

# Machine Learning Inference of the Dense Matter EoS and Quark-Hadron Continuity

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in collaboration with: K. Fukushima, K. Murase & W. Weise

YF, K. Fukushima, K. Murase, Phys. Rev. D **98**, 023019 (2018)

YF, K. Fukushima, K. Murase, arXiv:1903.03400 [nucl-th]

YF, K. Fukushima, W. Weise, arXiv:1908.09360 [hep-ph].

22 Sep. 2019, Erice 2019

# Phases of Dense Matter

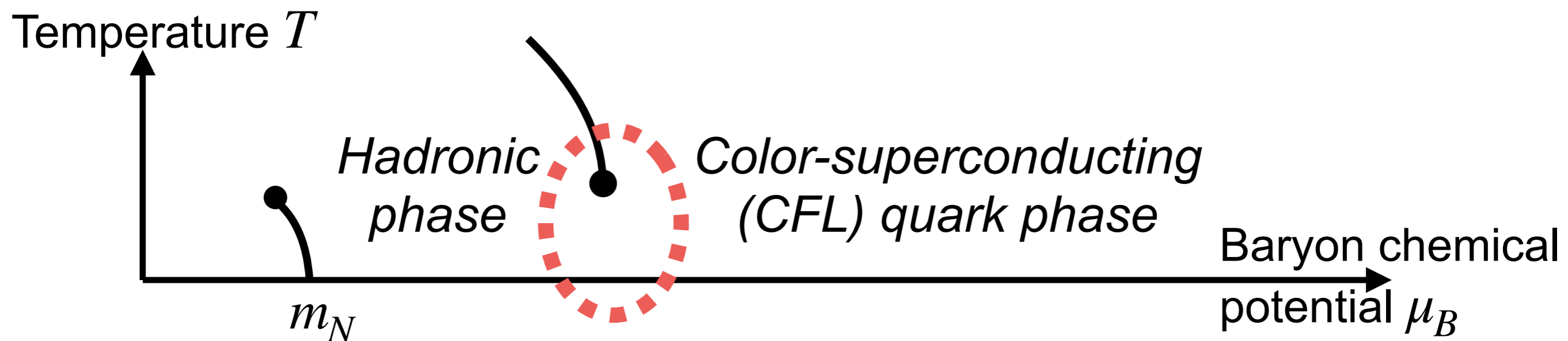
- Degrees of freedom of dense matter:

## Hadron & Quarks

- One possibility from neutron star phenomenology:

## Quark-hadron continuity

- both phases are indistinguishable (no order parameter)
- connected by **crossover** transition



# Equation of State of Dense Matter

- Dense matter properties are characterized by **Equation of State (EoS):**

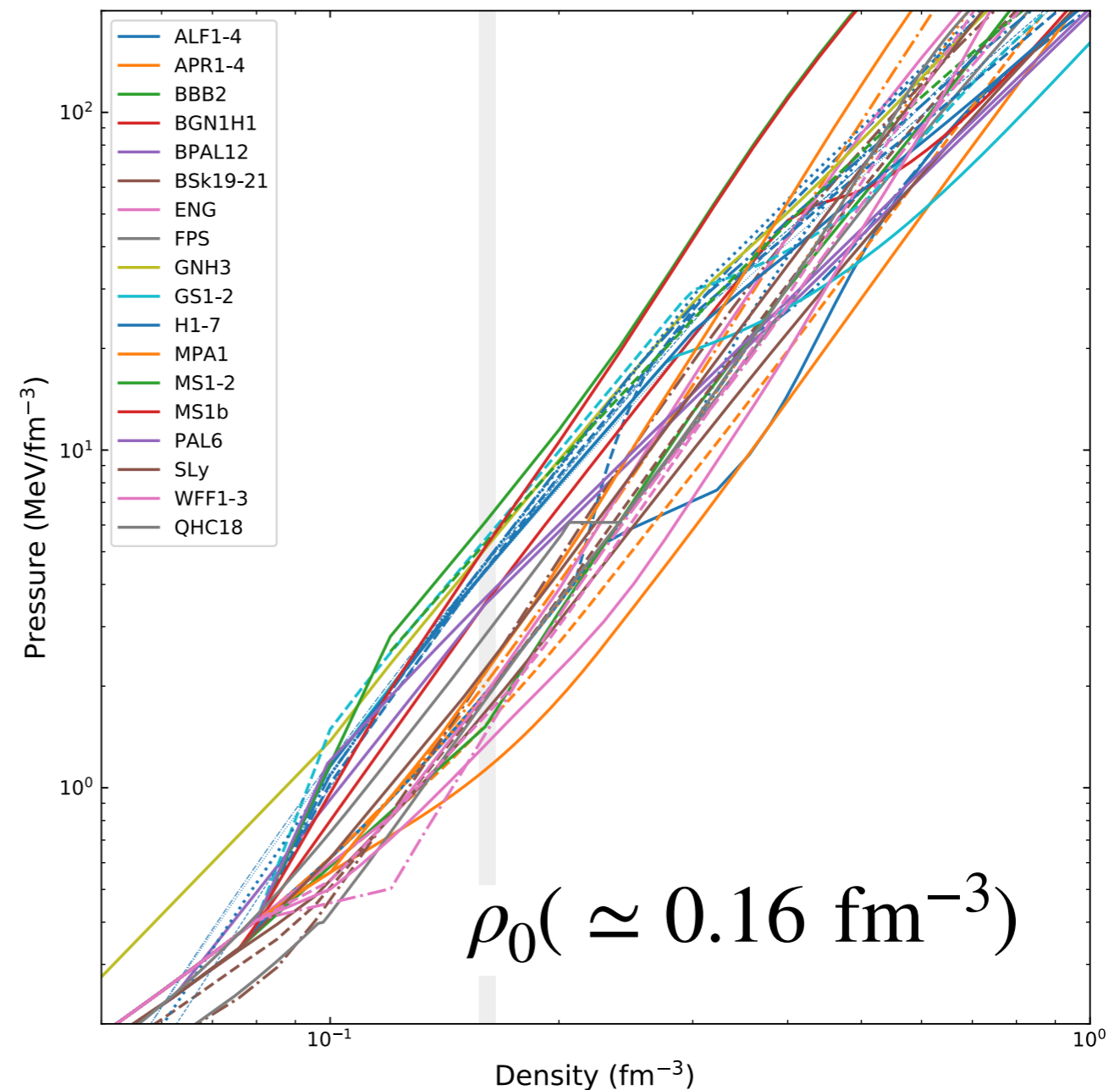
relation between pressure  $p$  and (mass) density  $\rho$

$$p = p(\rho)$$

- Essential ingredients for neutron star theory; characterizes neutron star structure
- Should be derived from QCD in principle, but many difficulties such as...
  - sign problem; no lattice data
  - renormalization scale dependence in pQCD

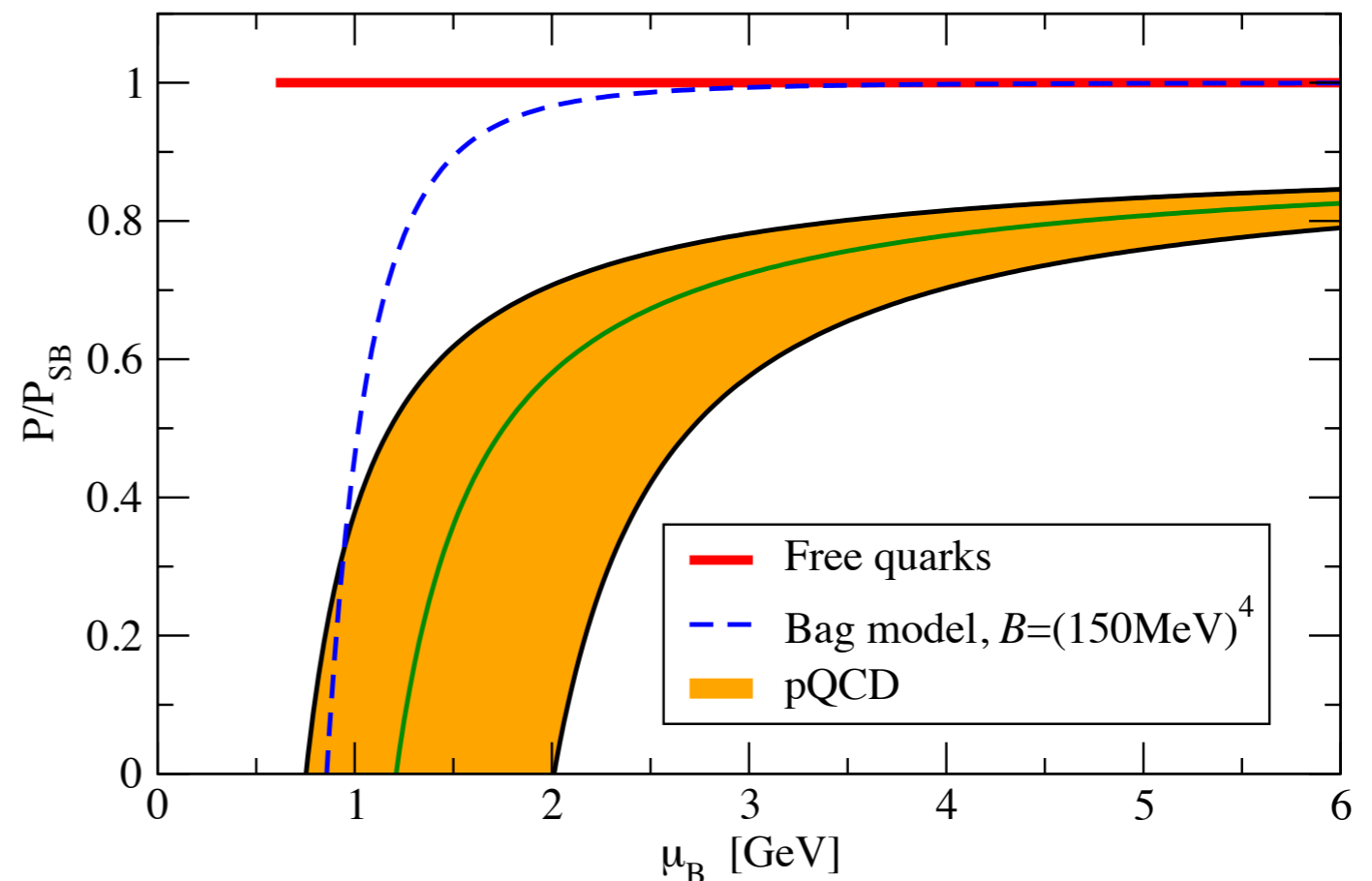
# Current Status of the EoS

- Many nuclear theory calculations  
...but reliability of these models decline with growing  $\rho$



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- Many nuclear theory calculations  
...but reliability of these models decline with growing  $\rho$
- Perturbative QCD calculation also has large uncertainty

