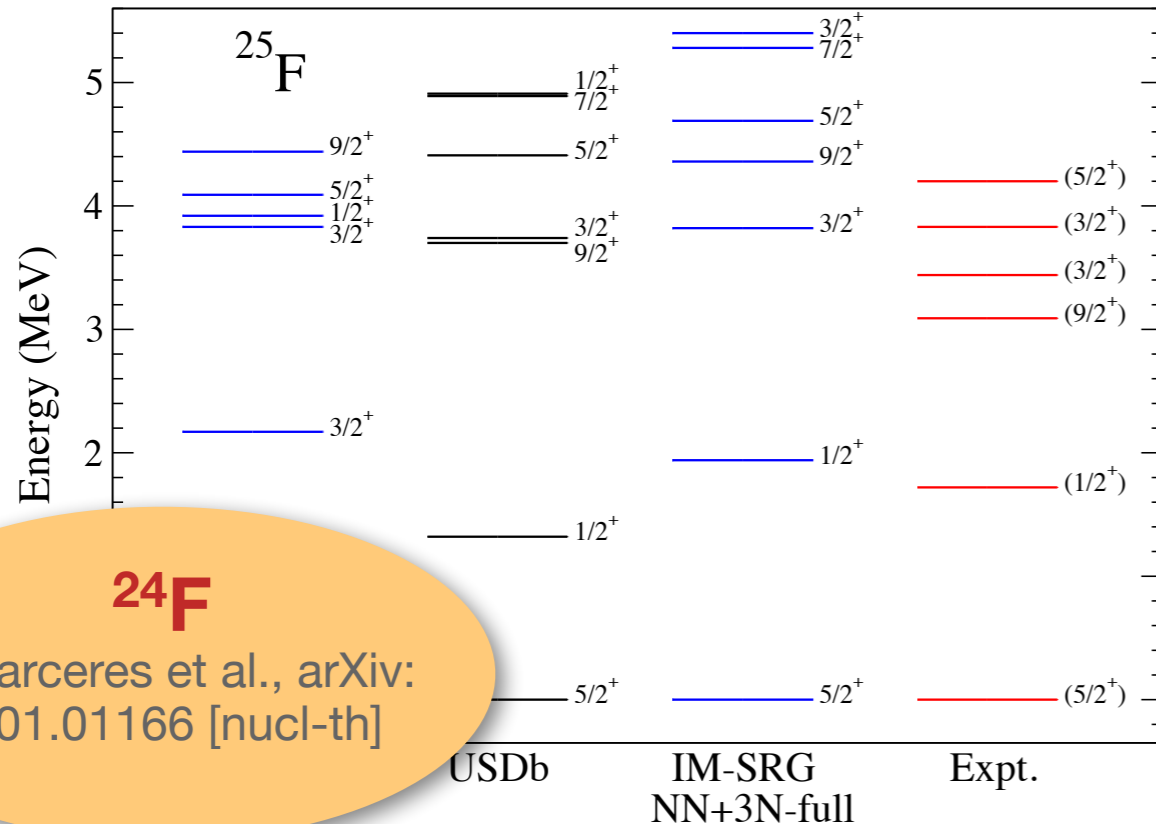
shading:  $\hbar\Omega$  variation

Phys. Rev. Lett. **113**, 142501 (2014)

**continuum lowers states by <1 MeV**

- **3N forces crucial**
- IM-SRG improves on finite-order MBPT effective interaction
- competitive with phenomenological calculations

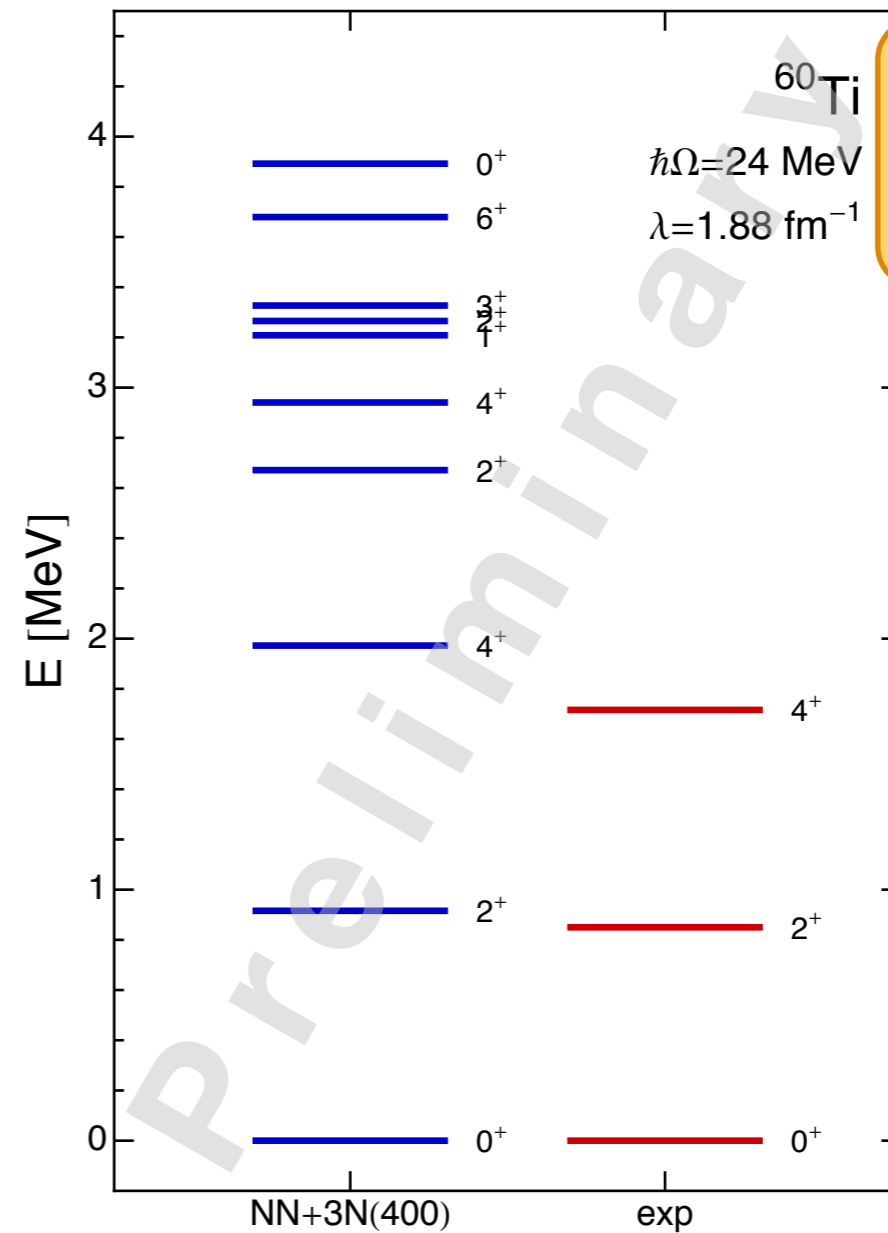
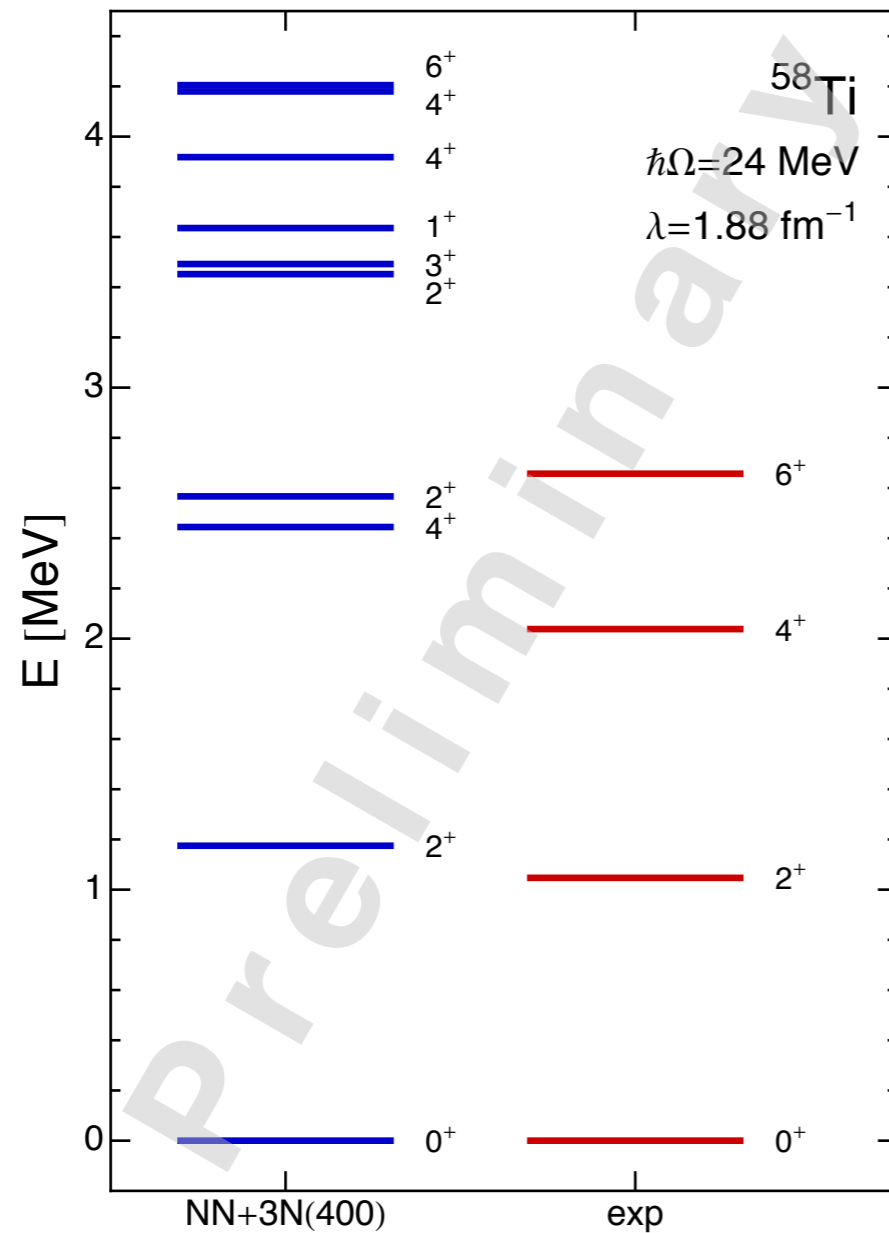
# ... Into the *sd*-Shell



