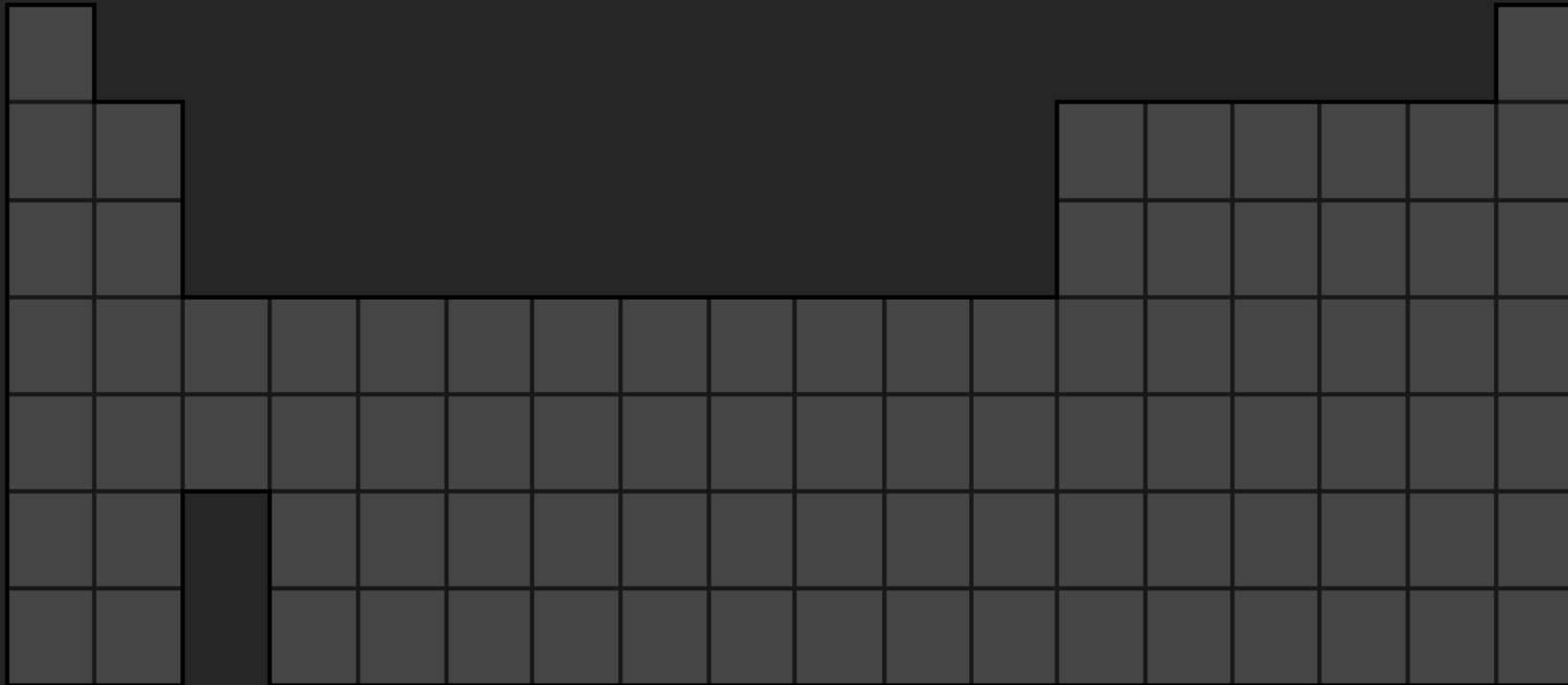


Kilonova Modelling: Nuclear Physics, Magnetic Fields, Neutrinos

Kelsey Lund

20 September 2022

Introduction: High Opacity Material



Lanthanides

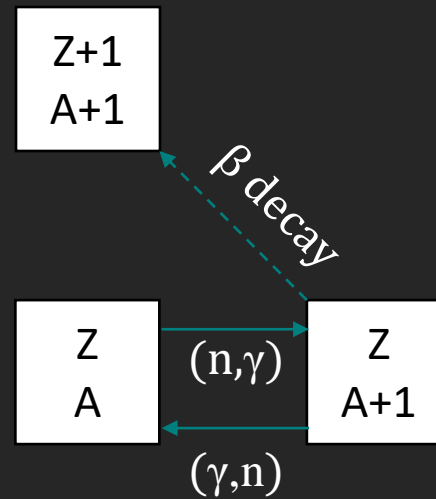
57	58	59	60	61	62	63	64	65	66	67	68	69	70	71
La	Ce	Pr	Nd	Pm	Sm	Eu	Gd	Tb	Dy	Ho	Er	Tm	Yb	Lu

Actinides

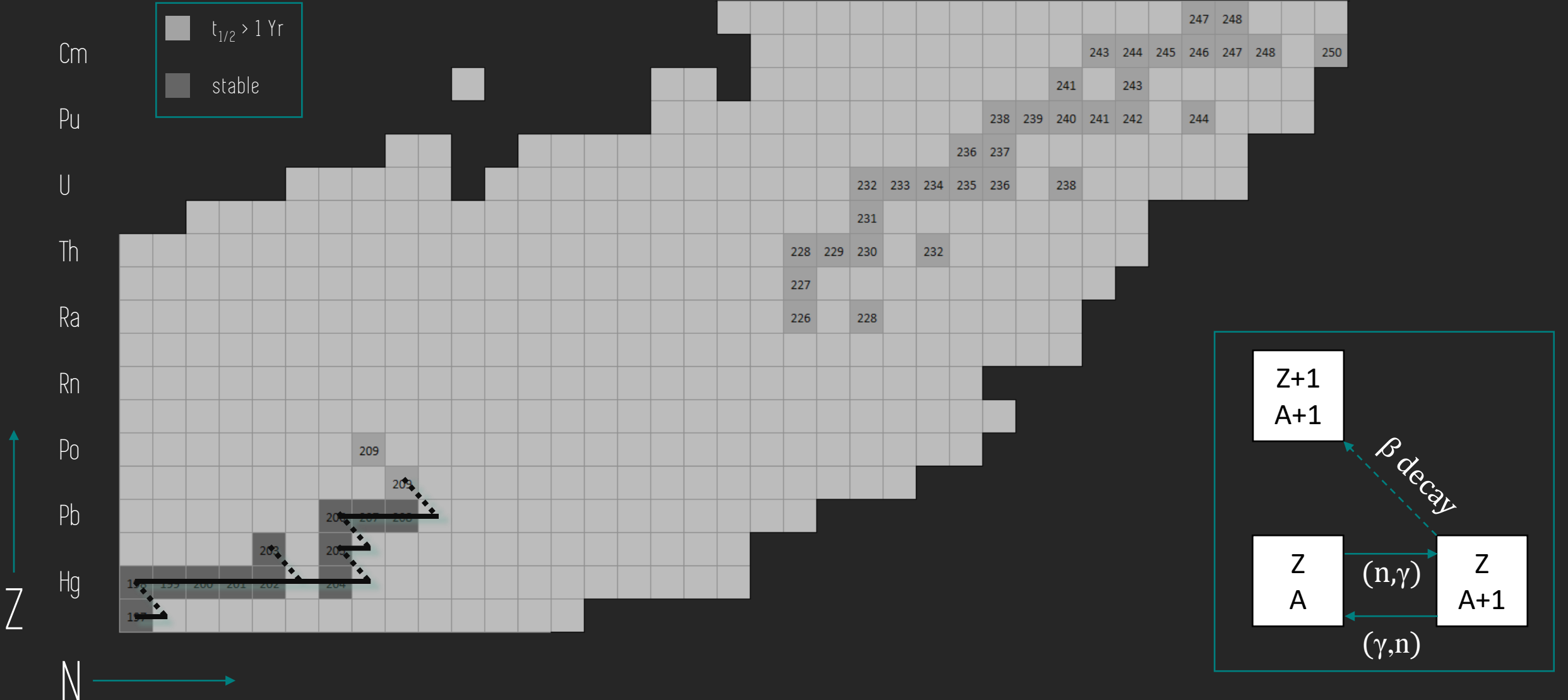
89	90	91	92	93	94	95	96	97	98	99	100	101	102	103
Ac	Th	Pa	U	Np	Pu	Am	Cm	Bk	Cf	Es	Fm	Md	No	Lr

(recall
weekend talks
by
Gail McLaughlin
&
Jonah Miller)