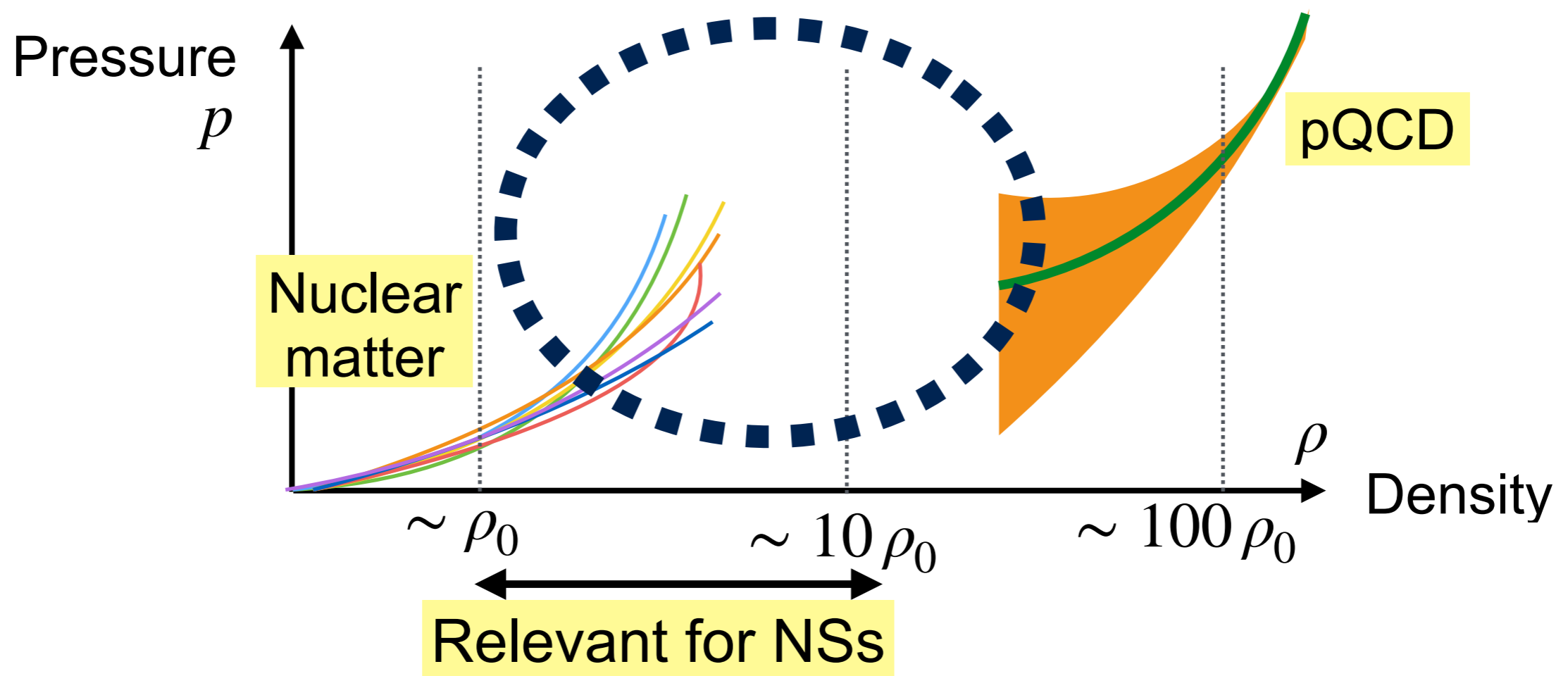


Freedman-McLerran (1977); Baluni (1978)

Fraga-Kurkela-Vuorinen (2014)

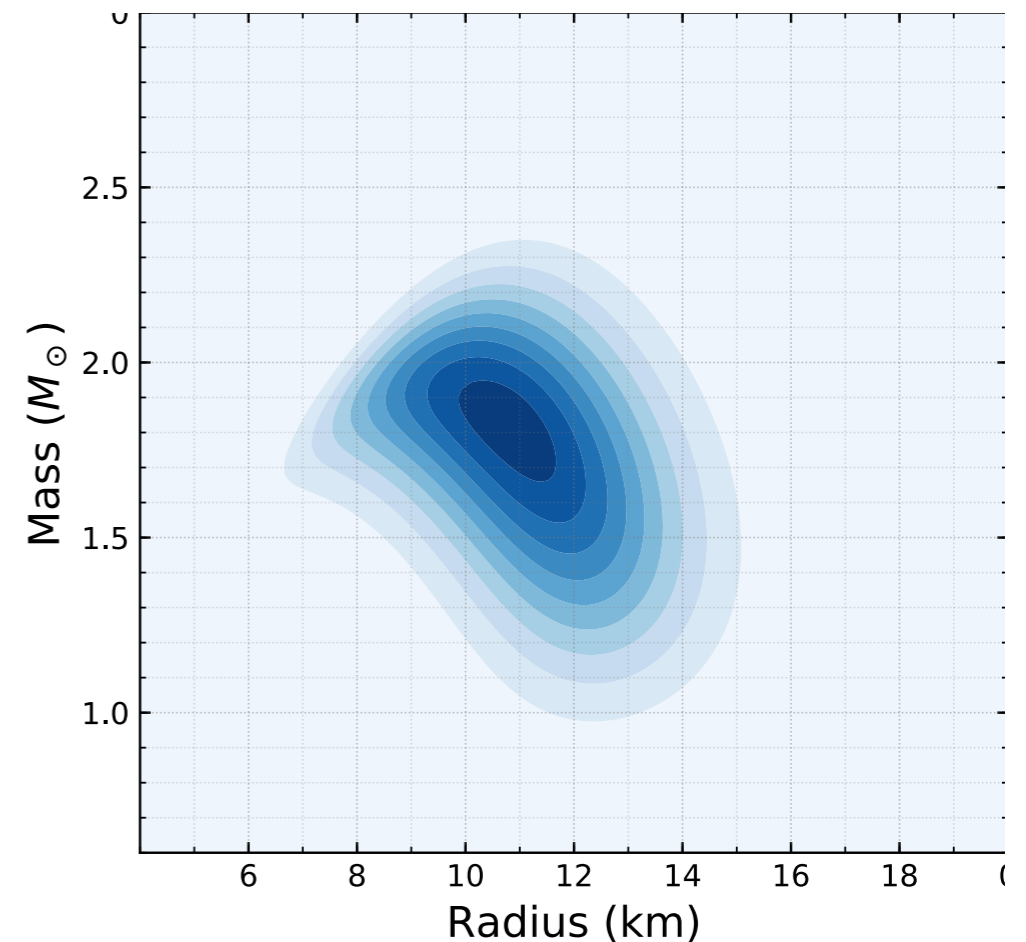
# Current Status of the EoS

- Many nuclear theory calculations  
...but reliability of these models decline with growing  $\rho$
- Perturbative QCD calculation also has large uncertainty
- Need systematic way to interpolate these two regions



# Systematic Way of EoS Inference

- Growing number of observables of NSs:  
e.g.) LIGO-Virgo, NICER experiment
- Utilize NS masses ( $M$ ) and radii ( $R$ ) Özel-Freire (2016)  
14 simultaneous measurements are publicly available;  
given in terms of likelihood distribution in  $(R, M)$  plane
- Infer EoS from  $M$ - $R$ :  
Use **machine learning (ML)**  
with **deep neural network (NN)**



