

The High Density Symmetry Energy in Relativistic Heavy Ion Collisions

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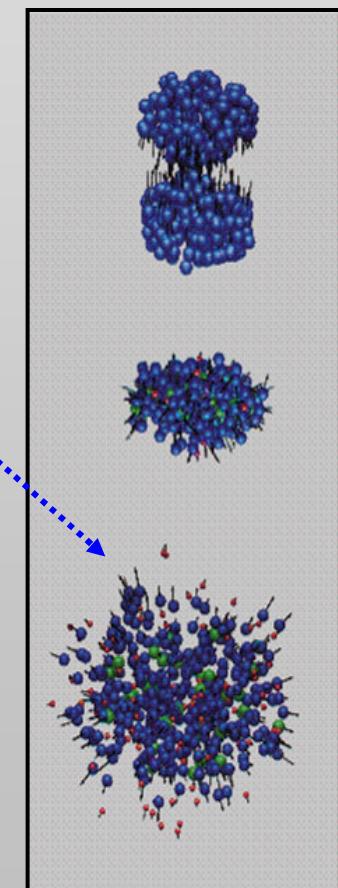
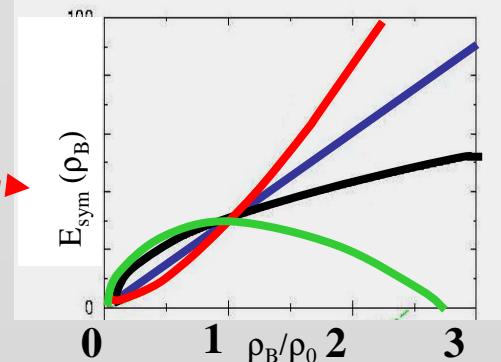
G. Ferini, M. Di Toro, V. Greco, Lab. Naz. del Sud, Catania

Hermann Wolter, LMU München



discuss:

1. Uncertainty of symmetry energy
2. test of symmetry energy in heavy ion collisions
 - low densities (Fermi energies) → talk by Massimo Di Toro
 - high density (rel. energies) → here
3. models for the symmetry energy
4. Transport calculations and search for observables sensitive to the symmetry energy:
 - proton-neutron differential flow,
 - meson production: ratios π^+/π^- and K^0/K^+
 - threshold and mean field effects, models for K-potential
5. Present status: high density symmetry energy almost not constrained by data,
but planned experiments should improve situation:
CHIMERA@GSI, RIKEN, FAIR (R^3B)

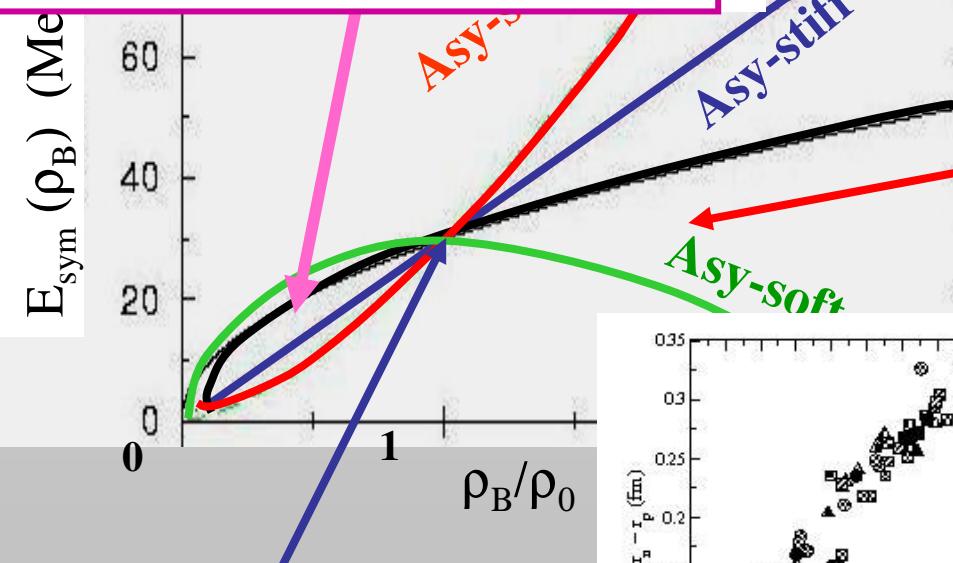
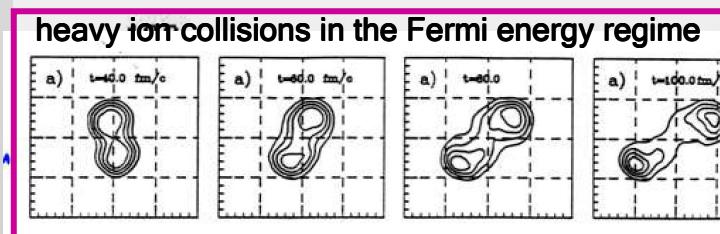


Symmetry Energy: Bethe-Weizsäcker Massenformel

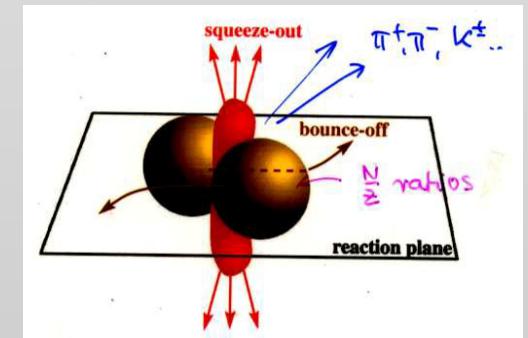
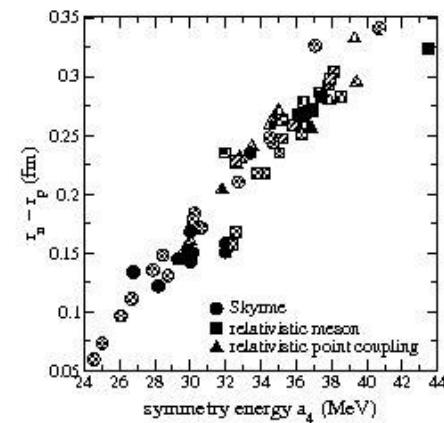
$$E(A, Z) = a_v A - a_s A^{2/3} - a_c Z(Z-1)A^{-1/3} - \text{circled term} - a_I (N-Z)^2 / A + \delta_{pair}$$

$$E(\rho_B, I)/A = E(\rho_B) + E_{sym}(\rho_B)I^2 + O(I^4) + \dots$$

$$I = \frac{N-Z}{N+Z}$$

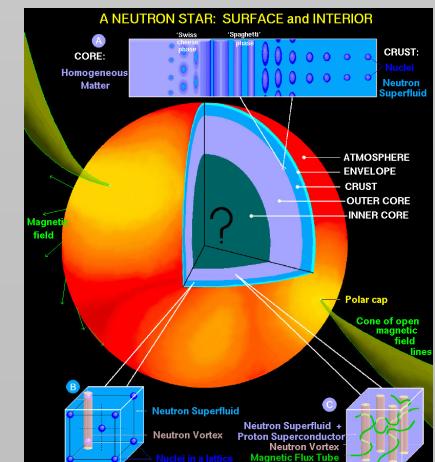


• Around normal density:
Structure, neutron skins



High density: HIC at relativistic energies (differential flow, particle production)

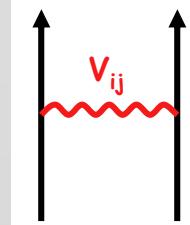
neutron stars



Theoretical Description of Nuclear Matter

Non-relativistic:

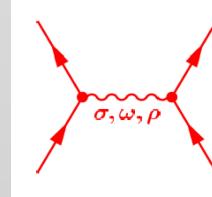
Hamiltonian $H = \sum T_i + \sum V_{ij}$; V nucleon-nucleon interaction



Relativistic:

Hadronic Lagrangian $L(\psi; \sigma, \omega, \pi, \eta, \delta, \dots)$

ψ , nucleon, resonances; $\sigma, \omega, \pi, \dots$ mesons



	phenomenological (fitted to nucl. matter)	microscopic (based on realistic NN interactions)
non-relativistic (Schrödinger)	Skyrme-type	Brueckner-HF (BHF) (+ 3-body forces)
Relativistic (Quantumhadrodyn.)	RMF	Dirac-Brueckner HF (DB)

Density functional
theory; EFT

Hadronic field theory → Quantenhadrodynamics (QHD)

Extensions of simplest $\sigma\omega$ -model necessary; many choices

$$L = \overline{\Psi} \left[i\gamma_\mu \left(\partial^\mu + ig_\omega \omega^\mu + ig_\rho \frac{\vec{\tau} \cdot \vec{b}^\mu}{2} \right) - \left(m - g_\sigma \sigma - g_\delta \frac{\vec{\tau} \cdot \vec{\delta}}{2} \right) \right] \Psi + L^{mes}$$

isovector mesons: symmetry energy

T. Gaitanos, et al., NPA732 (2004) 24

non-linear meson self-interactions

G. Lalazissis et al., PRC55 (1997) 540

density dependent coupling vertices

S. Typel, HHW, NPA656 (1999) 331

$$L^{mes} = \frac{1}{2} \left(\partial^\mu \sigma \partial_\mu \sigma - m_\sigma^2 \sigma^2 \right) - \frac{b_3}{3} \sigma^3 - \frac{b_4}{4} \sigma^4$$

$$\Gamma_\omega(\hat{\rho})$$

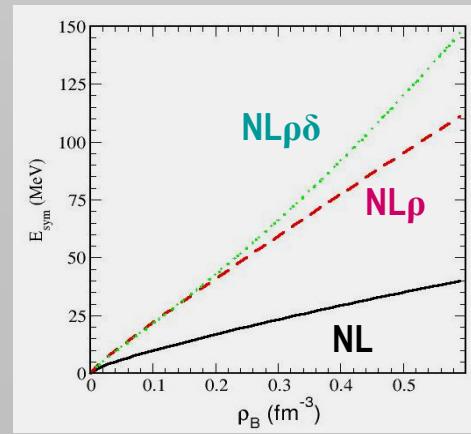
$$\Gamma_\rho(\hat{\rho})$$

...