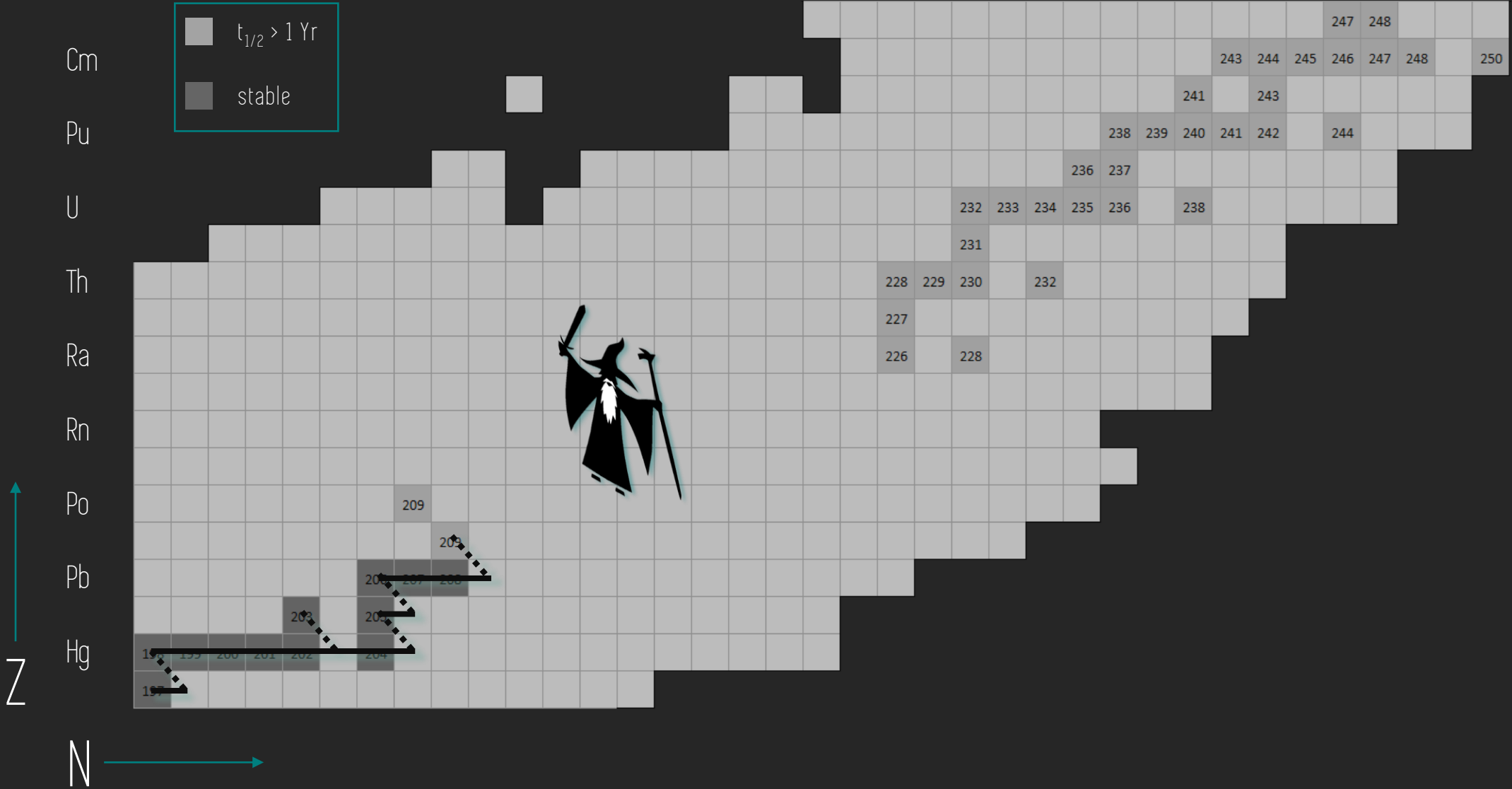
Introduction: Neutron Capture



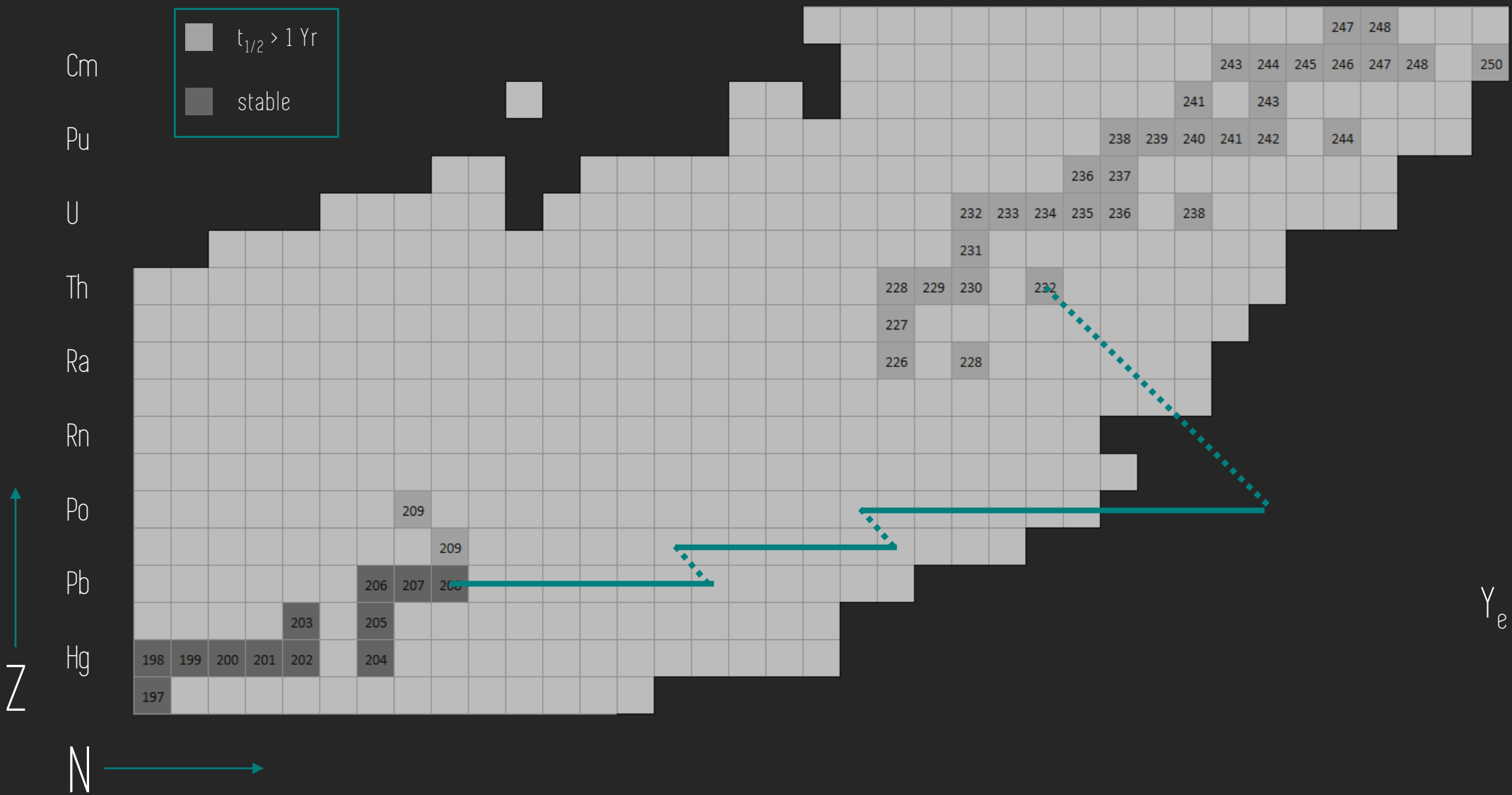
Introduction: (slow) Neutron Capture



Introduction: (slow) Neutron Capture



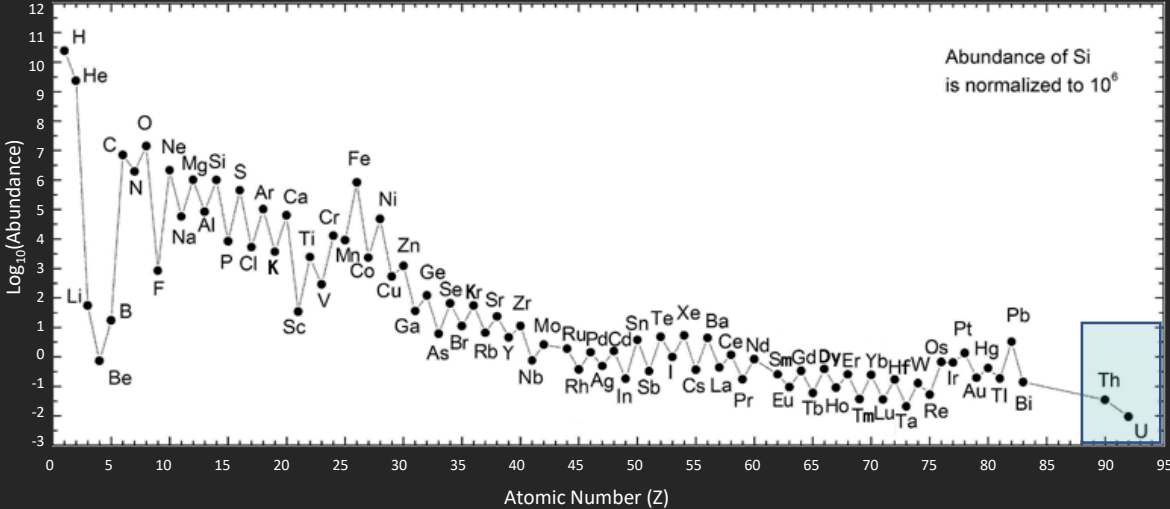
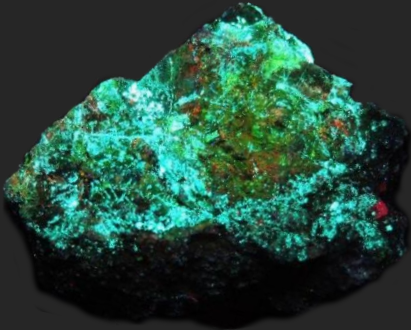
Introduction: Rapid Neutron Capture