# Method of Machine Learning

Input

$x_1^{(0)}$

⋮

$x_{N_0}^{(0)}$

Answer

$y_1$

⋮

$y_{N_L}$

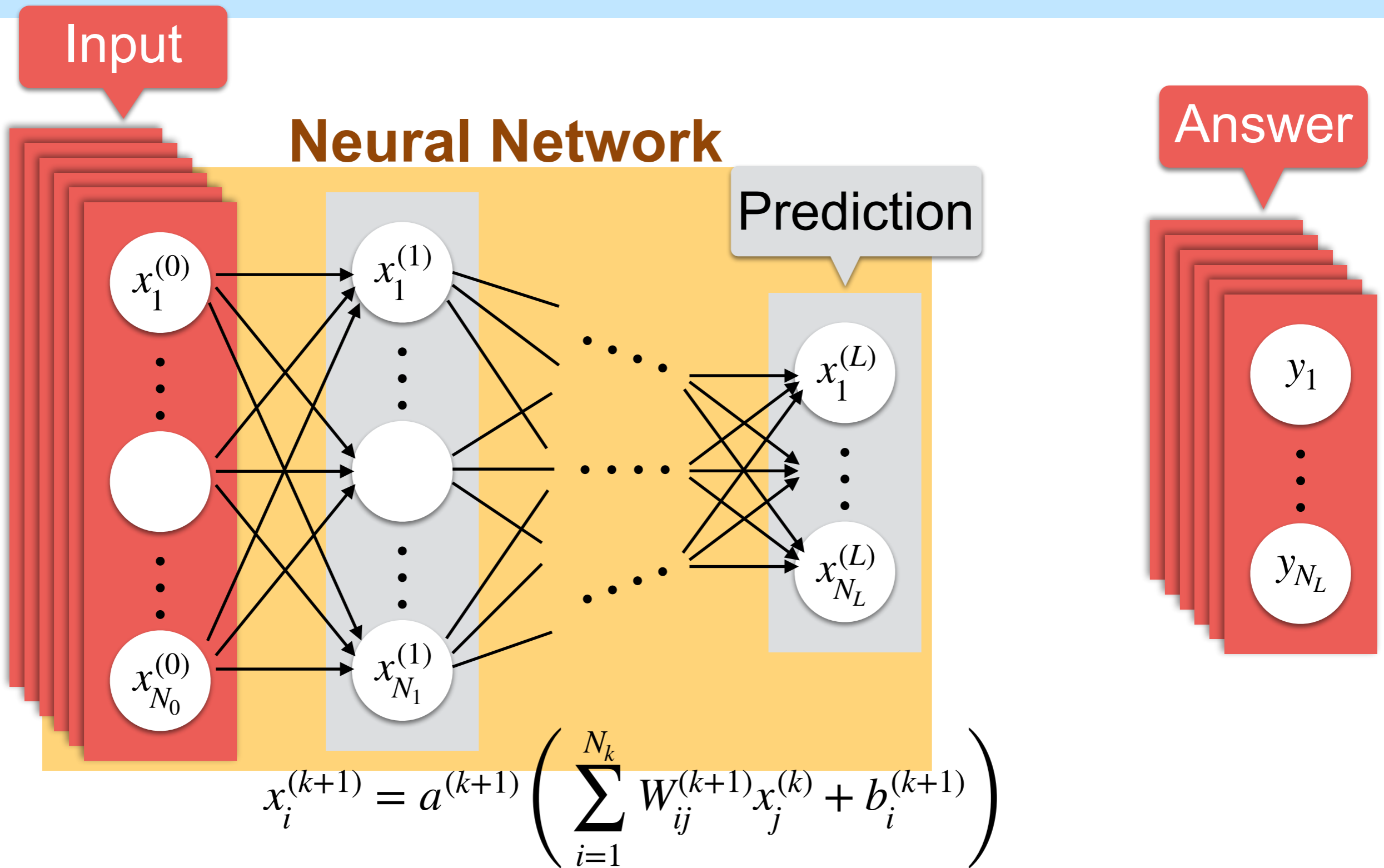
**Training data**



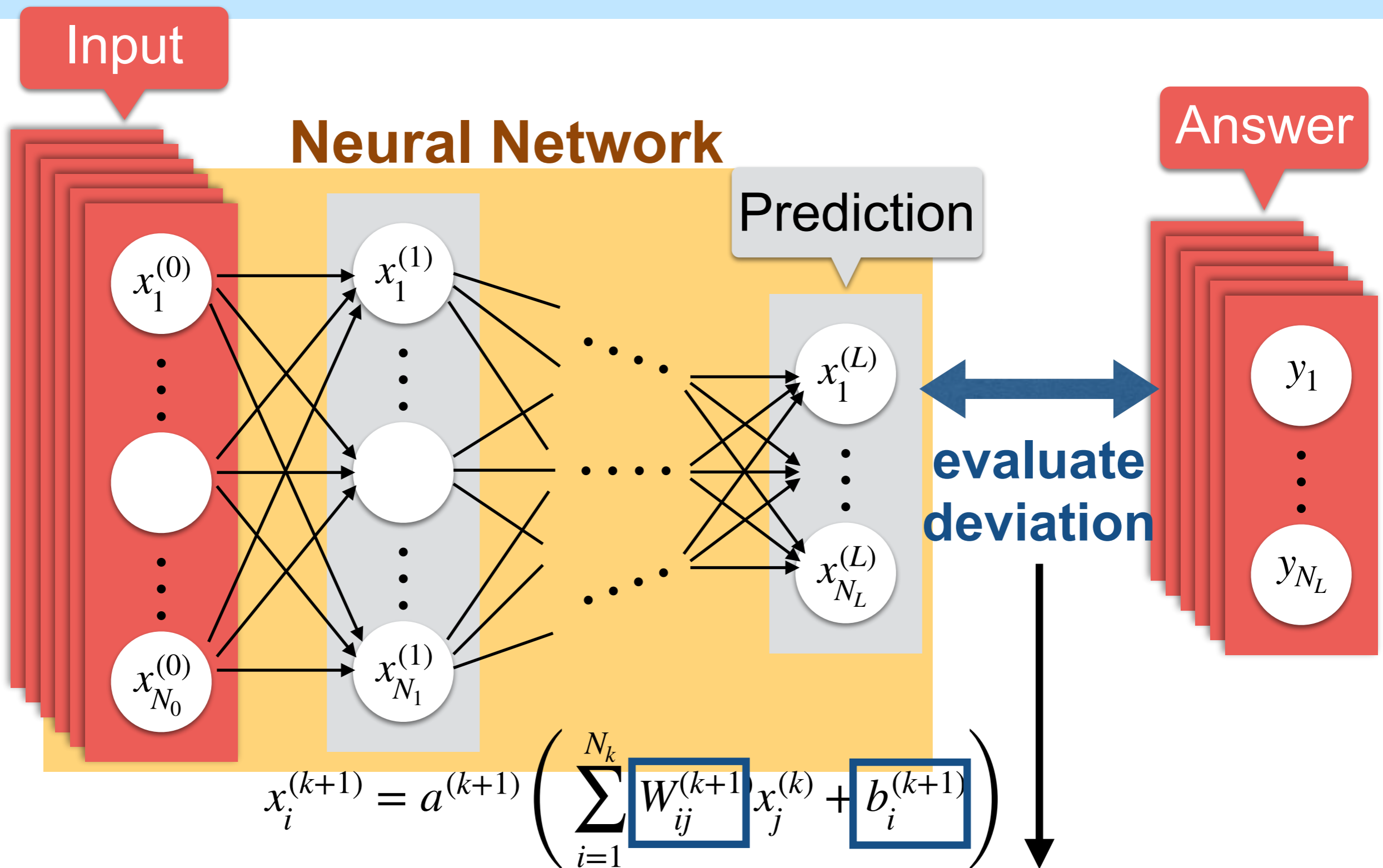
**By “training”  
the neural network,  
we extract the relation  
between input and answer**



# Method of Machine Learning



# Method of Machine Learning



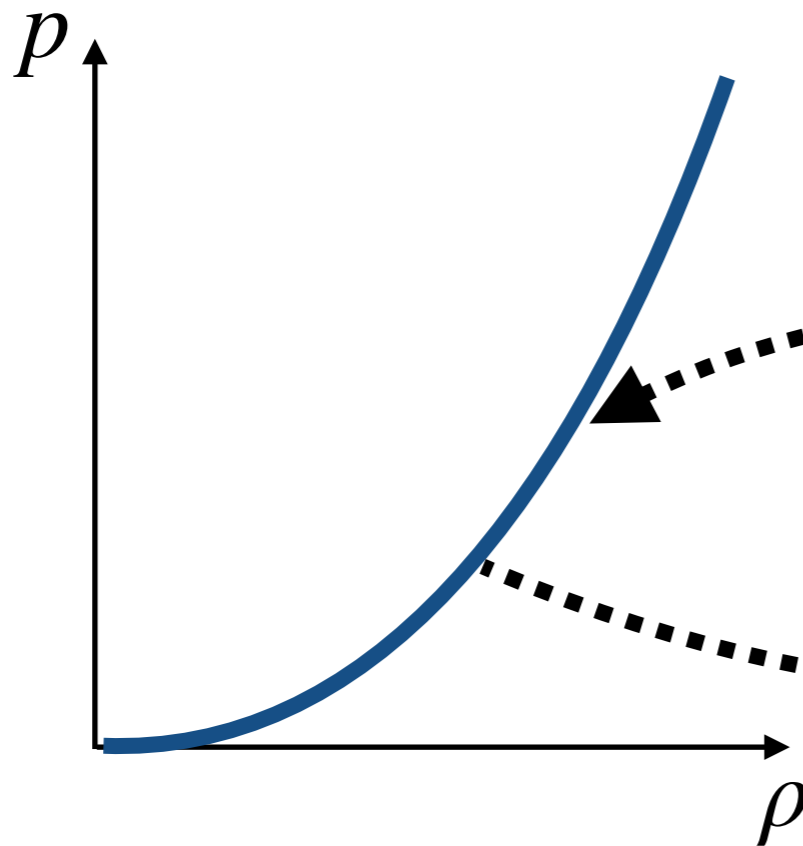
**Tune these parameters inside network to minimize the deviation**

# EoS & $M$ - $R$ : TOV Mapping $\Psi_{\text{TOV}}$

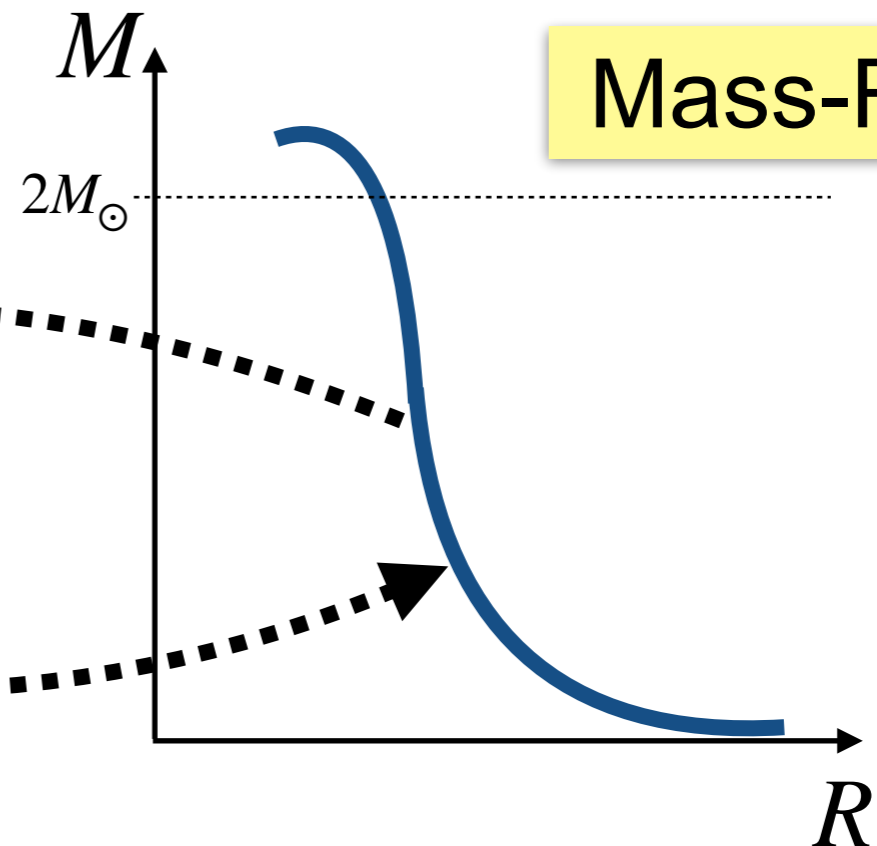
If  $M$ - $R$  curve is given,  
there is **one-to-one correspondence**

Lindblom (1992)