**$^{24}\text{F}$**   
 L. Carceres et al., arXiv: 1501.01166 [nucl-th]



# Heavier Cores



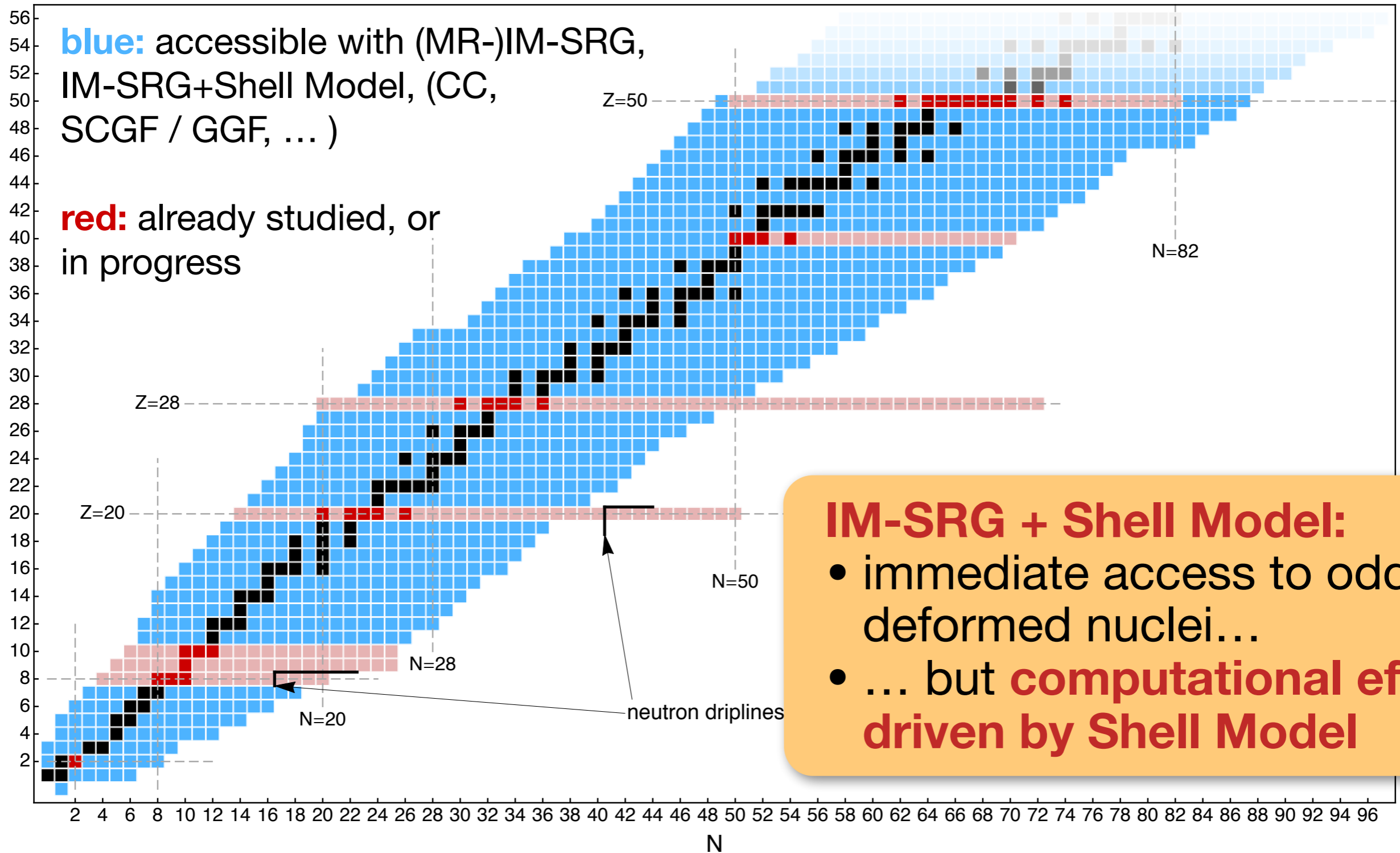
$^{48}\text{Ca}$  core,  
 $\nu\text{pf}_{5/2}$ ,  $\pi\text{pf}$   
 valence space

experimental data: A. Gade et al., Phys. Rev. Lett. **112**, 112503 (2014) and NNDC

➔ theoretical level scheme similar to empirical interactions (LNPS, GXPF1A)

# Next Steps

# Reach of Ab Initio Methods



**IM-SRG + Shell Model:**

- immediate access to odd & deformed nuclei...
- ... but **computational effort driven by Shell Model**

- describe “excited states” based on reference state:

$$|\psi_k\rangle \equiv R_k |\psi_0\rangle$$

- **(MR-)IM-SRG effective Hamiltonian** in EOM approach:

$$[H(\infty), R_k] = \omega_k R_k, \quad \omega_k = E_k - E_0$$

- computational effort scales **polynomially**, vs. factorial scaling of Shell Model
- can exploit Multi-Reference capabilities (commutator formulation identical to flow equations)

➔ **complementary** to Shell Model

- particle-hole excitations (TDA, RPA, Second RPA, ...)

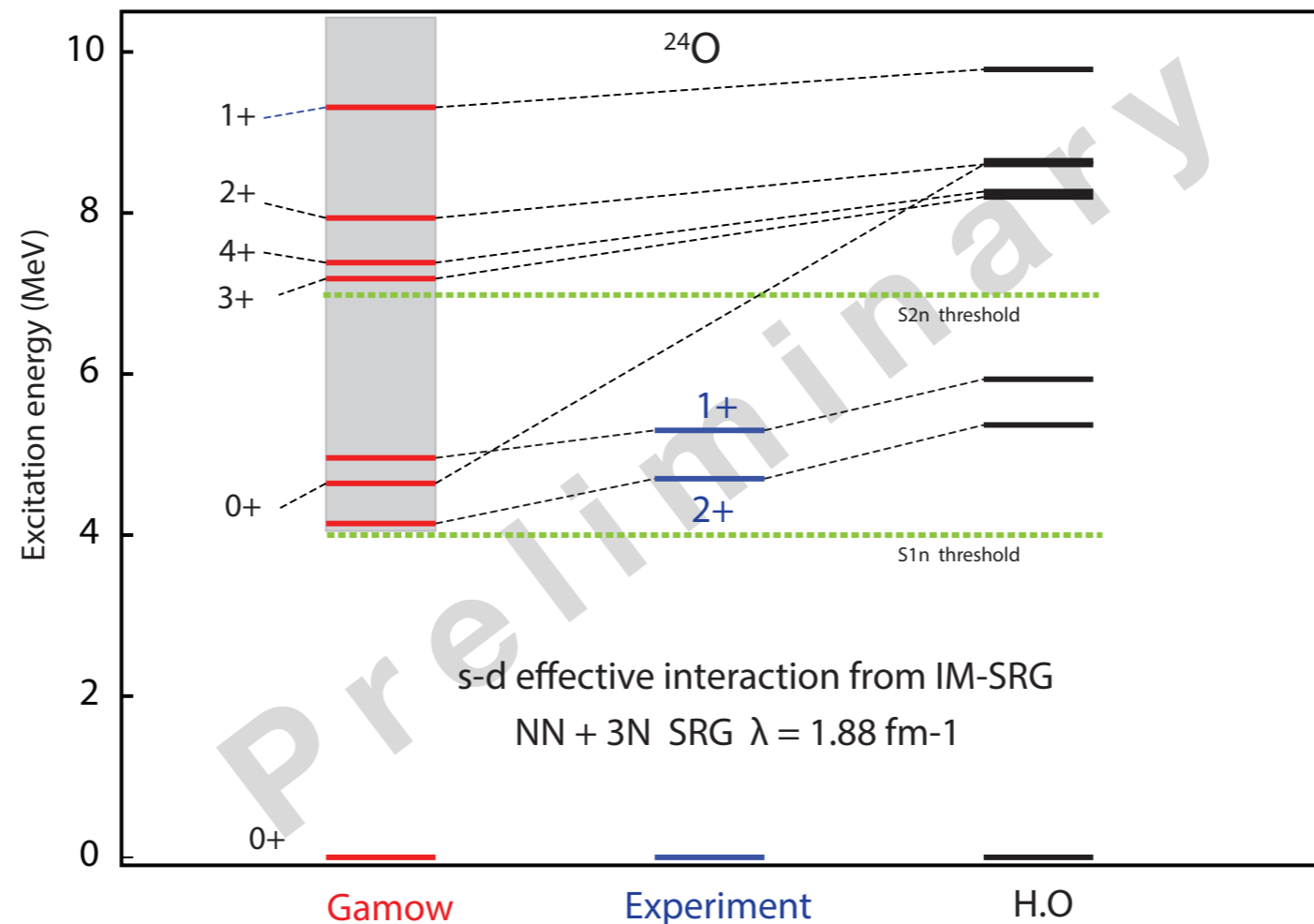
$$R_k = \sum_{ph} R_{ph}^{(k)} : a_p^\dagger a_h : + \sum_{pp'hh'} R_{pp'hh'}^{(k)} : a_p^\dagger a_{p'}^\dagger a_{h'} a_h : + \dots$$