$$Y_e = \frac{n_p}{n_p + n_n}$$

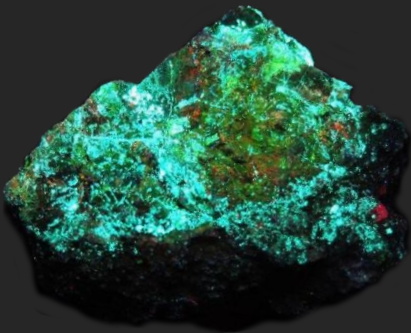
Introduction: Motivation

Natural Abundance Production in Solar System

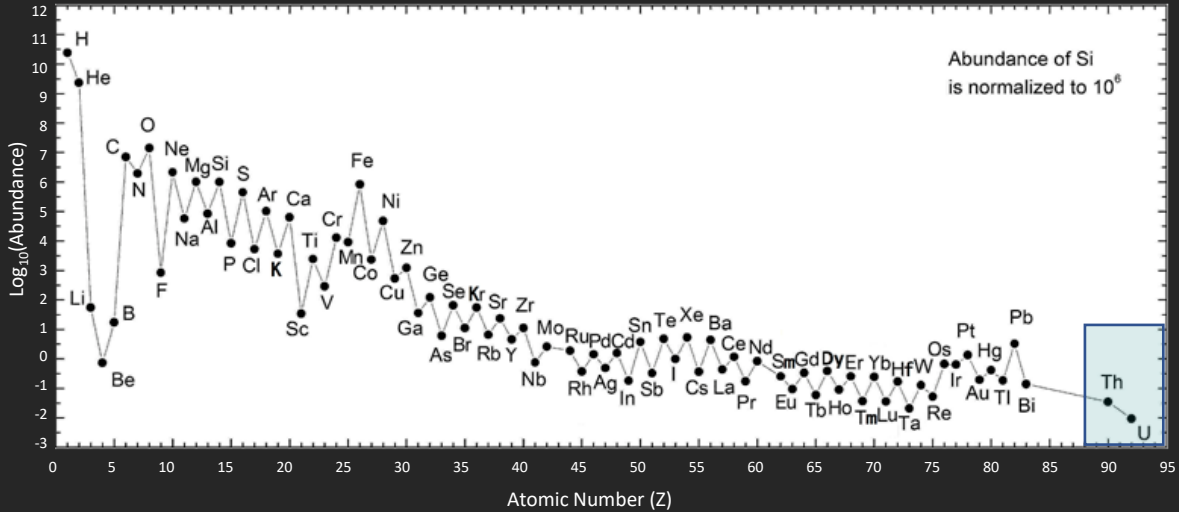


Introduction: Motivation

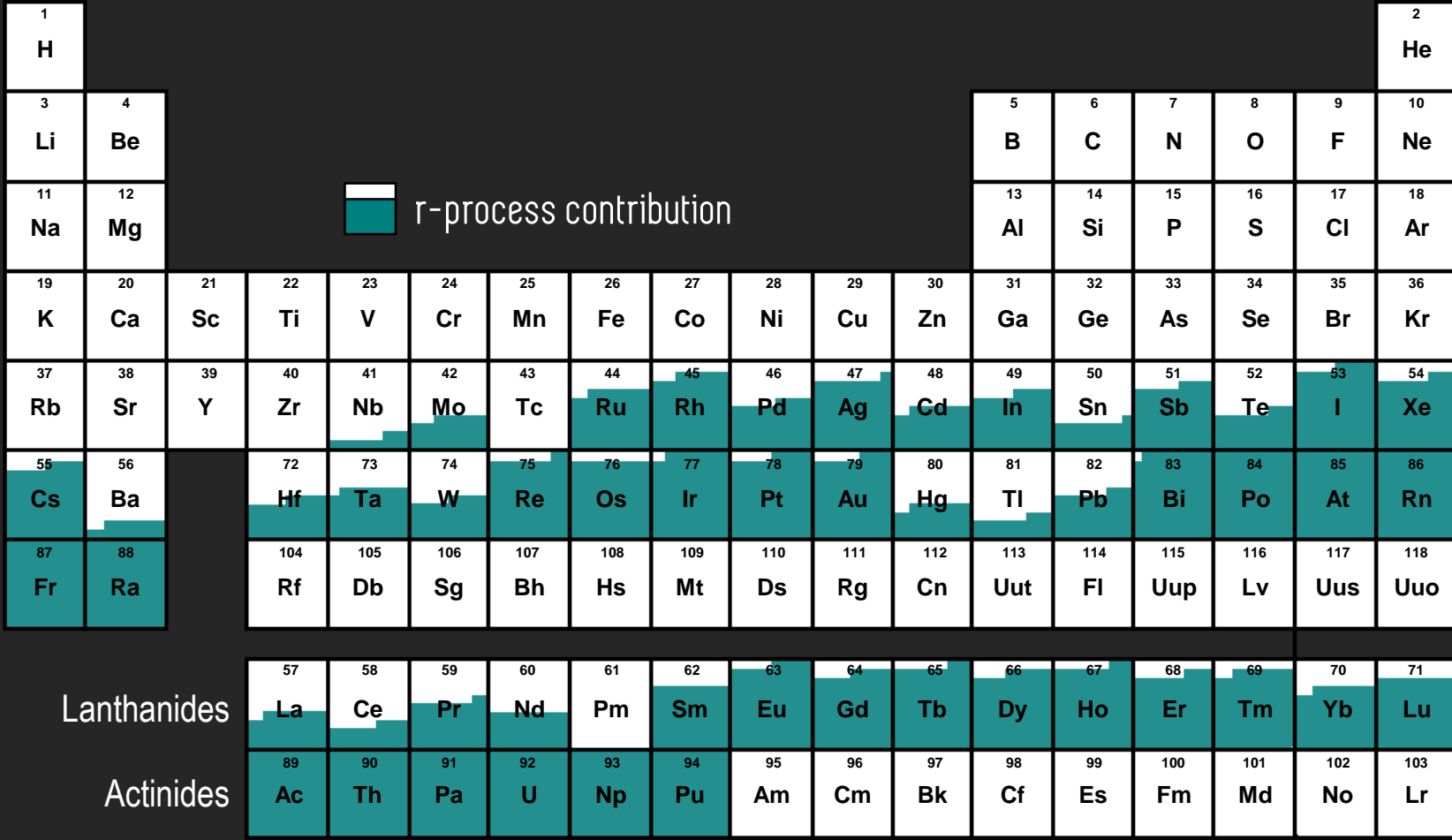
Natural Abundance Production in Solar System



Cosmic Chronometers

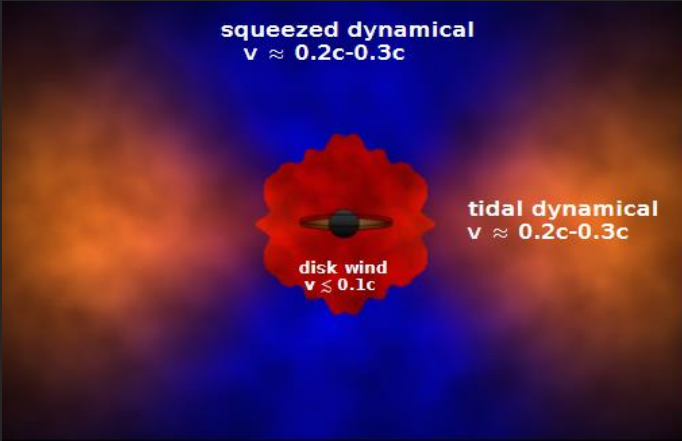


Introduction: Motivation



r-Process Site: Compact Object Mergers

NS+NS



Prompt Collapse

NS+NS



Central Engine

NS+BH



Prompt Collapse

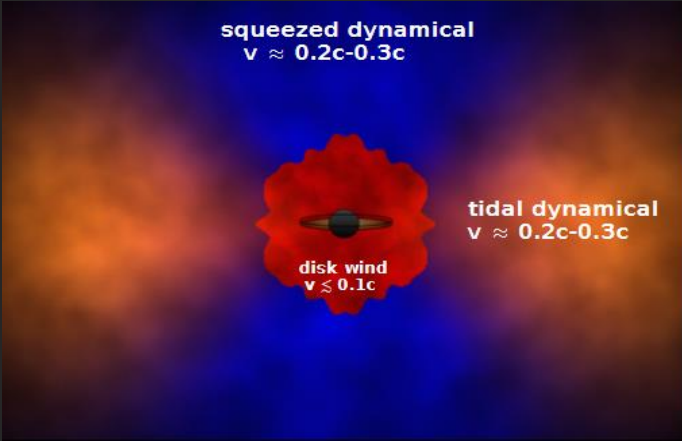
■ high opacity

■ low opacity

(Adapted from Kasen+ 2017)

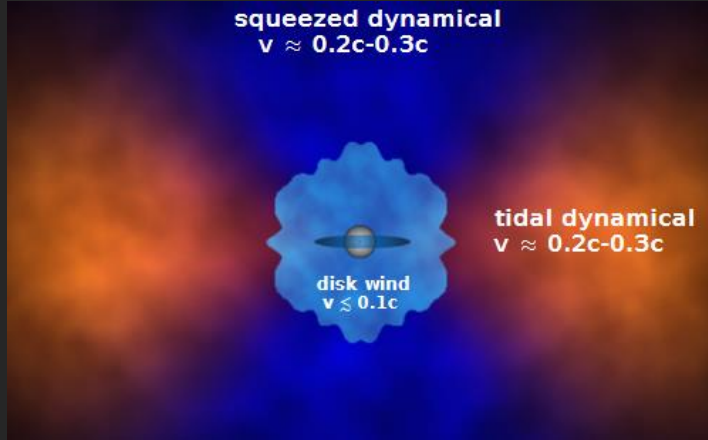
r-Process Site: Compact Object Mergers

NS+NS



Prompt Collapse

NS+NS



Central Engine

NS+BH



Prompt Collapse

Charged current interactions change composition of material:

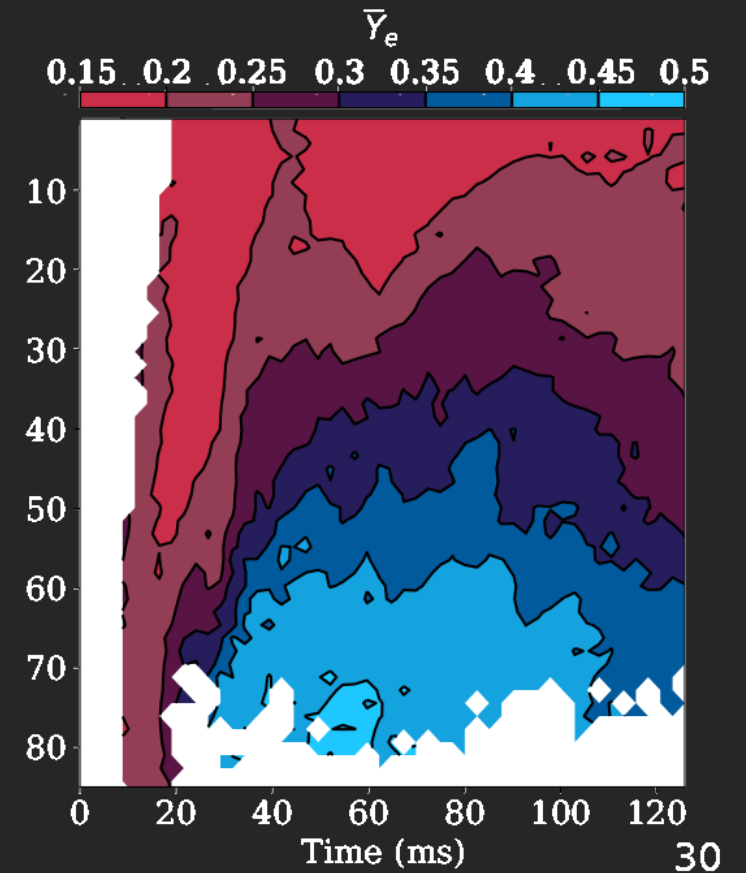
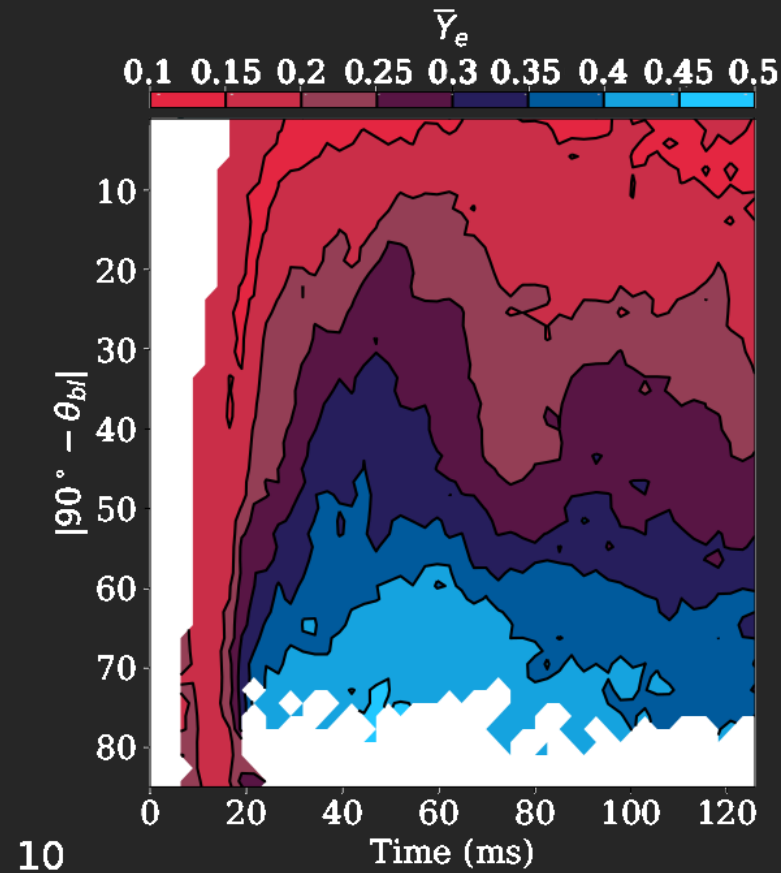


