

Gluonic Spin Orbit Correlations in Photon Pair Production

Marc Schlegel
Tuebingen University

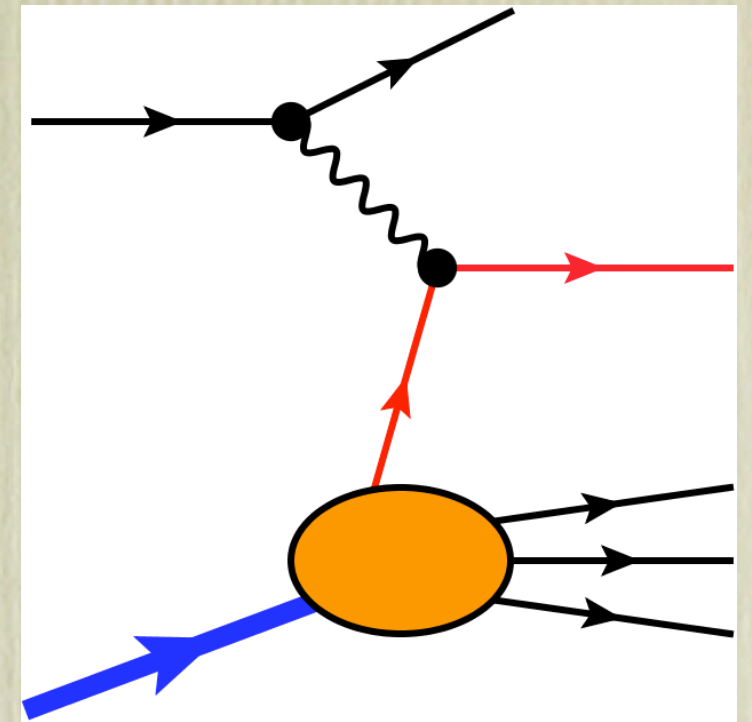
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Transverse Momentum Dependent (TMD) Parton Distributions

Collinear Factorization in pQCD:

- Cross Section at high energies \rightarrow (hard part) \times (soft part)
- Hard Part \rightarrow pQCD ; Soft Part \rightarrow Universal
- applicable to many (integrated) observables, e.g. incl. DIS
- Soft Part \rightarrow collinear Parton Distributions (1-dimensional)

$$f_1(x, \mu), g_1(x, \mu), h_1(x, \mu)$$

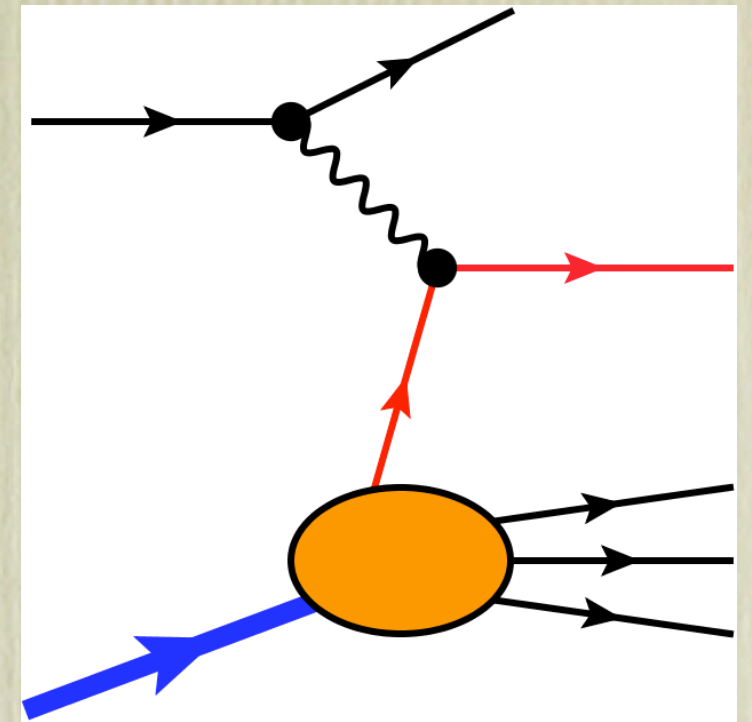


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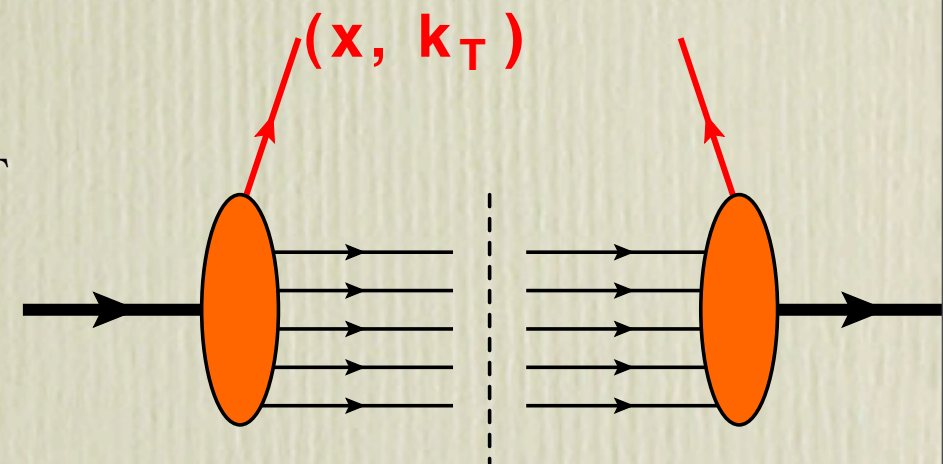
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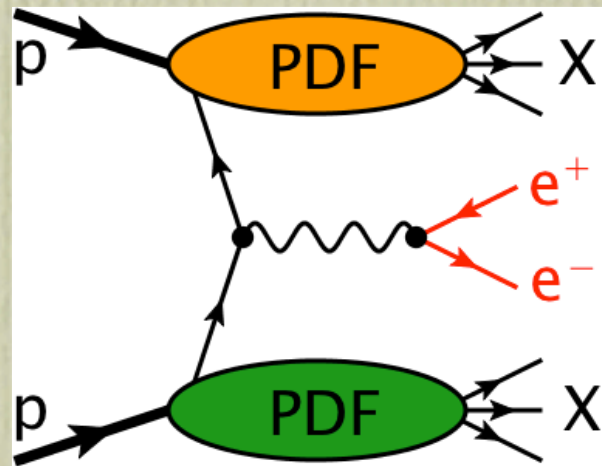
Idea of TMDs:

- Implement “intrinsic” transverse parton momentum k_T
 \rightarrow different kind of factorization
 \rightarrow opportunity to study different aspects of hadron structure

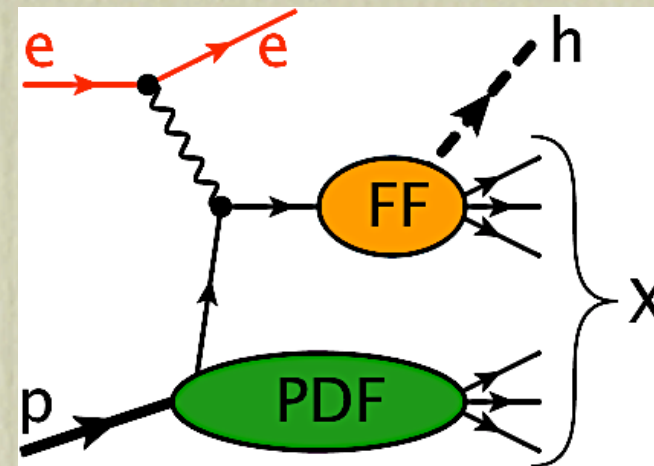


Quark TMDs in Drell-Yan & SIDIS

“intrinsic” transverse parton momentum through small final state transverse momenta



$$q_T \ll Q$$



$$P_{hT} \ll Q$$

TMD factorization (tree-level)

Drell-Yan

$$W^{\mu\nu} \sim \int d^2 k_{aT} d^2 k_{bT} \delta^{(2)}(\vec{k}_{aT} + \vec{k}_{bT} - \vec{q}_T) \text{Tr}[\gamma^\mu \Phi(x_a, \vec{k}_{aT}) \gamma^\nu \bar{\Phi}(x_b, \vec{k}_{bT})]$$

SIDIS

$$W^{\mu\nu} \sim \int d^2 k_T d^2 p_T \delta^{(2)}(\vec{k}_T - \vec{p}_T - \vec{P}_{hT}/z) \text{Tr}[\gamma^\mu \Phi(x, \vec{k}_T) \gamma^\nu \Delta(z, \vec{p}_T)]$$

(Naive) definition of the quark TMD correlator

$$\Phi_{ij}(x, \vec{k}_T; S) = \int \frac{dz^- d^2 z_T}{2(2\pi)^3} e^{ik \cdot z} \langle P, S | \bar{\psi}_j(0) \mathcal{W}_{\text{SIDIS/DY}}[0, z] \psi_i(z) | P, S \rangle \Big|_{z^+=0}$$

→ correct definition (incl. soft gluon resum., rapidity cut-offs, etc.)