➔ **giant resonances**

- particle attachment (analogous for removal):

$$R_k = \sum_{ph} R_p^{(k)} : a_p^\dagger : + \sum_{pp'h} R_{pp'h}^{(k)} : a_p^\dagger a_{p'}^\dagger a_h : + \dots$$

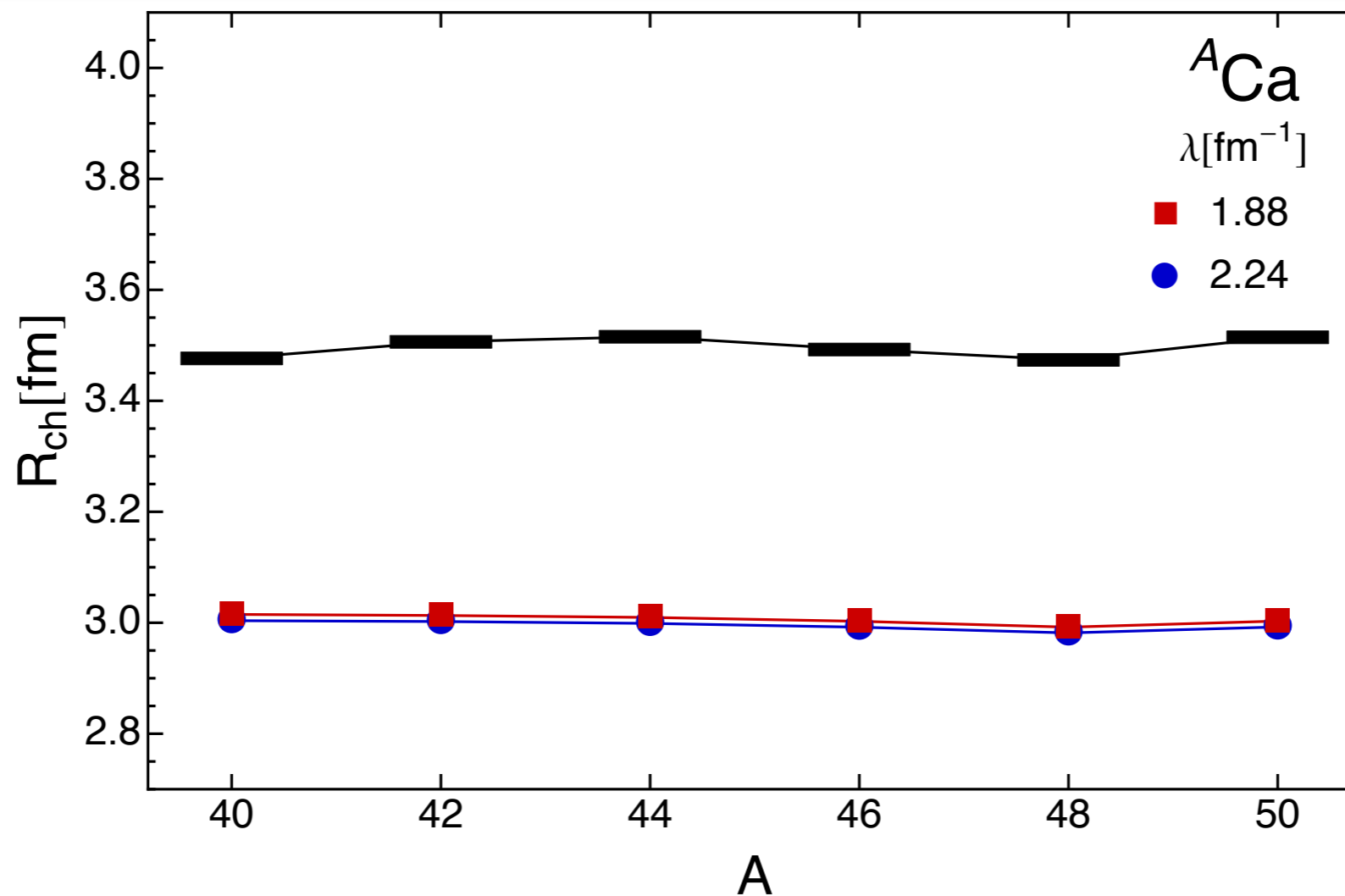
➔ **ground and excited states in odd nuclei**



(with G. Papadimitriou, ISU)

- **Gamow Shell Model:** use Berggren s.p. basis containing bound, resonance, and scattering states
- first step: GSM calculations with IM-SRG interaction
- **future:** include continuum in IM-SRG evolution

# Effective Operators



- small radii: **interaction issue** (power counting, regulators, LECs, ...)? importance of **currents**?
- implementation of electromagnetic & weak **transition operators in progress**; aim for **consistent treatment: chiral EFT, SRG, IM-SRG ( & Shell Model code !)**

# Conclusions



- enormous progress in *ab initio* nuclear structure and reactions, driven by (S)RG and (chiral) EFT
  - stringent link to QCD
  - enhanced control over many-body methods, **uncertainty quantification**
- IM-SRG is a powerful *ab initio* framework for closed- and open-shell, medium-mass & (heavy) nuclei
  - allows derivation of **Shell-Model interactions**
  - ➔ immediate access to **spectra, odd nuclei, intrinsic deformation** (at Shell Model numerical cost)
- new perspectives for old (?) problems: evolution of long-range correlations, construction of density functionals...

# Acknowledgments



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# Supplements

- SRG is a **unitary transformation** in **A-body space**
- up to **A-body interactions** are **induced** during the flow:

$$\frac{dH}{d\lambda} = \left[ \left[ \sum a^\dagger a, \underbrace{\sum a^\dagger a^\dagger aa}_{2\text{-body}} \right], \underbrace{\sum a^\dagger a^\dagger aa}_{2\text{-body}} \right] = \dots + \underbrace{\sum a^\dagger a^\dagger a^\dagger aaa}_{3\text{-body}} + \dots$$

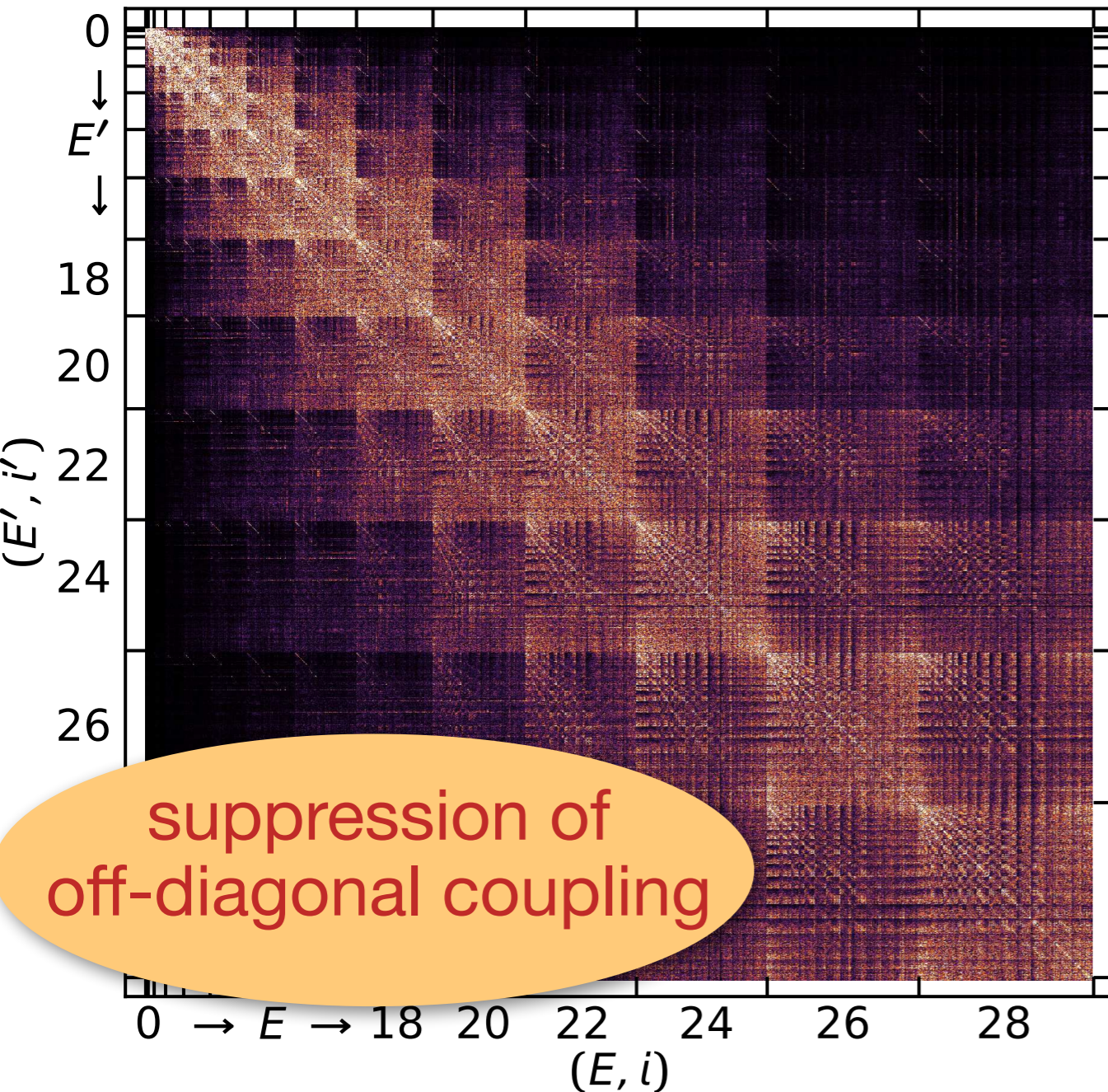
- state-of-the-art: evolve in three-body space, truncate induced four- and higher many-body forces  
(Jurgenson, Furnstahl, Navratil, PRL 103, 082501; Hebeler, PRC 85, 021002; Wendt, PRC 87, 061001 )
- **$\lambda$ -dependence** of eigenvalues is a **diagnostic** for size of omitted induced interactions