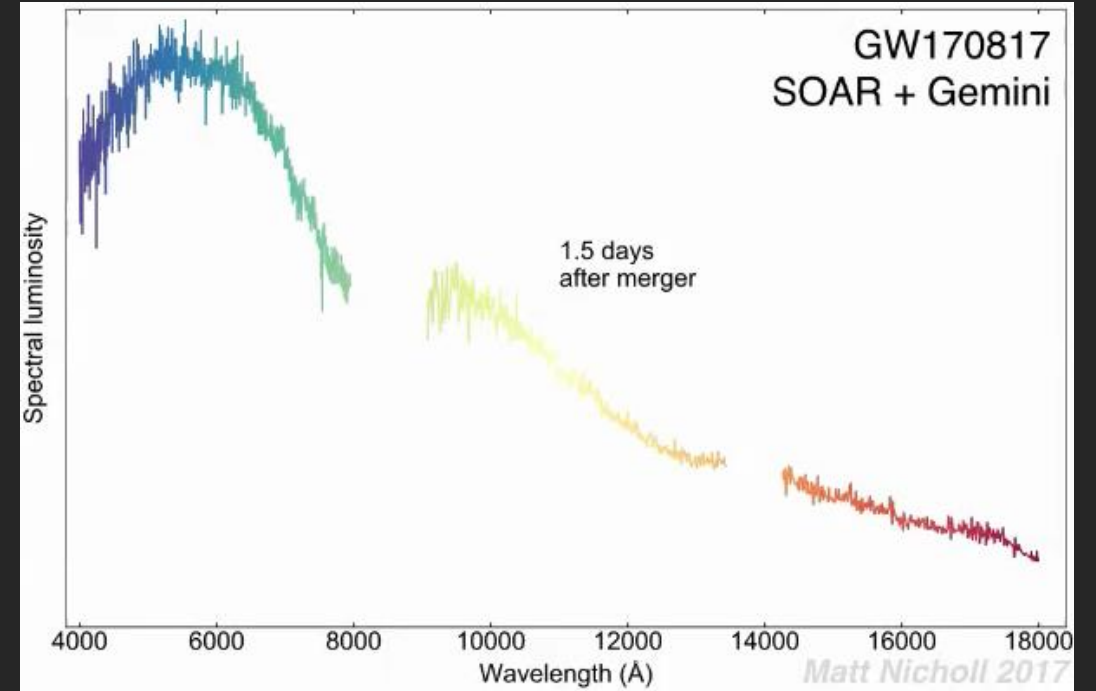
Important quantity for r-process:
$$Y_e = \frac{n_p}{n_p + n_n}$$

(Adapted from Kasen+ 2017)

r-Process Site: Evolution of Post-Merger Disk

Neutrino effects on Y_e affected by geometry and evolution of magnetic fields in the post-merger disk.
(wip)







Nuclear Physics Sources of Uncertainties

Nucleosynthesis calculations probe:

- Nuclear Energy Generation
- Light Curve Evolution
- Final Abundance Patterns
- Cosmochronometry

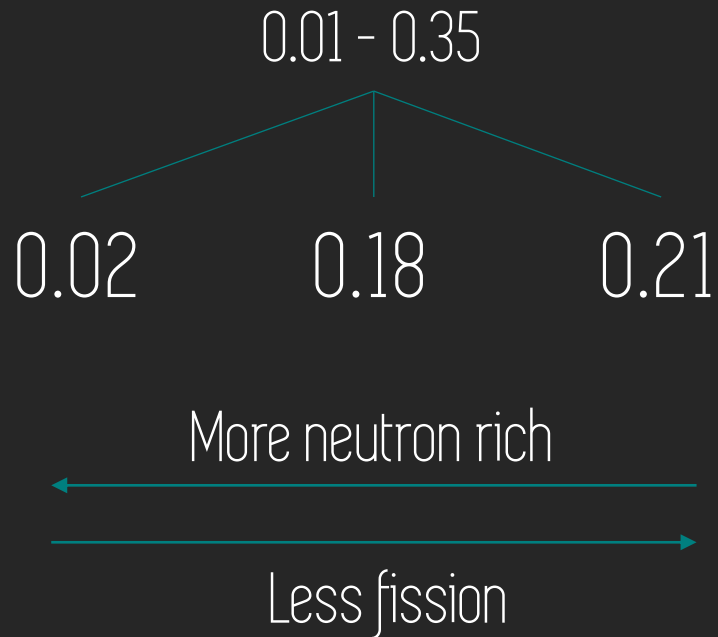
arXiv: 2208.06373

2010.03668

2010.11182

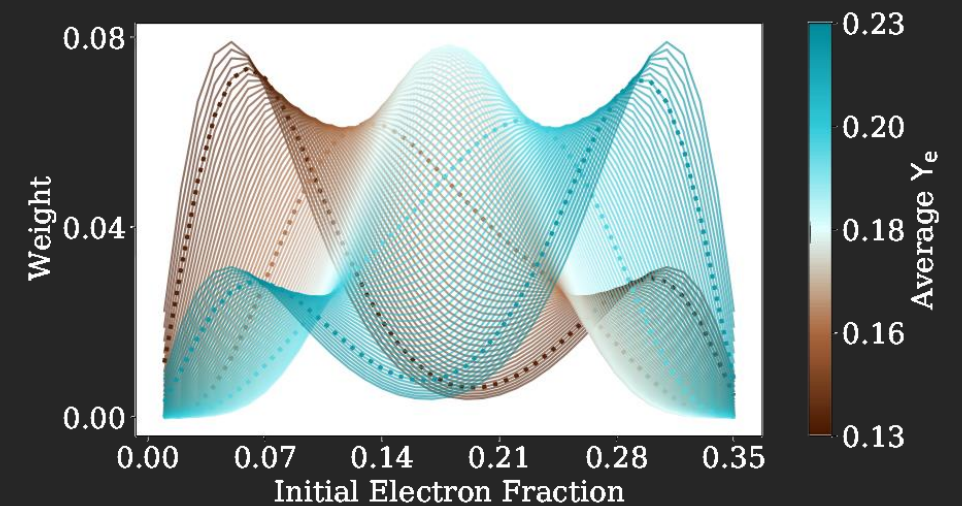
Building a Dataset: Initial Y_e

Single Trajectory: (all nuclear datasets)

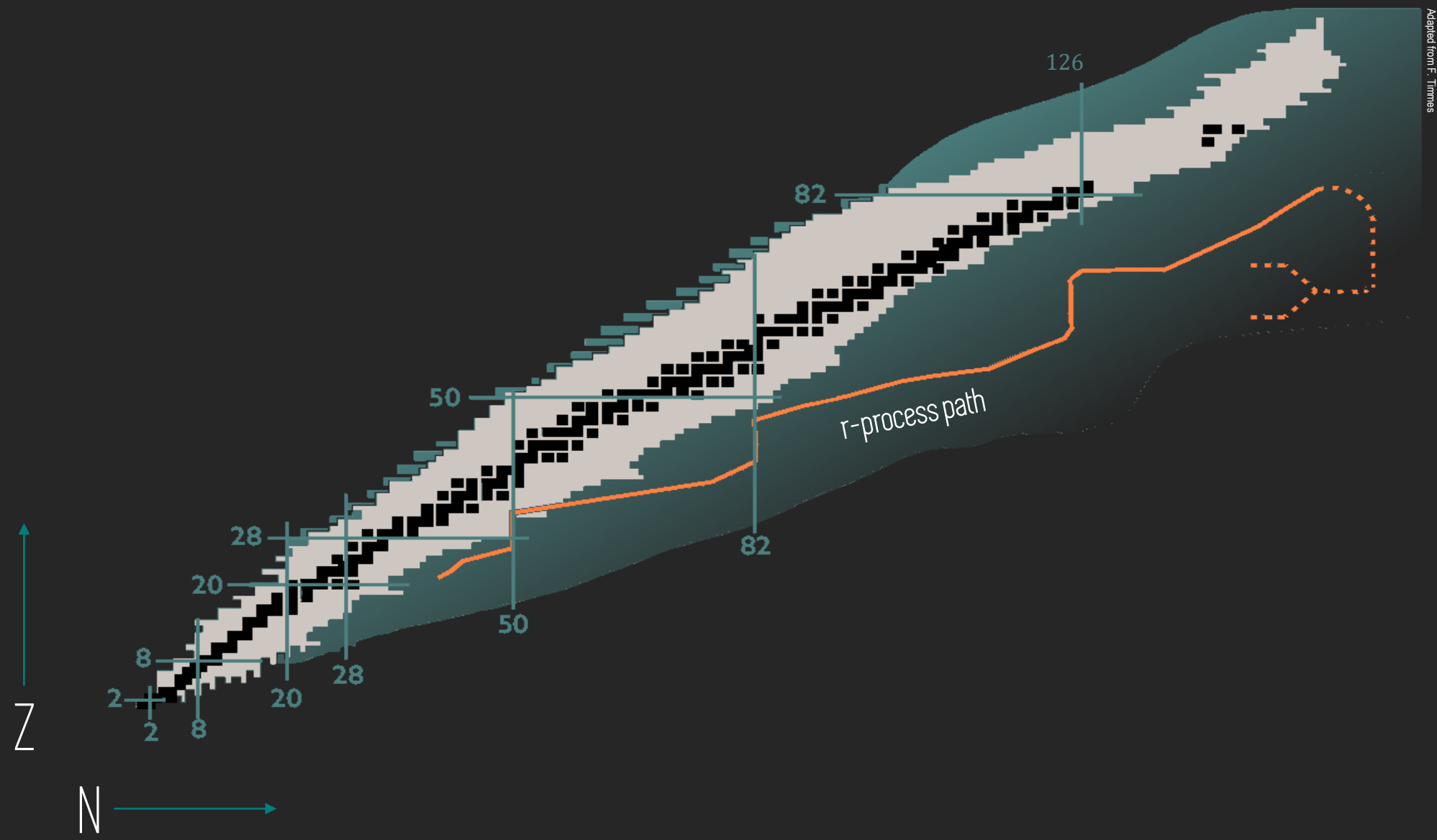


Zhu, Lund+(2021); Barnes, Zhu, Lund +(2021); Lund+ (2022)

Linear Combinations: (limited nuclear dataset)