Connection to DBHF

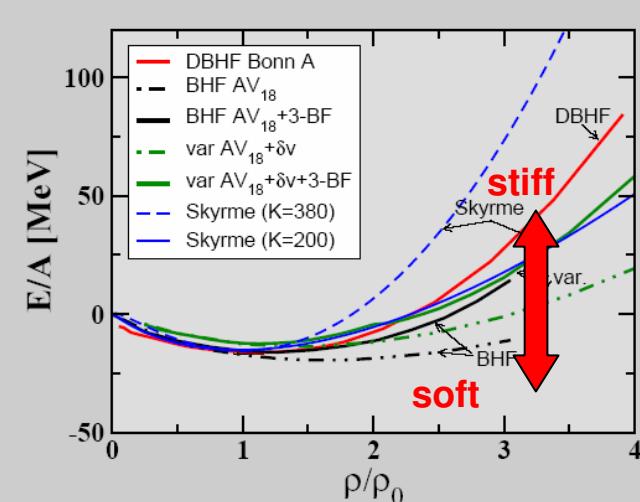
Full Lorentz structure:

	isoscalar	isovector
scalar	σ	δ
vector	ω	ρ
	↑ cancellation!	



$$E_{sym} = \frac{1}{6} \frac{k_F^2}{E_F^{*2}} + \frac{1}{2} \left[f_\rho - f_\delta \left(\frac{M^*}{E^*} \right)^2 \right] \rho_B$$

EOS for Symmetric Nuclear Matter



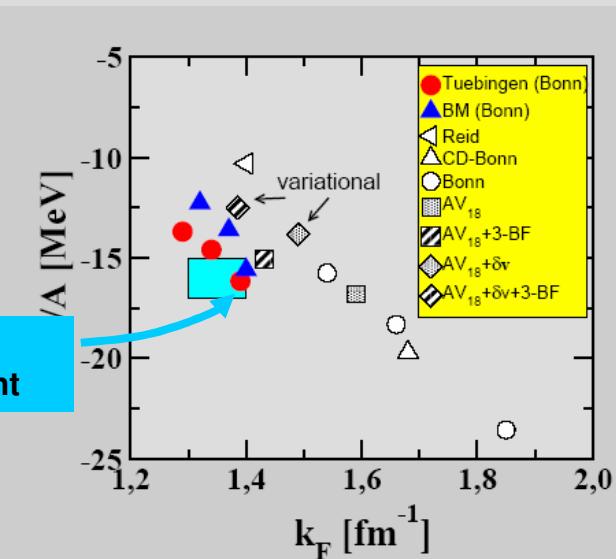
Chr. Fuchs, H.H. Wolter EPJ A30 (2006) 5

→ for BHF to describe saturation, 3-body forces are necessary

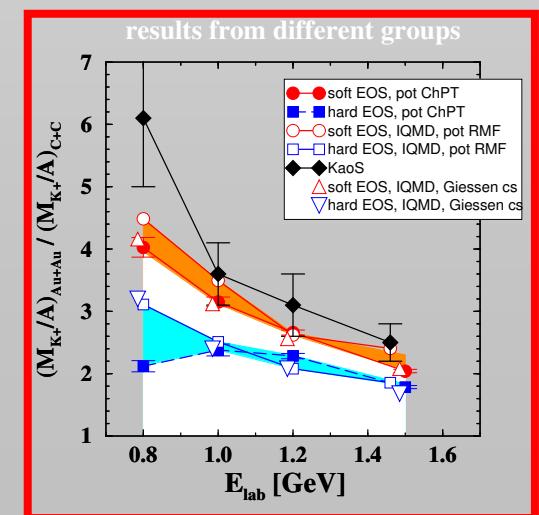
→ Microscopic EOS's are soft!

Fairly well fixed:
what about the symmetry energy?

Empirical
saturation point



Constraint
from ratio
of kaon
production
in heavy
and light
system



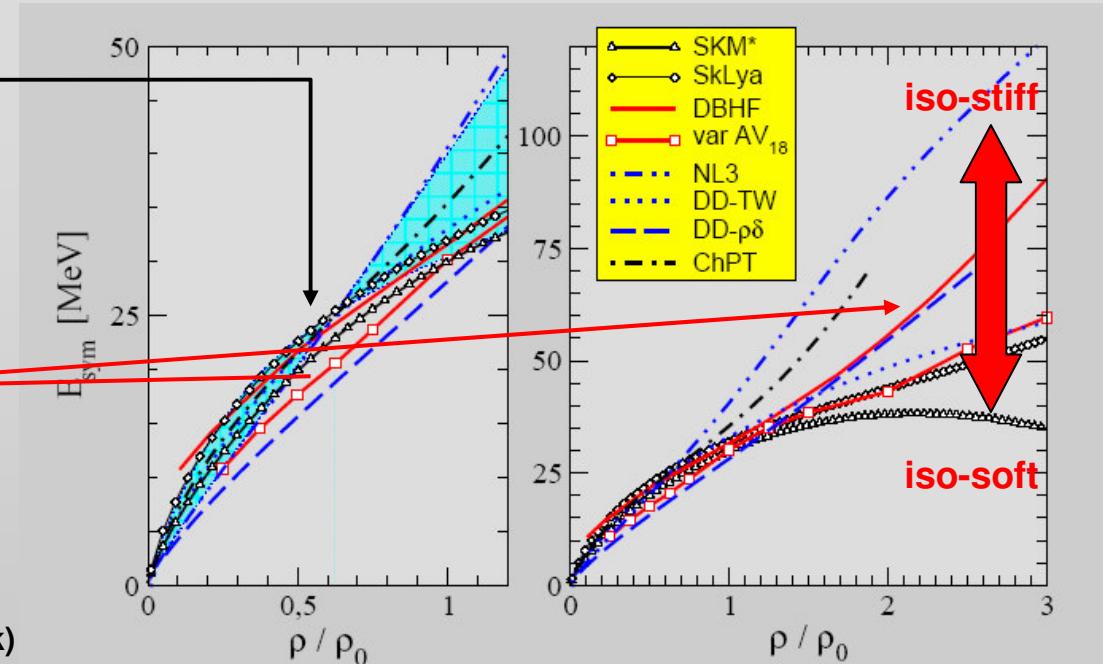
Chr. Fuchs, et al., PRL 86 (2001)

The Nuclear Symmetry Energy in different Models

empirical asy-EOS's cross at about

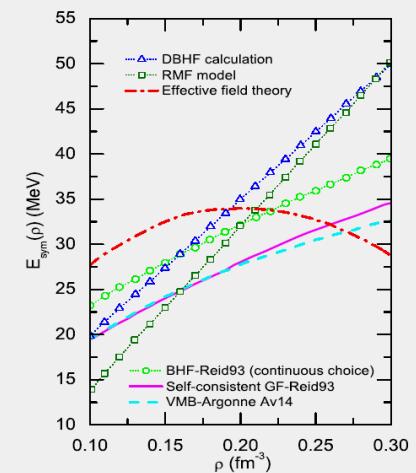
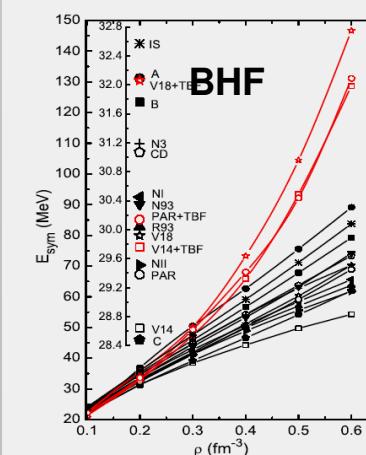
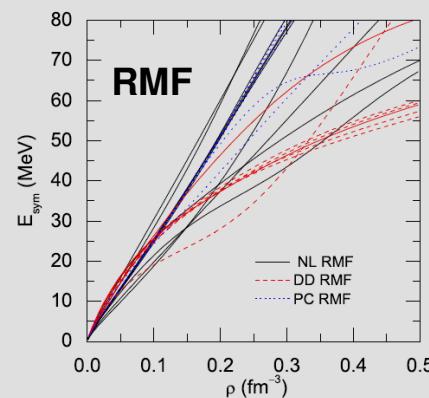
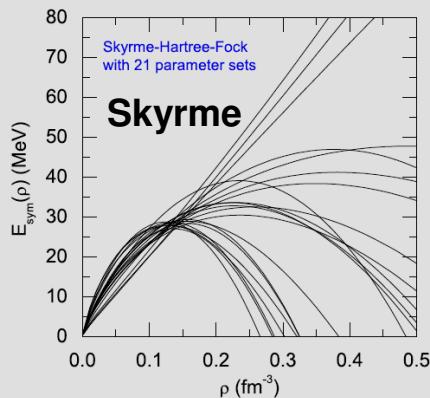
$$\rho \approx 0.6 \rho_0$$