EoS



Mass-Radius

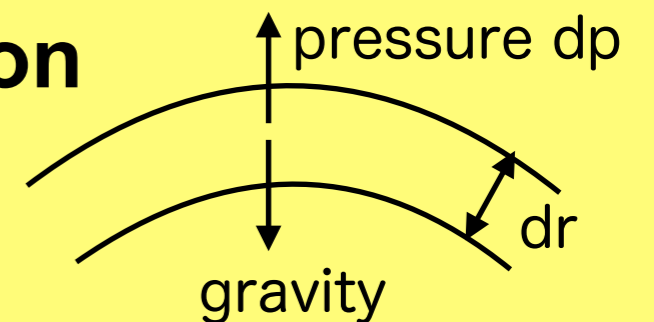


$\Psi_{\text{TOV}}^{-1}$

$\Psi_{\text{TOV}}$

**Tolman-Oppenheimer-Volkoff (TOV) equation**

General relativistic structural equation  
for stars at hydrostatic equilibrium

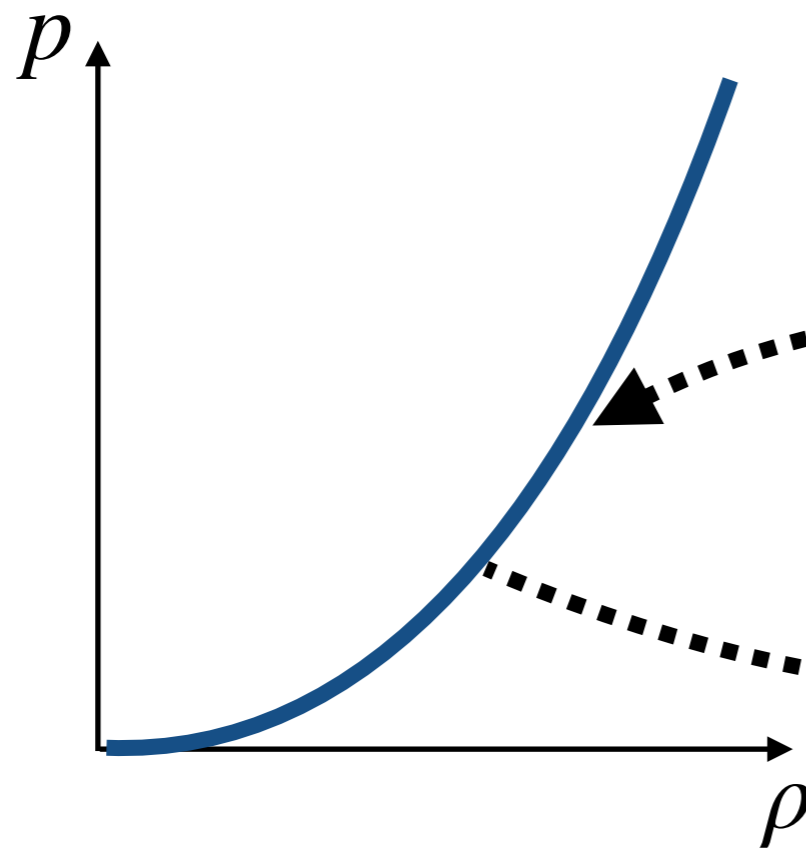


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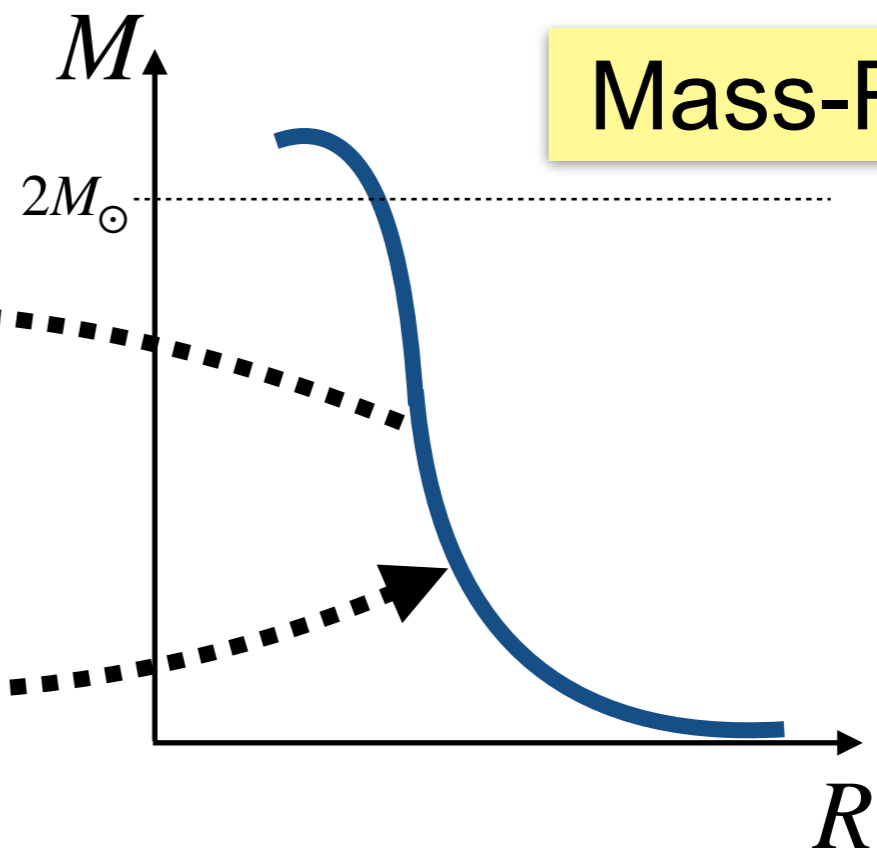
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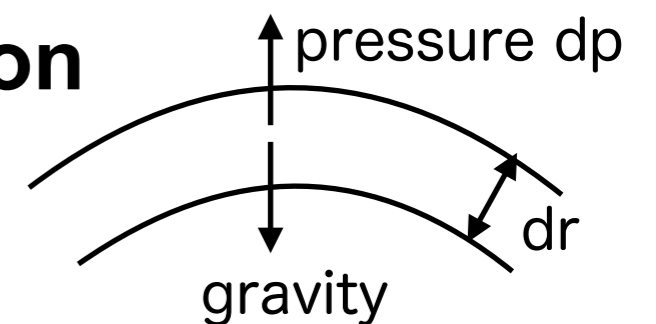


$\Psi_{\text{TOV}}^{-1}$

$\Psi_{\text{TOV}}$

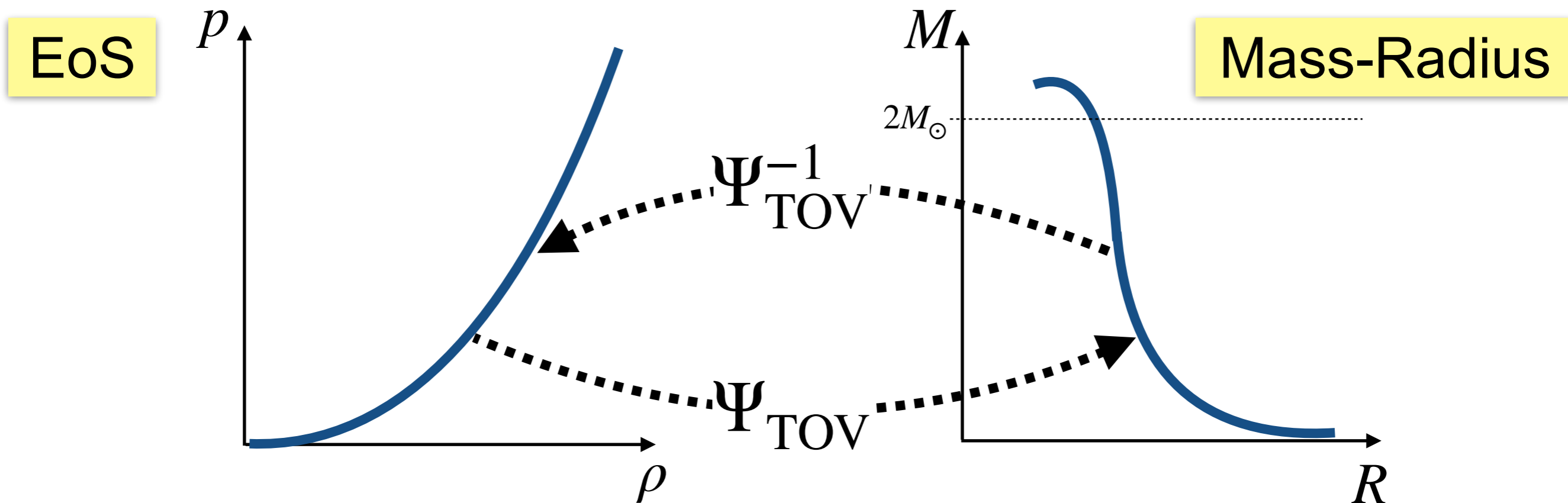
**Tolman-Oppenheimer-Volkoff (TOV) equation**

General relativistic equation for  
hydrostatic equilibrium



# EoS & $M$ - $R$ : TOV Mapping $\Psi_{\text{TOV}}$

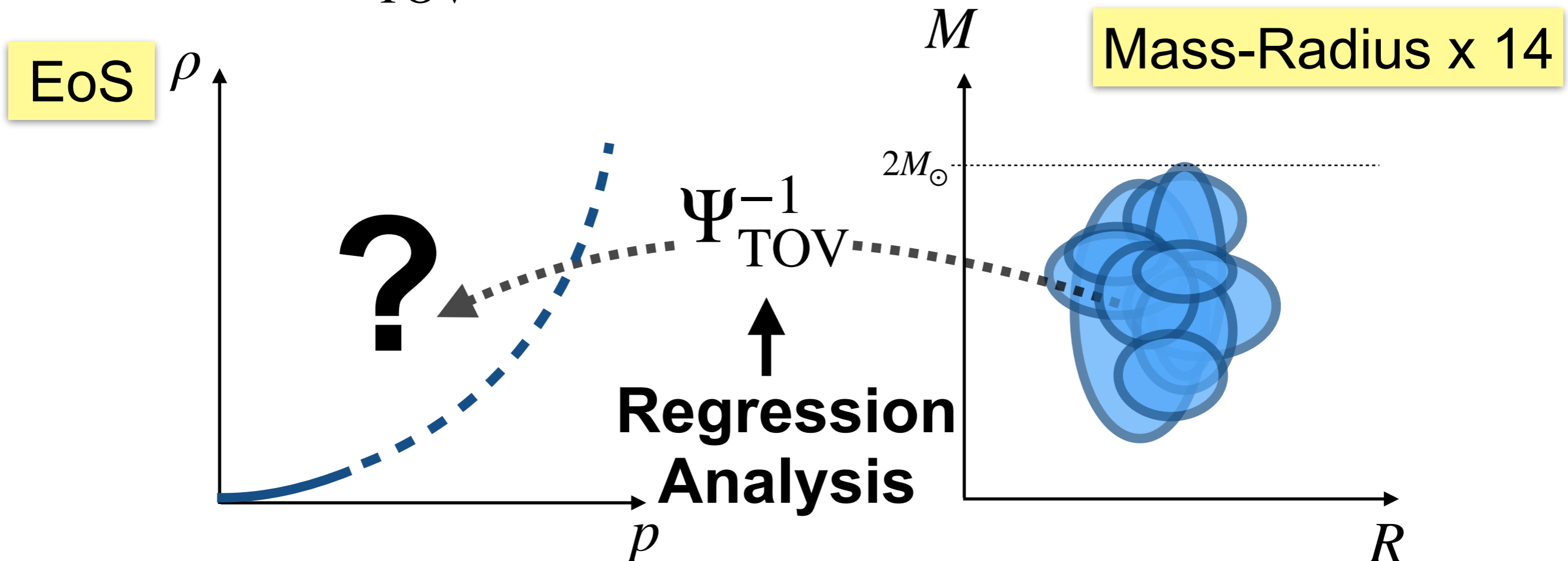
The operation of the TOV equation can be regarded as one-to-one mapping: **TOV mapping**  $\Psi_{\text{TOV}}$



# Regression Analysis of TOV Mapping

In reality...  $M$ - $R$  is point-like and has uncertainties

→ finding  $\Psi_{\text{TOV}}^{-1}$  becomes non-trivial!