Sources of Nuclear Uncertainty



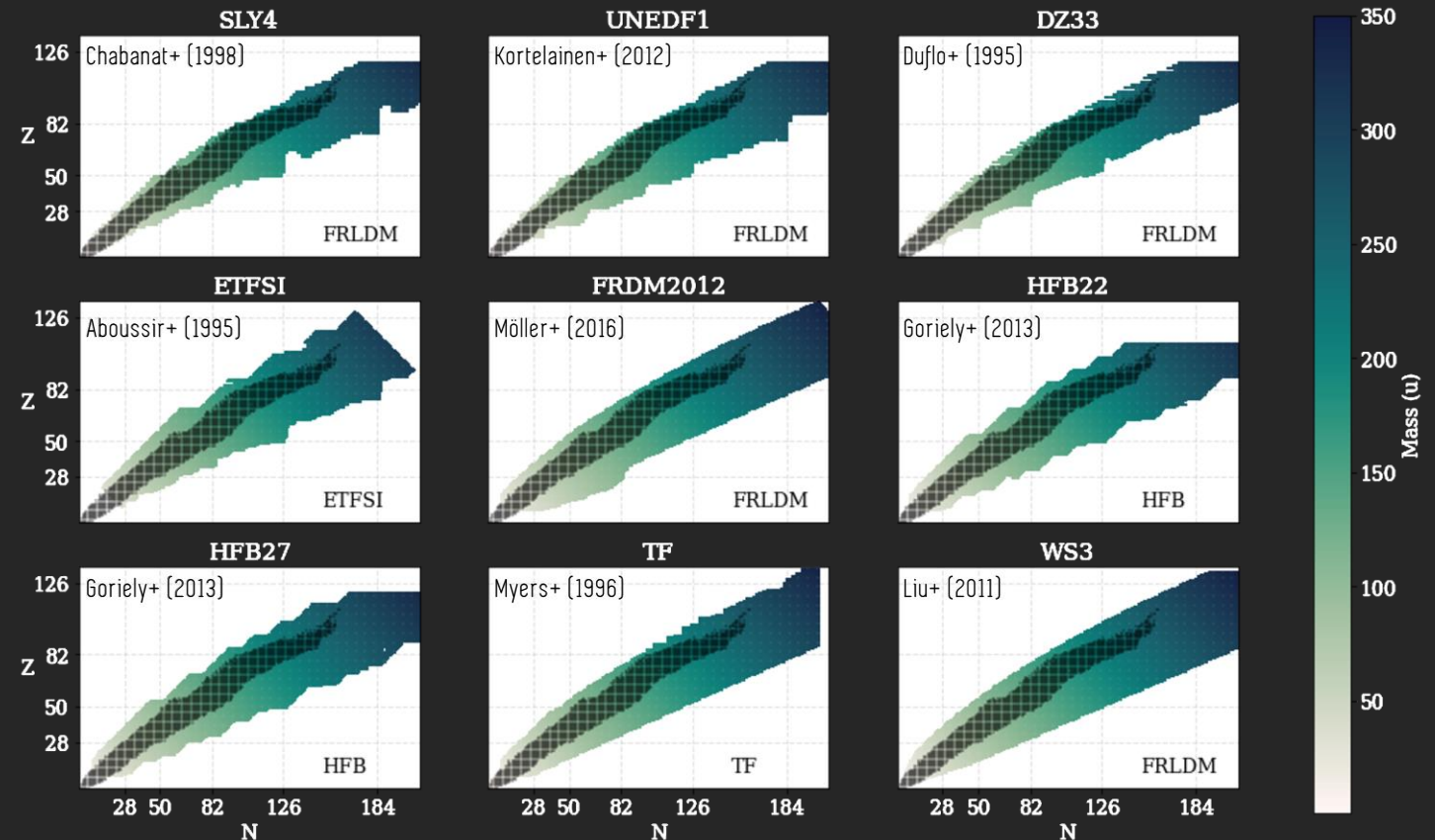
Adapted from F. Timmes

Building a Dataset: Mass Model

Most basic nuclear property: mass

Common approach: fit parameters to experimental data, extrapolate to make predictions about unknown nuclei

Each mass model associated with fission barrier height model



*Experimental data from AME2016 (Wang+2017, Audi+2017)

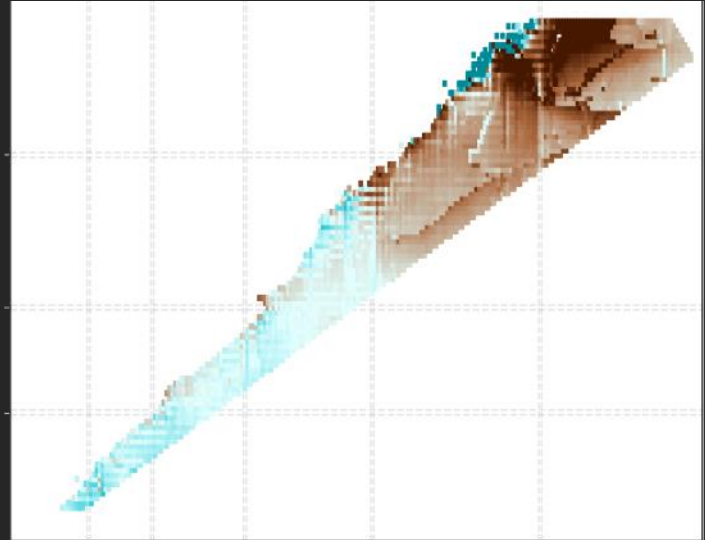
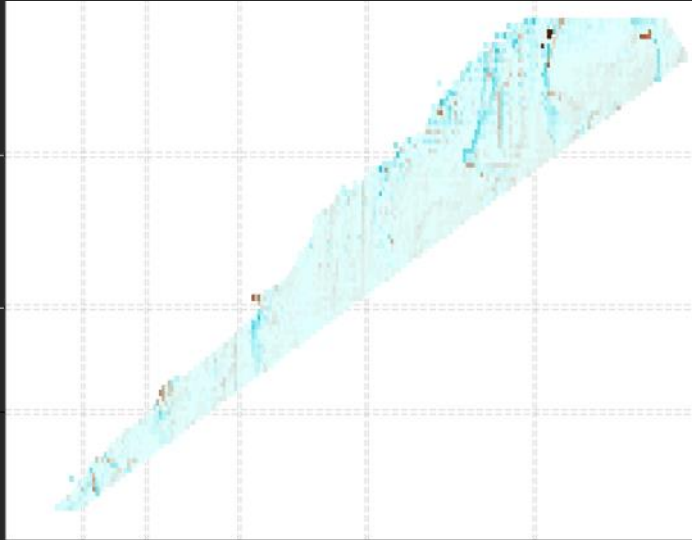
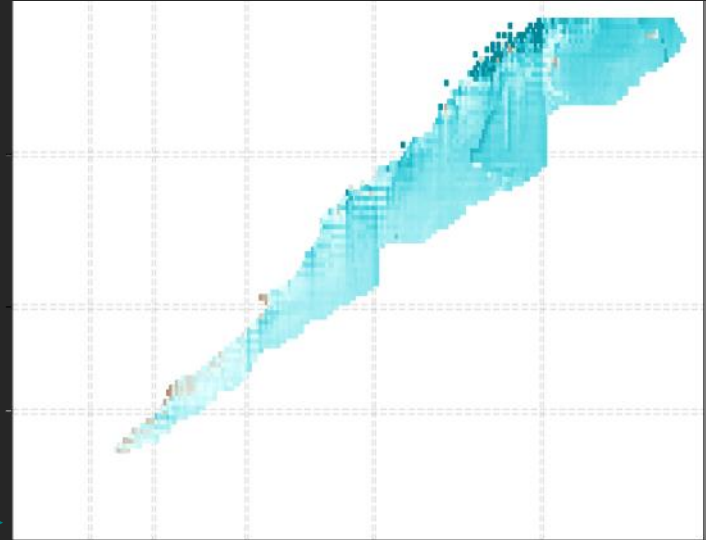
Building a Dataset: Beta Decay Rates

← ~ slower rate* →

NES/MLR03

MLR/MLR03

MKT/MLR03



FAM+QRPA

QRPA+FRDM

Covariant DFT

Ney (NES)

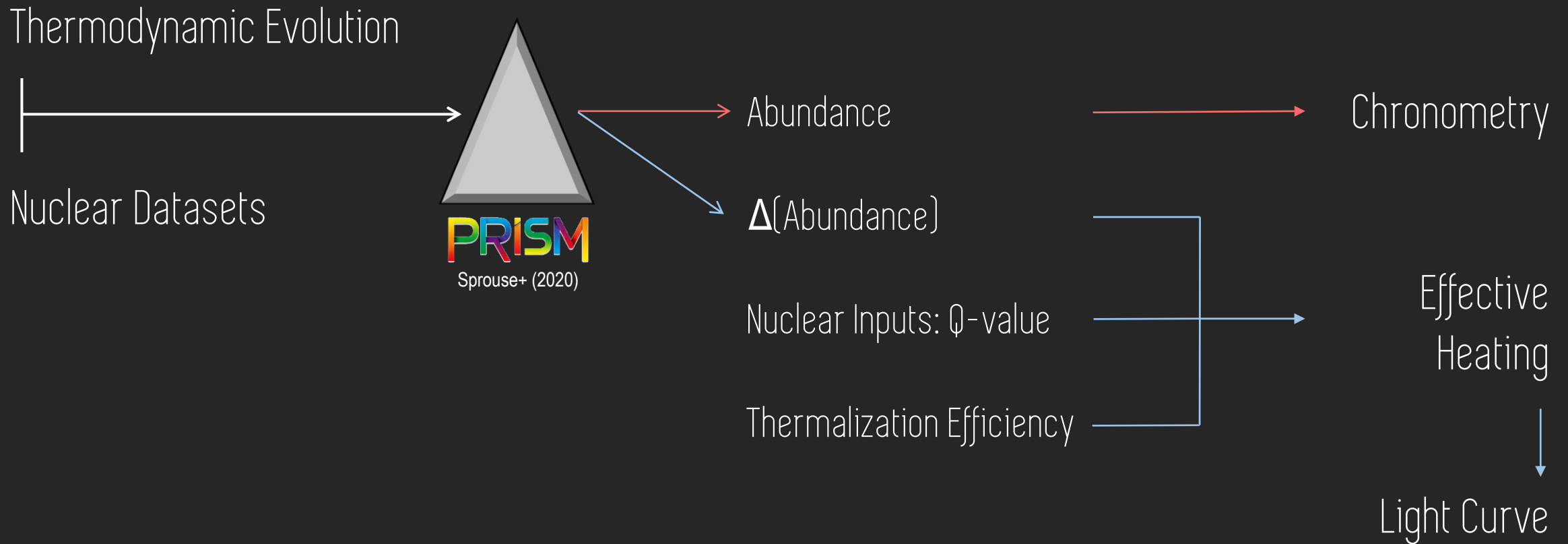
Ney+ 2020

Möller (MLR)