microscopic asy-EOS's soft
at low densities but stiff at
high densities



C. Fuchs, H.H. Wolter, EPJA 30(2006)5,(WCI book)

many more in: B.A. Li, Phys. Rep. 08

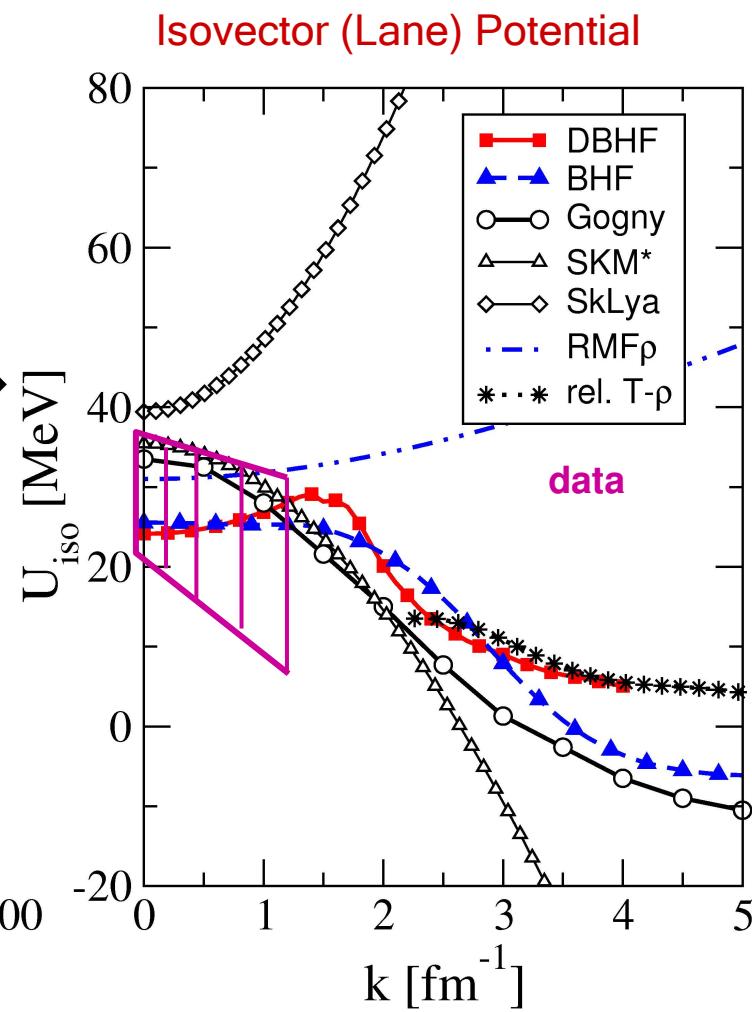
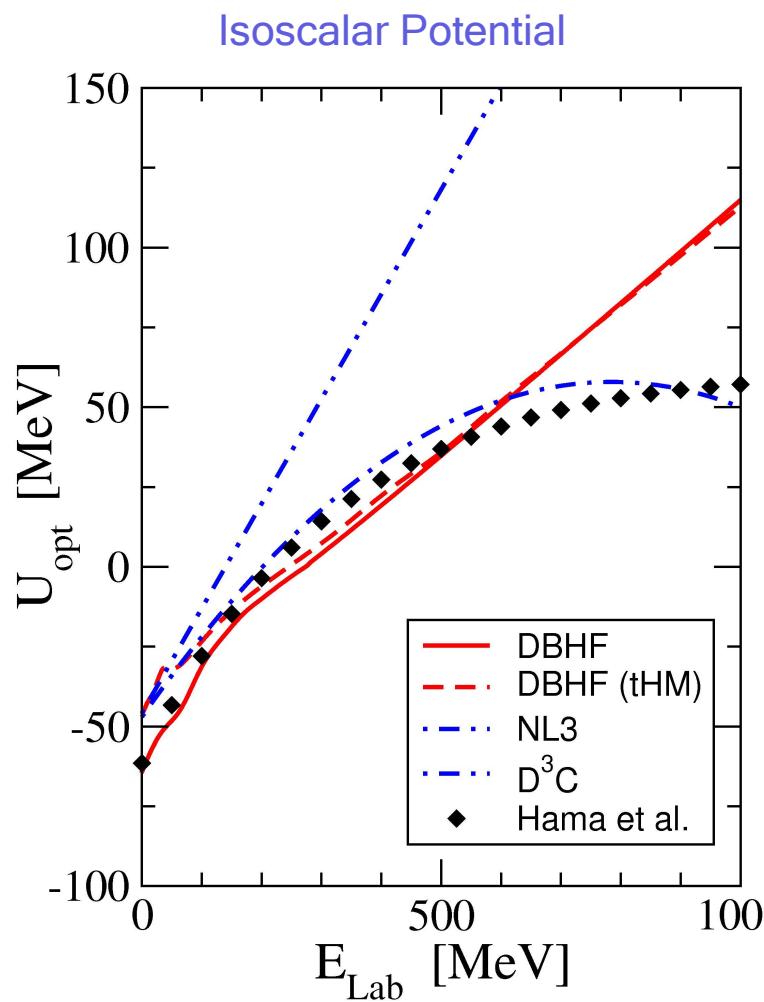


B.A. Li, Phys. Rep. (2008) in print

Momentum dependence of mean fields

$$U_{\text{opt}} = -\Sigma_S + \frac{E}{M} \Sigma_V + \frac{\Sigma_S^2 - \Sigma_V^2}{2M} .$$

Schrödinger Equivalent Optical Potential



Connected to splitting of
proton/neutron effective masses

Heavy Ion Collisions with particle production:

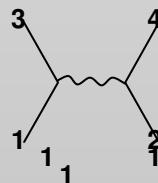
→ Coupled relativistic transport eqs. with elastic and inelastic collision terms

Rel. Transport eq. (RBUU, RLV, RQMD,..)

$$\left[\frac{p_i^{*\mu}}{m_i^*} \partial_\mu + \left(\frac{p_{vi}^*}{m_i^*} F_i^{\mu\nu} + \partial^\mu m_i^* \right) \partial_\mu^{(p^*)} \right] f_i(x, p^*) = I_i^{coll}$$

$i = p, n, \Delta^{(-,0,+,\text{++})}, \pi^\pm (\text{free}), K^{(0,+)} (\text{see below})$

	isoscalar	isovector
Scalar	$\sigma(\text{NL})$	δ
vector	ω	ρ
		cancellation!



$$I_i^{(coll)} = \sum_{j=\text{elast,inelast}} \left[\frac{1}{(2\pi)^{3/2}} \iiint dp_2 dp_3 dp_4 v_{12} \frac{d\sigma^{(ij)}}{d\Omega}_{12 \rightarrow 34} \delta(\text{energy conserv.}) \right] [(1-f)(1-f_2)f_3f_4 - f f_2(1-f_3)(1-f_4)]$$

Elastic baryon-baryon coll.: $\text{NN} \leftrightarrow \text{NN}$ (in-medium σ_{NN}), $\text{N}\Delta \leftrightarrow \text{N}\Delta$, $\Delta\Delta \leftrightarrow \Delta\Delta$

Inelastic baryon-baryon coll, (hard Δ -production): $\text{NN} \leftrightarrow \text{N}\Delta$, $\text{NN} \leftrightarrow \Delta\Delta$
 Inelastic baryon-meson coll. (soft Δ -production): $\text{N}\pi \leftrightarrow \Delta$

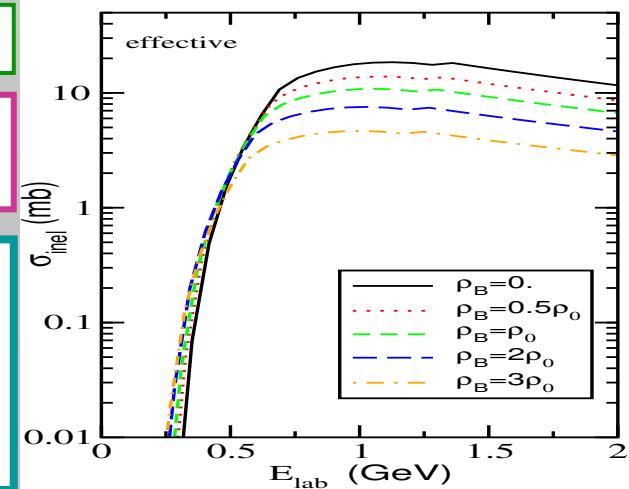
Channels with strangeness (perturbative kaon production):

Baryon-Baryon : $\text{BB} \rightarrow \text{BYK}$ ($\text{B} = \text{N}, \Delta^{\pm,0,\text{++}}$, $\text{Y} = \Lambda, \Sigma^{\pm,0}$, $\text{K} = \text{K}^{0,+}$)

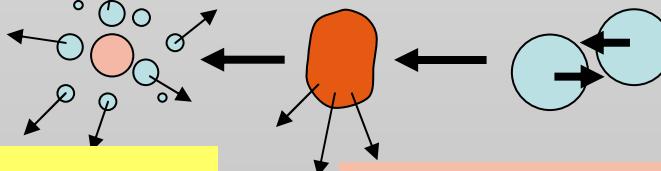
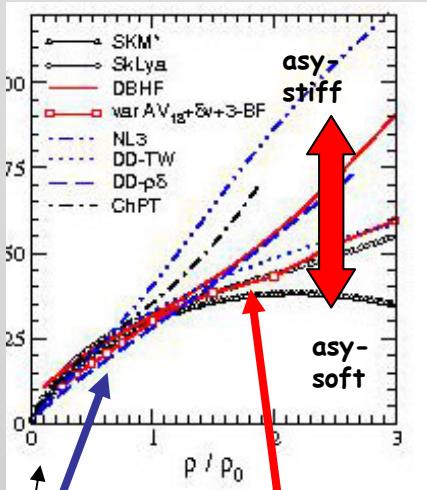
Pion-Baryon : $\pi\text{B} \rightarrow \text{YK}$

Kaon-Baryon : $\text{BK} \rightarrow \text{BK}$ (elastic, isospin exchange)

No channels with antistrangeness (K^-)



Probes of the ASY-EOS in Heavy Ion Collisions



Low density/energy

- fragments, p/n ratios
- isospin diffusion
- isoscaling
- migration/fractionat.
- collective excitations
- phase transitions

Covered
in talk by
M. Di Toro

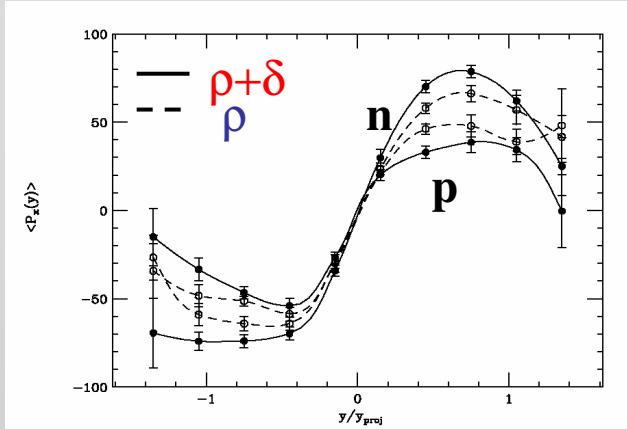
High density/energy

- n/p ratios
- p/n differential flow
- pions ratios
- kaon ratios
- neutron stars