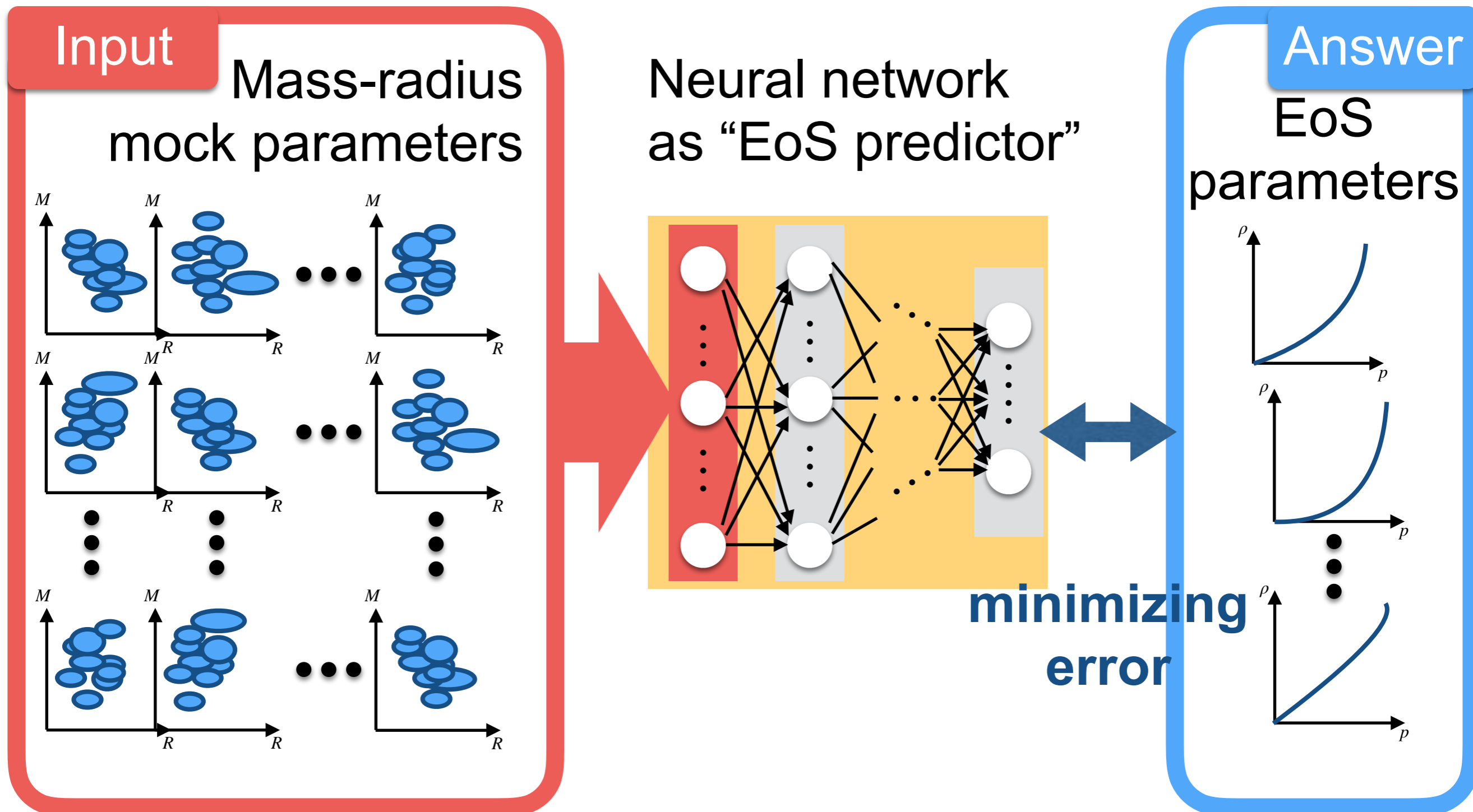
Express  $\Psi_{\text{TOV}}^{-1}$  in terms of Neural Network

→ problem becomes regression analysis finding

$$\text{EoS} = \Psi_{\text{TOV}}^{-1}(M-R)$$

# Expressing TOV by Neural Network

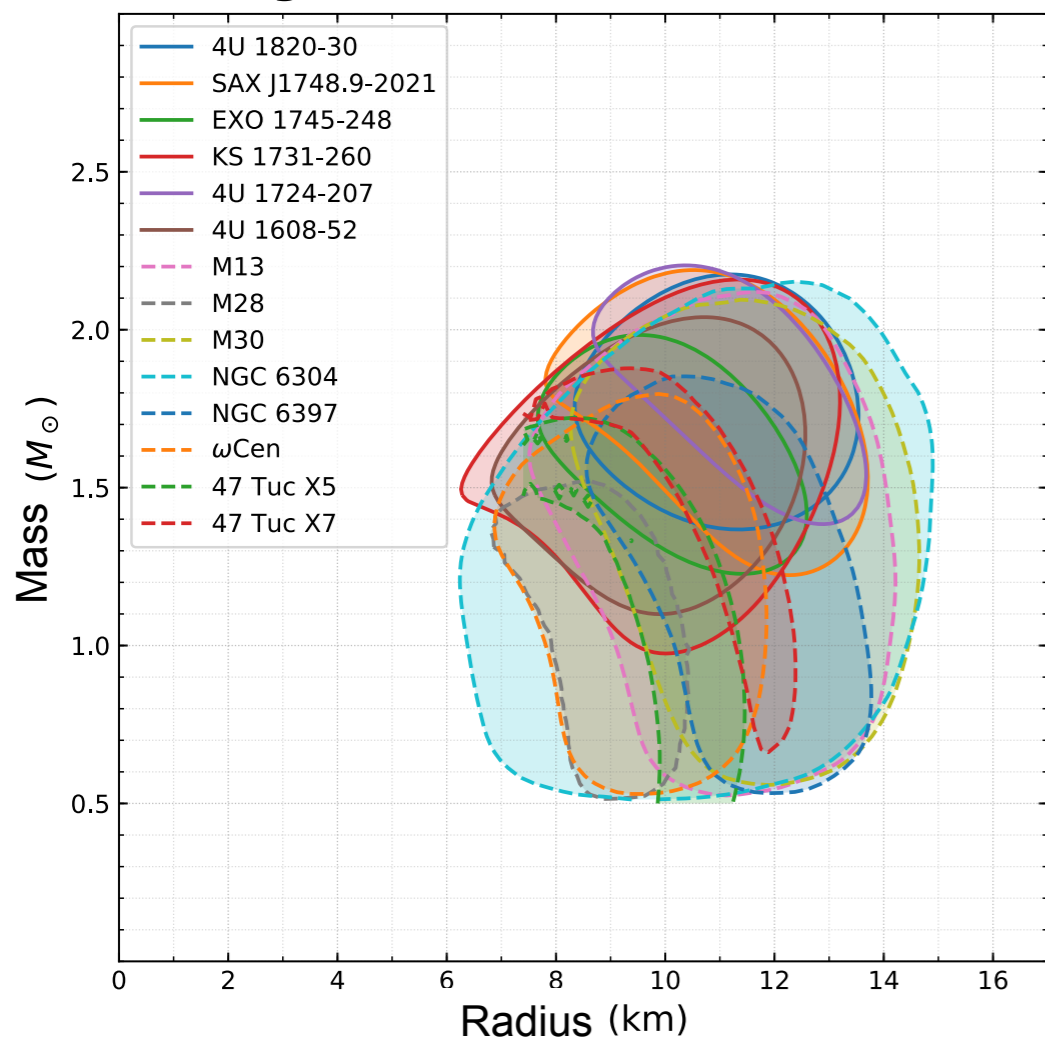
Find the “EoS predictor” using machine learning



# Result: Inferred EoS

Inferred by machine learning

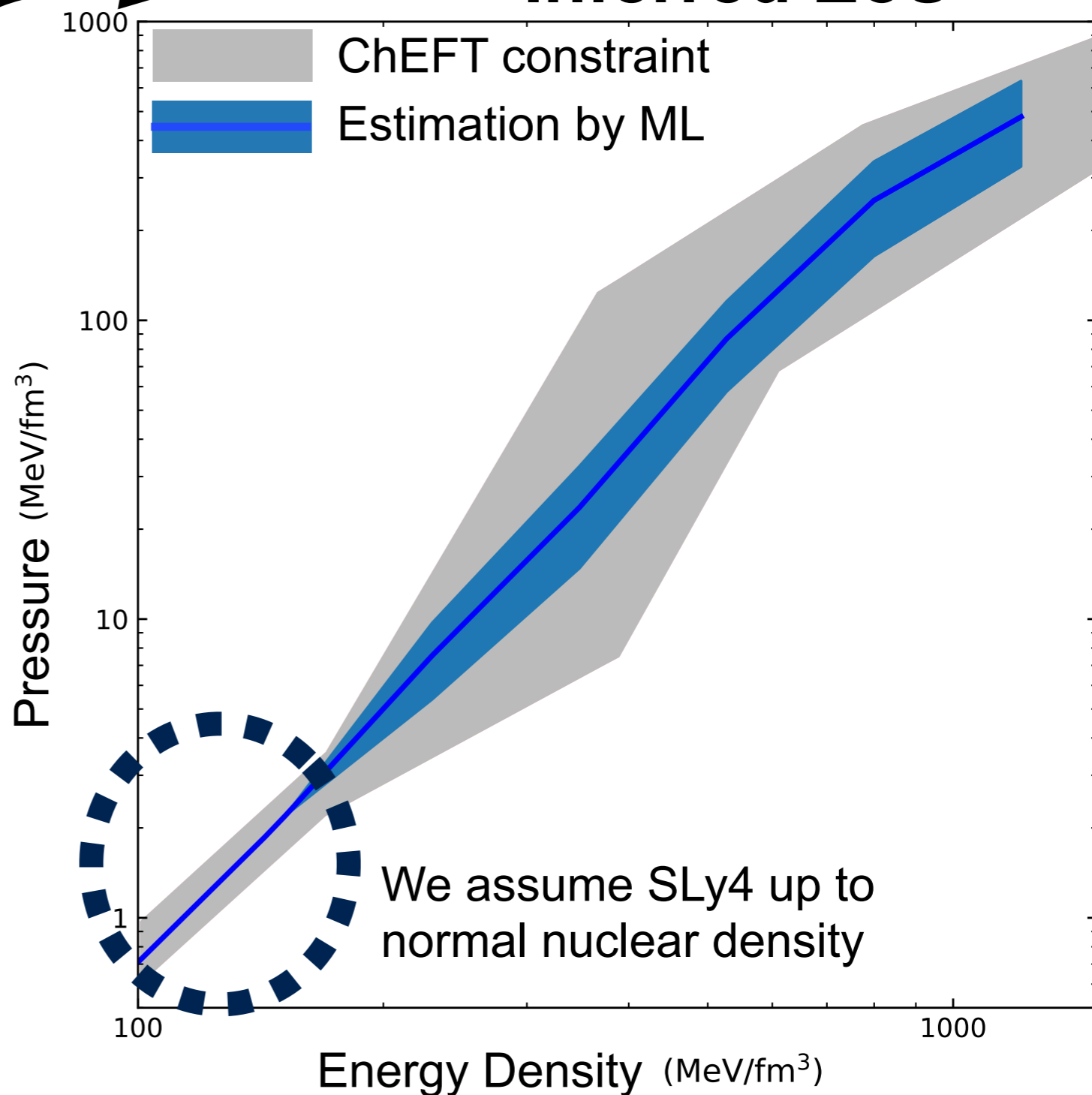
X-ray observation



Özel-Freire (2016)

YF-Fukushima-Murase (2019)