Marketin (MKT)

Marketin+ 2016

PRISM: A Sparse Matrix Solver

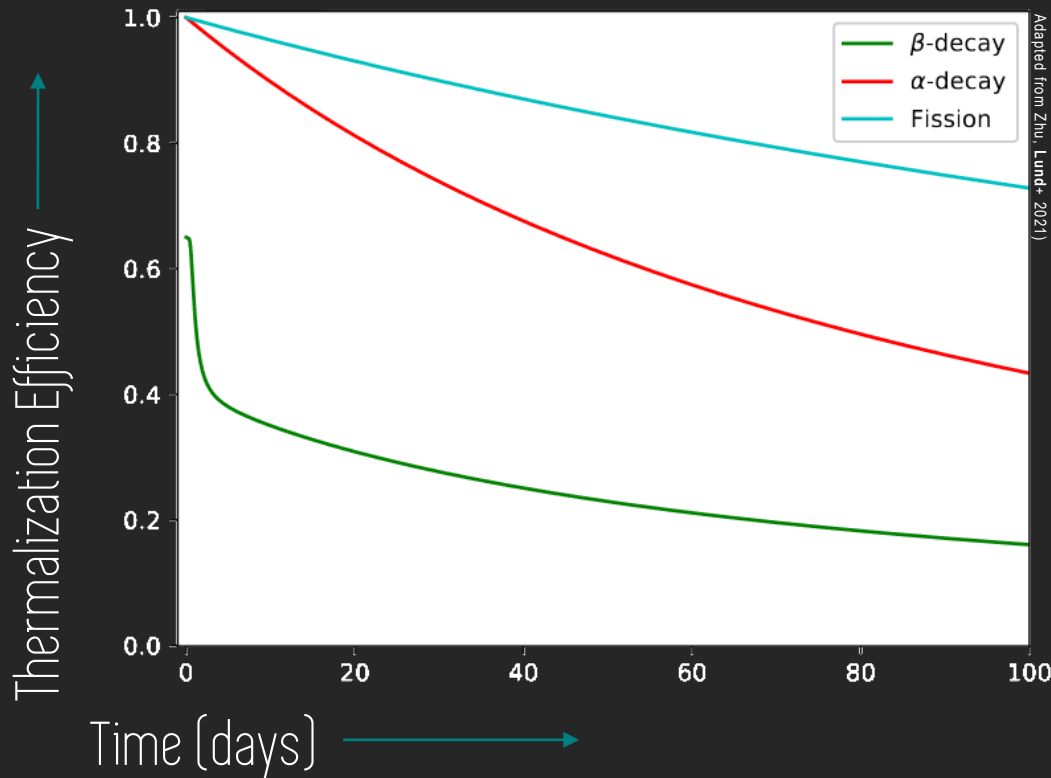


Nuclear Heating



Nuclear Heating Calculation

PRISM flow: specific contribution of total time-derivative of an abundance due to a single nuclear process [s⁻¹]



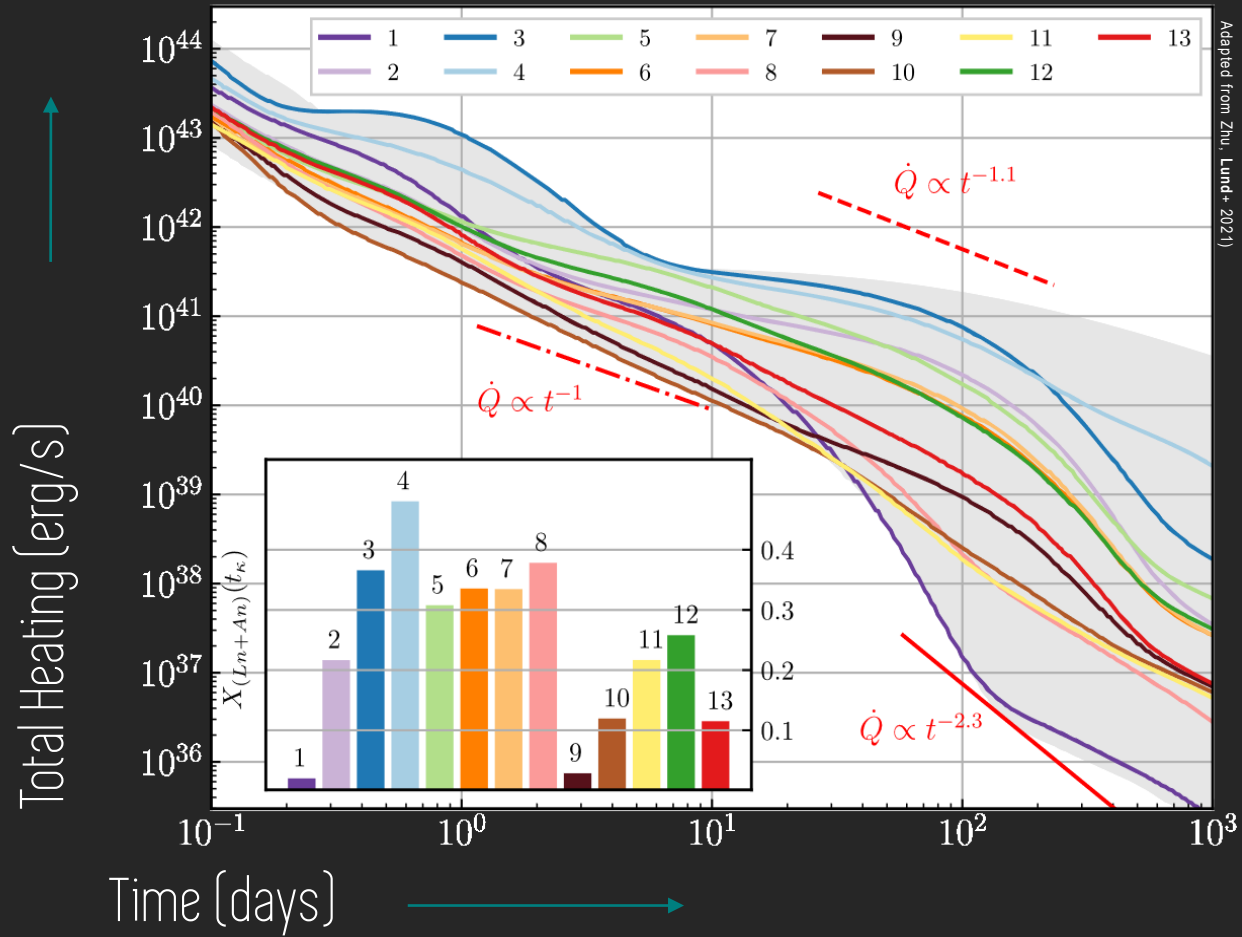
Thermalization efficiency: how effectively decay products can heat ejecta (function of time, ejecta mass, and characteristic velocity)

$$\dot{Q}(t) = \sum_i f_i(M_{ej}, v_{ej}, t) \dot{q}_i(t) M_{ej}$$

Labels for the equation components:

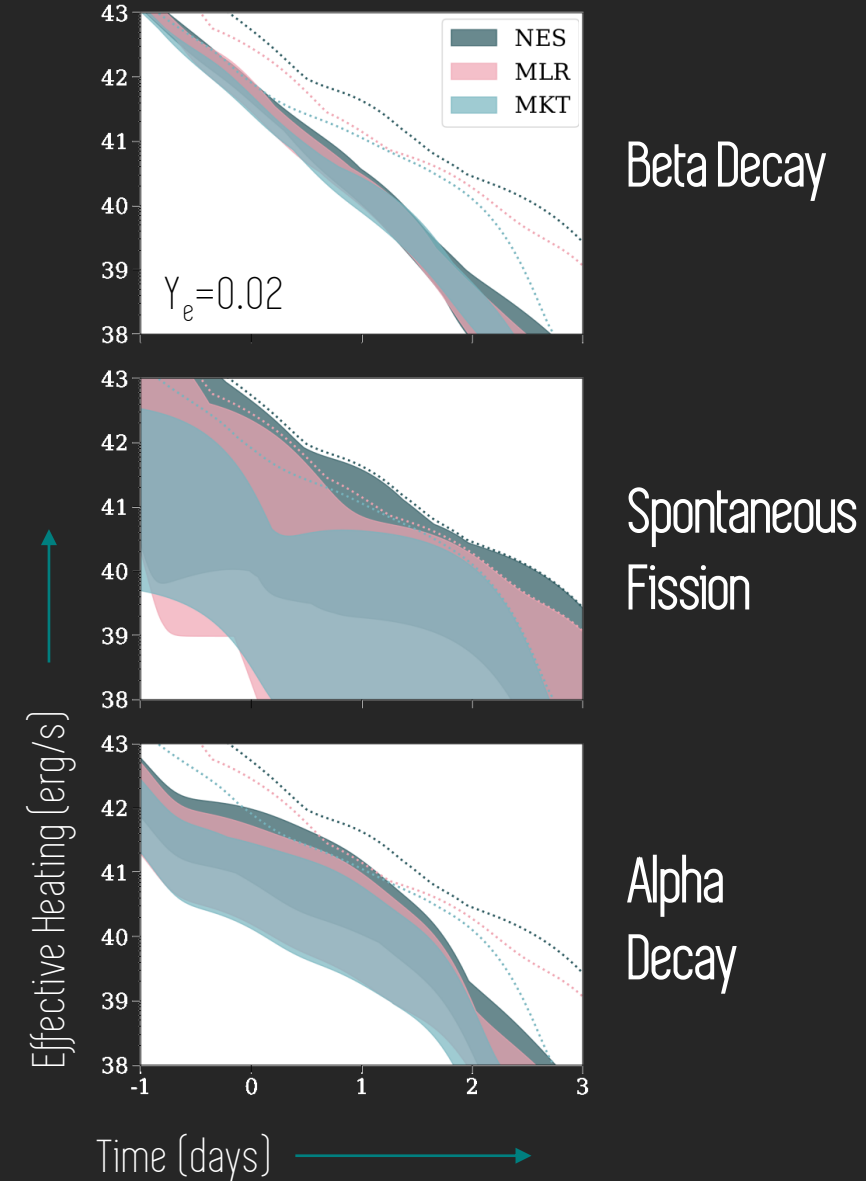
- $\dot{Q}(t)$: Total effective heating
- $f_i(M_{ej}, v_{ej}, t)$: Thermalization efficiency
- $\dot{q}_i(t)$: Heating (Q-value * flow)
- M_{ej} : Ejecta mass