[Aybat, Rogers, PRD83, 114042; Collins' new book]

$$\Phi_{ij}(x, \vec{k}_T; S; \xi, \mu)$$

Projection of correlator on $\gamma^+, \gamma^+ \gamma_5, \gamma^+ \gamma^\perp \gamma_5 \rightarrow 8$ quark TMDs

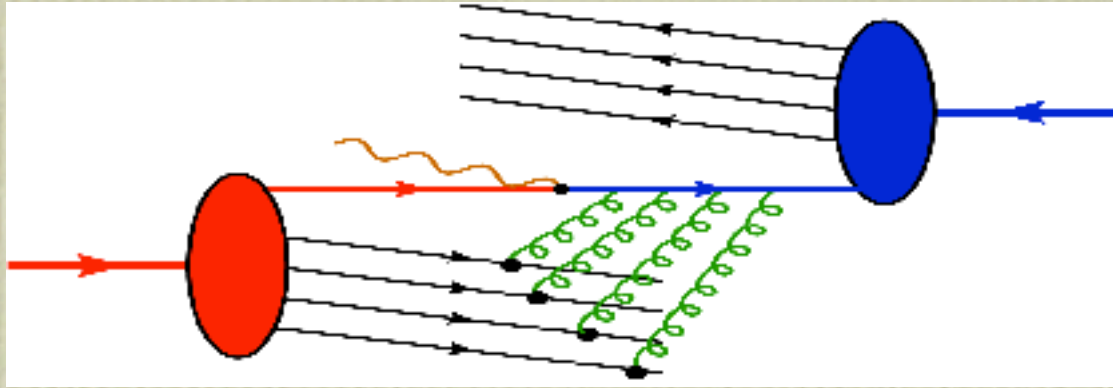
N \ q	U	L	T
U	f_1		h_1^\perp
L		g_1	h_{1L}^\perp
T	f_{1T}^\perp	g_{1T}	h_1, h_{1T}^\perp

time-reversal odd

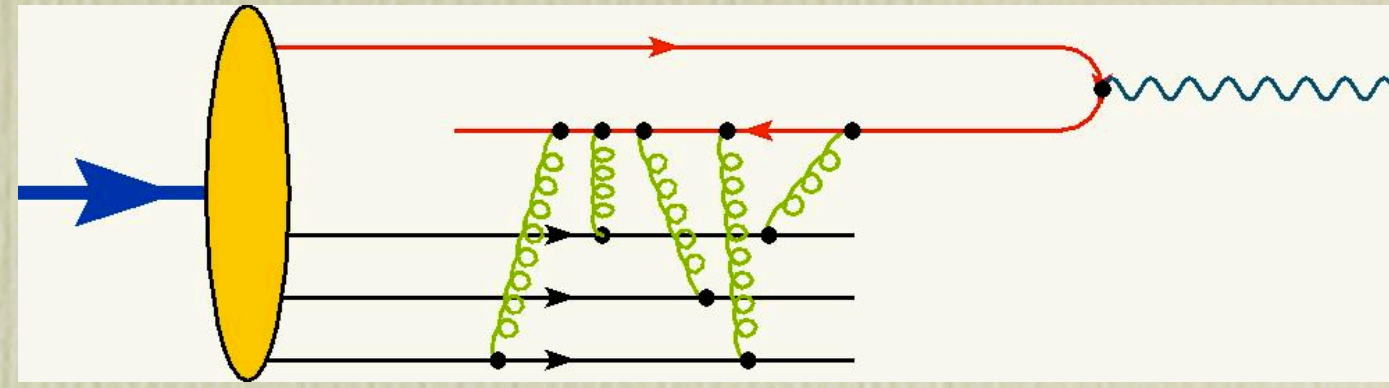
Plot courtesy of B. Musch

Physics of the Wilson line

Initial State Interactions: Drell-Yan

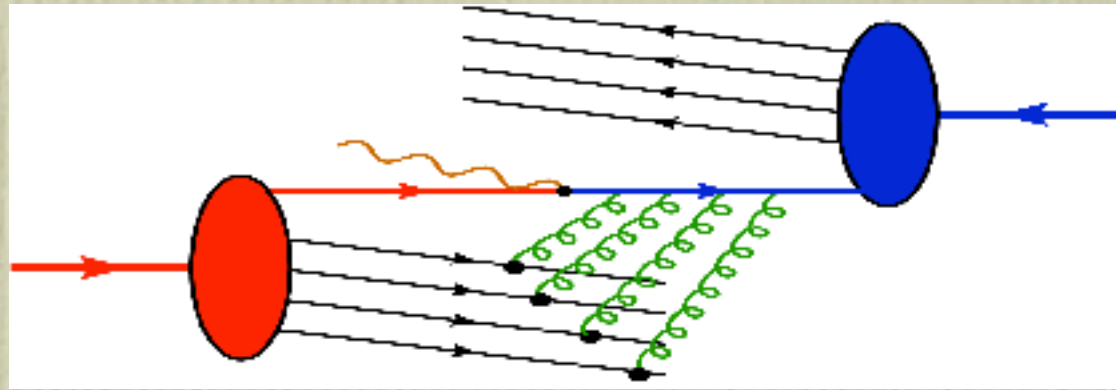


Final State Interactions: SIDIS

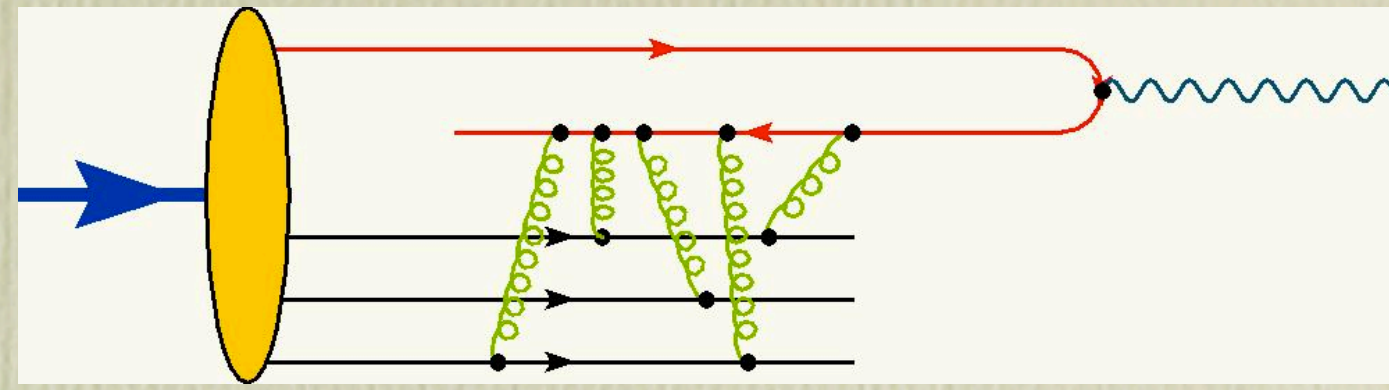


Physics of the Wilson line

Initial State Interactions: Drell-Yan



Final State Interactions: SIDIS



\mathcal{PT} - Transformation on the quark correlator \rightarrow ISI \Leftrightarrow FSI

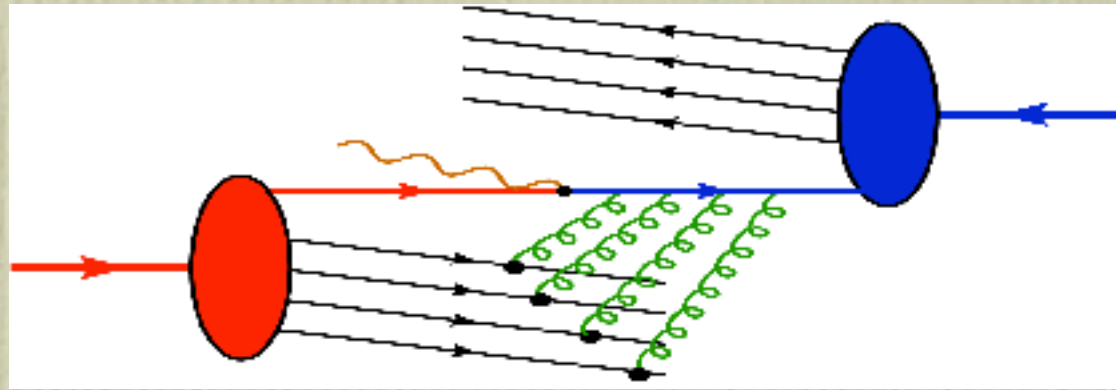
\rightarrow sign switch of Sivers and Boer-Mulder function “T-odd”

$$f_{1T}^{\perp} |_{DIS} = -f_{1T}^{\perp} |_{DY}$$

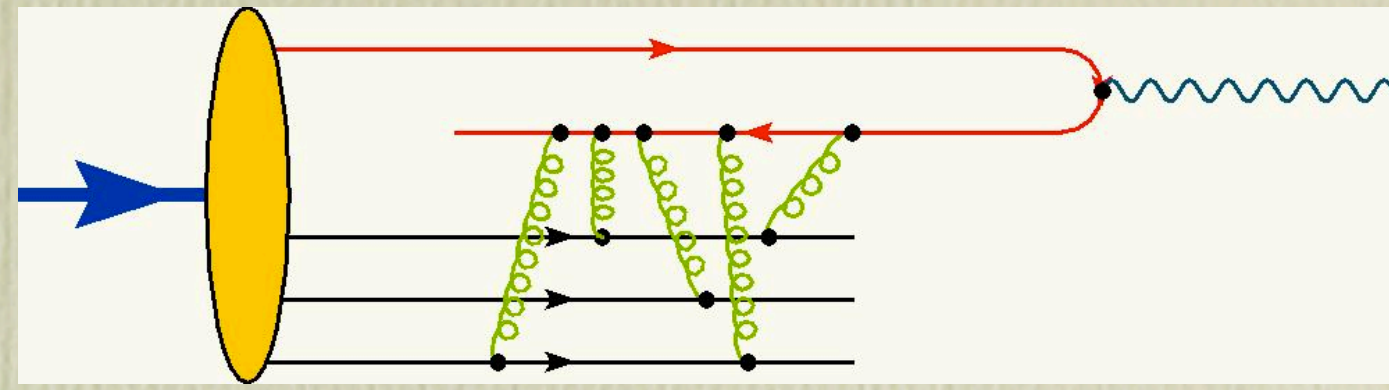
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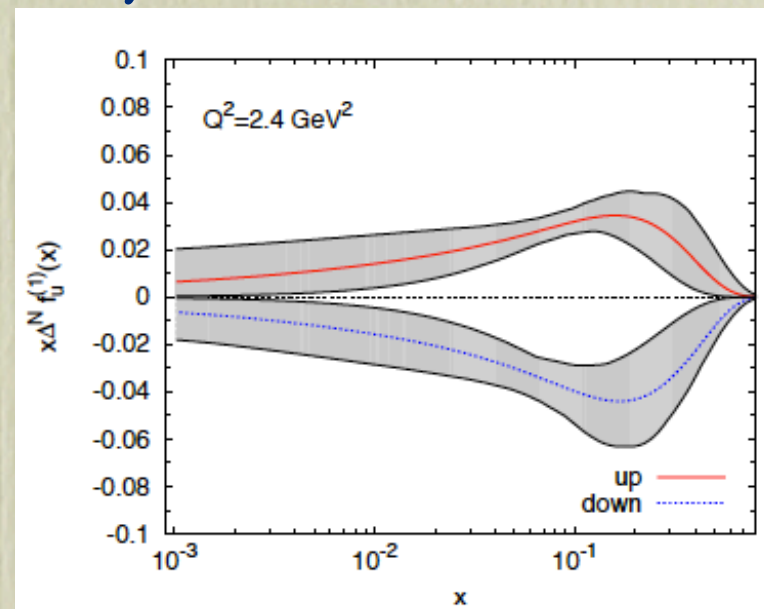
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“T-odd” TMDs generate / appear in observables