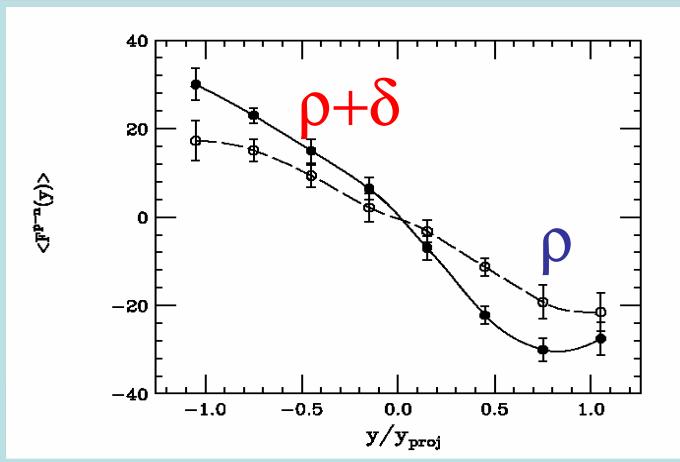
To be
discussed
(partially)
here

Asymmetric matter: Differential directed and elliptic flow

$^{132}\text{Sn} + ^{132}\text{Sn}$ @ 1.5 AGeV b=6fm



differential directed flow

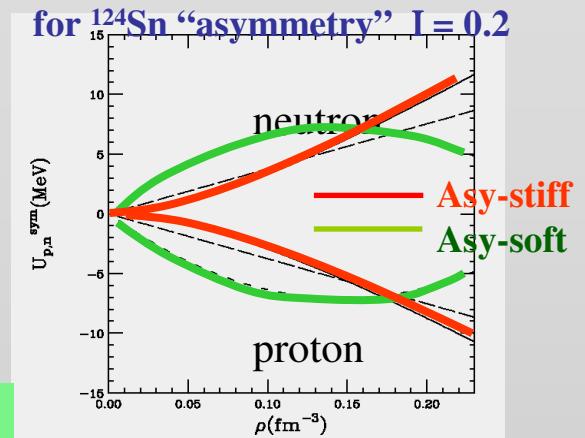


Proton-neutron differential flow

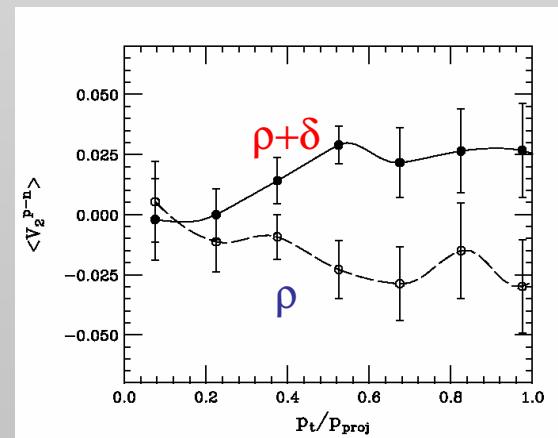
$$F_{n-p}^x(y) = \frac{1}{N(y)} \sum_{i=1}^{N(y)} (p_i^x w_i),$$

$w_i = +1(-1)$ for neutron (proton)

and analogously for elliptic flow



differential elliptic flow

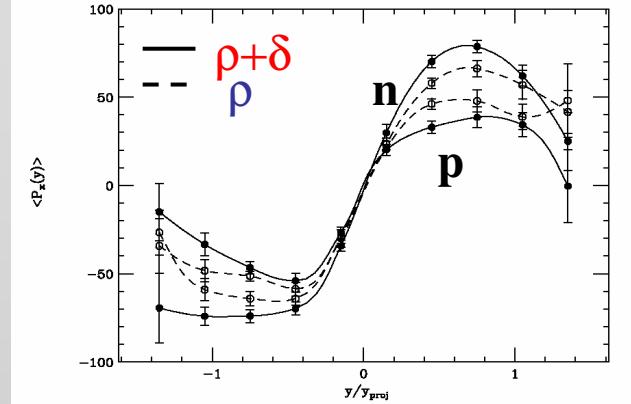


✿ Difference at high $p_t \leftrightarrow$ first stage , dynamical boosting of vector contribution

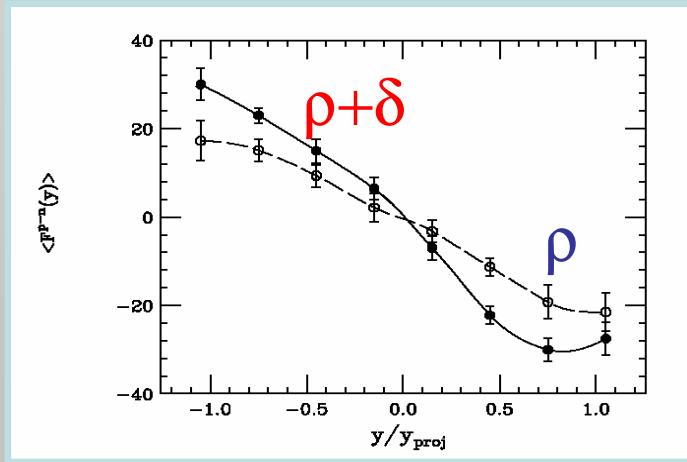
$$\frac{d\vec{p}_p^*}{d\tau} - \frac{d\vec{p}_n^*}{d\tau} \simeq 2 \left[\gamma f_\rho - \frac{f_\delta}{\gamma} \right] \vec{\nabla} \rho_3 = \frac{4}{\rho_B} E_{\text{sym}}^* \vec{\nabla} \rho_3$$

T. Gaitanos, M. Di Toro, et al., PLB562(2003)

Asymmetric matter: Differential directed and elliptic flow



differential directed flow

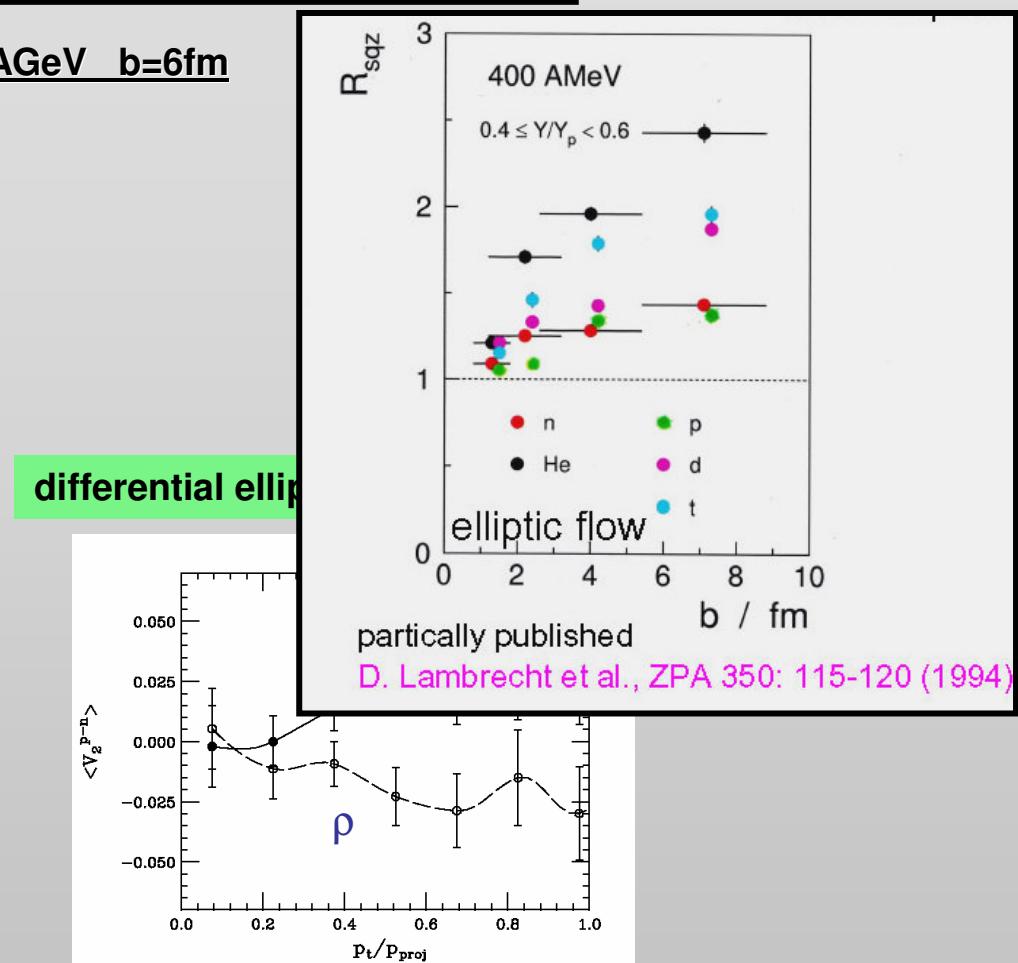


Proton-neutron differential flow

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$w_i = +1(-1)$ for neutron (proton)

and analogously for elliptic flow

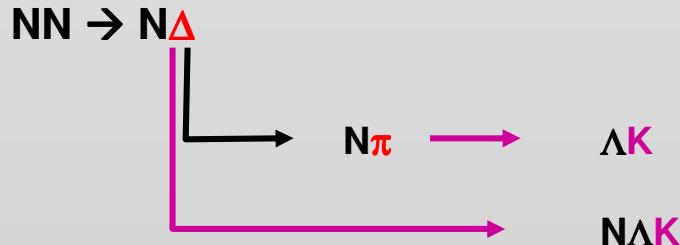


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T. Gaitanos, M. Di Toro, et al., PLB562(2003)

Sensitivity of particle production to the symmetry energy



consider only K^+ and K^0 , since antistrange kaons have a very different dynamics (talk by H. Oeschler)