Total Heating for Different Nuclear Physics



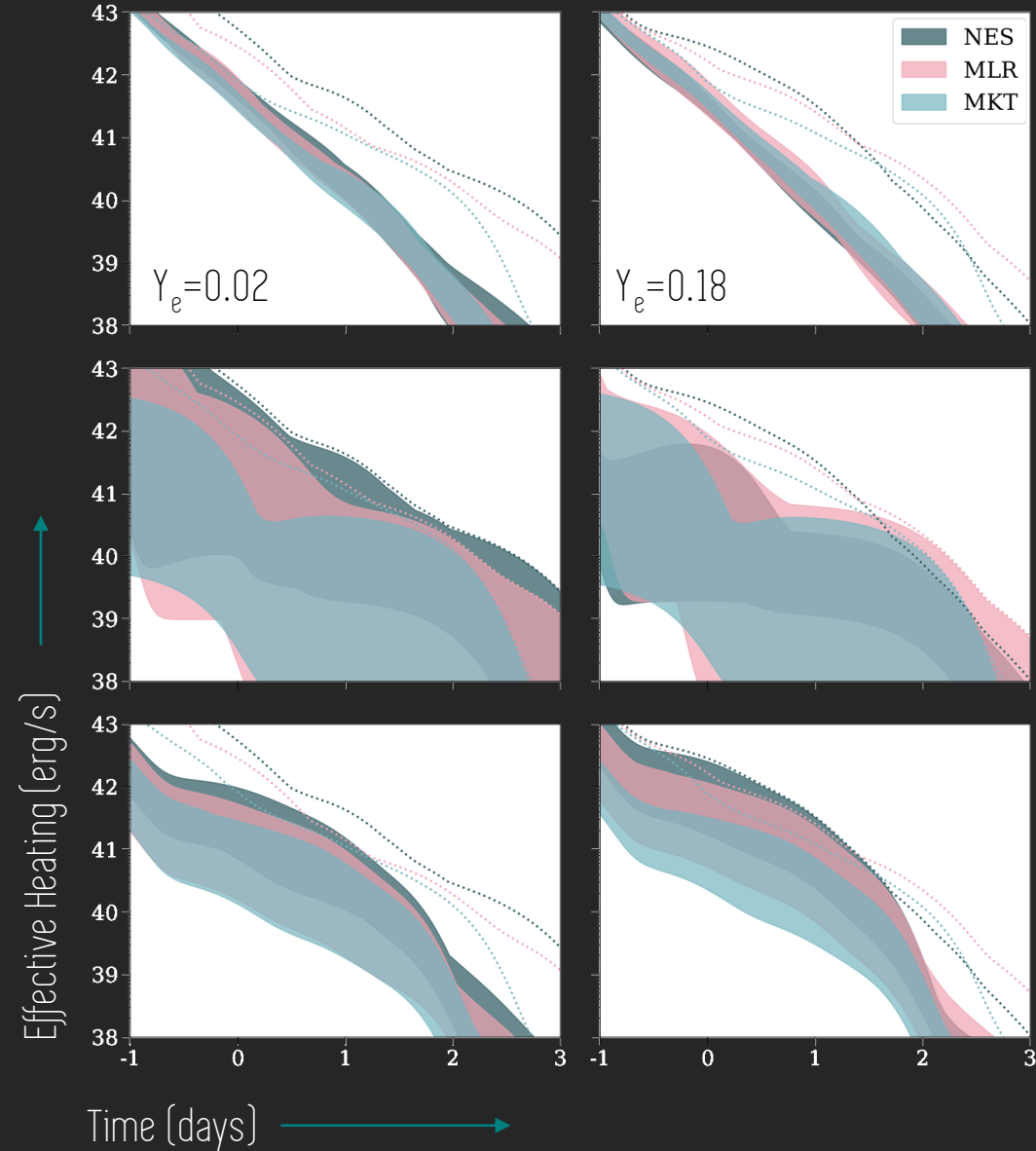
Index	Mass Model	Y_e
1	FRDM12	0.28
2	FRDM12	0.16
3	HFB22	0.16
4	HFB27	0.16
5	DZ33	0.16
6	UNEDF1	0.16
7	UNEDF1	0.16
8	UNEDF1	0.24
9	SLY4	0.18
10	SLY4	0.21
11	TF*	0.16
12	DZ33	Solar mix
13	UNEDF1	Solar mix

Nuclear Heating by Reaction Type



- Upper limit of heating uncertainty set by fission of few mass models
- Beta models differ in behavior of dominating fission heating

Nuclear Heating by Reaction Type



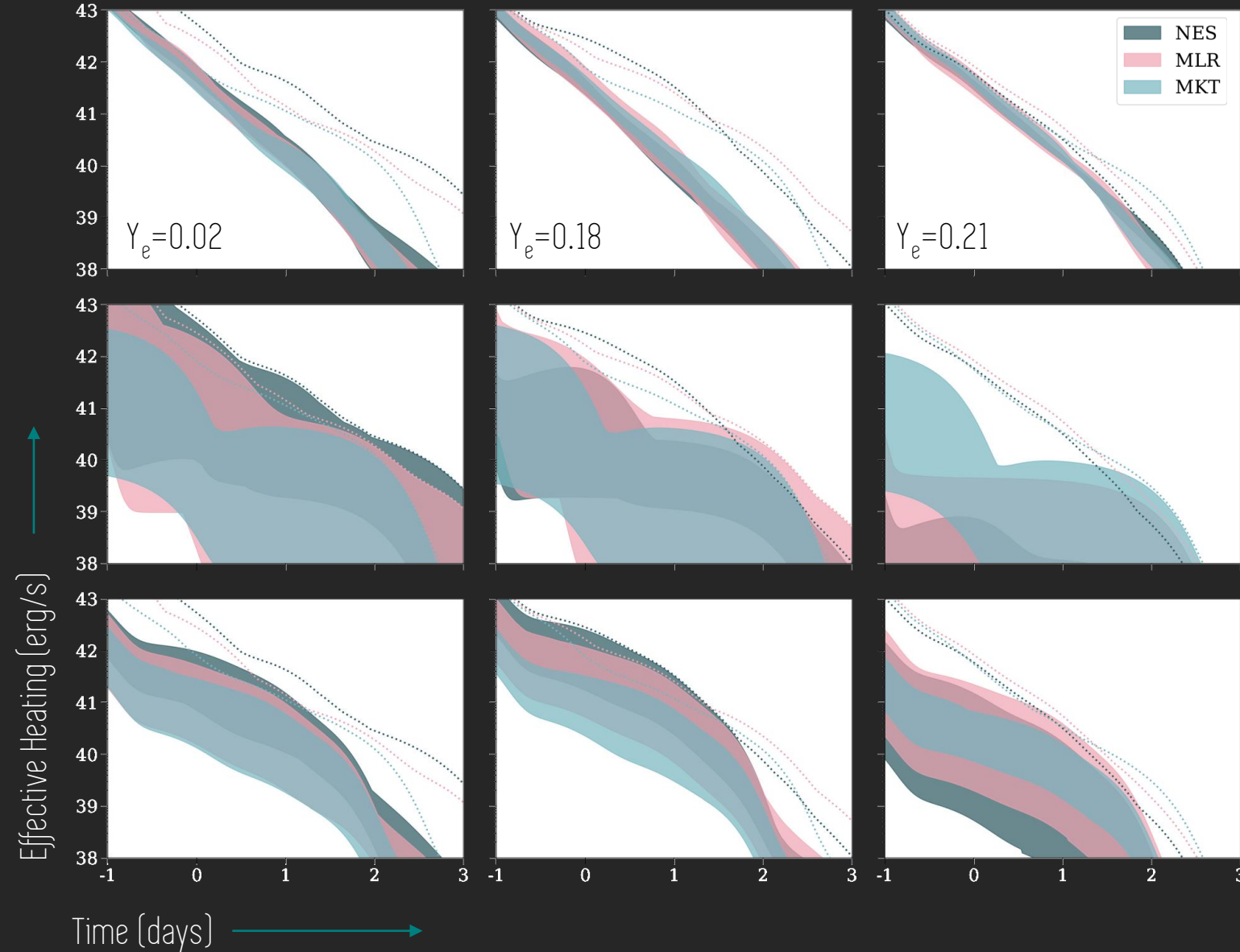
Beta Decay

Spontaneous Fission

Alpha Decay

- Alpha heating becomes more important <100 days
- Beta models differ in predicting when alpha tends to dominate + late-time tail shape of fission heating

Nuclear Heating by Reaction Type



Beta Decay

- Much more overlap, total heating tends to be set by beta (and some alpha) decay

Spontaneous Fission

- Overall effect on beta decay heating is small

Alpha Decay

Light Curve



Light Curve Shell Model

Similar procedure as effective heating calculation, but computationally more intensive (ref Metzger 2017)

Shell model for ejecta: the mass of each shell, M_v , depends on the velocity, v , of that shell (100 shells evenly distributed between 0.1c and 0.4c)

Time evolution of the energy of a shell:

Luminosity (ultimately want to plot this!)

$$\frac{dE_v}{dt} = \frac{M_v}{M_{ej}} \dot{Q}(t, v) - \frac{E_v}{t} - \frac{E_v}{t_{d,v} + t_{lc,v}}$$