\rightarrow e.g. transverse SSA in SIDIS (measured by HERMES, COMPASS etc.) or DY

$$A_{UT}^{SIDIS} \propto \frac{[f_{1T}^{\perp,q} \otimes D_1^q]}{[f_1^q \otimes D_1^q]}$$

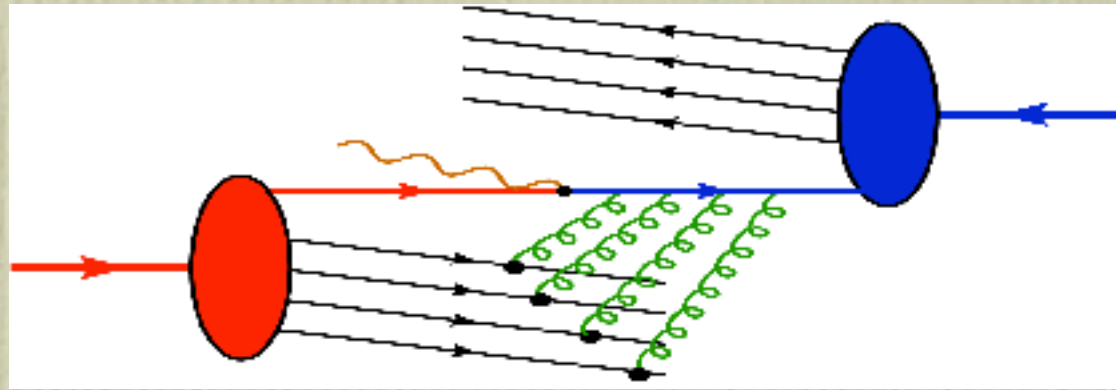
$$A_{UT}^{DY} \propto \frac{[f_{1T}^{\perp,q} \otimes f_1^{\bar{q}}]}{[f_1^q \otimes f_1^{\bar{q}}]}$$



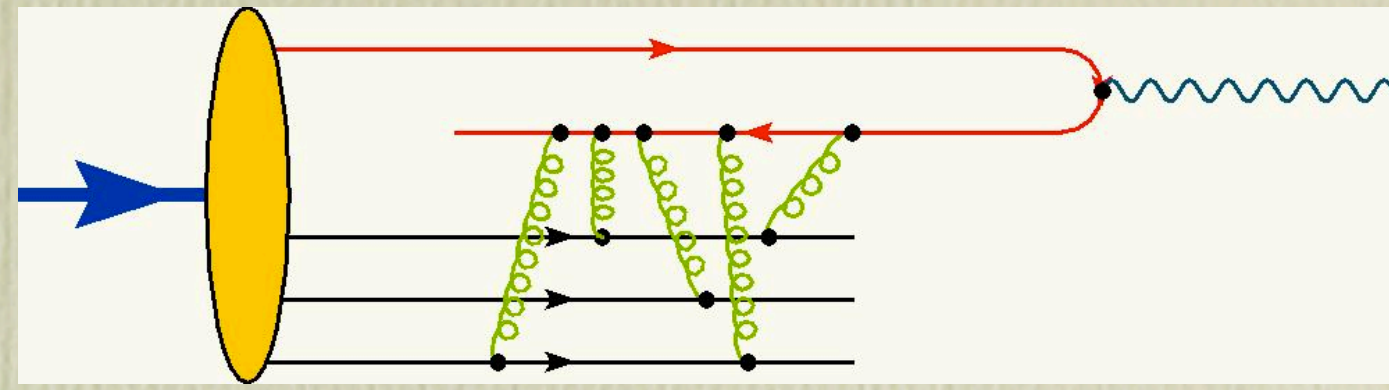
most recent param. from [Anselmino et al., 1107.4446]

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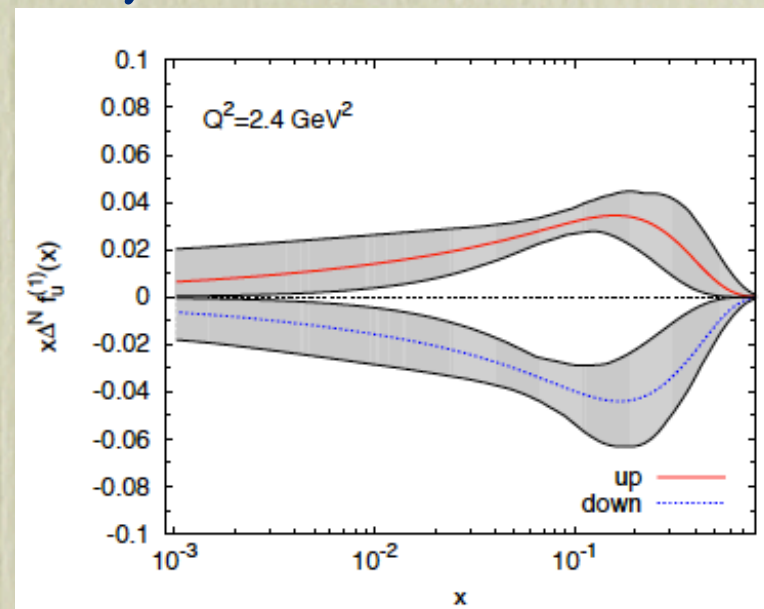
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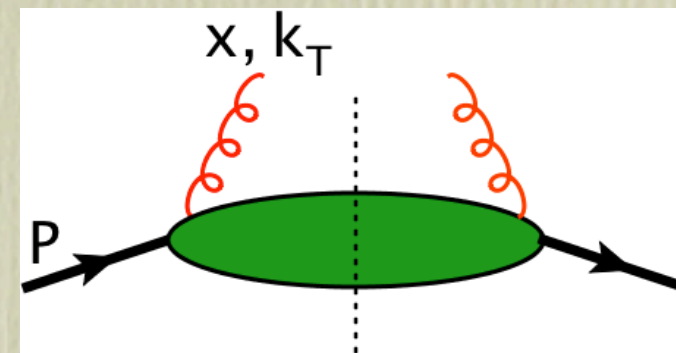


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Eight Gluon TMDs

$$\Gamma^{ij}(x, \vec{k}_T) = \frac{1}{xP^+} \int \frac{dz^- d^2z_T}{(2\pi)^3} e^{ik \cdot z} \langle P, S | F^{+i}(0) \mathcal{W}[0; z] F^{+j}(z) | P, S \rangle \Big|_{z^+=0}$$

$\Gamma^{[T-even]}(x, \vec{k}_T)$		$\Gamma^{[T-odd]}(x, \vec{k}_T)$	
		flip	flip
U	f_1^g $h_1^{\perp g}$		
L	$g_{1L}^{\perp g}$		$h_{1L}^{\perp g}$
T	$g_{1T}^{\perp g}$	$f_{1T}^{\perp g}$	h_1^g $h_{1T}^{\perp g}$



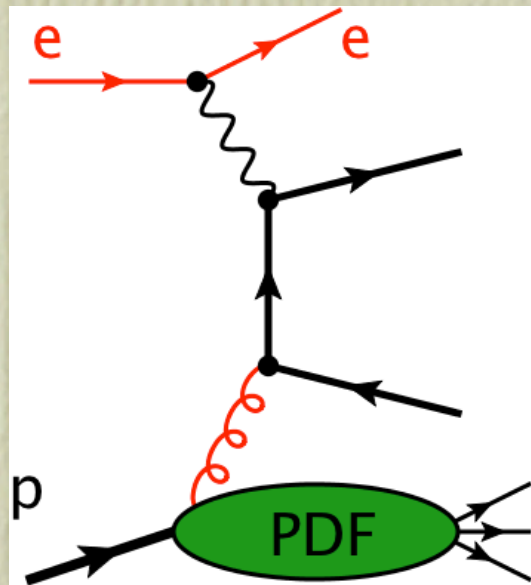
- * linearly polarized gluons: T-even
- * unpolarized gluons in transversely pol. proton:
gluon Sivers function
- * gluonic transversity / pretzelosity / wormgears:
T-odd
- * no chirality
- * two collinear PDFs

Processes sensitive to gluon TMDs

Gluon TMDs do not appear in Drell-Yan or SIDIS

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Heavy Quark production in ep - collisions

[Boer, Brodsky, Mulders, Pisano, PRL 106, 132001]

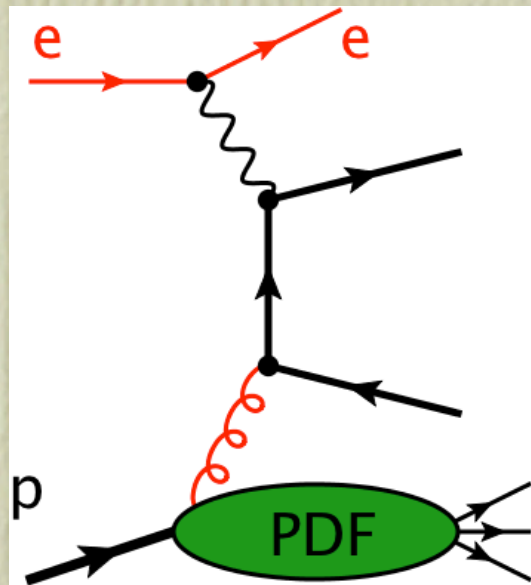
TMD factorization ok!

Spin dependent gluon TMDs: EIC

(Nucleon) spin independent gluon TMDs: EIC / HERA(?)