1. Mean field effect: U_{sym} more repulsive for neutrons, and more for asystiff

→ pre-equilibrium emmission of neutron,
reduction of asymmetry of residue

2. Threshold effect, in medium effective masses:

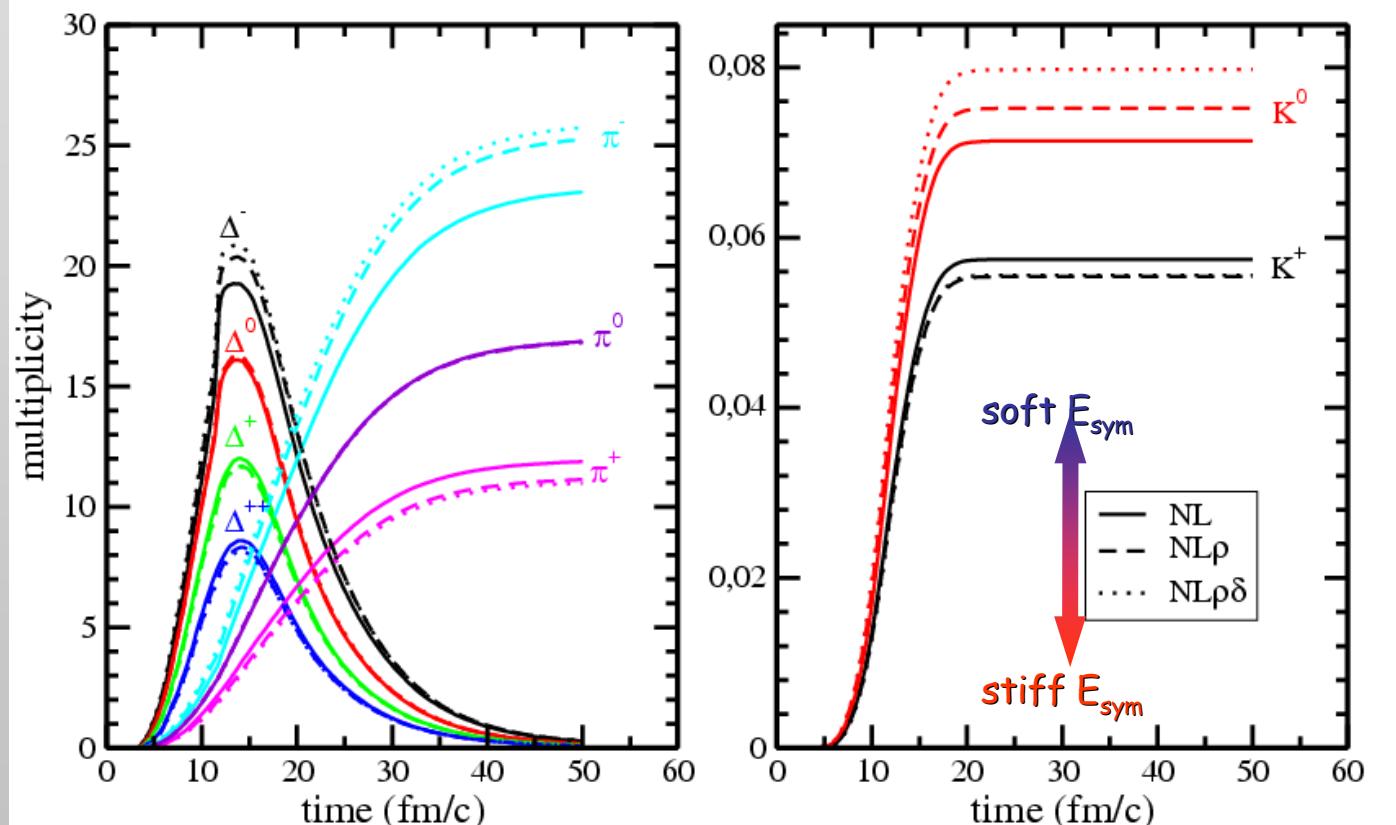
→ m^*_N , m^*_Δ , contribution of symmetry energy
 m^*_K , models for K-potentials

3. Discuss in particular ratios: π^-/π^+ and K^0/K^+ , to enhance effect of symmetry energy

Pion and Kaon production in “open” system (HIC)...



mean field effects: n/p repulsion
 threshold effects: Δ and K self energies



Pions: compensation

Kaons:

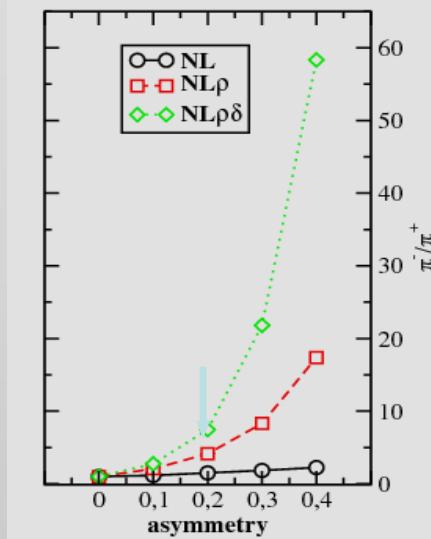
- direct early production: high density phase
- isovector channel effects

Pion production: Au+Au, semicentral

Equilibrium production
(box results)

$$\frac{\pi^-}{\pi^+}$$

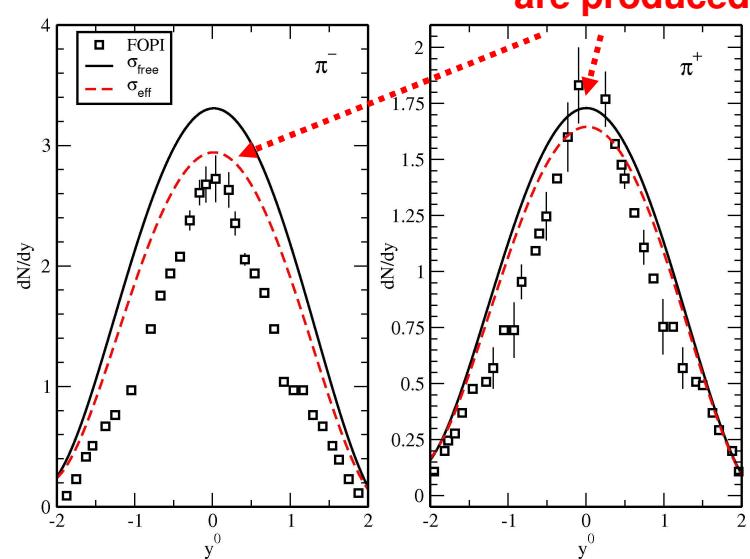
Ferini, NPA762(2005) 147



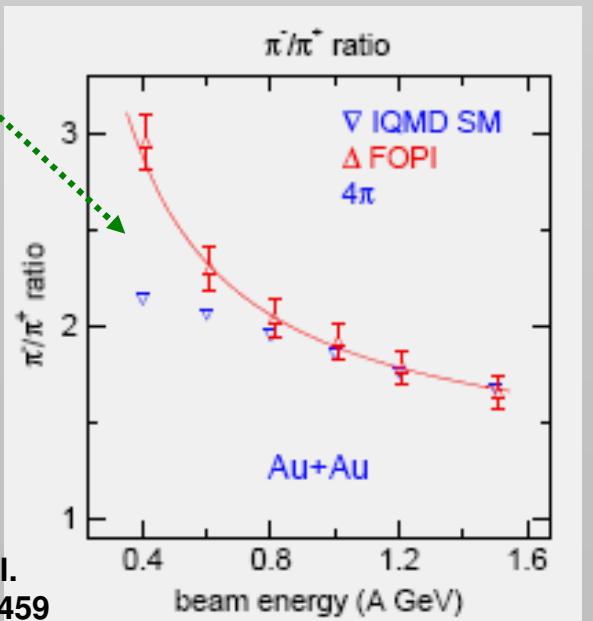
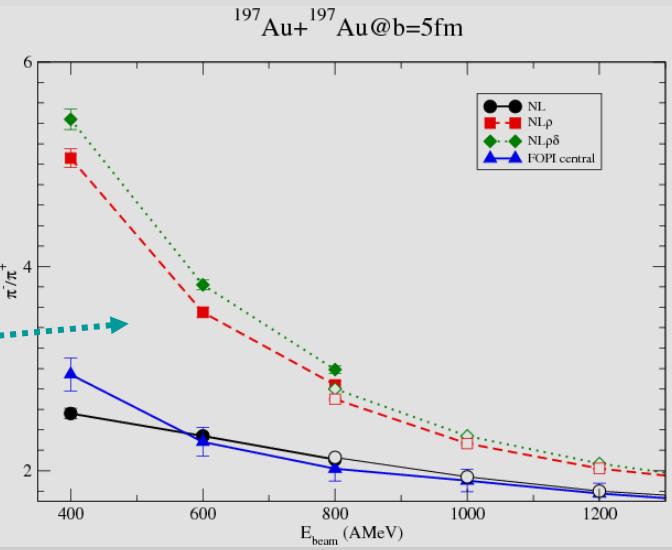
$$\frac{\pi^-}{\pi^+} = \frac{\sigma_{\pi^-}^{abs}}{\sigma_{\pi^+}^{abs}} \exp[2(\mu_n - \mu_p)/T]$$

~ 5 (NL ρ) to 10 (NL $\rho\delta$)

disagreement in magnitude,
particularly at low energies,
(also in other calc.),
but better at midrapidity
(high density), where Kaons
are produced.