Effective heating

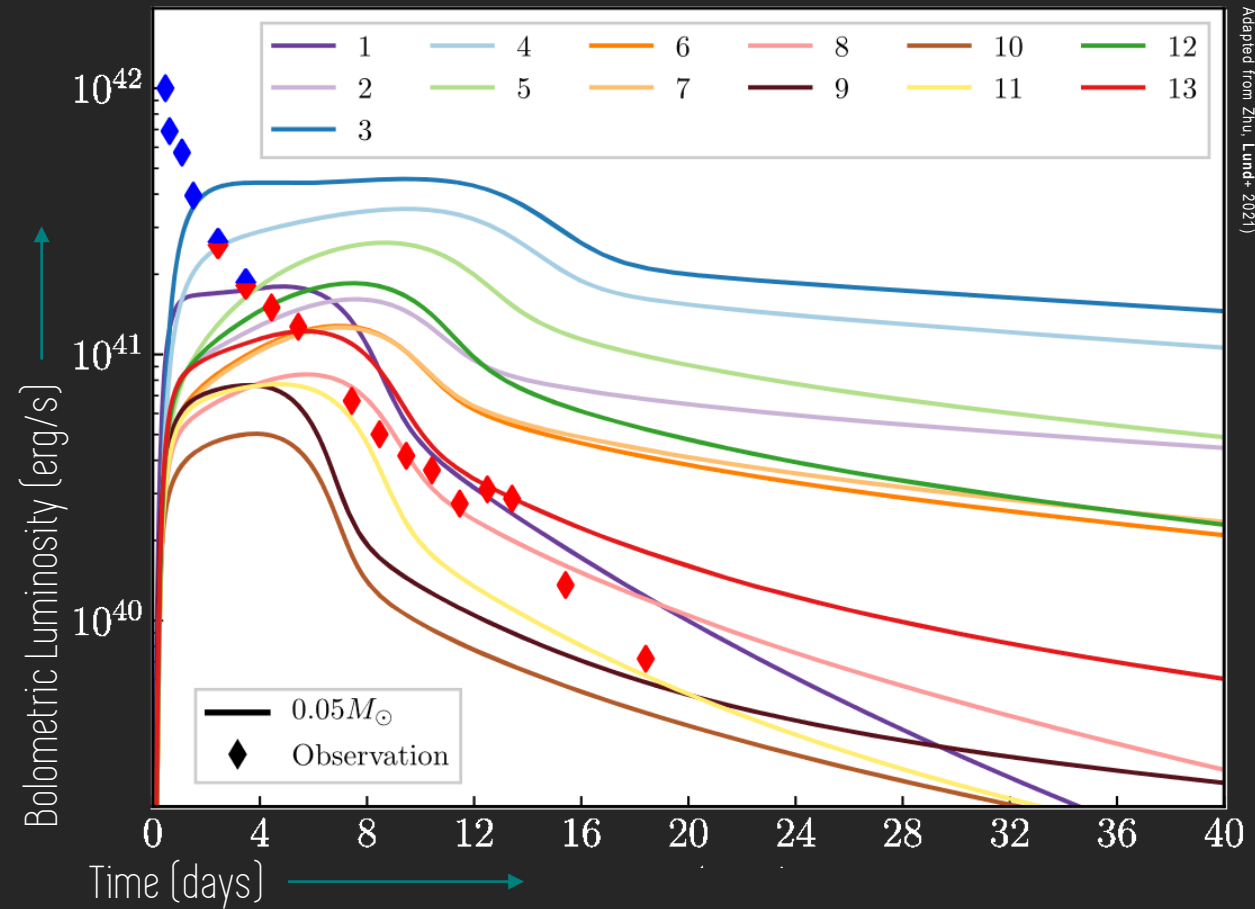
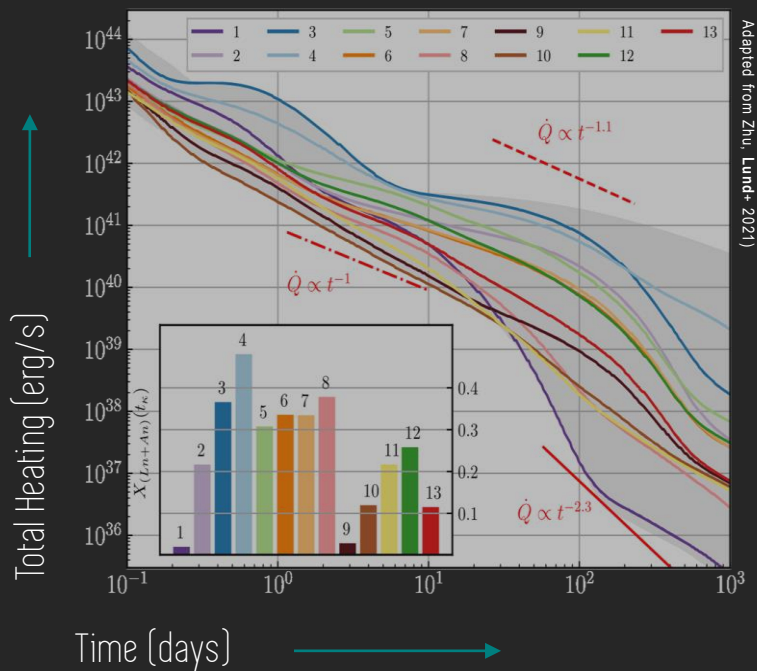
Adiabatic expansion

Diffusion timescale
(depends on opacity)

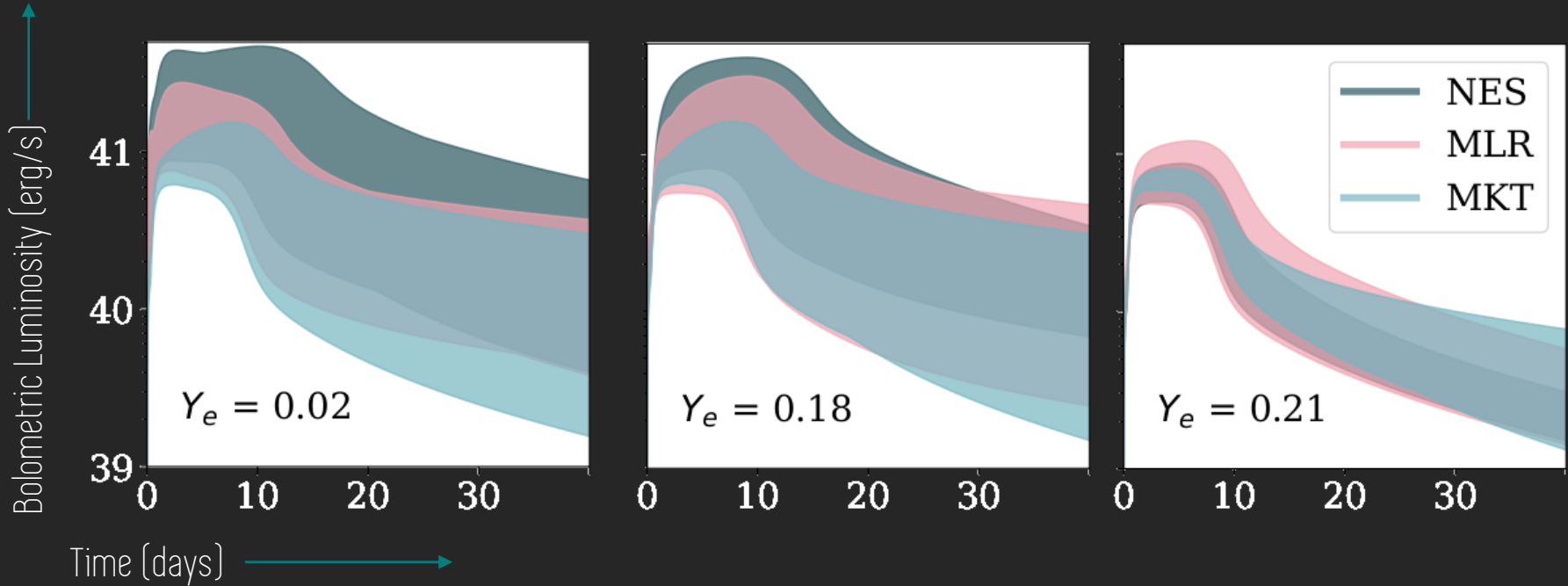
Light-crossing
timescale

Light Curves

- Differences in nuclear heating predictions propagate through to light curve
- Overall magnitude of luminosity ~ overall heating BUT peak depends on ejecta opacity
- For a given effective heating rate, higher χ_{Ln+An} elongates curve



Light Curves



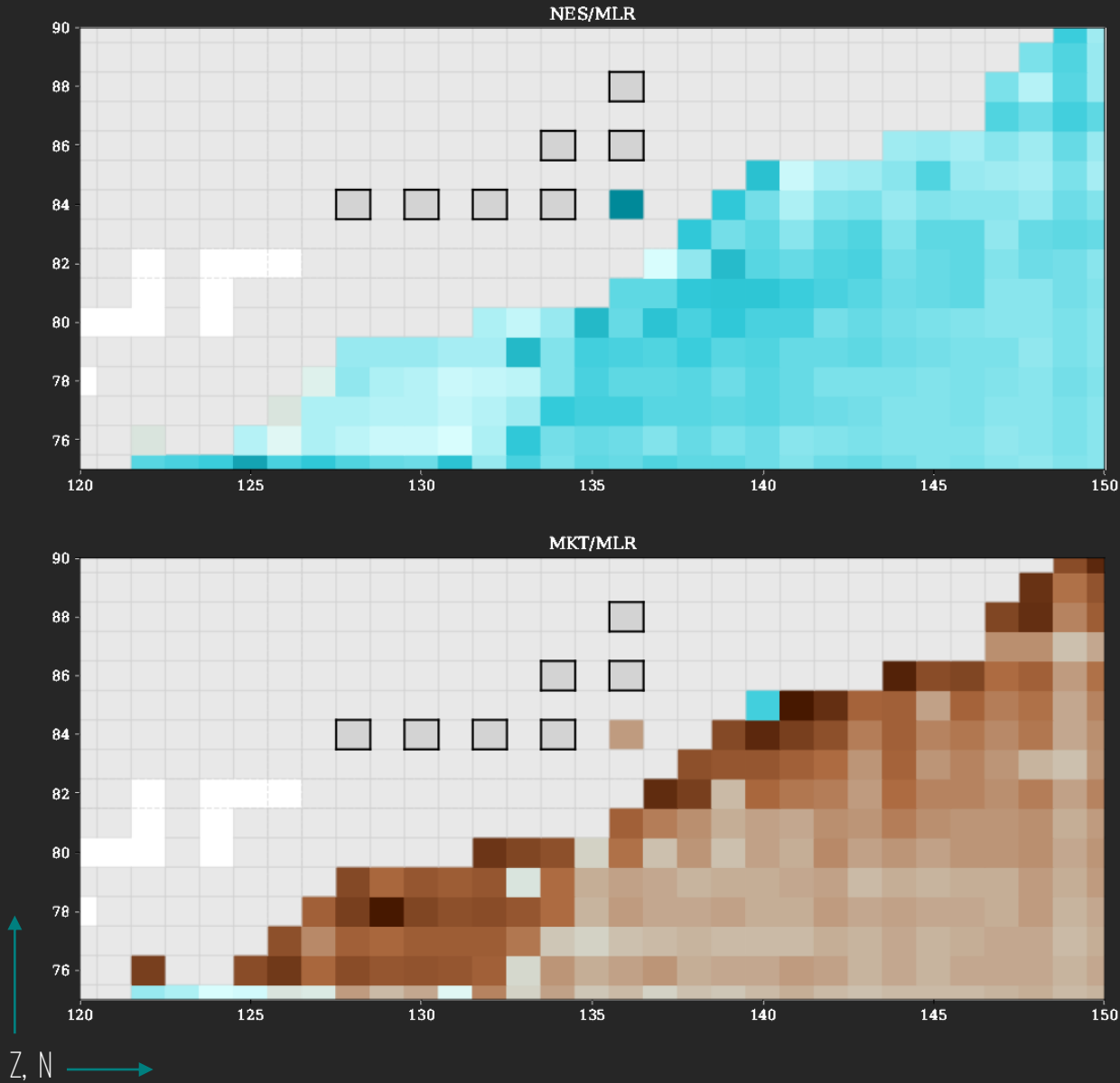
~70% more heating
can yield ~50%
brighter light curve
(NES:MLR)

~40% less heating can
yield ~50% dimmer
light curve
(MKT:MLR)

Nuclear Heating (revisited)



Alpha Decay Heating



Differences in beta decay rates affect heating from alpha heaters with measured decay times, especially:



Spontaneous Fission (et al.) Heating

Theoretical branching ratios affect spontaneous fission heating