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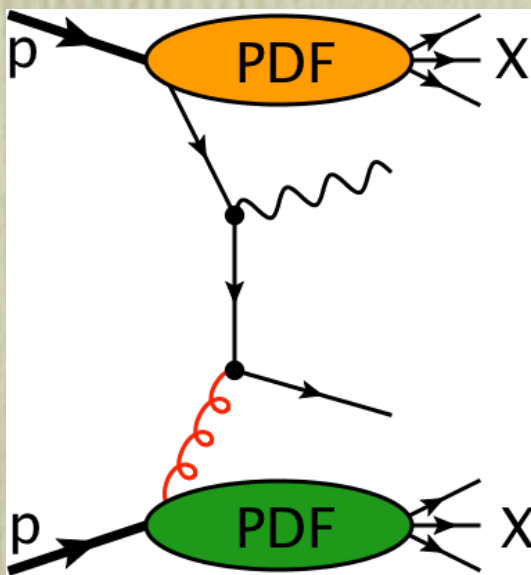
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Jet / Hadron production in pp - collisions

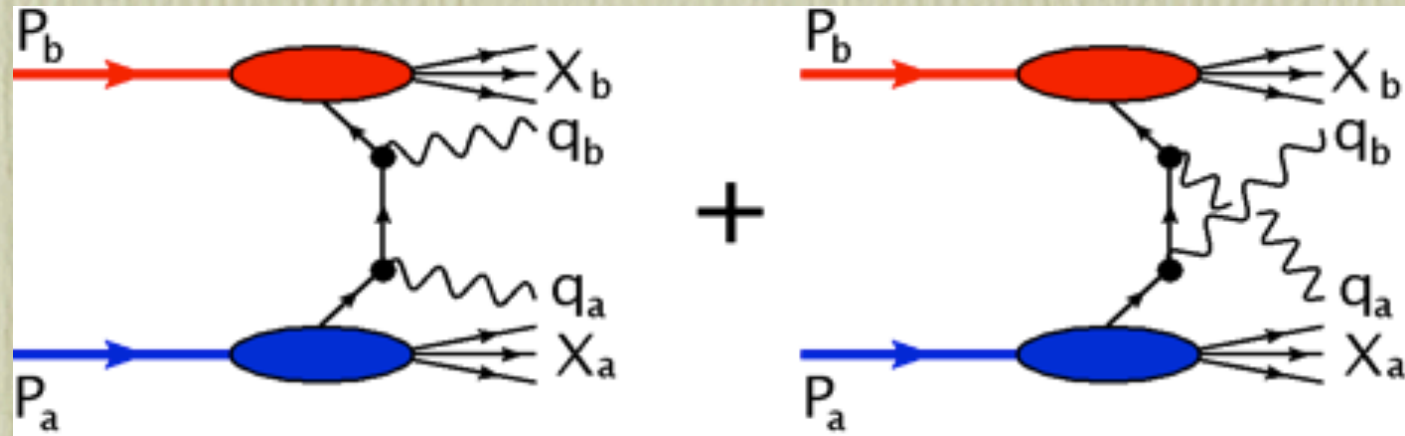
Spin dependent processes feasible at RHIC

colored final states: problems with TMD factorization

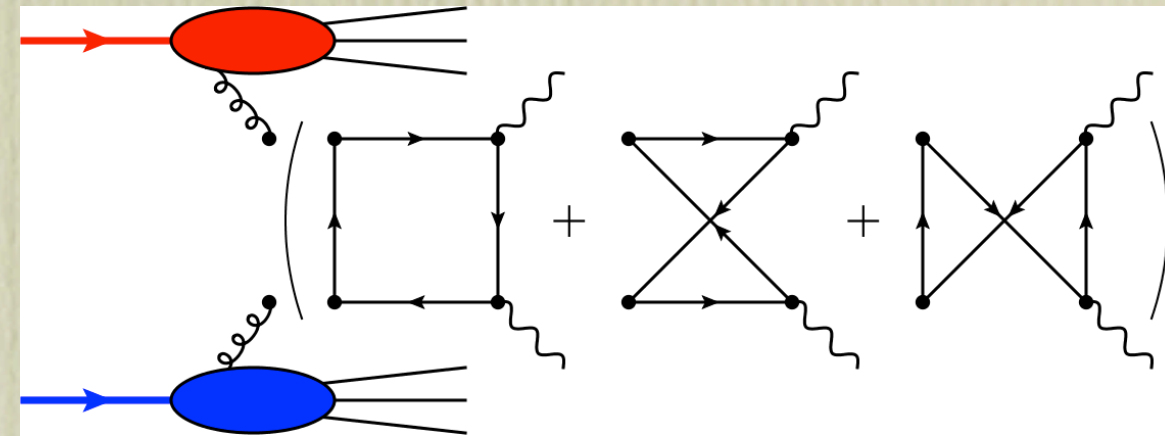
Photon pair production

[Qiu, M.S., Vogelsang, PRL 107, 062001 (2011)]

quark TMDs



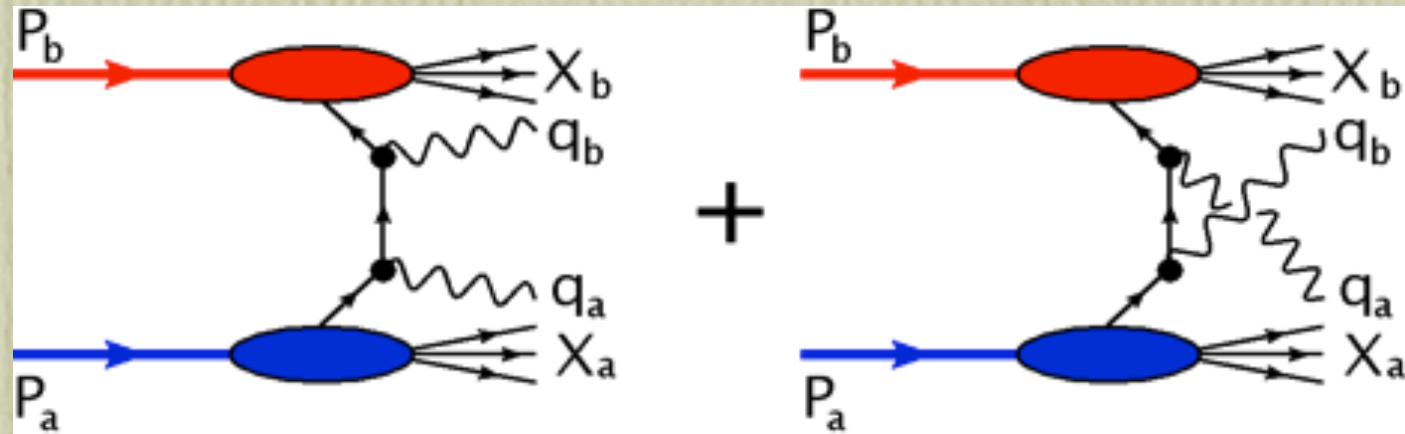
gluon TMDs at $O(\alpha_s^2)$



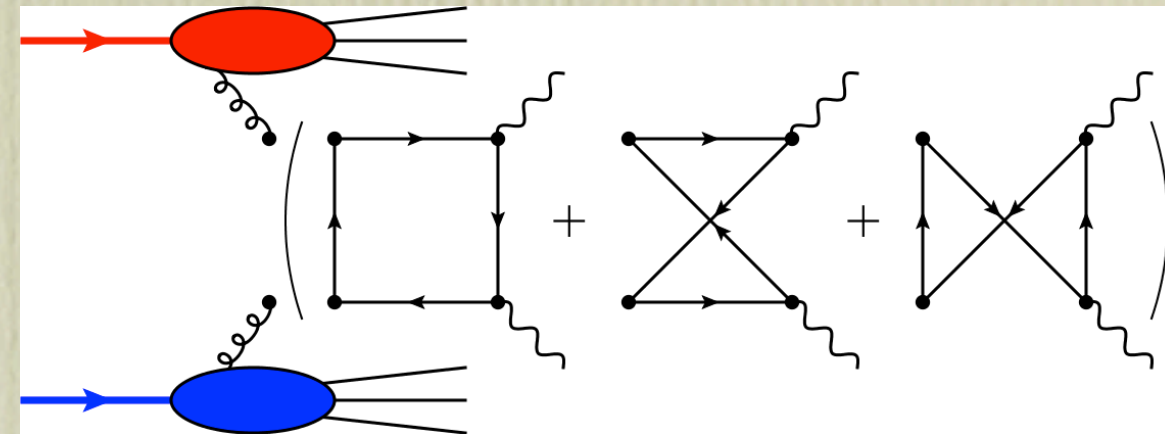
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quark TMDs



gluon TMDs at $O(\alpha_s^2)$



- * no colored final state \Rightarrow TMD factorization ok
- * gauge invariance \Rightarrow box finite \Rightarrow effectively tree-level
- * potentially large gluon distributions
- * new azimuthal observables

Unpolarized $pp \rightarrow \gamma\gamma X$ Cross-Section at $q_T \ll Q$

$$\frac{d\sigma_{UU}}{d^4q d\Omega} \sim \left(\frac{2}{\sin^2 \theta} \right) \left((1 + \cos^2 \theta) [f_1^q \otimes f_1^{\bar{q}}] + \cos(2\phi) \sin(2\theta) [h_1^{\perp q} \otimes h_1^{\perp \bar{q}}] \right)$$

$$+ \left(\frac{\alpha_s}{2\pi} \right)^2 \left(\mathcal{F}_1 [f_1^g \otimes f_1^g] + \mathcal{F}_2 [h_1^{\perp g} \otimes h_1^{\perp g}] + \cos(2\phi) \mathcal{F}_3 [h_1^{\perp g} \otimes f_1^g + f_1^g \otimes h_1^{\perp g}] + \cos(4\phi) \mathcal{F}_4 [h_1^{\perp g} \otimes h_1^{\perp g}] \right)$$

$\mathcal{F}_i \rightarrow$ non-trivial functions of $\cos(\theta)$ and $\sin(\theta)$