# SRG in Three-Body Space



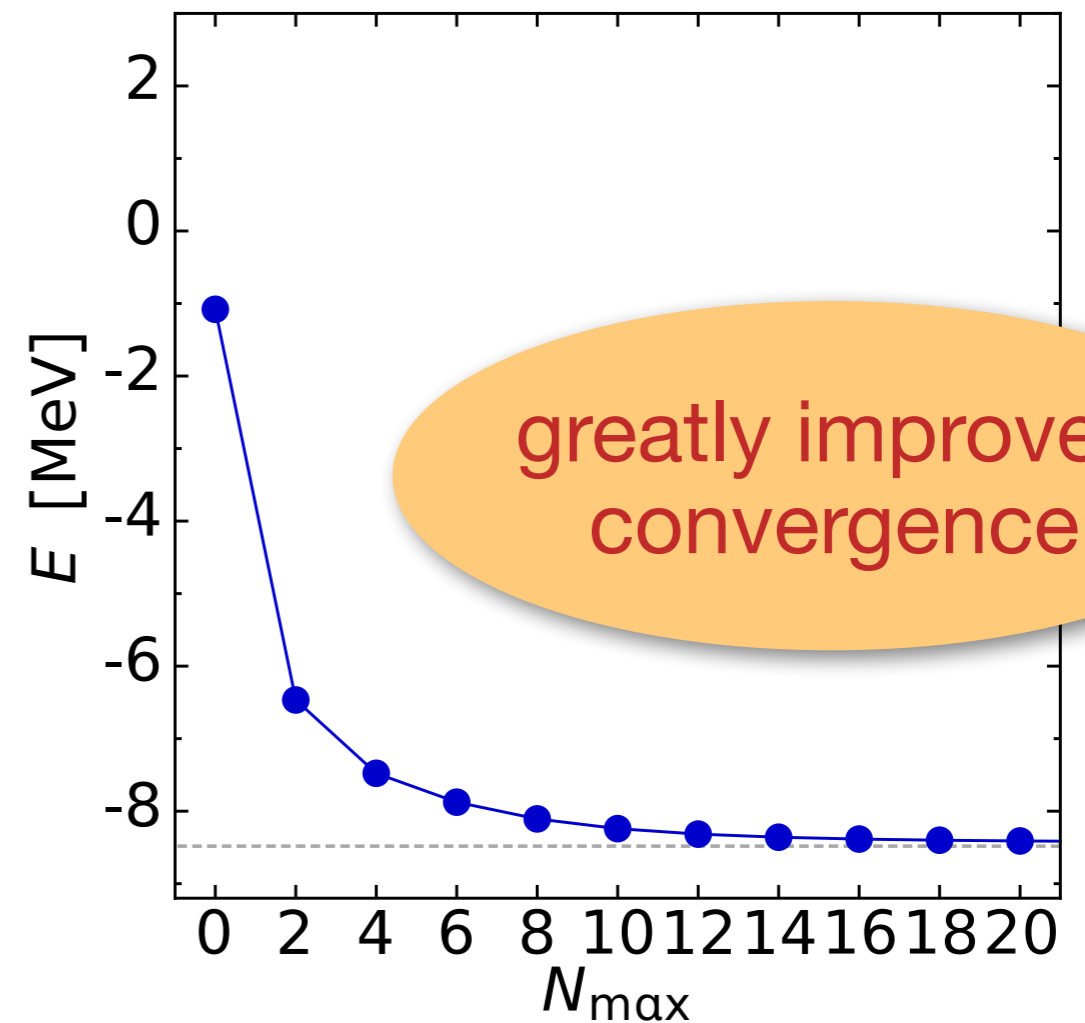
## 3B Jacobi-HO Matrix Elements

$$J^\pi = \frac{1}{2}^+, T = \frac{1}{2}, \hbar\Omega = 28 \text{ MeV}$$



$$\lambda = 1.33 \text{ fm}^{-1}$$

## ${}^3\text{H}$ ground-state (NCSM)



[figures by R. Roth, A. Calci, J. Langhammer]

# Choice of Generator



- **Wegner:**  $\eta^I = [H^d, H^{od}]$

- **White:** (J. Chem. Phys. 117, 7472)

$$\eta^{II} = \sum_{ph} \frac{f_h^p}{\Delta_h^p} : A_h^p : + \sum_{pp'hh'} \frac{\Gamma_{hh'}^{pp'}}{\Delta_{hh'}^{pp'}} : A_{hh'}^{pp'} : + \text{H.c.}$$

$\Delta_h^p, \Delta_{hh'}^{pp'}$  : approx. 1p1h, 2p2h excitation energies

- **“imaginary time”:** (Morris, Bogner)

$$\eta^{III} = \sum_{ph} \text{sgn}(\Delta_h^p) f_h^p : A_h^p : + \sum_{pp'hh'} \text{sgn}(\Delta_{hh'}^{pp'}) \Gamma_{hh'}^{pp'} : A_{hh'}^{pp'} : + \text{H.c.}$$

- off-diagonal matrix elements are suppressed like  $e^{-\Delta^2 s}$  (Wegner),  $e^{-s}$  (White), and  $e^{-|\Delta|s}$  (imaginary time)
- g.s. energies ( $s \rightarrow \infty$ ) differ by  $\ll 1\%$

## 0-body Flow

$$\frac{dE}{ds} = \sum_{ab} (n_a - n_b) \left( \eta_b^a f_a^b - f_b^a \eta_a^b \right) + \frac{1}{4} \sum_{abcd} \left( \eta_{cd}^{ab} \Gamma_{ab}^{cd} - \Gamma_{cd}^{ab} \eta_{ab}^{cd} \right) n_a n_b \bar{n}_c \bar{n}_d$$

~ 2nd order MBPT for  $H(s)$

## 1-body Flow

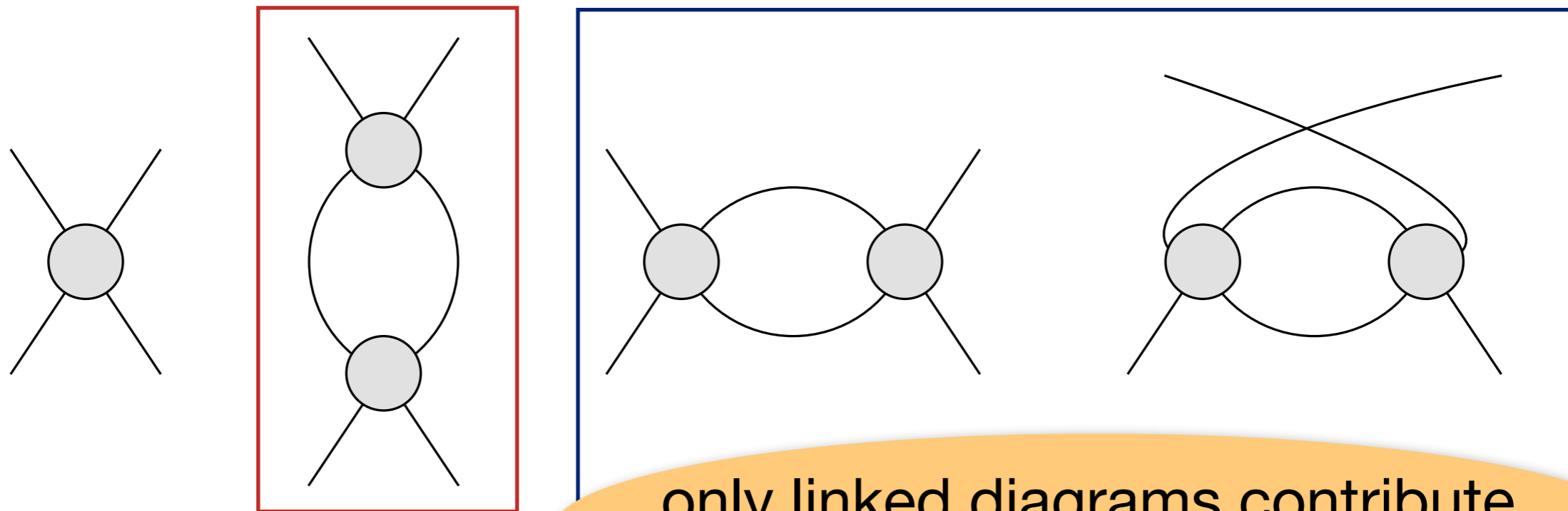
$$\begin{aligned} \frac{d}{ds} f_2^1 = & \sum_a \left( \eta_a^1 f_2^a - f_a^1 \eta_2^a \right) + \sum_{ab} \left( \eta_b^a \Gamma_{a2}^{b1} - f_b^a \eta_{a2}^{b1} \right) (n_a - n_b) \\ & + \frac{1}{2} \sum_{abcdef} \left( \eta_{bc}^{1a} \Gamma_{2a}^{bc} - \Gamma_{bc}^{1a} \eta_{2a}^{bc} \right) (n_a \bar{n}_b \bar{n}_c + \bar{n}_a n_b n_c) \end{aligned}$$