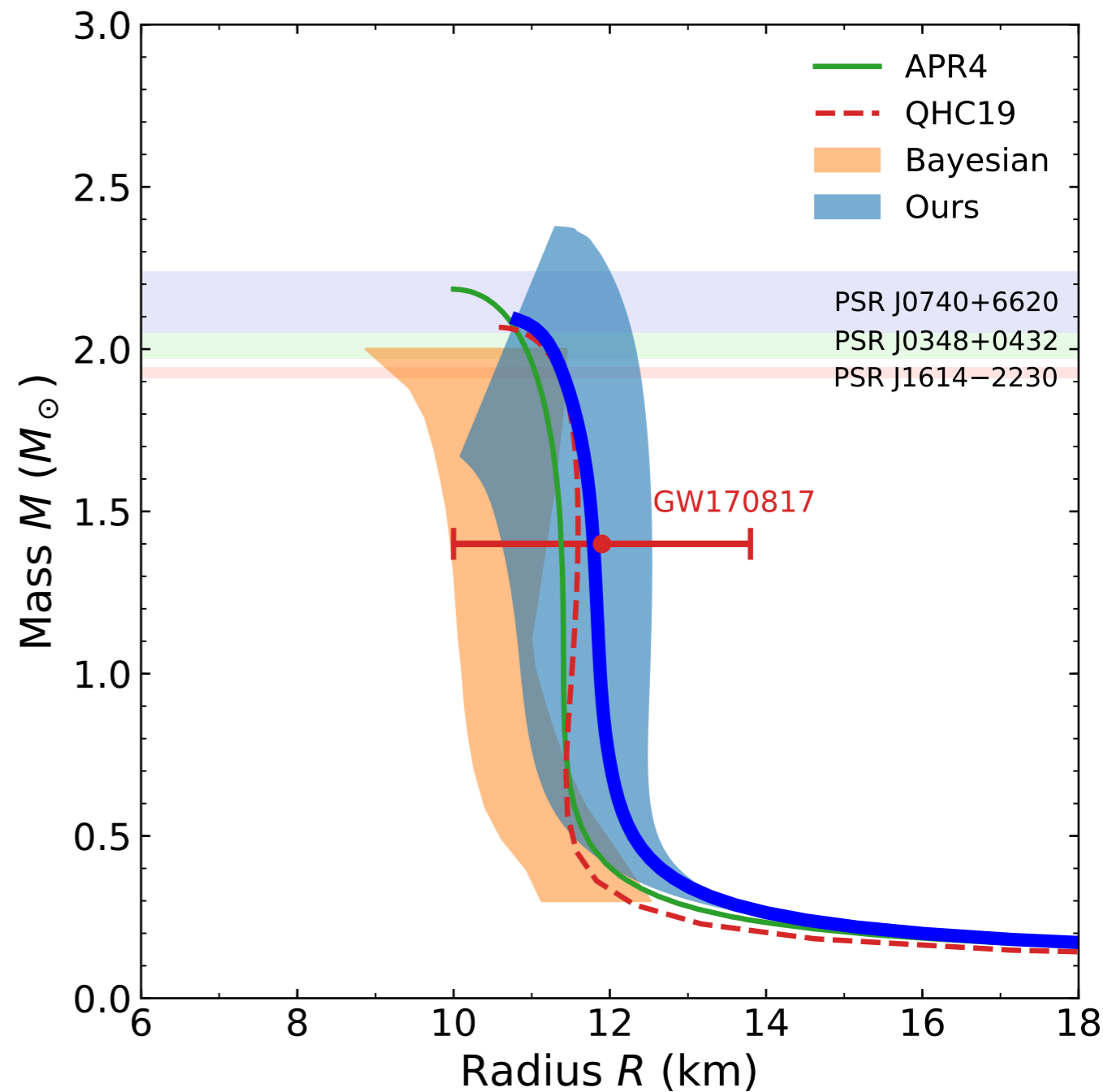
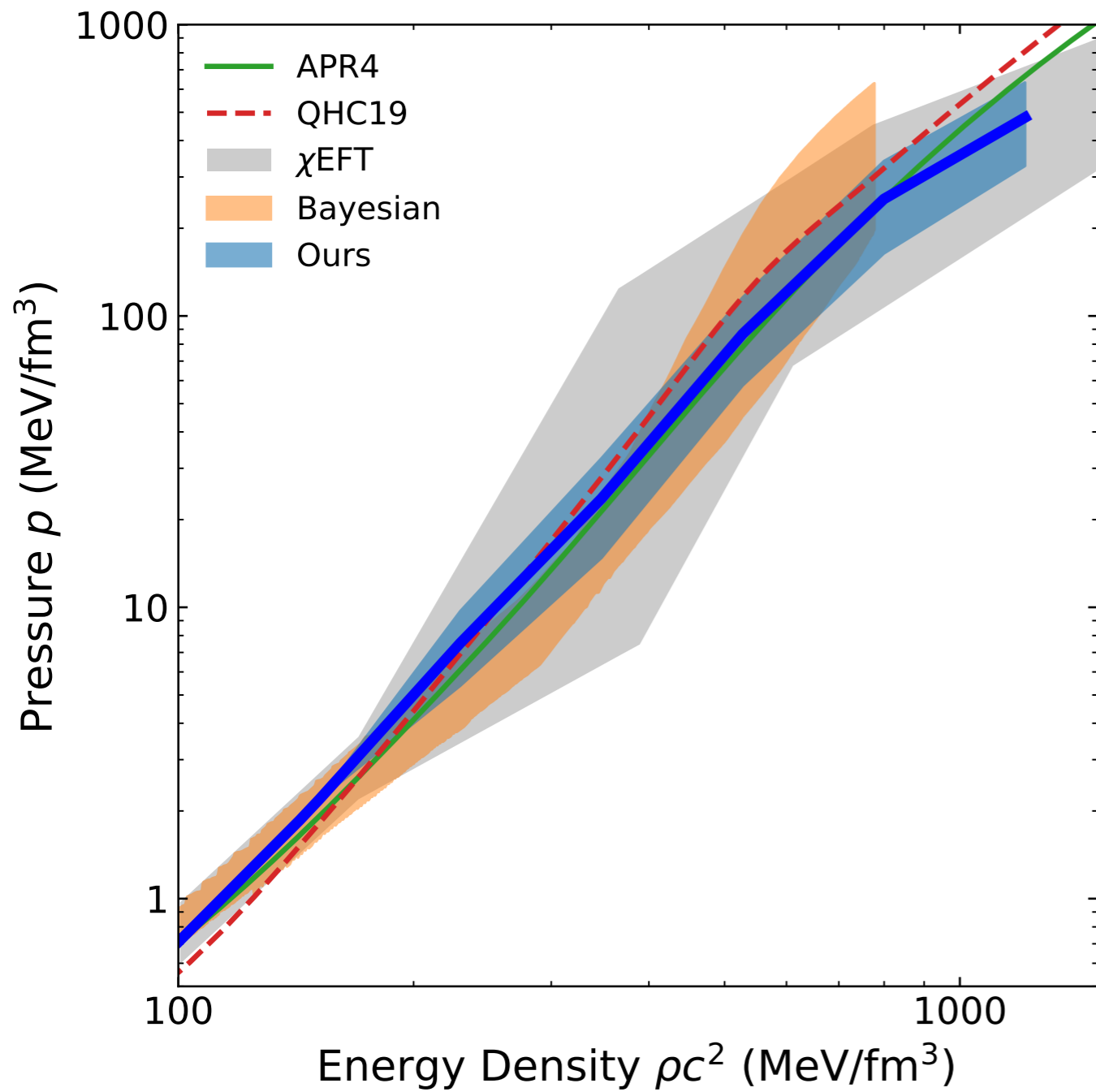
Inferred EoS