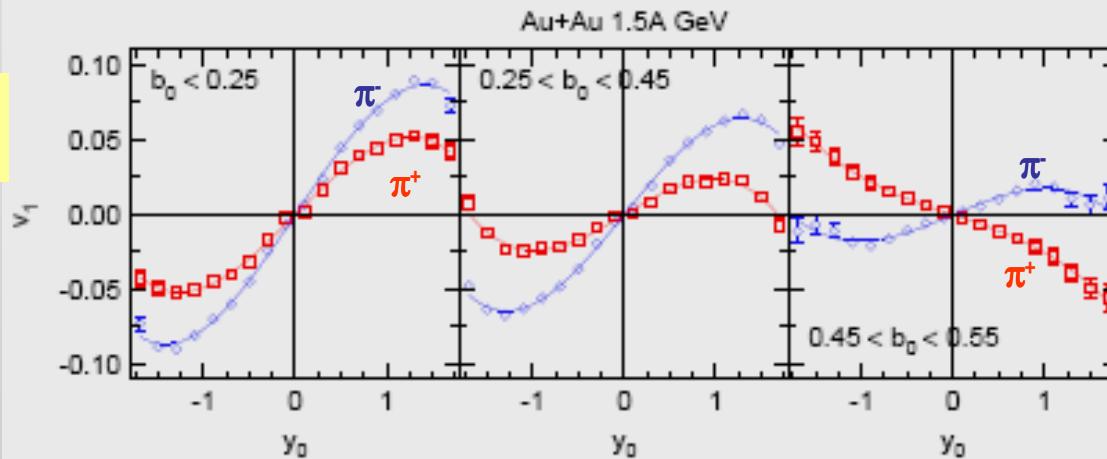
Finite nucleus simulation:



W.Reisdorf et al.
NPA781 (2007) 459

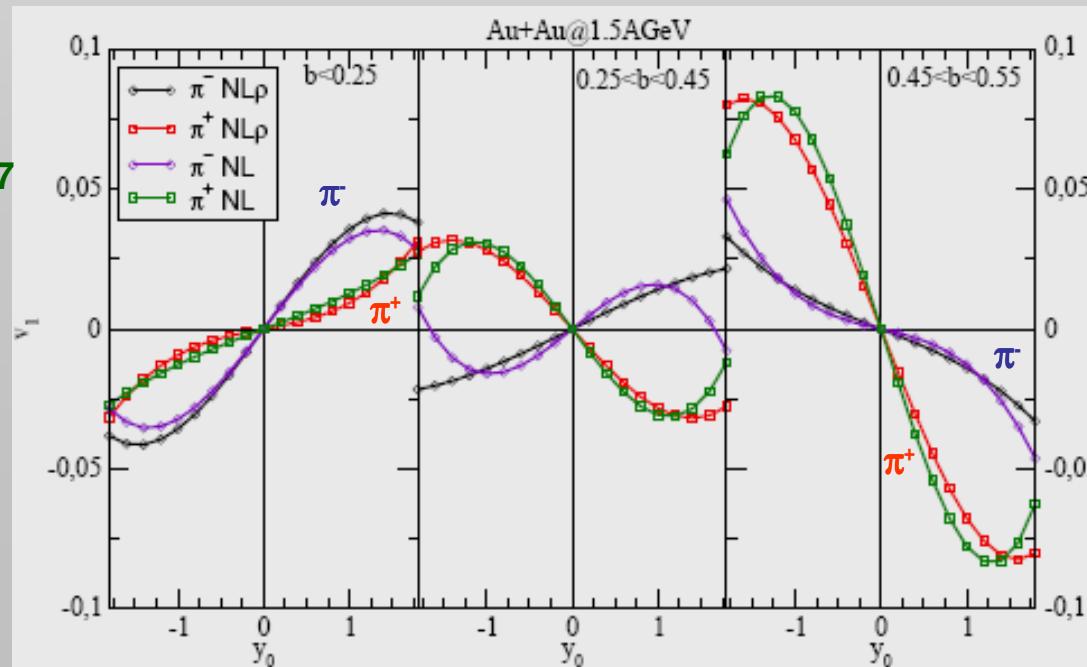
Transverse Pion Flows

W.Reisdorf et al.
NPA781 (2007) 459



Antiflow:
Decoupling of the
Pion/Nucleon flows

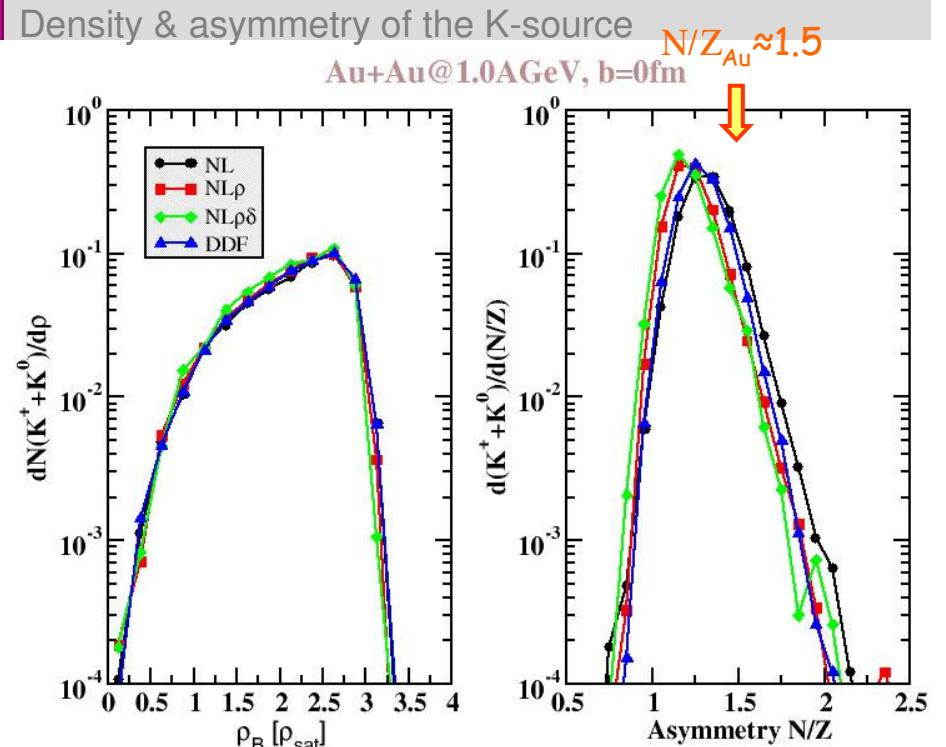
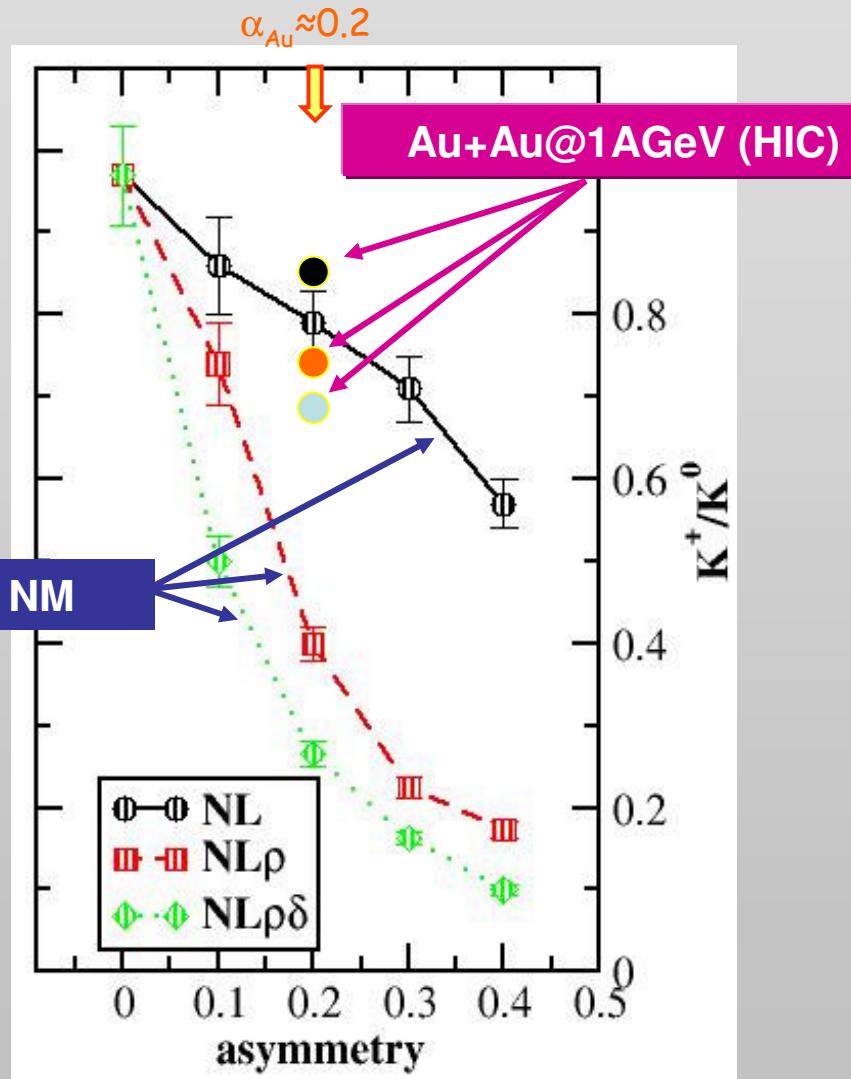
Simulations:
V.Prassa Sept.07



OK general trend.
but:
- smaller flow for
both
 π^- and π^+
- not much
dependent
on iso-EoS