# In-Medium SRG Flow Equations



## 2-body Flow

$$\begin{aligned} \frac{d}{ds} \Gamma_{34}^{12} = & \sum_a \left( \eta_a^1 \Gamma_{34}^{a2} + \eta_a^2 \Gamma_{34}^{1a} - \eta_3^a \Gamma_{a4}^{12} - \eta_4^a \Gamma_{3a}^{12} - f_a^1 \eta_{34}^{a2} - f_a^2 \eta_{34}^{1a} + f_3^a \eta_{a4}^{12} + f_4^a \eta_{3a}^{12} \right) \\ & + \frac{1}{2} \sum_{ab} \left( \eta_{ab}^{12} \Gamma_{34}^{ab} - \Gamma_{ab}^{12} \eta_{34}^{ab} \right) (1 - n_a - n_b) \\ & + \sum_{ab} (n_a - n_b) \left( \left( \eta_{3b}^{1a} \Gamma_{4a}^{2b} - \Gamma_{3b}^{1a} \eta_{4a}^{2b} \right) - \left( \eta_{3b}^{2a} \Gamma_{4a}^{1b} - \Gamma_{3b}^{2a} \eta_{4a}^{1b} \right) \right) \end{aligned}$$



only linked diagrams contribute,  
IM-SRG **size-extensive**



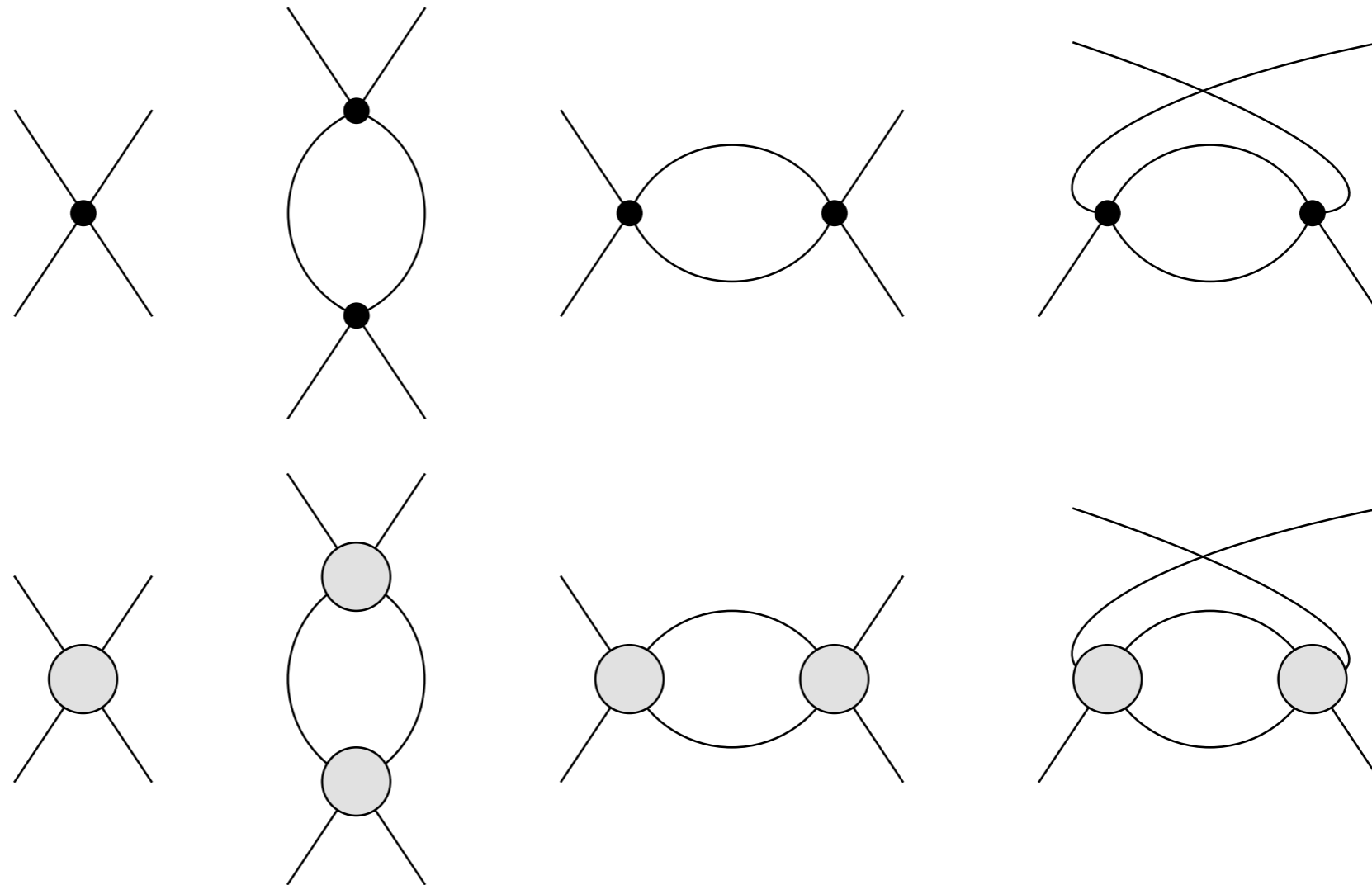
# In-Medium SRG Flow: Diagrams



$\Gamma(\delta s) \sim$



$\Gamma(2\delta s) \sim$



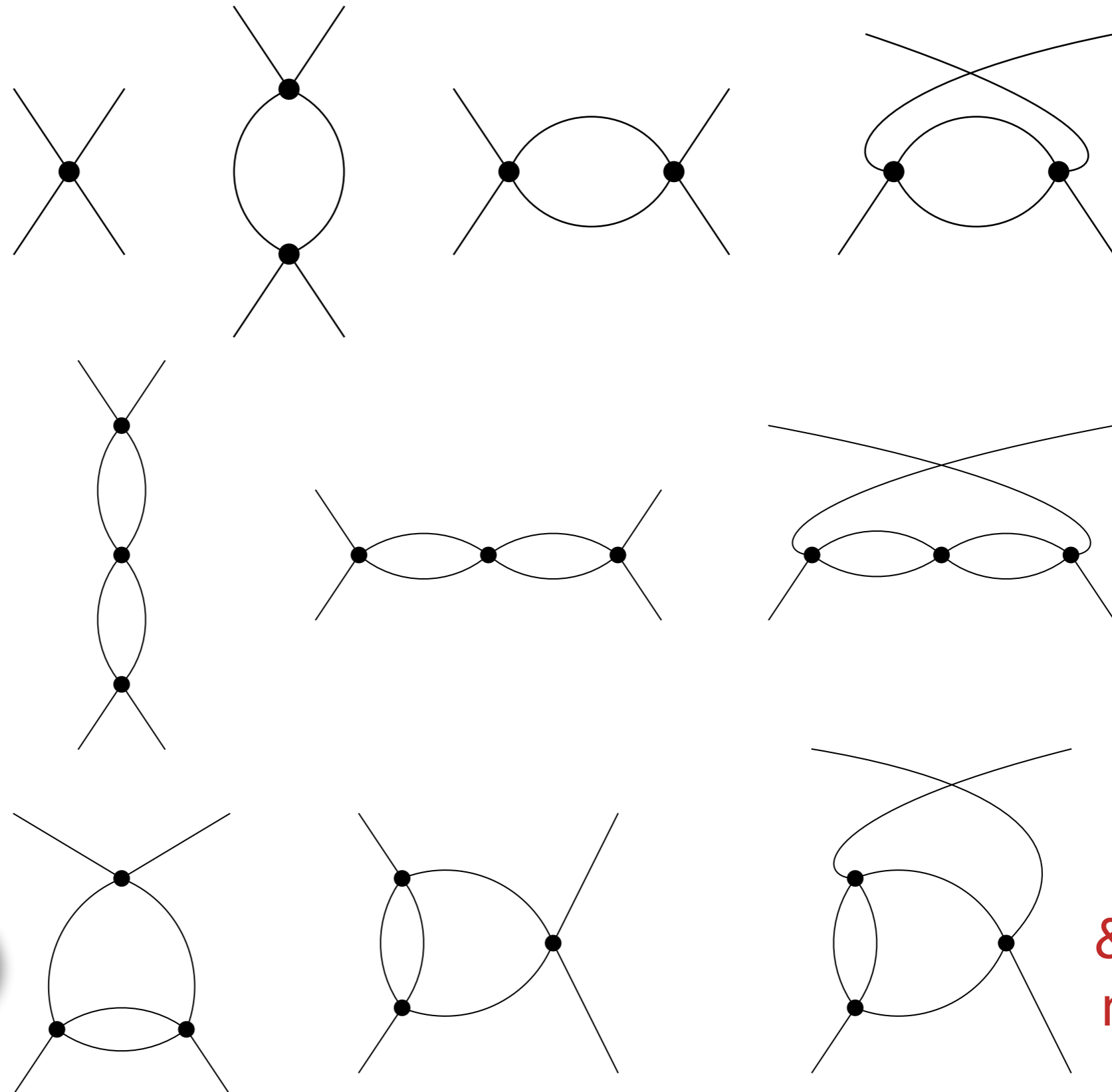
# In-Medium SRG Flow: Diagrams



$$\Gamma(\delta s) \sim$$



$$\Gamma(2\delta s) \sim$$



non-  
perturbative  
resummation

& many  
more...