Hebeler-Lattimer-Pethick-Schwenk (2013) 16

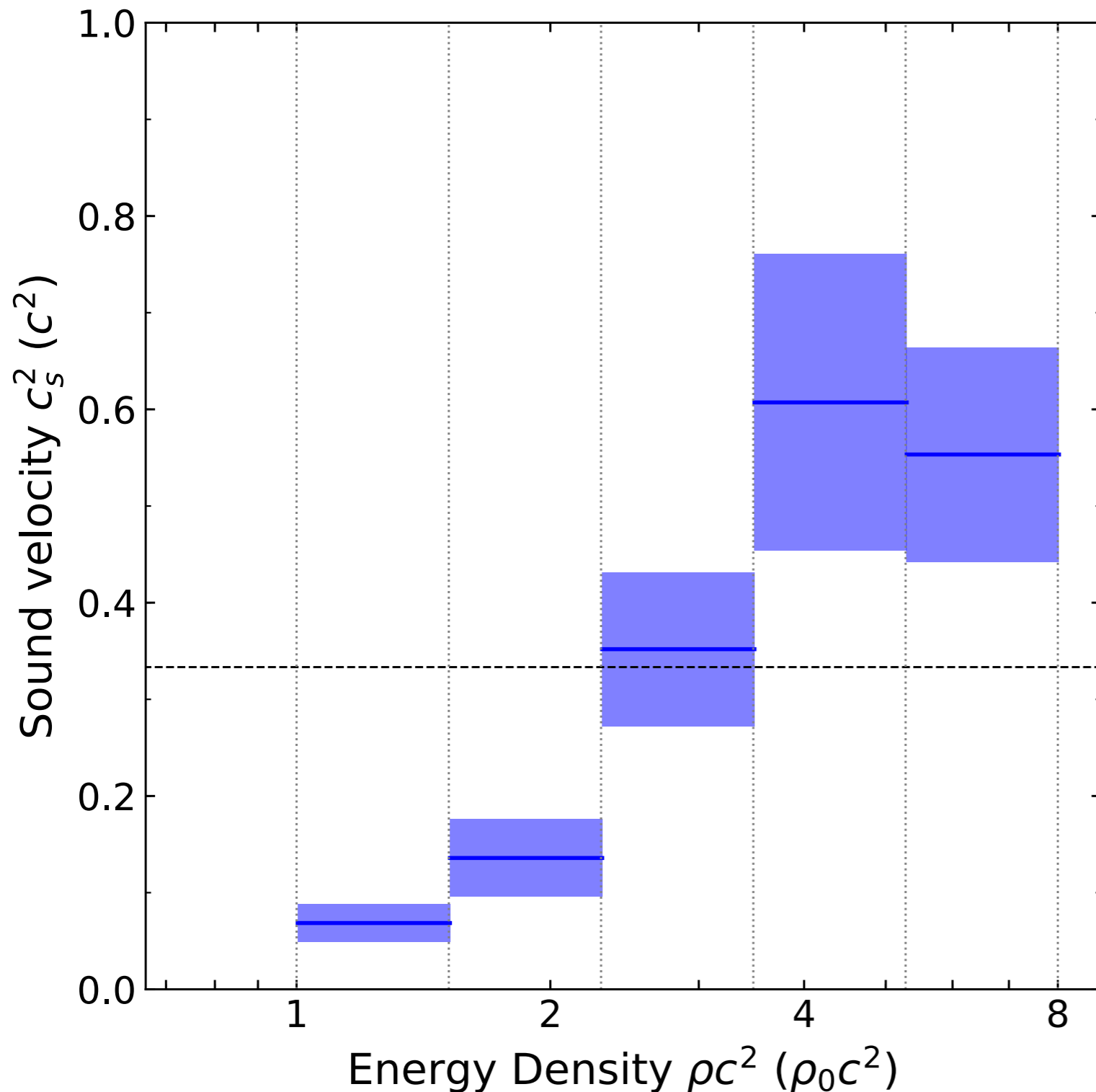


# EoS with M-R relation



# Result: Sound velocity

## Average sound velocity:



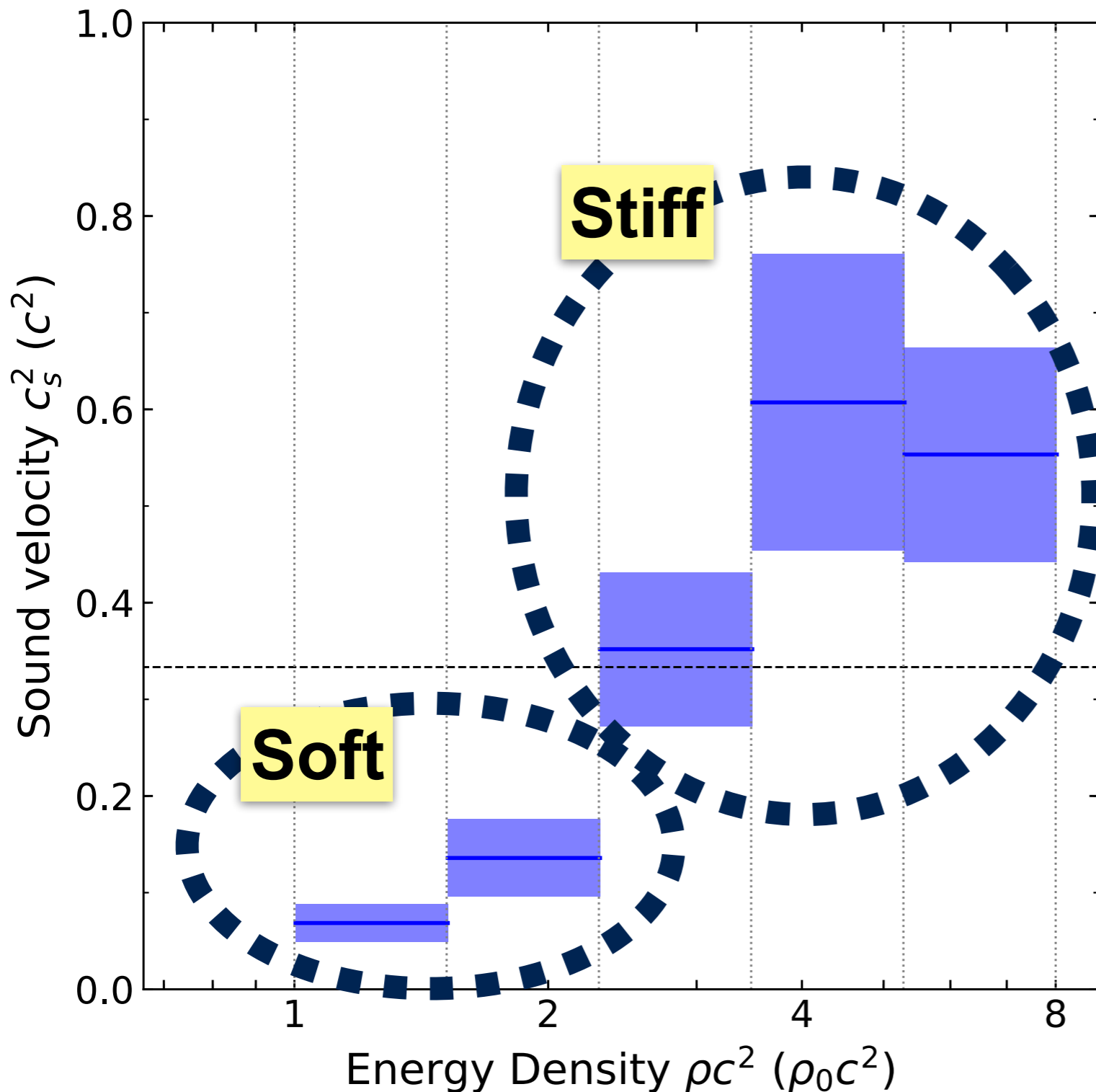
$$c_s^2 = \frac{\partial p}{\partial \rho}$$

- Measure of **stiffness** for EoS
- Our inferred result shows **Soft-to-Stiff** behavior
- This property may imply **Quarkyonic** picture; Nuclear matter seems as if it is quark matter

Fukushima-Kojo (2015),  
McLerran-Reddy (2018)

# Result: Sound velocity

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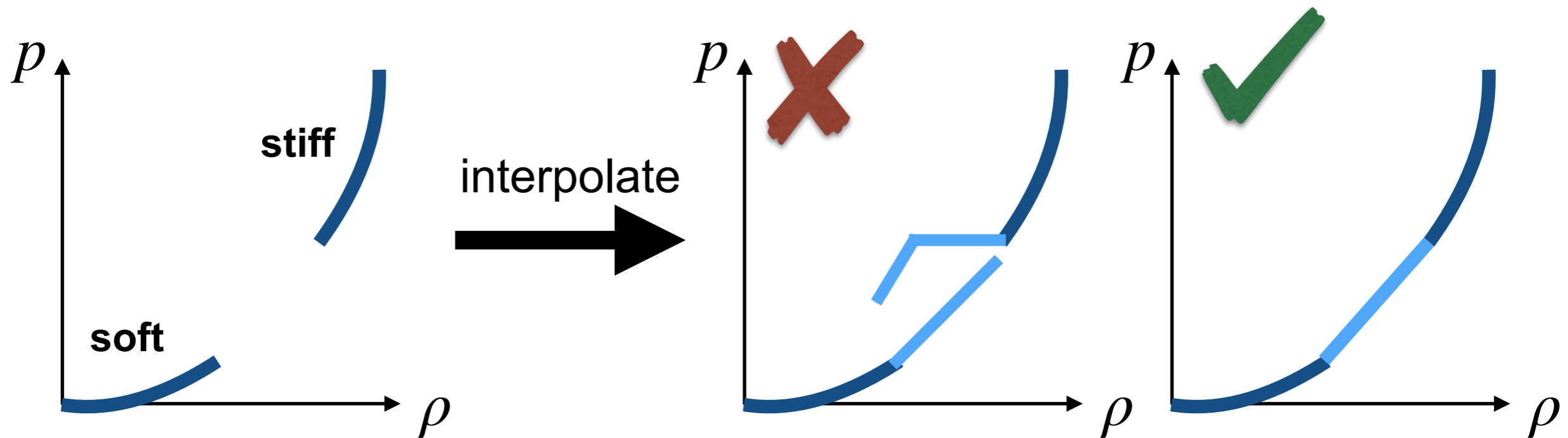
# Soft-to-Stiff EoS

- Our inferred EoS shows soft-to-stiff behavior

- Model study of an EoS shows:

Han et al. (2019)

If EoS is soft-to-stiff, strong 1st-order QCD transition cannot exist... Only the **crossover** transition is allowed!



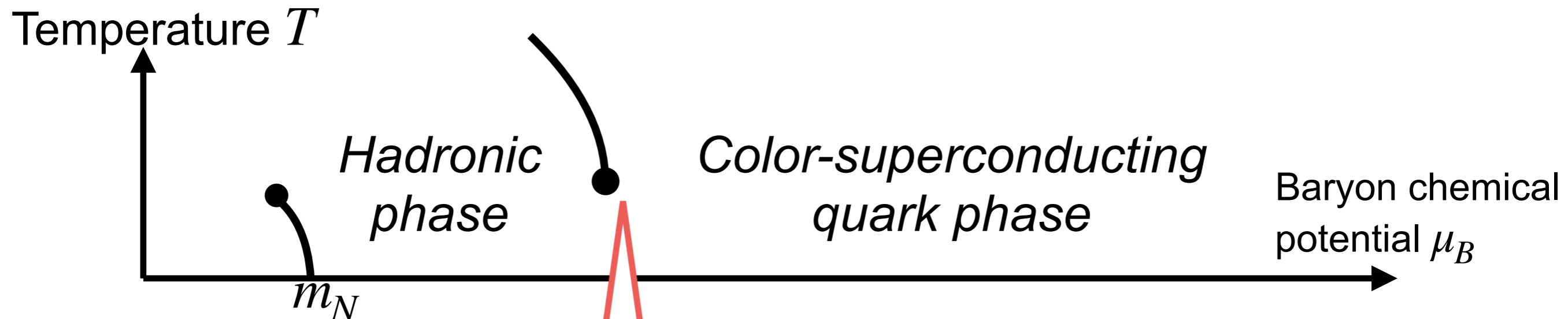
- What is the physical foundation for the crossover?:

**Quark-Hadron Continuity**

# Quark-Hadron Continuity

Hatsuda et al. (2006)

## Phase diagram of QCD at low temperature:



Both phases are continuously connected  
(**Quark-Hadron Continuity**)

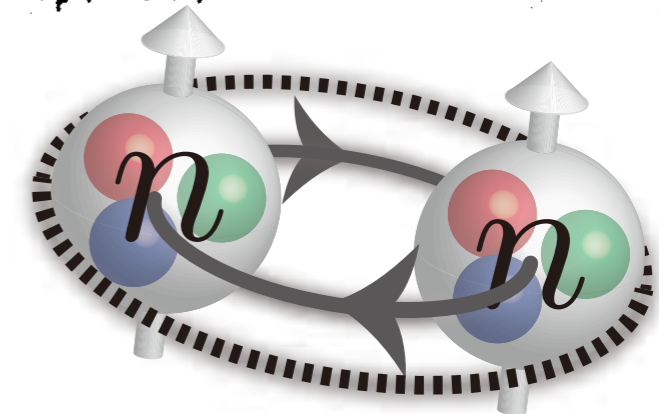
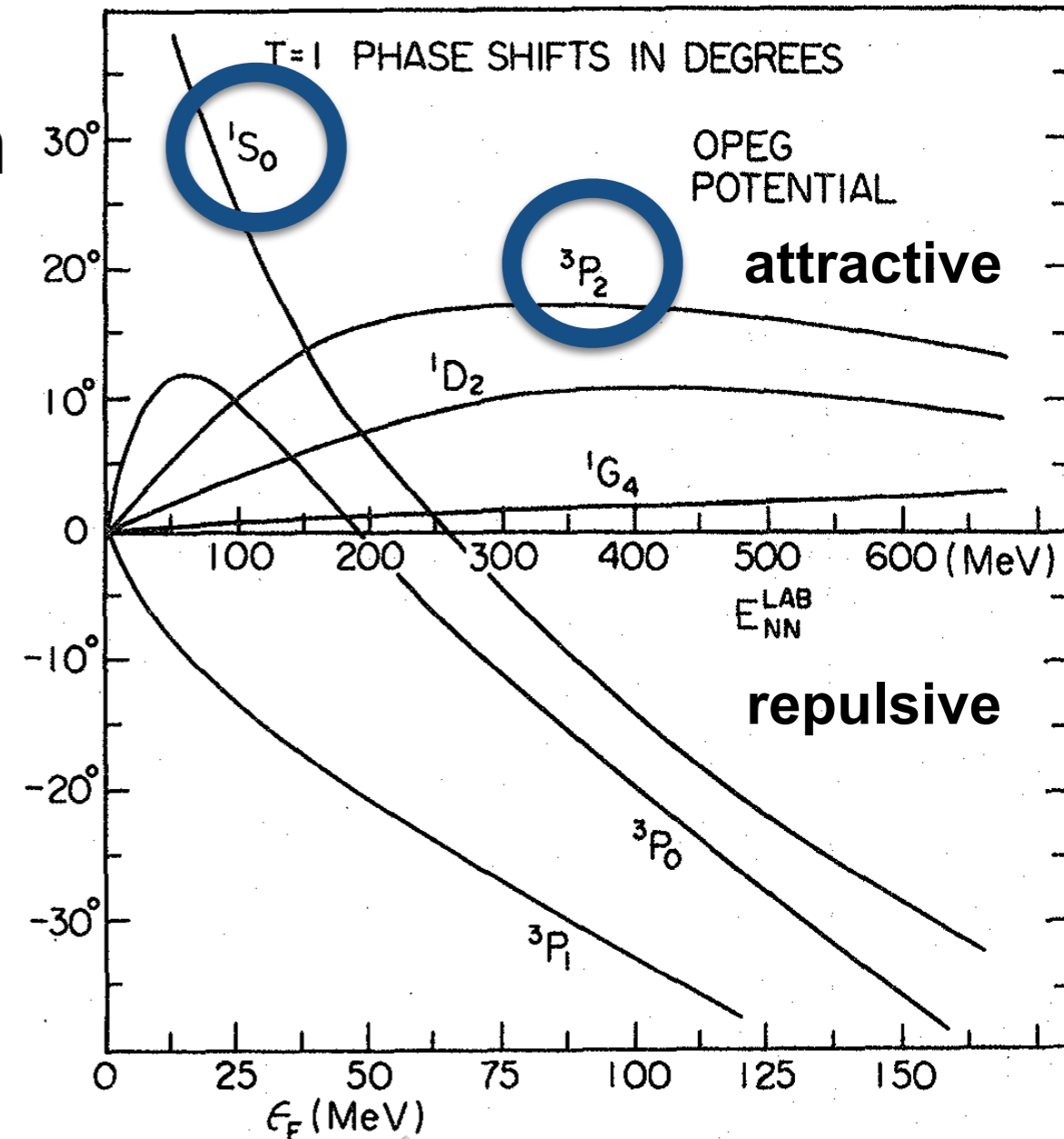
Schäfer-Wilczek (1999)