Strangeness ratio : Infinite Nuclear Matter vs. HIC

G. Ferini, et al., NPA762(2005) 147 and nucl-th/0607005

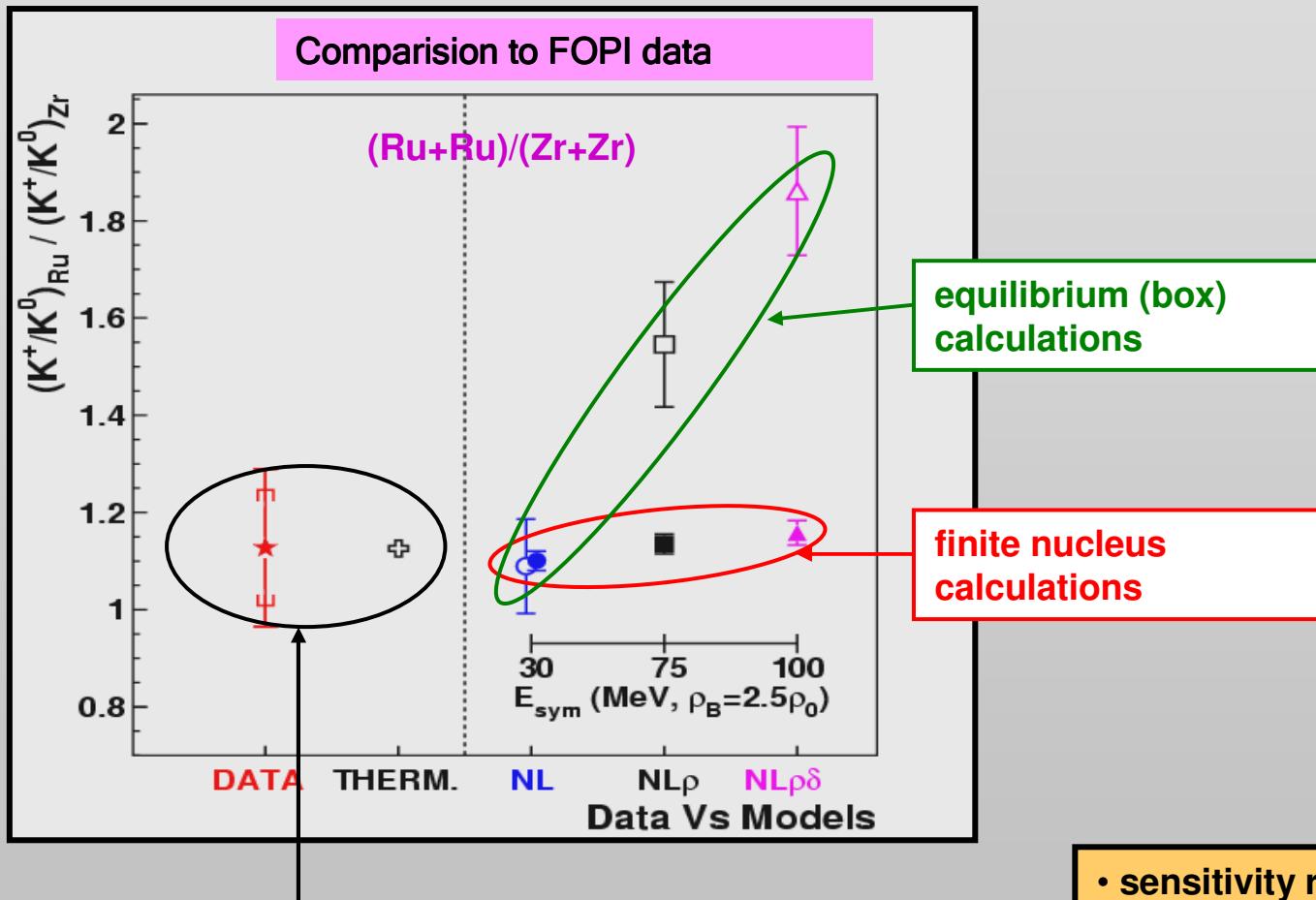


$NL \rightarrow DDF \rightarrow NL\rho \rightarrow NL\rho\delta$:
more neutron escape and more $n \rightarrow p$ transformation
(less asymmetry in the source)

Pre-equilibrium emission (mainly of neutrons) reduced asymmetry of source for kaon production → reduces sensitivity relative to equilibrium (box) calculation

Kaon ratios: comparison with experiment

G. Ferini, et al., NPA762(2005) 147 and nucl-th/0607005



Data (Fopi)

X. Lopez, et al. (FOPI), PRC 75
(2007)

- sensitivity reduced in collisions of finite nuclei
- single ratios more sensitive
- enhanced in larger systems

Effect of kaon potentials

In-medium Klein-Gordon eq. for Kaon propagation:

$$[(\partial_\mu + iV_\mu)^2 + m_K^{*2}] \phi_K(x) = 0$$

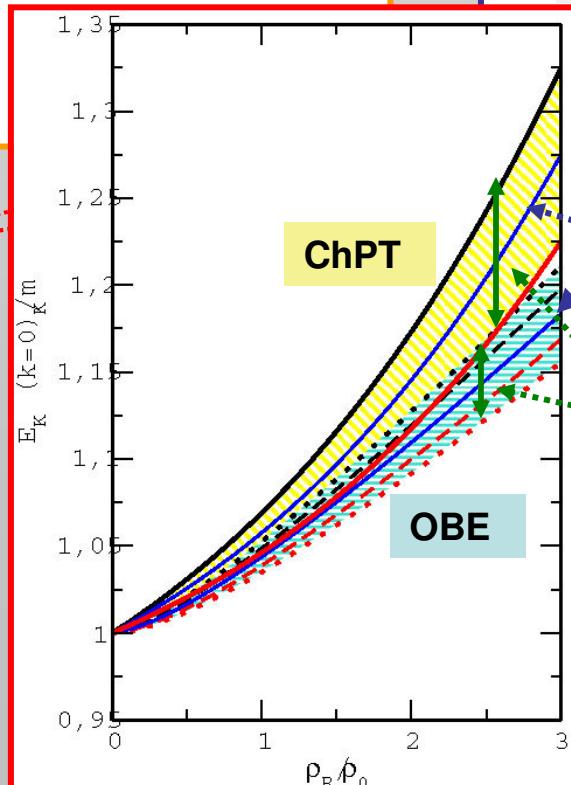
Two models for medium effects tested:

Chiral perturbation (Kaplan, Nelson et al.)

$$m_K^* = \sqrt{m_K^2 - \frac{\Sigma_{KN}}{f_\pi^2} \rho_s \mp \frac{C}{f_\pi^2} \rho_{s3} + V_\mu V^\mu} \quad (\text{upper sign, } K^+)$$

$$V_\mu = \frac{3}{8f_\pi^{*2}} j_\mu \pm \frac{1}{8f_\pi^{*2}} j_{\mu 3}$$

Isospin-dependence



One-Boson Exchange (Schaffner-Bielich et al.)

$$m_K^* = \sqrt{m_K^2 + \frac{m_K}{3} (g_{\sigma N} \sigma \mp f_\delta \rho_{s3})}$$

$$V^\mu = \frac{1}{3} (f_\omega^* j^\mu \pm f_\rho j_3^\mu)$$

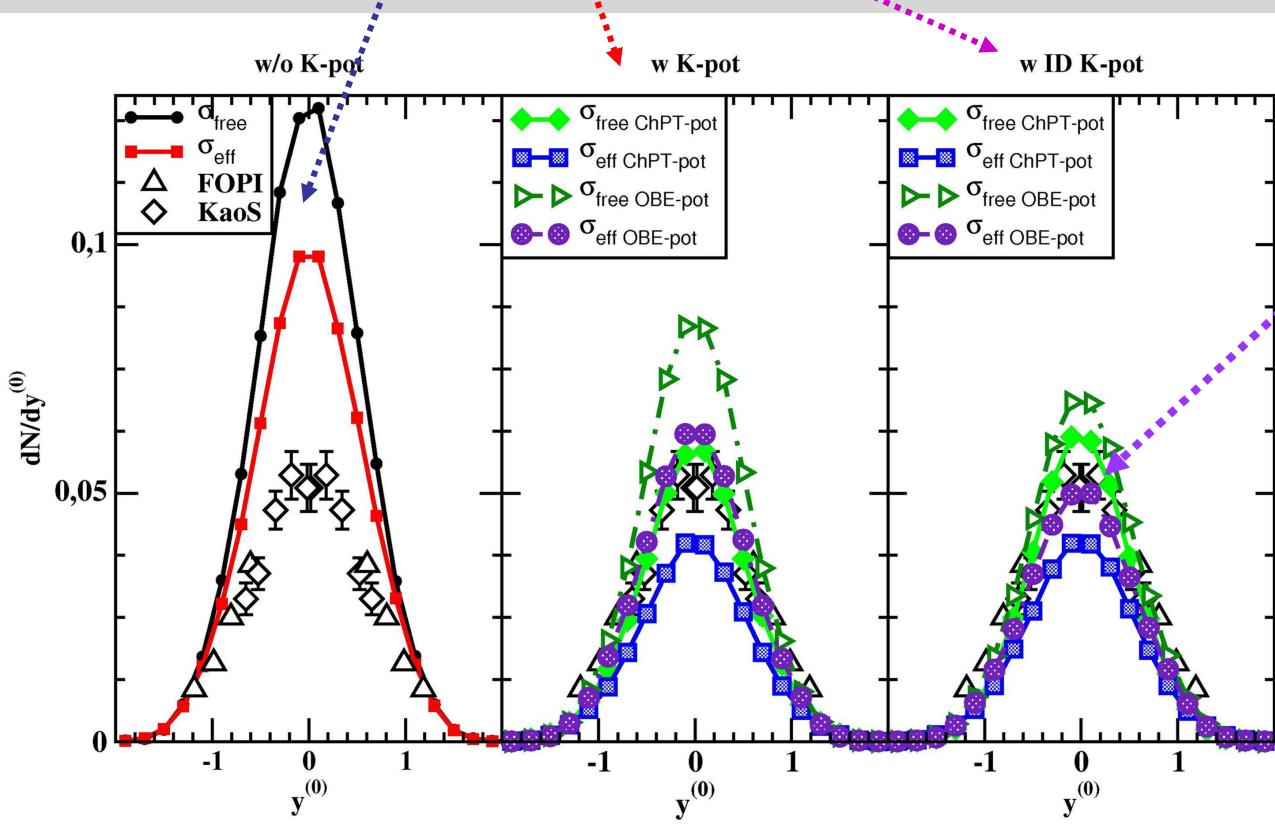
In-medium K energy ($k=0$)

$$E_K(k) = k_0 = \sqrt{k^2 + m_K^{*2}} + V_0$$

Splitting for $K^{0,+}$
for NLρ and NLρδ

Kaon production (absolute yields)