# Generalized Normal Ordering



- generalized Wick's theorem for **arbitrary reference states** (Kutzelnigg & Mukherjee)
- define **irreducible n-body density matrices** of reference state:

$$\rho_{mn}^{kl} = \lambda_{mn}^{kl} + \lambda_m^k \lambda_n^l - \lambda_n^k \lambda_m^l$$

$$\rho_{lmn}^{ijk} = \lambda_{lmn}^{ijk} + \lambda_l^i \lambda_{mn}^{jk} + \lambda_l^j \lambda_{mn}^{ik} + \lambda_l^k \lambda_{mn}^{ij} + \text{permutations}$$

...

- irreducible densities give rise to **additional contractions**:

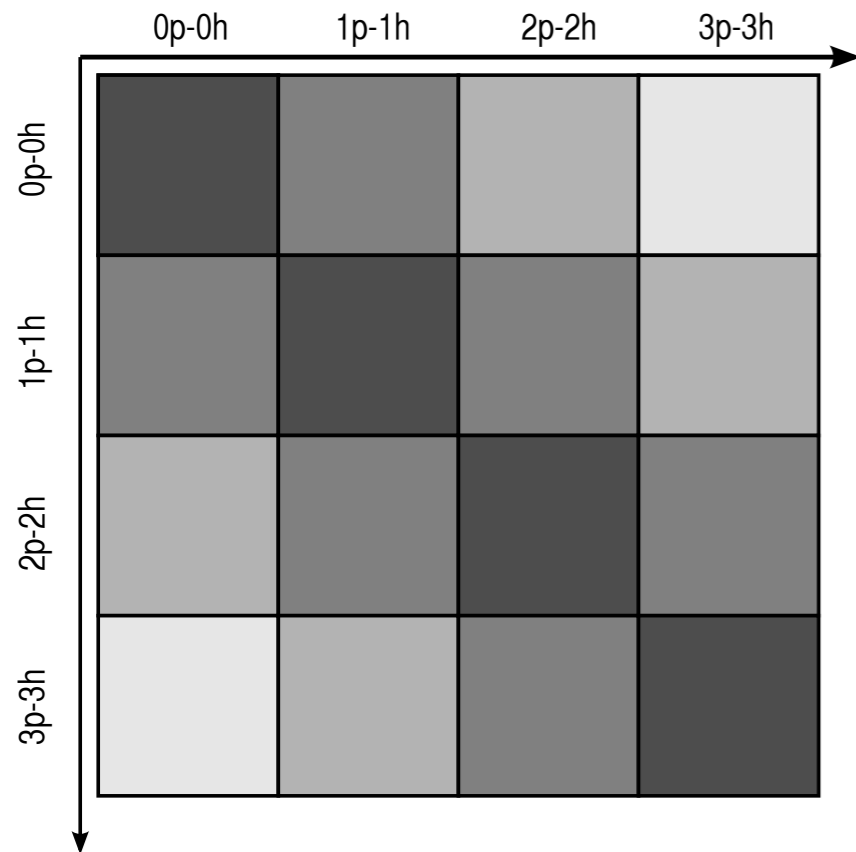
$$: A_{cd\dots}^{ab\dots} :: A_{mn\dots}^{kl\dots} : \longrightarrow \lambda_{mn}^{ab}$$

$$: A_{cd\dots}^{ab\dots} :: A_{mn\dots}^{kl\dots} : \longrightarrow$$

...

**two-body flow unchanged,  
O(N<sup>6</sup>) scaling preserved**

# Decoupling



$$\langle \begin{smallmatrix} p \\ h \end{smallmatrix} | H | \Psi \rangle \sim f_h^p, \sum_{kl} f_l^k \lambda_{pl}^{hk}, \sum_{klmn} \Gamma_{mn}^{kl} \lambda_{pmn}^{hkl}, \dots$$

$$\langle \begin{smallmatrix} pp' \\ hh' \end{smallmatrix} | H | \Psi \rangle \sim \Gamma_{hh'}^{pp'}, \sum_{km} \Gamma_{hm}^{pk} \lambda_{p'm}^{h'k}, \sum_{kl} f_l^k \lambda_{pp'l}^{hh'k}, \sum_{klmn} \Gamma_{mn}^{kl} \lambda_{pp'mn}^{hh'kl}, \dots$$

$$\langle \begin{smallmatrix} pp'p'' \\ hh'hh' \end{smallmatrix} | H | \Psi \rangle \sim \dots$$

- truncation in irreducible density matrices
- number of **correlated vs. total** pairs, triples, ... (**caveat:** highly collective reference states)
- perturbative analysis (e.g. for shell-model like states)
- **verify for chosen multi-reference state when possible**

- HFB ground state is a **superposition** of states with **different particle number**:

$$|\Psi\rangle = \sum_{A=N, N\pm 2, \dots} c_A |\Psi_A\rangle, \quad |\Psi_N\rangle \equiv P_N |\Psi\rangle \equiv \frac{1}{2\pi} \int_0^{2\pi} d\phi e^{i\phi(\hat{N}-N)} |\Psi\rangle$$

- calculate one- and two-body densities (**project only once**):

$$\lambda_i^k = \frac{\langle \Psi | A_i^k P_N | \Psi \rangle}{\langle \Psi | \Psi \rangle}, \quad \lambda_{mn}^{kl} = \frac{\langle \Psi | A_{mn}^{kl} P_N | \Psi \rangle}{\langle \Psi | \Psi \rangle} - \lambda_m^k \lambda_m^l + \lambda_n^k \lambda_m^l$$

- work in natural orbitals (= HFB **canonical basis**):

$$\lambda_i^k = n_k \delta_i^k (= v_k^2 \delta_i^k), \quad 0 \leq n_k \leq 1$$

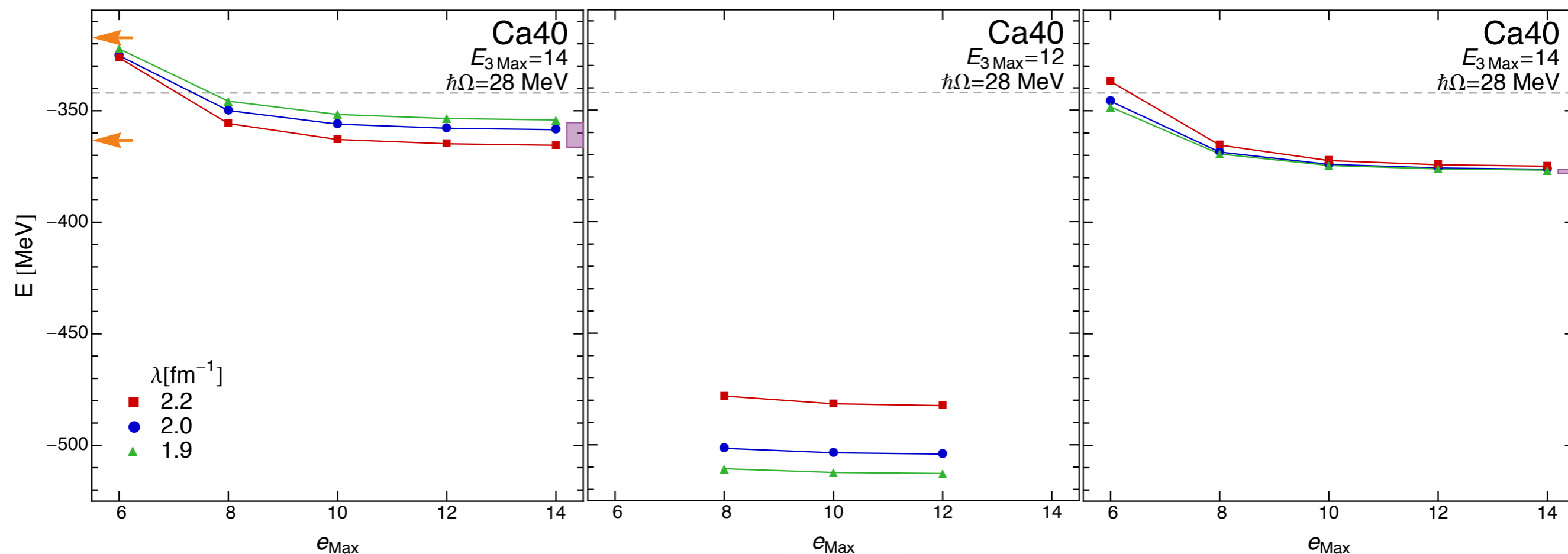
# Results: Closed-Shell Nuclei



NN + 3N-ind.