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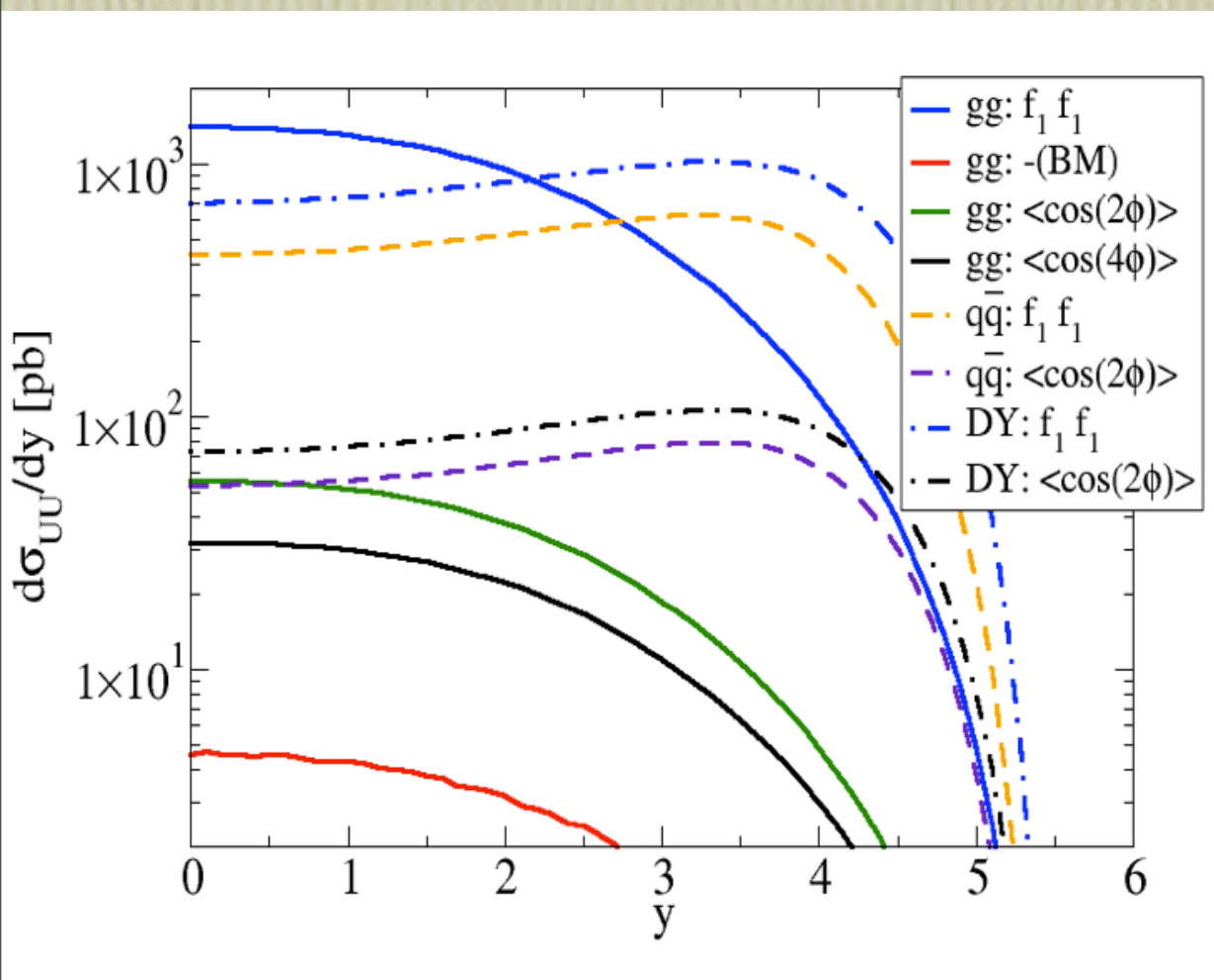
- Quark contribution similar to DY
→ only ISI / past-pointing Wilson line
- requires p_T / isolation cuts for the photons
- powerful in combination with DY
map out quark TMDs in DY → gluon TMDs in $\gamma\gamma$

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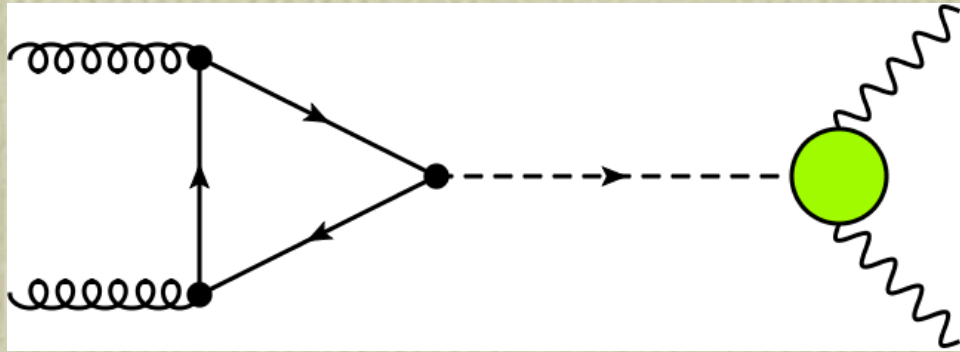
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map out quark TMDs in DY → gluon TMDs in $\gamma\gamma$
- numerical estimates →
Gaussian Ansatz + saturation of positivity bound
- $\cos(4\phi) \rightarrow$ only due to lin. pol. gluons → clean, ~1%
- $\cos(2\phi) \rightarrow$ determination of sign of $h_{1\perp}^g$

Linearly polarized gluons and Higgs production

[Boer, den Dunnen, Pisano, M.S., Vogelsang, arXiv:1109.1444]



ϕ - integrated cross section of Higgs + box:

$$\int d\phi \frac{d\sigma^{gg}}{d^4q d\Omega} \propto \bar{\mathcal{F}}_1 [f_1^g \otimes f_1^g] + \bar{\mathcal{F}}_2 [h_1^{\perp g} \otimes h_1^{\perp g}]$$

$$Q \neq m_H: \quad \bar{\mathcal{F}}_1 \gg \bar{\mathcal{F}}_2$$

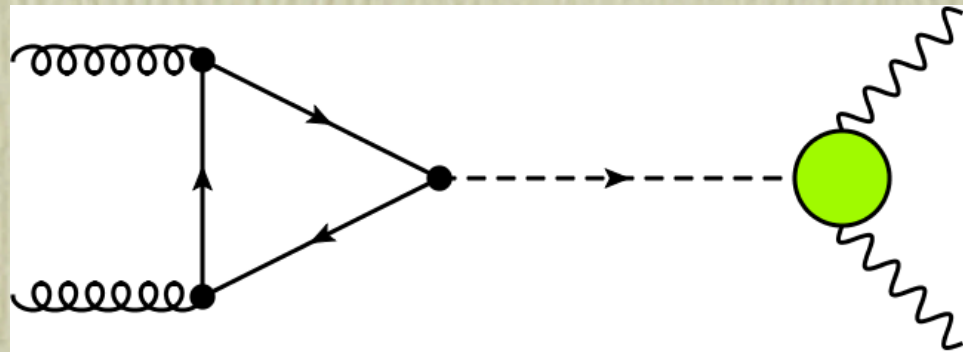
box dominant

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Higgs dominant (pole of the propagator)

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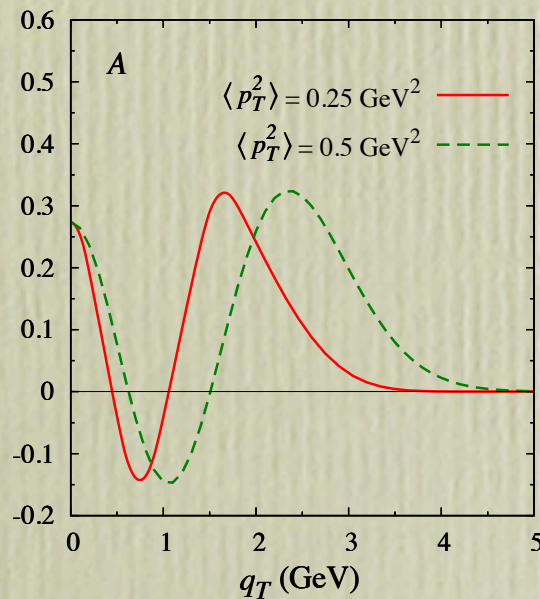
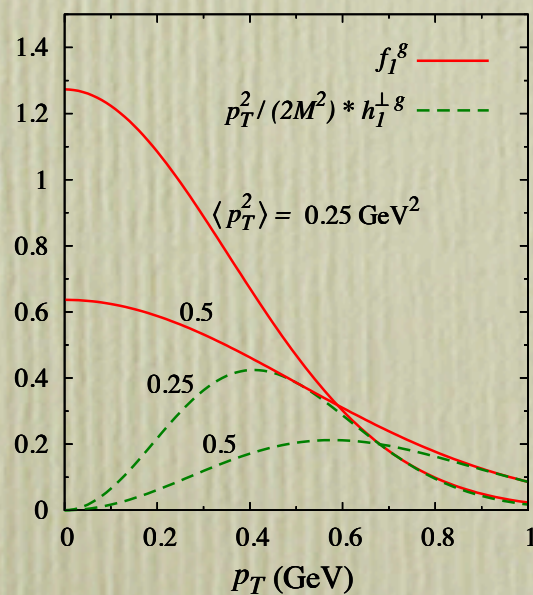
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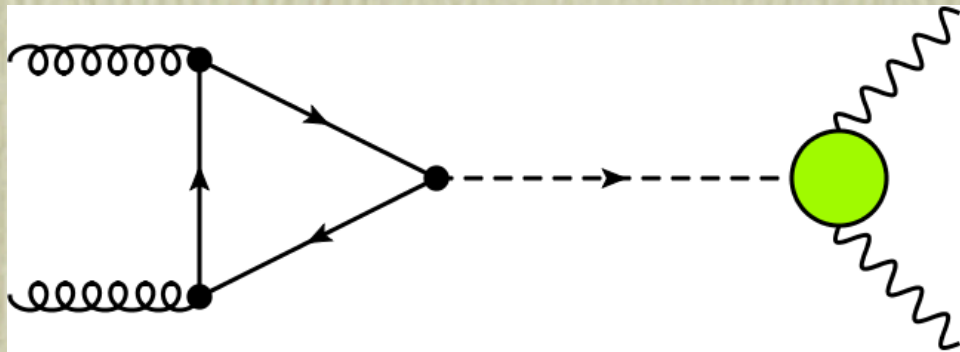
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- sizeable contribution of lin. pol. gluons
- characteristic double node in q_T