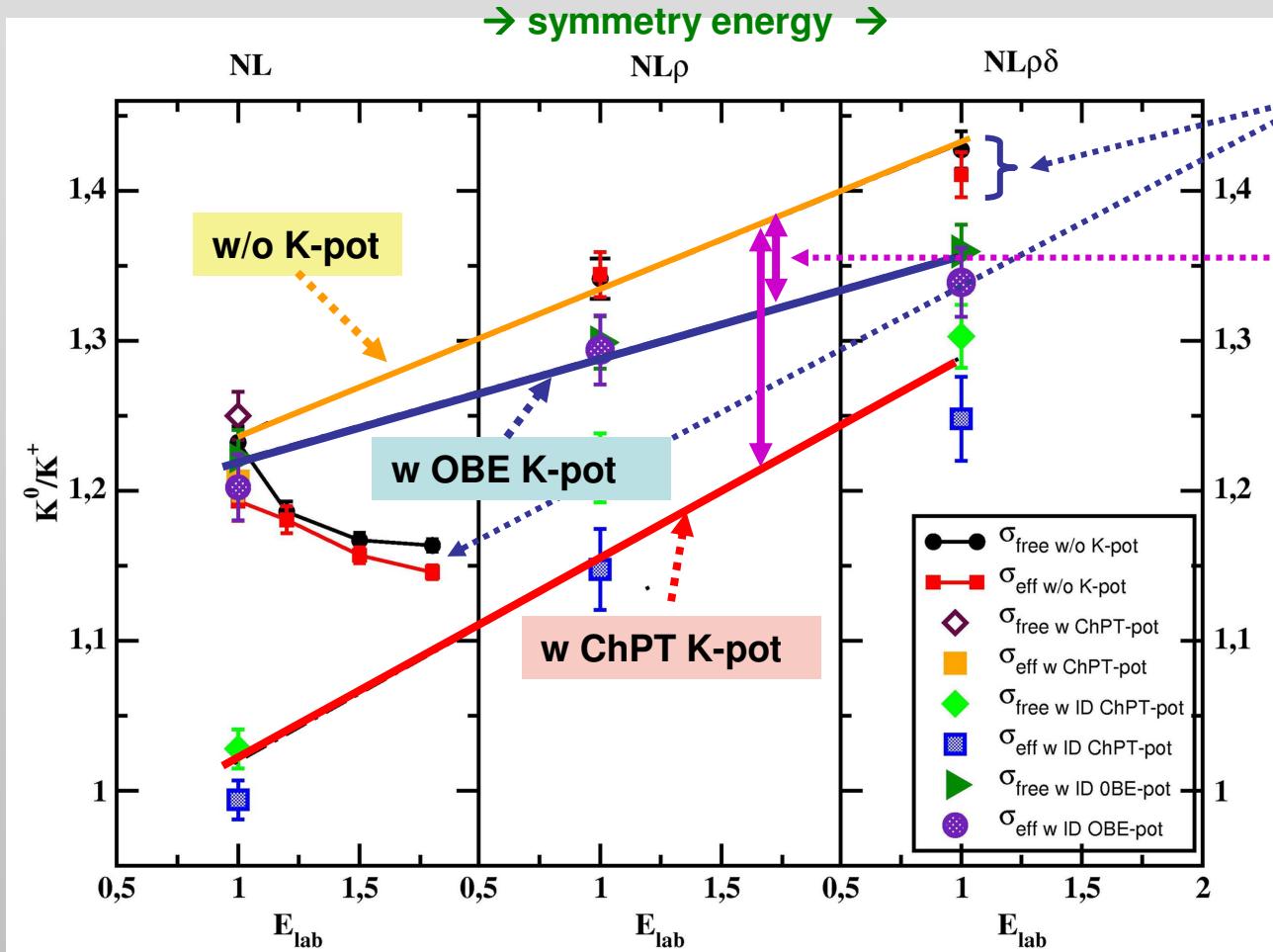
- Sensitivity to:
1. in-medium cross section: σ_{free} vs. σ_{eff}
 2. K-potential: ChPT vs. OBE
 3. Isospin-dep- (ID) K-potential



- Influence of all ingredients substantial
 - rather good description with OBE potential and σ_{eff}
- consider ratios to minimize influence of ingredients

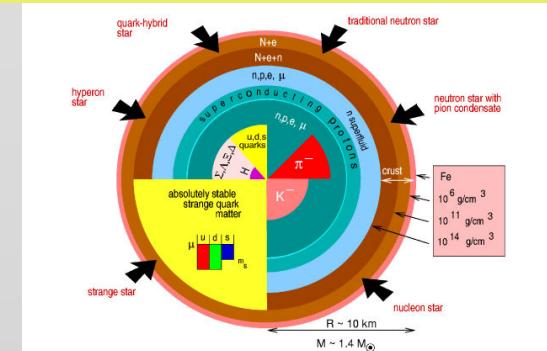
Ni+Ni, $E=1.93$ AGeV, $b < 4$ fm, rapidity distrib.

Kaon ratios: test of „robustness“ against variation of K-potential and σ_{inmed}

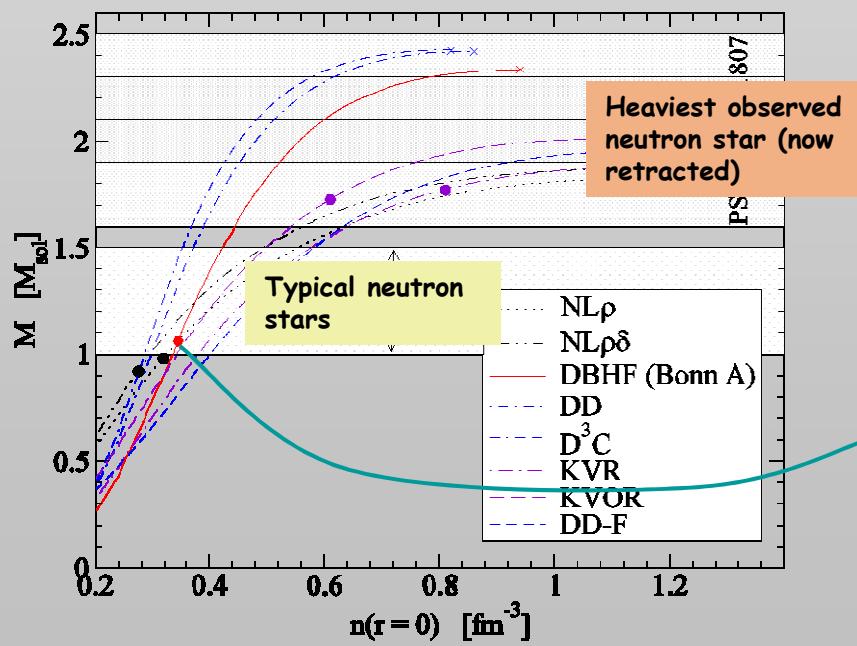


A \bar{u} +Au, b=0 fm

Neutron star masses and cooling and the symmetry energy



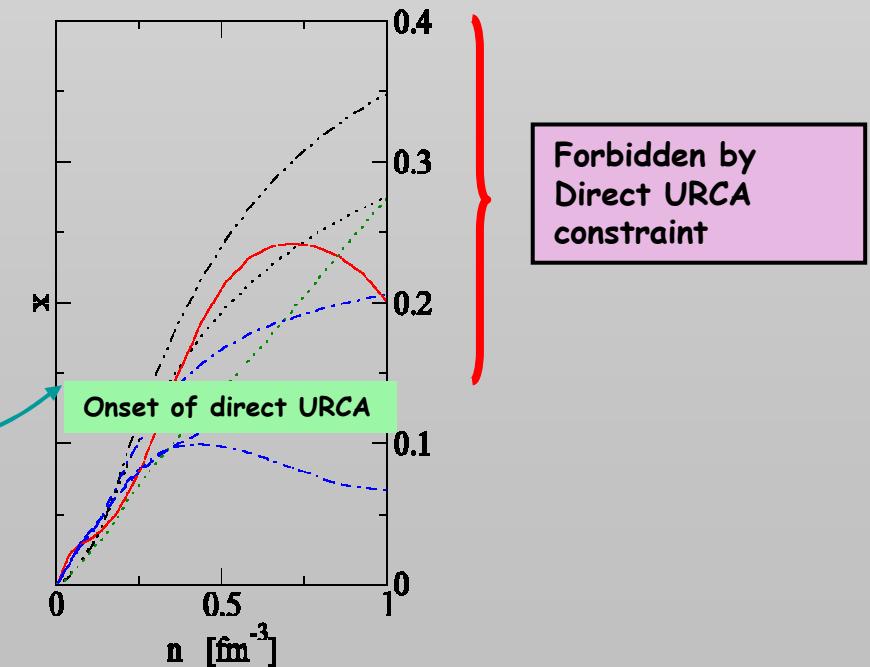
Tolman-Oppenheimer-Volkov equation to determine mass of neutron star



Klähn, Blaschke, Typel, Faessler, Fuchs, Gaitanos, Gregorian, Trümper, Weber, Wolter, Phys. Rev. C74 (2006) 035802

Proton fraction and direct URCA

- β -equilibrium and charge neutrality: $y = \frac{N}{Z} = y(\epsilon_{\text{sym}})$
- direct URCA process: $p \rightarrow n + e^+ + \nu_e$
threshold: $y \approx 11\%$, fast neutrino cooling



Summary and Conclusions:

- While the EOS of symmetric NM is fairly well determined, the density (and momentum) dependence of the symmetry energy is still rather uncertain (but important for exotic nuclei, neutron stars and supernovae)
- Can be investigated in HIC both at low densities (Fermi energy regime, isospin transport) and high densities (relativistic collisions, flow, particle production)
- Nucleon and light cluster pre-equilibrium emission and flow are directly sensitive to the symmetry energy.
- Pion and Subthreshold Kaon production are a promising signal, in particular when considering ratios π^-/π^+ and K^0/K^+
The kaon signal is robust with respect to the in-medium cross sections, but influenced by the model for the Kaon potential
- Effects scale with the asymmetry - thus reactions with RIB are very important