NN + 3N-full (500)

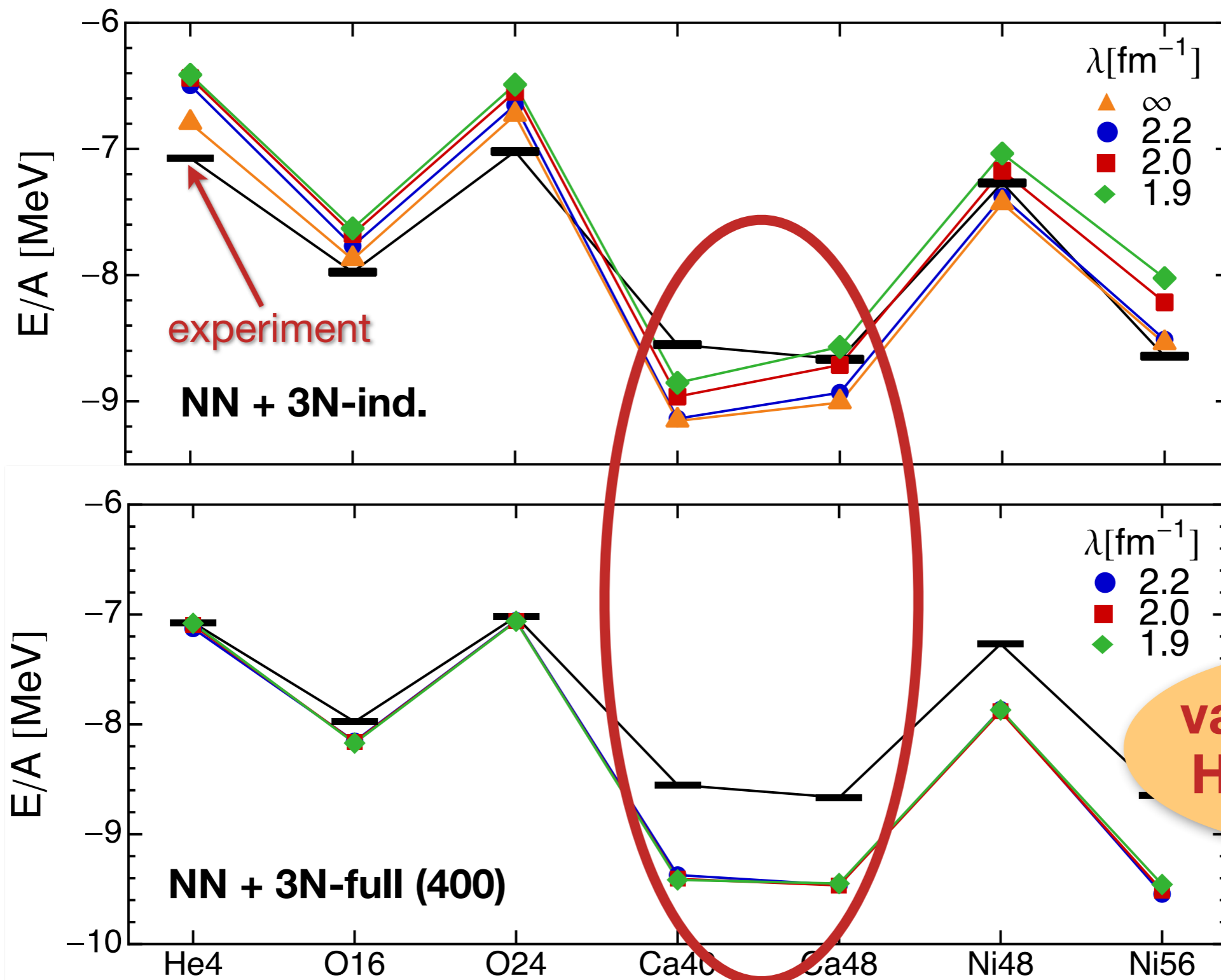
NN + 3N-full (400)



validate chiral  
Hamiltonians

- ← CCSD/ $\Lambda$ -CCSD(T),  $\lambda = \infty$ , G. Hagen et al., PRL 109, 032502 (2012)
- ▣  $\Lambda$ -CCSD(T),  $\lambda = 1.9 - 2.2 \text{ fm}^{-1}$ , S. Binder et al., arXiv:1211.4748 [nucl-th] & PRL 109, 052501 (2012)

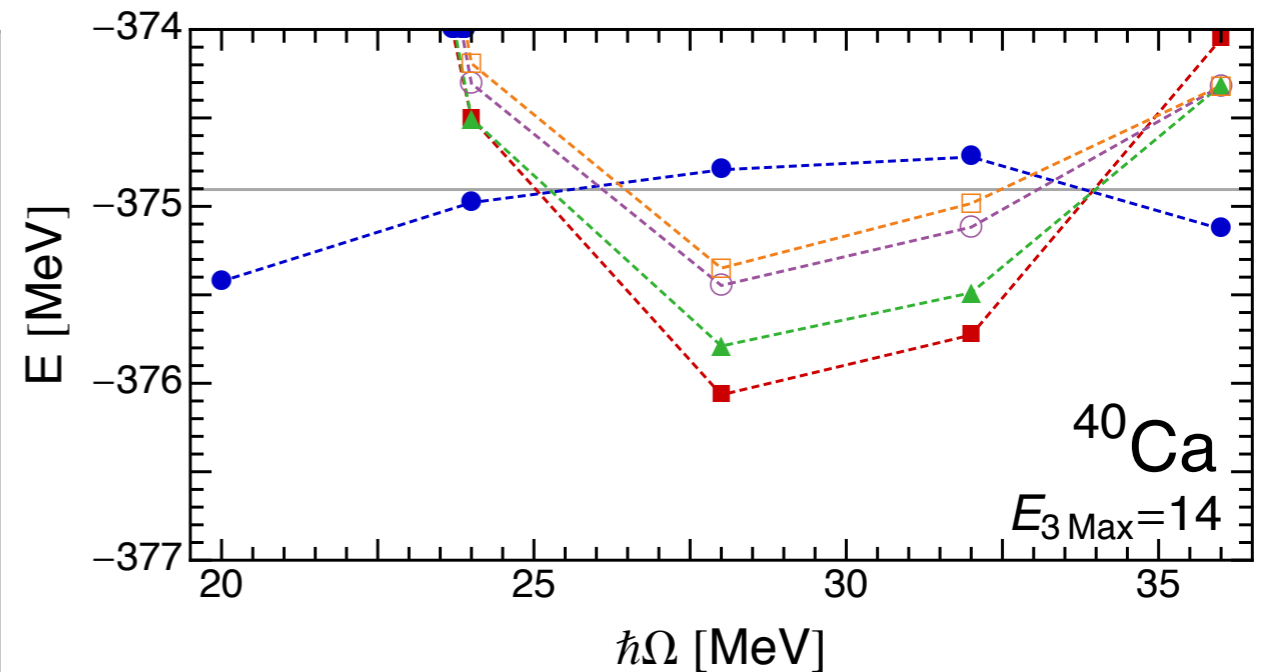
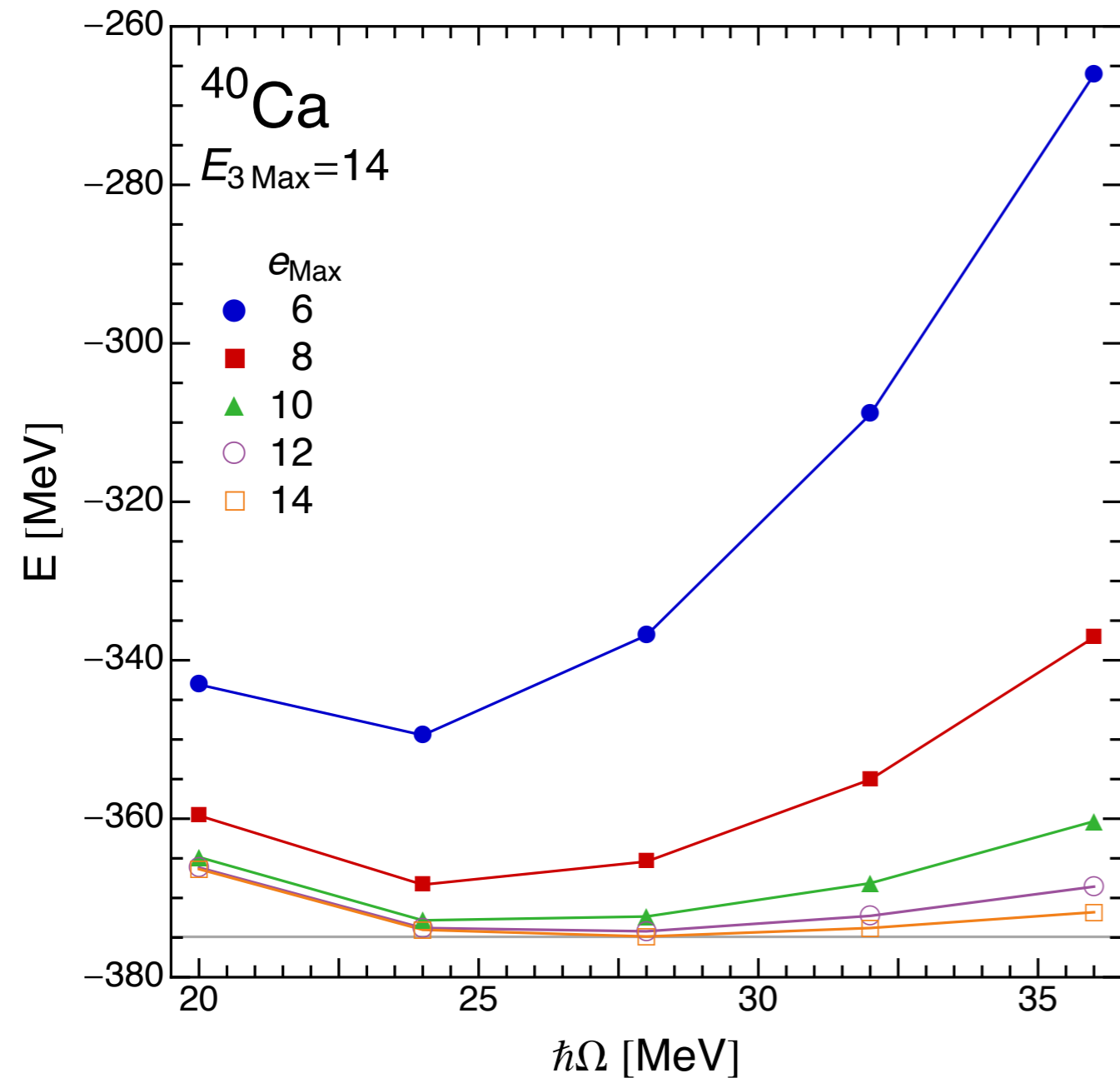
# Results: Closed-Shell Nuclei



validate chiral Hamiltonians

Phys. Rev. C **87**, 034307 (2013), arXiv: 1212.1190 [nucl-th]