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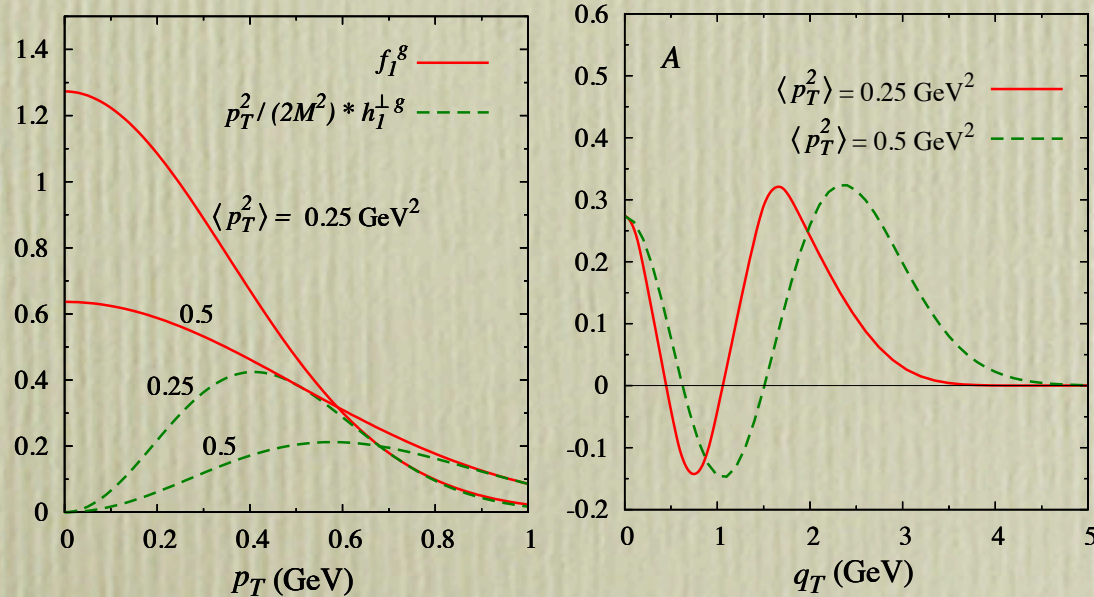
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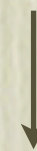
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linearly polarized gluons sensitive to Higgs parity

$$[f_1^g \otimes f_1^g] \pm [h_1^{\perp g} \otimes h_1^{\perp g}]$$

+: scalar Higgs -: pseudoscalar Higgs



precise q_T measurement may offer a way to determine Higgs parity

- sizeable contribution of lin. pol. gluons
- characteristic double node in q_T

Gluon Sivers Effect

(Transverse) Spin dependent photon pair cross section:

$$\frac{d\sigma_{TU}}{d^4q d\Omega} \sim S_T \sin \phi_S \left[\frac{2}{\sin^2 \theta} (1 + \cos^2 \theta) [f_{1T}^{\perp,g} \otimes f_1^{\bar{g}}] \right. \\ \left. + \left(\frac{\alpha_s}{2\pi} \right)^2 \left(\mathcal{F}_1 [f_{1T}^{\perp,g} \otimes f_1^g] + \mathcal{F}_2 [h_1^g \otimes h_1^{\perp,g}] + \mathcal{F}_2 [h_{1T}^{\perp,g} \otimes h_1^{\perp,g}] \right) \right] + \dots$$

Gluon Sivers Effect

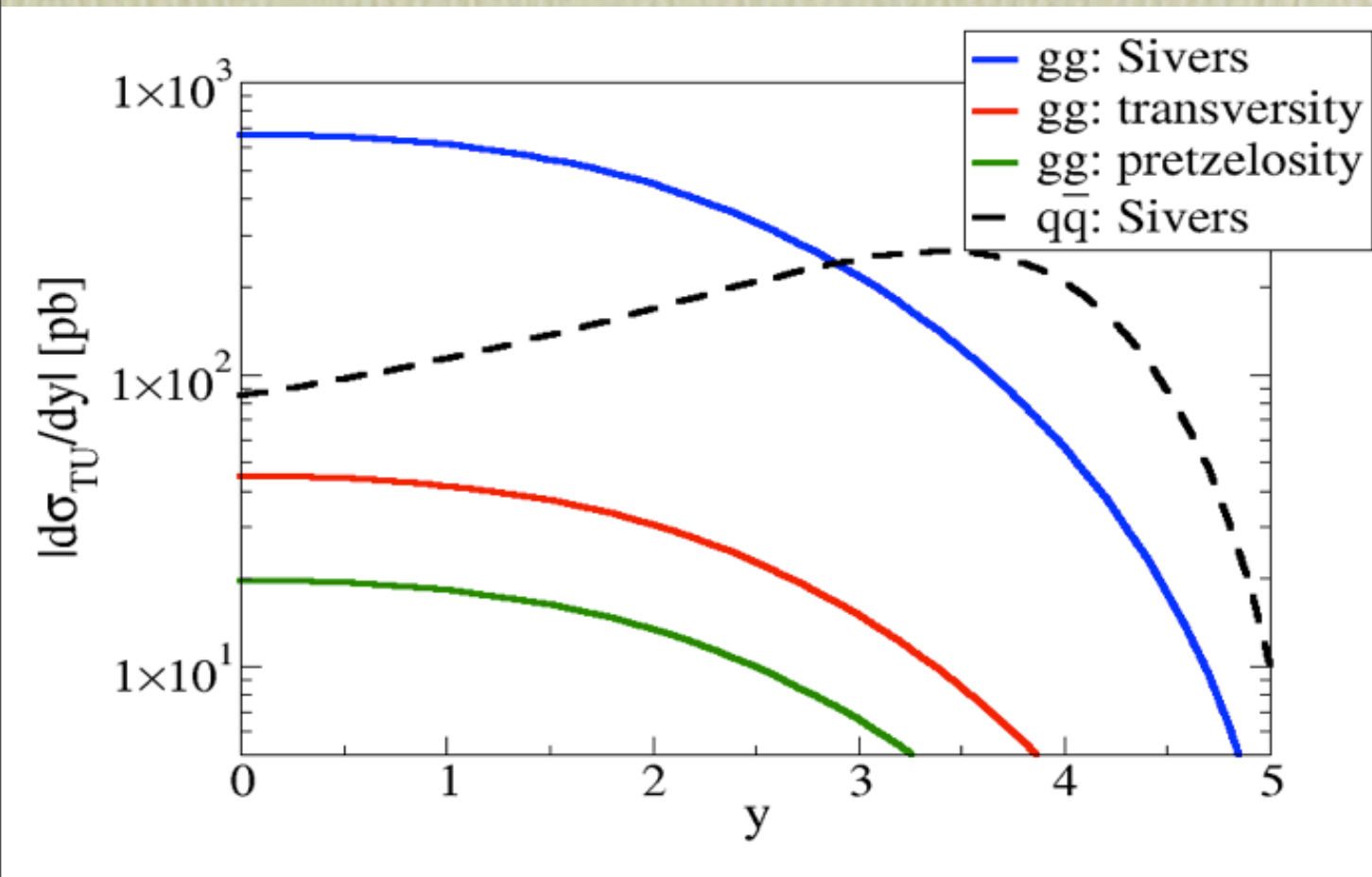
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Estimates for RHIC 500 GeV

- Gaussian Ansatz + positivity bound for gluon and quark TMDs
- Flavor cancellation for quark Sivers func.:

$$f_{1T}^{\perp,u} \simeq -f_{1T}^{\perp,d}$$
 → bound only for u-quarks
- Sign not fixed by bound
 → quark and gluon Sivers effect could add.
- Gluons dominate at mid-rapidity, quarks at large rapidity
- Effects by gluon “transv. / pretzel.” small

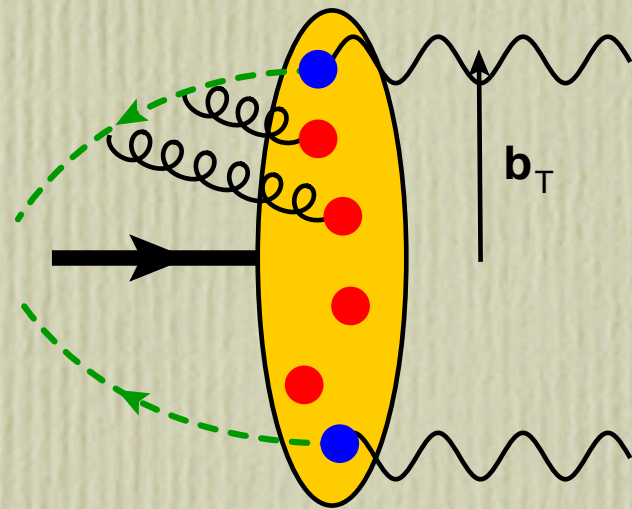


Summary

- T-odd TMDs generated by ISI/FSI \rightarrow sign switch
- Gluon TMDs from Photon Pair Production at RHIC
- Lin. polarized Gluons may offer a way to determine parity of the Higgs boson
- Gluon Sivers function may be feasible at RHIC (if $\neq 0$)

Intuitive picture of Sivers effect (M. Burkardt)