- Preceding analysis has been made to **three-flavor symmetric** matter with masses  $m_u = m_d = m_s$
- Natural starting point: **two-flavor symmetric matter** because in reality  $m_s/m_{u,d} \sim 30$

# Neutron Superfluidity

- Neutrons show superfluidity because effective  $nn$ -interaction is attractive in some channel
- At higher densities, pairing in  $^1S_0$  state gets weakened by repulsive core;  $\rightarrow$   $^3P_2$  state takes over
- Is this neutron superfluid continuously connected to quark phase?  $\rightarrow$  YES!

Hoffberg et al. (1970), Tamagaki (1970)



# Order Parameter Rearrangement

YF-Fukushima-Weise (2019)

Order parameter of neutron superfluid:  $\langle \hat{n} \gamma^i \nabla^j \hat{n} \rangle$

Neutron operator:  $\hat{n} = (\hat{u} \gamma_5 \hat{d}) \hat{d}$

Order parameter can be rearranged (in MFA) as ...

$$\langle \hat{n} \gamma^i \nabla^j \hat{n} \rangle \simeq \underbrace{\langle \hat{u} \gamma_5 \hat{d} \rangle}_{\text{2SC diquark condensate}} \underbrace{\langle \hat{d} \gamma^i \nabla^j \hat{d} \rangle}_{\text{Additional } \langle dd \rangle \text{ arises!}}$$

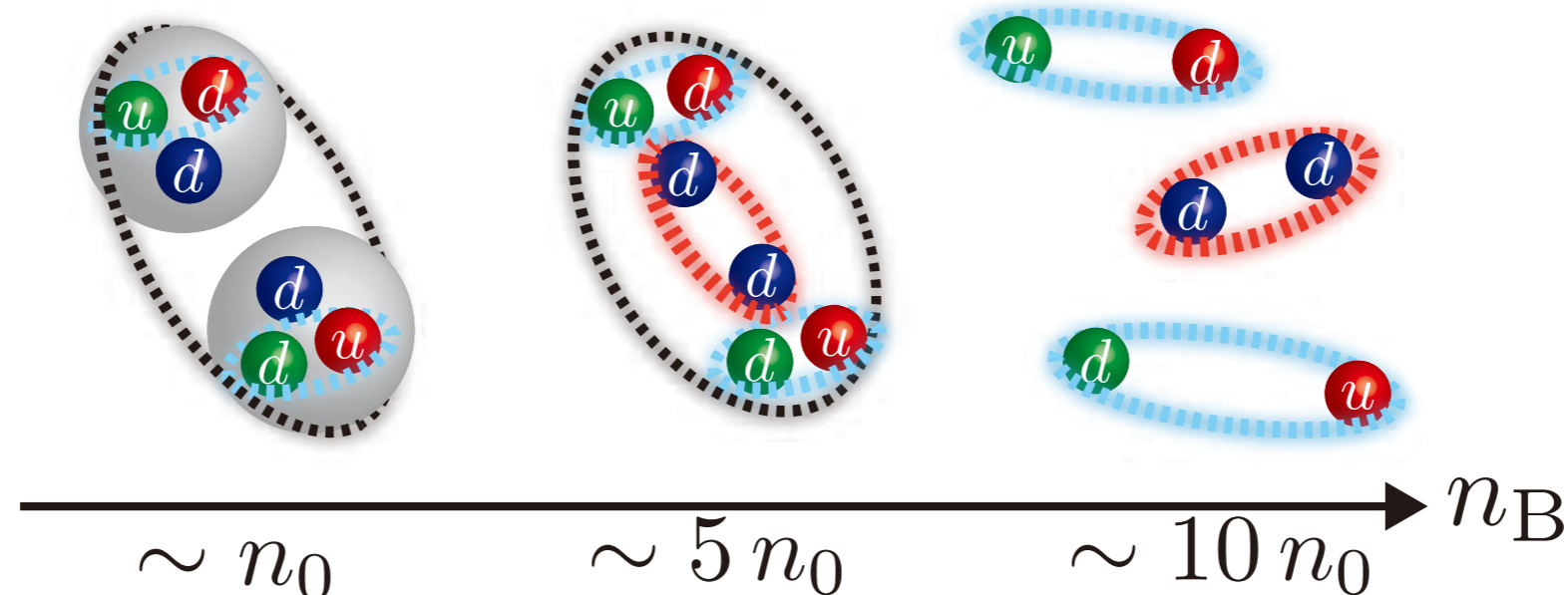
2SC diquark

Additional  $\langle dd \rangle$  arises!

condensate

Neutron superfluid

Color superconductor



# Two-flavor Continuity: 2SC+ $\langle dd \rangle$

YF-Fukushima-Weise (2019)

- Two diquark condensates:

$$\langle ud \rangle + \langle dd \rangle$$

- **2SC +  $\langle dd \rangle$**  induces...

- Chiral symmetry breaking

- Baryon  $U(1)_B$  breaking (leads baryon superfluidity)

→ Same as the neutron superfluid, **continuity set in!**

- For dynamical aspects and supportive evidences, please refer to:

**YF**, K. Fukushima & W. Weise, arXiv:1908.09360



# Summary

- Established the method of machine learning to infer EoS from mass and radius observations
- Our EoS result is consistent with quark-hadron continuity scenario
- Established the quark-hadron continuity in more realistic context of two-flavor than the previous one; neutron superfluid  $\rightarrow$  2SC +  $\langle dd \rangle$  condensate

# **Supplementary materials**

# Dynamical Aspects of $2SC+\langle dd \rangle$

To have  ${}^3P_2$  pairing at high density, one needs:

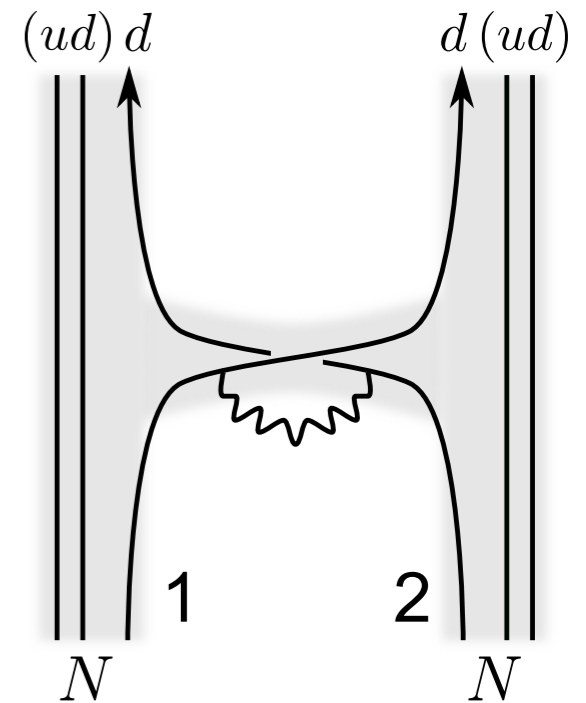
- **short-range repulsive core**  
→ disfavors pairing in  ${}^1S_0$  state
- **attractive spin-orbit potential**  
→ favors  ${}^3P_2$  state among  ${}^3P_{J=0,1,2}$  states

# Dynamical Aspects - Repulsive Core

## Dynamics that disfavors $^1S_0$ pairing

- Short-range interaction between d-quarks can be understood with this diagram
- Potential between d-quarks reads:

$$V_{12}^{\text{OGE}} = \left( \sum_A T_1^A T_2^A \right) \frac{\alpha_s}{4} \left[ \frac{1}{r_{12}} - \frac{2\pi}{3m_d^2} (\mathbf{s}_1 \cdot \mathbf{s}_2) \delta(\mathbf{r}_{12}) \right]$$



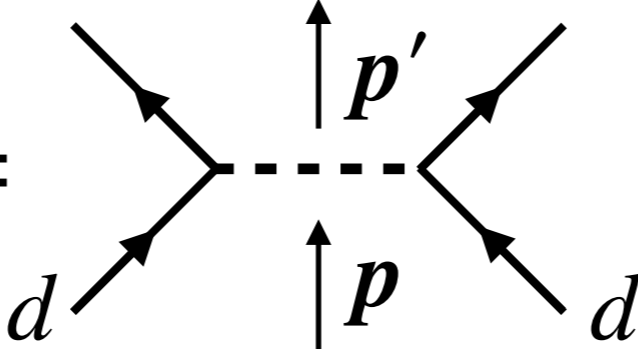
- $\langle dd \rangle$  condensate is in color-sextet channel  
 $\rightarrow$  Color factor above becomes repulsive!

# Dynamical Aspects - LS Potential

## Dynamics that favors ${}^3P_2$ pairing

- NJL-type model with scalar and vector couplings + diquark correlations:

$$\mathcal{L}_{\text{int}} = G(\bar{\psi}\psi)^2 + H(\bar{\psi}\bar{\psi})(\psi\psi) - G_V(\bar{\psi}\gamma^\mu\psi)^2$$

$$\langle \mathbf{p}' | V_{LS} | \mathbf{p} \rangle =$$


$$\rightarrow \text{Fourier transform } V_{LS}(\mathbf{r}) \propto - (1 + mr) \frac{e^{-mr}}{r^2} \mathbf{L} \cdot \mathbf{S} \quad (m > 0)$$

$$\mathbf{L} \cdot \mathbf{S} = \begin{cases} -2 & ({}^3P_{J=0}) \\ -1 & ({}^3P_{J=1}) \\ 1 & ({}^3P_{J=2}) \end{cases}$$

$\rightarrow$  **Spin-orbit potential is attractive only in  ${}^3P_2$  state**