# Extrapolation



max./UV momentum:

$$\Lambda_{\text{UV}} = \sqrt{2m(e_{\text{Max}} + 3/2)\hbar\Omega}$$

radial extent:

$$L = \sqrt{2(e_{\text{Max}} + 3/2)\hbar/m\Omega}$$

simultaneous ultraviolet & infrared extrapolation:

$$E(\Lambda_{\text{UV}}, L) = E_{\infty} + A_0 \exp\left(-2\Lambda_{\text{UV}}^2/A_1^2\right) + A_2 \exp(-2k_{\infty}L)$$

(R. Furnstahl, G. Hagen & T. Papenbrock, PRC 86,031301 (2012))



# Multi-Reference Flow Equations



0-body flow:

$$\begin{aligned} \frac{dE}{ds} = & \sum_{ab} (n_a - n_b) \left( \eta_b^a f_a^b - f_b^a \eta_a^b \right) + \frac{1}{4} \sum_{abcd} \left( \eta_{cd}^{ab} \Gamma_{ab}^{cd} - \Gamma_{cd}^{ab} \eta_{ab}^{cd} \right) n_a n_b \bar{n}_c \bar{n}_d \\ & + \frac{1}{4} \sum_{abcd} \left( \frac{d}{ds} \Gamma_{cd}^{ab} \right) \lambda_{cd}^{ab} + \frac{1}{4} \sum_{abcdklm} \left( \eta_{cd}^{ab} \Gamma_{am}^{kl} - \Gamma_{cd}^{ab} \eta_{am}^{kl} \right) \lambda_{cdm}^{bkl} \end{aligned}$$

1-body flow:

$$\begin{aligned} \frac{d}{ds} f_2^1 = & \sum_a \left( \eta_a^1 f_2^a - f_a^1 \eta_2^a \right) + \sum_{ab} \left( \eta_b^a \Gamma_{a2}^{b1} - f_b^a \eta_{a2}^{b1} \right) (n_a - n_b) \\ & + \frac{1}{2} \sum_{abcdef} \left( \eta_{bc}^{1a} \Gamma_{2a}^{bc} - \Gamma_{bc}^{1a} \eta_{2a}^{bc} \right) (n_a \bar{n}_b \bar{n}_c + \bar{n}_a n_b n_c) \\ & + \frac{1}{4} \sum_{abcde} \left( \eta_{bc}^{1a} \Gamma_{2a}^{de} - \Gamma_{bc}^{1a} \eta_{2a}^{de} \right) \lambda_{bc}^{de} + \sum_{abcde} \left( \eta_{bc}^{1a} \Gamma_{2d}^{be} - \Gamma_{bc}^{1a} \eta_{2d}^{be} \right) \lambda_{cd}^{ae} \\ & - \frac{1}{2} \sum_{abcde} \left( \eta_{2b}^{1a} \Gamma_{ae}^{cd} - \Gamma_{2b}^{1a} \eta_{ae}^{cd} \right) \lambda_{be}^{cd} + \frac{1}{2} \sum_{abcde} \left( \eta_{2b}^{1a} \Gamma_{de}^{bc} - \Gamma_{2b}^{1a} \eta_{de}^{bc} \right) \lambda_{de}^{ac} \end{aligned}$$

## 2-body flow:

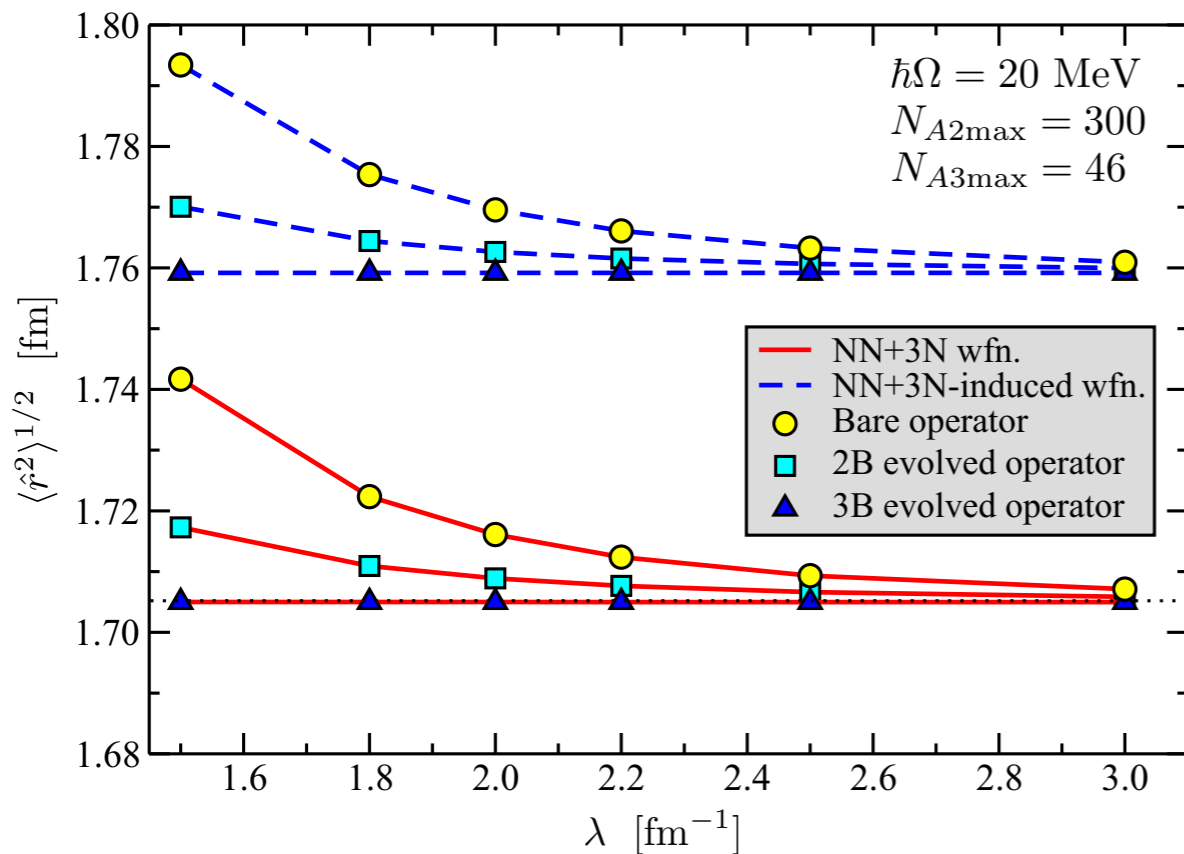
$$\begin{aligned} \frac{d}{ds} \Gamma_{34}^{12} = & \sum_a \left( \eta_a^1 \Gamma_{34}^{a2} + \eta_a^2 \Gamma_{34}^{1a} - \eta_3^a \Gamma_{a4}^{12} - \eta_4^a \Gamma_{3a}^{12} - f_a^1 \eta_{34}^{a2} - f_a^2 \eta_{34}^{1a} + f_3^a \eta_{a4}^{12} + f_4^a \eta_{3a}^{12} \right) \\ & + \frac{1}{2} \sum_{ab} \left( \eta_{ab}^{12} \Gamma_{34}^{ab} - \Gamma_{ab}^{12} \eta_{34}^{ab} \right) (1 - n_a - n_b) \\ & + \sum_{ab} (n_a - n_b) \left( \left( \eta_{3b}^{1a} \Gamma_{4a}^{2b} - \Gamma_{3b}^{1a} \eta_{4a}^{2b} \right) - \left( \eta_{3b}^{2a} \Gamma_{4a}^{1b} - \Gamma_{3b}^{2a} \eta_{4a}^{1b} \right) \right) \end{aligned}$$

2-body flow  
unchanged

# Effective Operators



${}^3\text{H}$  rms matter radius



from: Schuster et al., PRC90, 011301 (2014)

- derive operators from chiral EFT, **including currents**
- **optimize LECs together with interaction**
- **evolve** to desired resolution scale
- evaluate operator (1B+2B +...) in IM-SRG (and Shell Model)
- (most) existing ab initio & Shell model codes lack capabilities for many-body observables