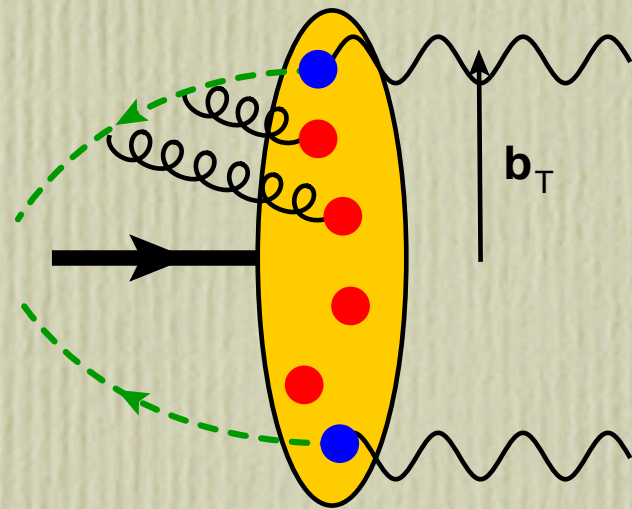
unpolarized



attr. FSI + eq. distributed quarks
= no net effect

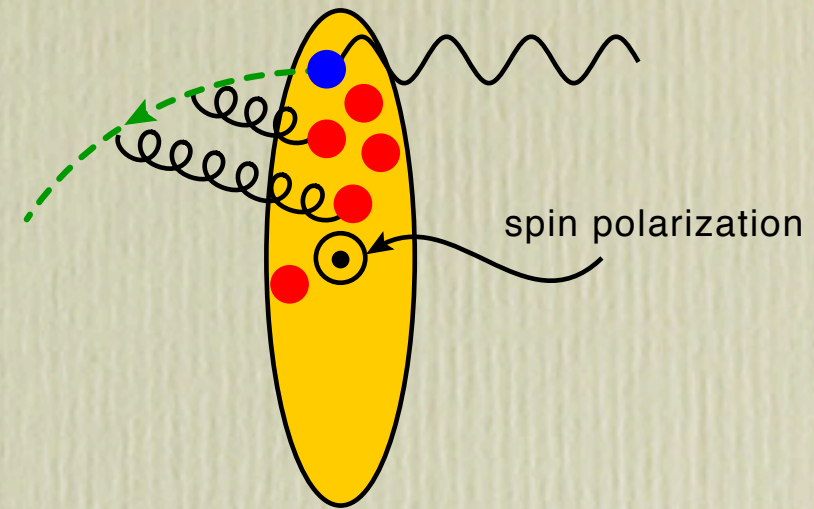
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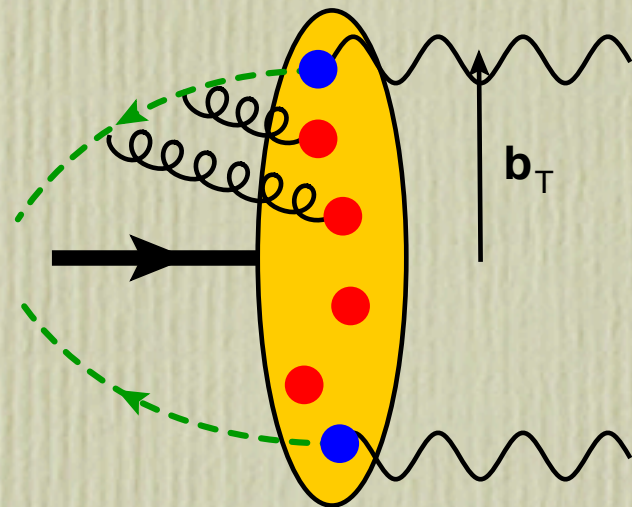
transversely polarized



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= Sivers effect

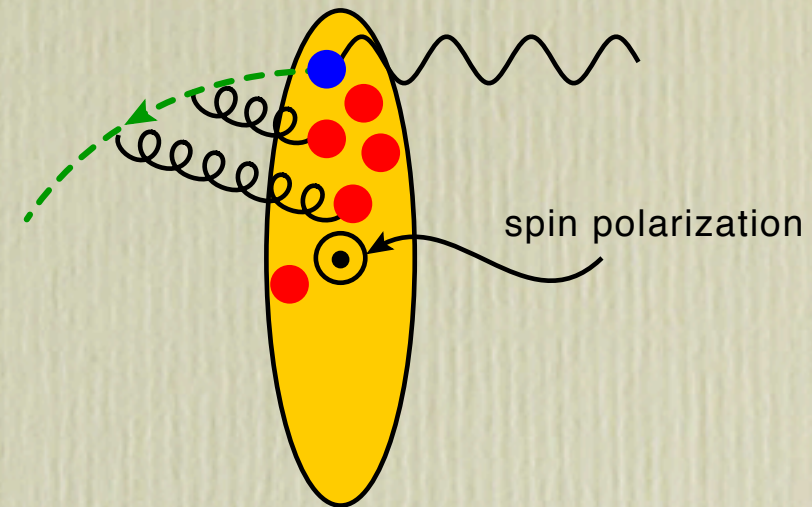
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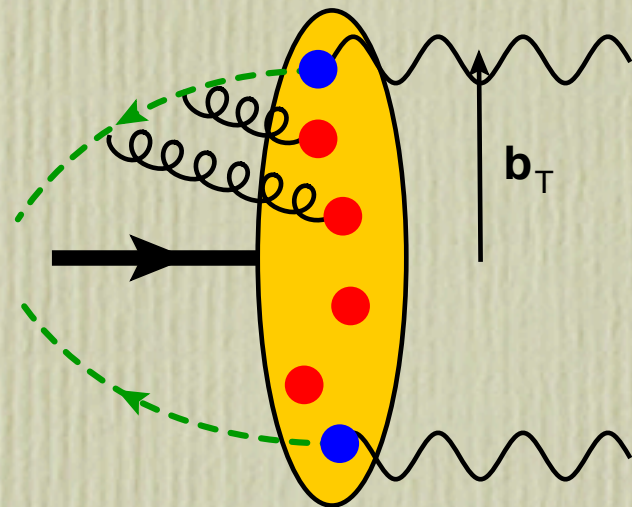
attr. FSI + distortion of distrib. (OAM)
= Sivers effect

$$\langle k_T^i \rangle = -M \epsilon_T^{ij} S_T^j f_{1T}^{\perp, (1)}(x) \simeq \int d^2 b_T \mathcal{I}^i(x, \vec{b}_T) \frac{\vec{b}_T \times \vec{S}_T}{M} \frac{\partial}{\partial b_T^2} \mathcal{E}(x, \vec{b}_T^2)$$

Not a model-independent relation (but valid in certain models...)

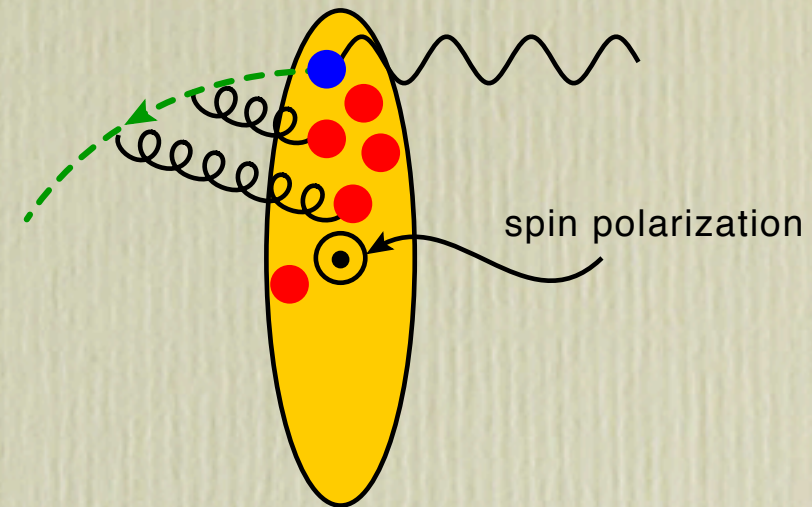
Intuitive picture of Sivers effect (M. Burkardt)

unpolarized



attr. FSI + eq. distributed quarks
= no net effect

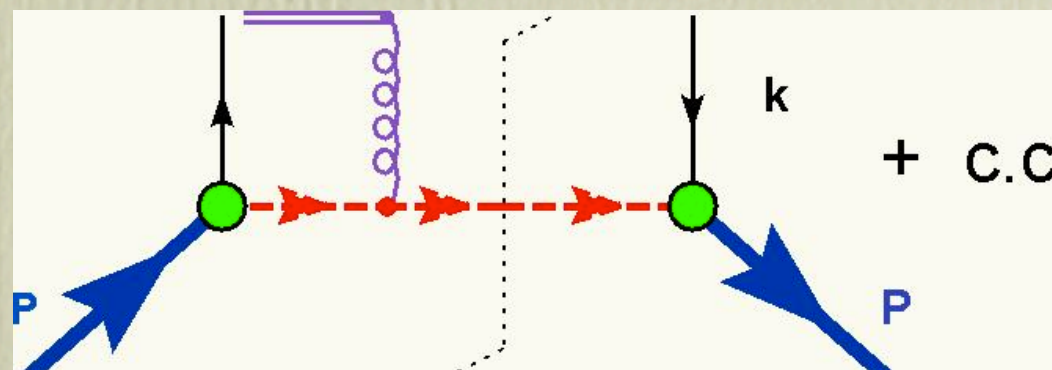
transversely polarized



attr. FSI + distortion of distrib. (OAM)
= Sivers effect

$$\langle k_T^i \rangle = -M \epsilon_T^{ij} S_T^j f_{1T}^{\perp, (1)}(x) \simeq \int d^2 b_T \mathcal{I}^i(x, \vec{b}_T) \frac{\vec{b}_T \times \vec{S}_T}{M} \frac{\partial}{\partial b_T^2} \mathcal{E}(x, \vec{b}_T^2)$$

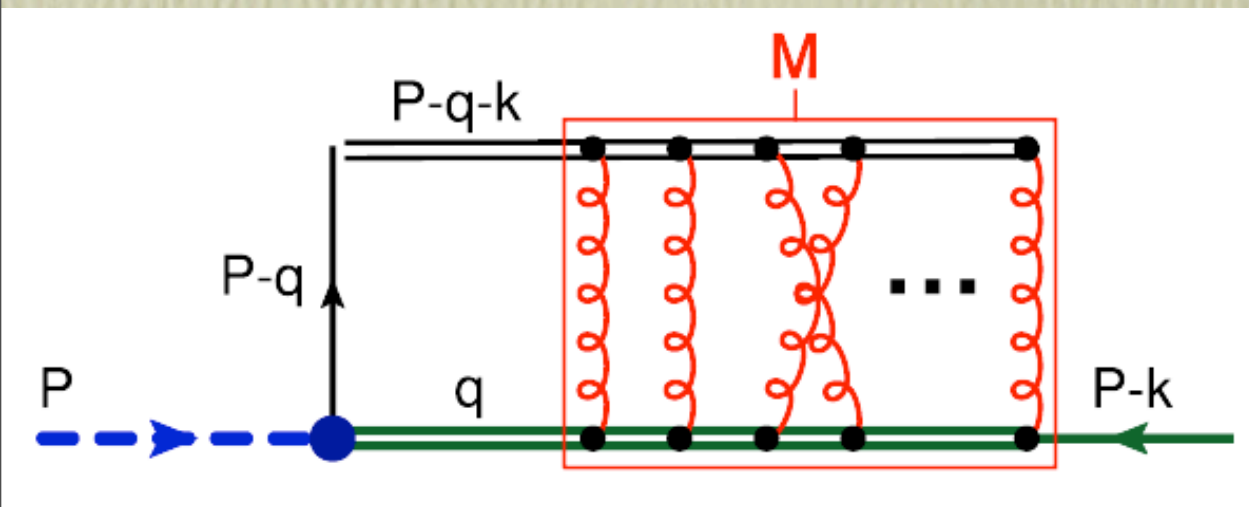
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→ Lensing function calculable

Beyond the one gluon exchange...

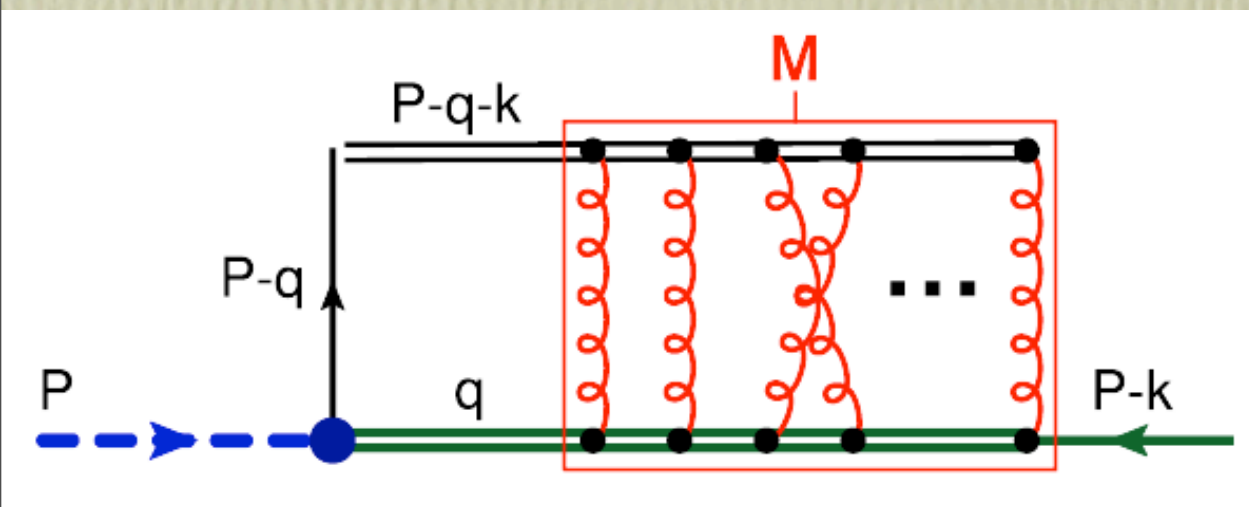
[Gamberg, M.S., PLB 685,95; arXiv:1012.3395; see also Courtoy, Scopetta, Vento, EPJA 47,49]



- Use spectator model to identify Lensing function
- Sum up soft gluons using non-perturbative eikonal methods
- Use input from Dyson-Schwinger eqs. for $SU(3)$ running coupling and non-pert. gluon propagator
→ **Lensing Function without free parameters!**
- Impact parameter GPD from models, parameterization

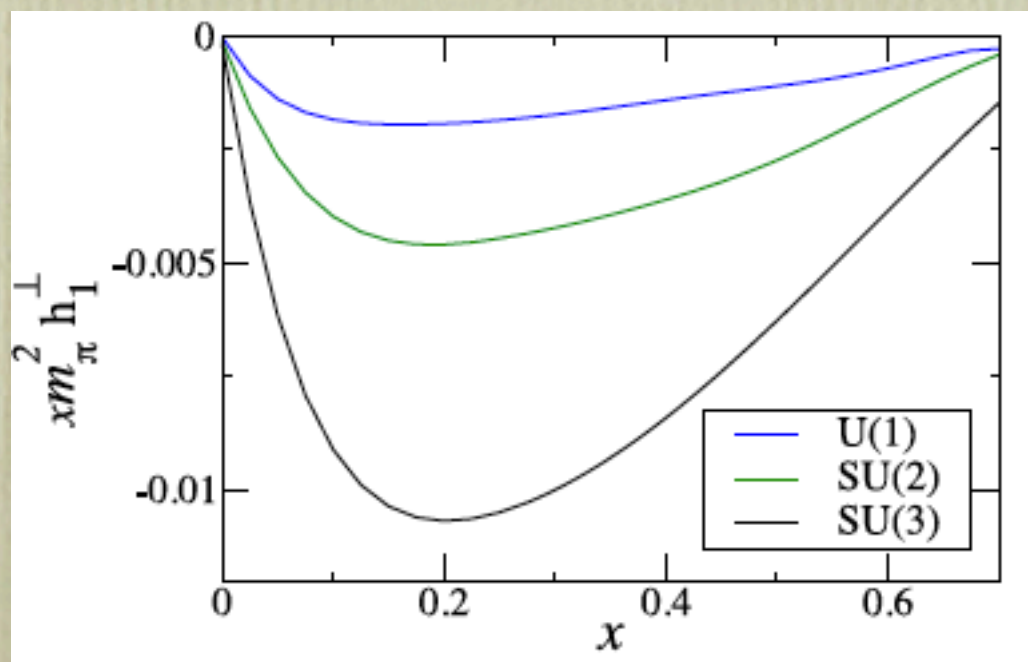
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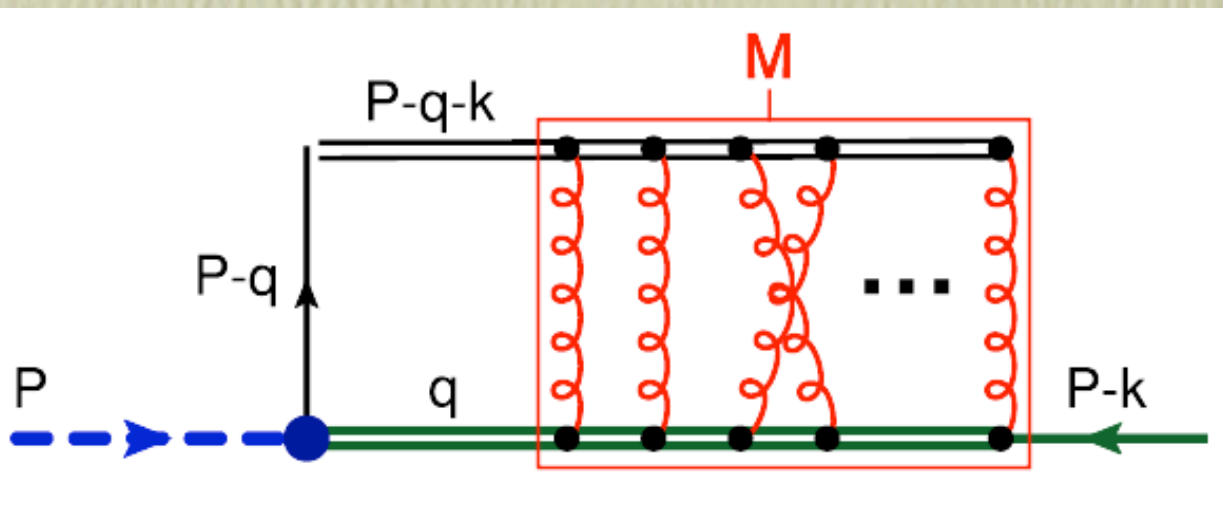
Pion Boer - Mulders



relevant for COMPASS DY

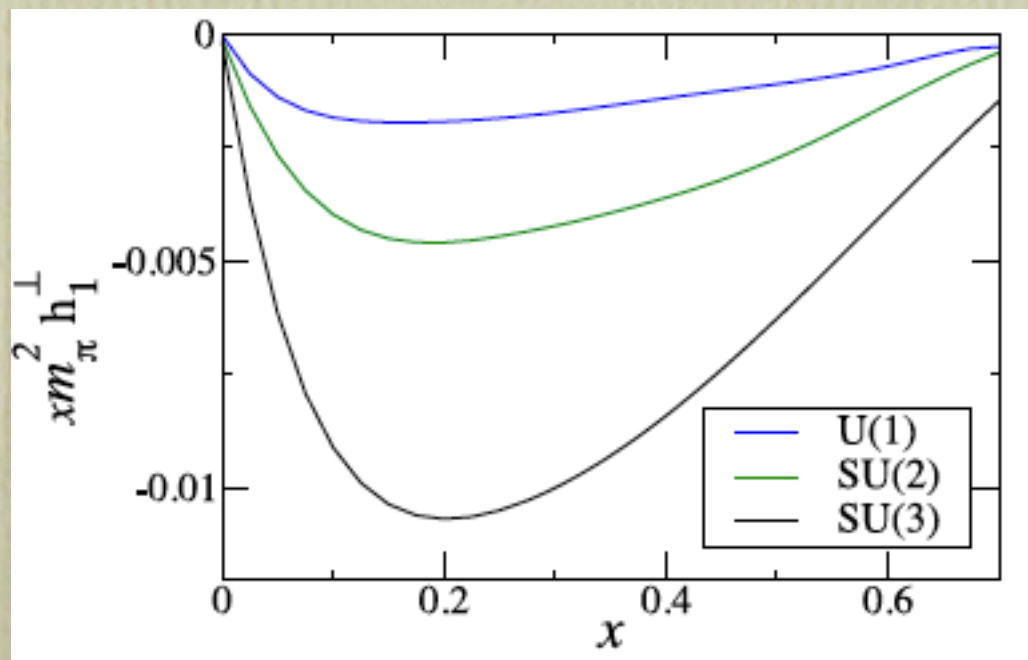
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relevant for COMPASS DY

Proton Sivers

